

Valence Bonds in Random Quantum Magnets

theory and application to YbMgGaO_4

Yukawa Institute, Kyoto, November 2017

Itamar Kimchi



I.K., Adam Nahum, T. Senthil, arXiv:1710.06860

Valence Bonds in Random Quantum Magnets

theory and application to YbMgGaO_4

Collaborators



Adam Nahum
(MIT -> Oxford)



T. Senthil
(MIT)

Spin-1/2 magnetic insulators are a playground for challenges in correlated quantum matter

Frustration: destabilizes classical magnetic order.

$T=0$ quantum paramagnets: valence bond liquids or solids

Spin-1/2 magnetic insulators are a playground for challenges in correlated quantum matter

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$T=0$ quantum paramagnets: valence bond liquids or solids

e.g. Kitaev honeycomb model (next talk) – already a challenge!

Aside: $K > 0$ chiral spin liquid has 10x stability to magnetic field
larger fields give intermediate gapless phase

Zheng Zhu, I.K., D.N. Sheng, Liang Fu, arxiv:1710.07595

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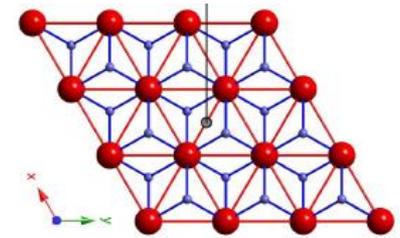
Quenched disorder: site impurities or bond-randomness

Spin glasses, irradiated high- T_c superconductors

Doped Mott insulators

This talk: interplay of quantum frustration and bond disorder

Experimental mystery: YbMgGaO_4



$S=1/2$ on triangular lattice – but no magnetic order

Strong spin-orbit-coupling ($\text{Yb}^{3+}: 4f^{13}$)

Exchanges: $J_1, J_2(?)$, XY, Kitaev... $\theta_{CW} = 4$ K

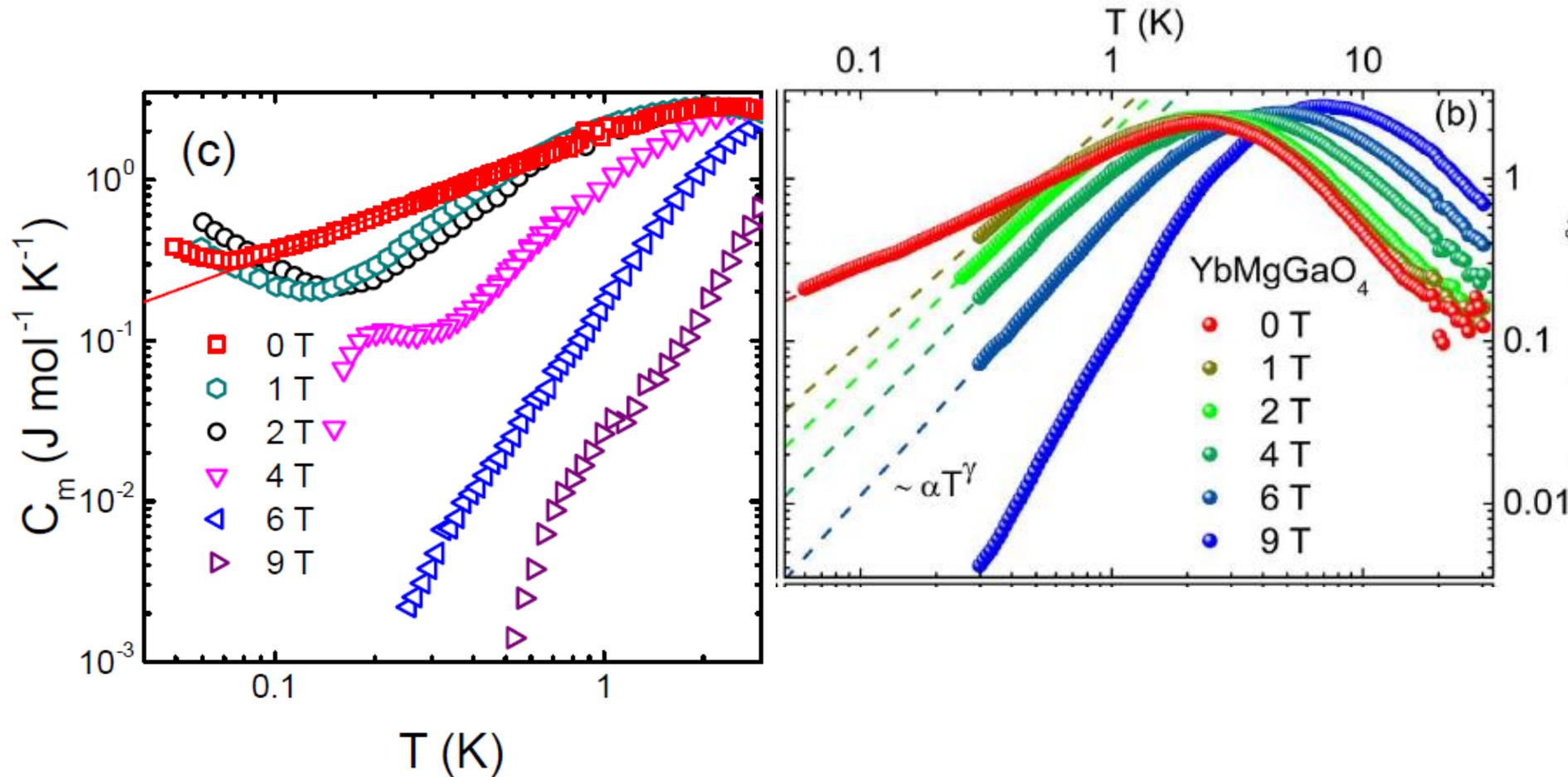
Unusual magnetic phenomenology:

No spin order or glass (50 mK neutrons & μSR)

Anomalous heat capacity $C(T) \sim T^{0.7}$

but no corresponding thermal conductivity

Low-temperature $C(T) \sim T^{0.7}$



Xu et al, PRL 117, 267202 (2016)

Li et al, Sci. Rep. 5, 16419 (2015)

How to understand this unusual phenomenology?

$$C(T) \sim T^{0.7}:$$

interpreted as “spinon Fermi surface” (Gang Chen *et al.*)

missing signatures of itinerant spinons

Ingredients for an alternative hypothesis:

Frustration: Geometrical & spin-orbit-coupling
capture via non-magnetic “valence bonds” basis

+

Disorder: Magnetic exchanges with random energies
due to Mg/Ga mixing in the non-magnetic layers

Say **frustration** prevents magnetic order:
describe clean magnet in **valence bond** basis,
then add **bond-randomness disorder**

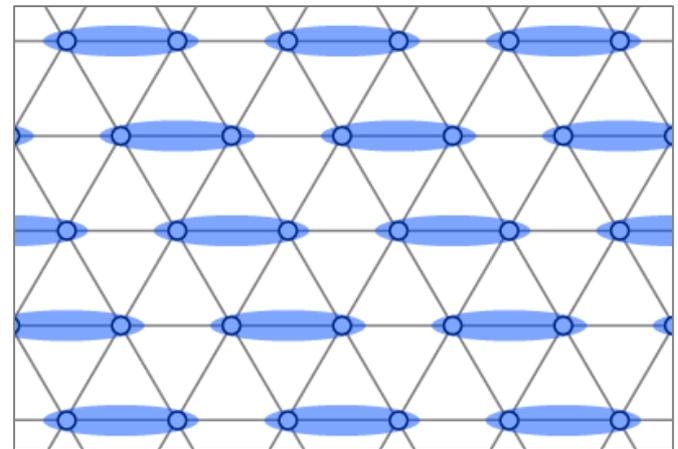
$$H = \sum J \vec{S} \cdot \vec{S} \quad J \rightarrow J + \Delta J_{ij}$$

Even in limit of weak disorder $\Delta J \ll J$

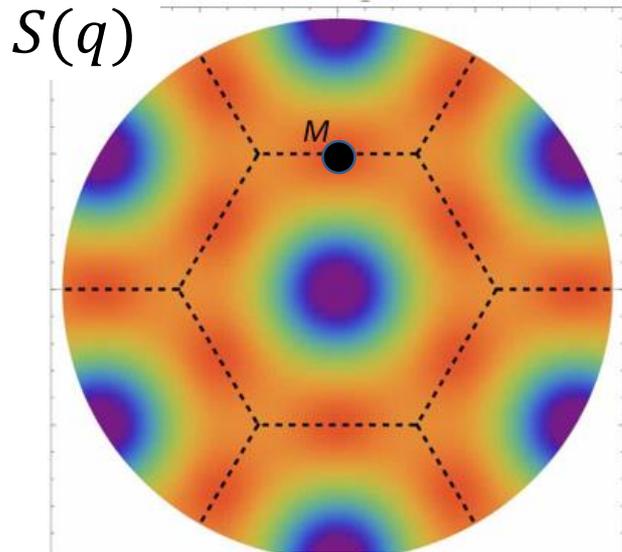
linear coupling to *valence-bond-solid* order (Imry-Ma)

→ splits VBS into domains
of short-ranged singlets

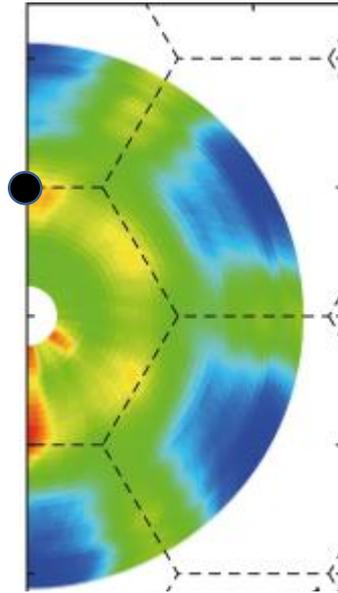
VBS: $T=0$ paramagnet
w/ broken lattice symms



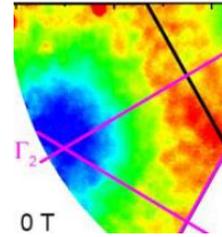
Random short-ranged singlets are a good starting point



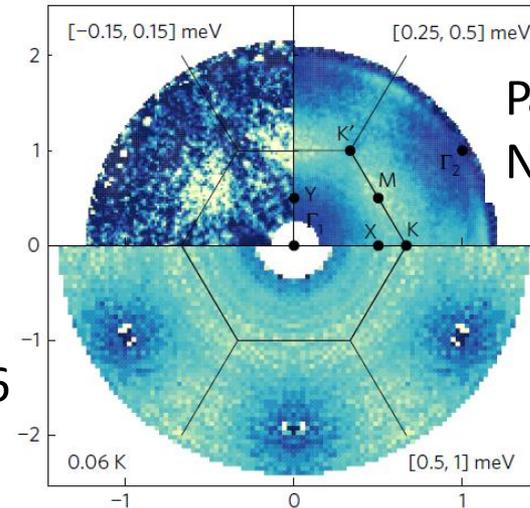
Randomly oriented short-ranged singlets



Shen *et al*, Nature 2016



Li *et al*,
Nat Comm 2017



Paddison *et al*,
Nat Phys 2017

But: frozen valence bonds \rightarrow gap.

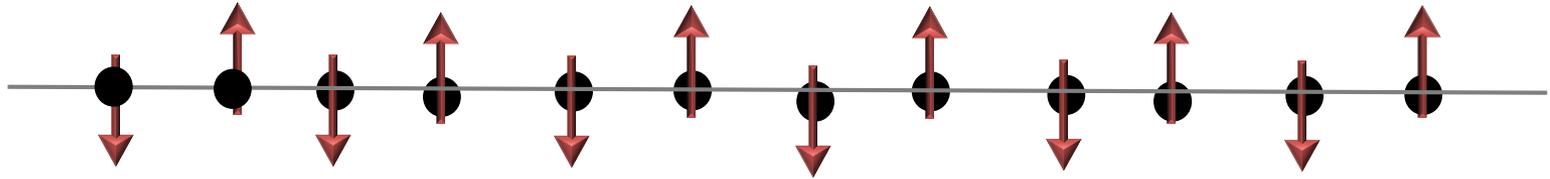
What gives gapless $C(T) \sim T^{0.7}$?

Competition between disorder and valence bonds
necessarily leads to low-energy spin excitations:
strong-randomness network of spin-1/2 emerges

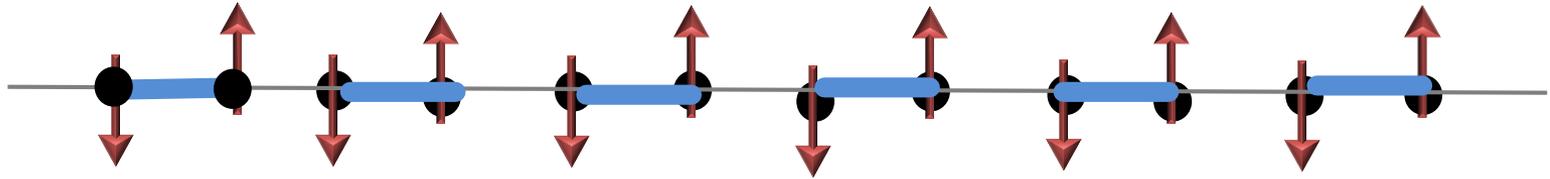
Rough sketch of the argument:

- (1) Weakly disordered VBS: vortices carrying spin-1/2
- (2) Stronger disorder: defect instability of pinned singlets
- (3) Disordered Lieb-Schultz-Mattis conjectures
- (4) Application to YbMgGaO_4 : $S(q, \omega), \kappa(T), C(T)$ & B -field

Warm up: 1D spin-1/2 chain
spontaneous dimerization + disorder

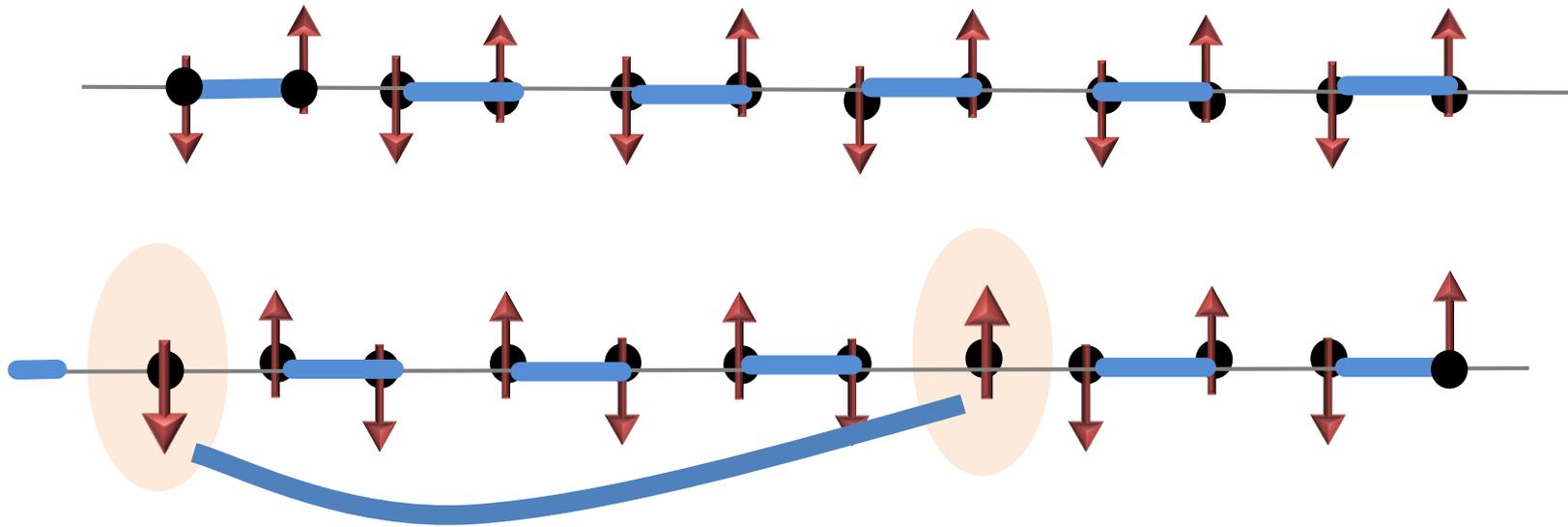


Warm up: 1D spin-1/2 chain
spontaneous dimerization + disorder



Clean system in dimerized (“VBS”) phase
Adding disorder breaks up domains

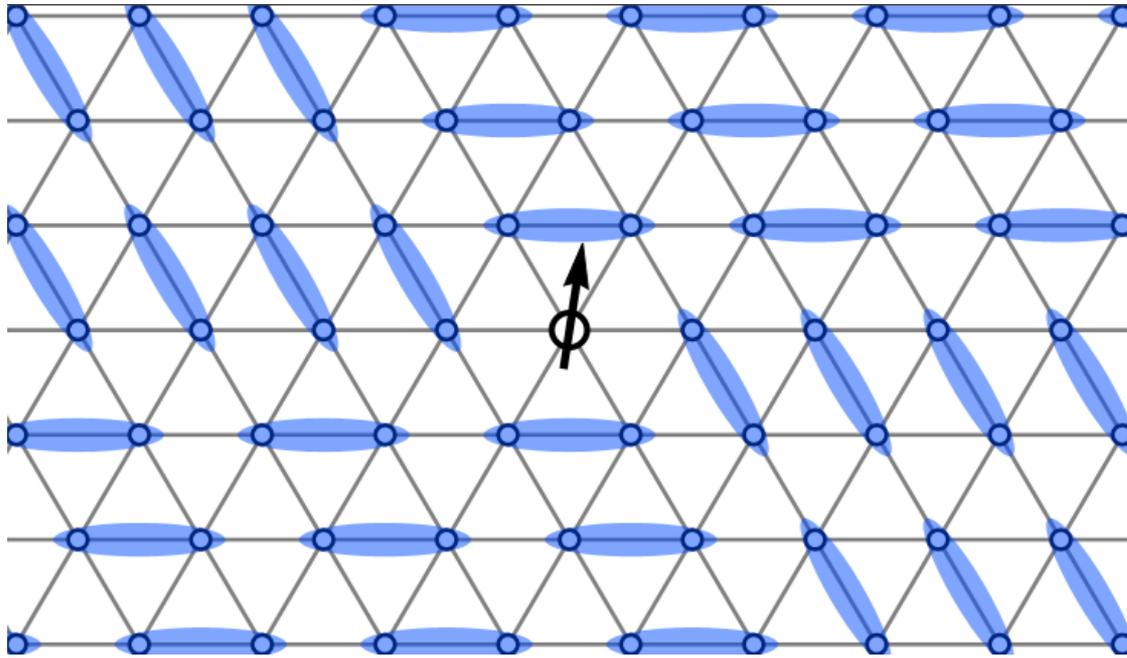
Warm up: 1D spin-1/2 chain
spontaneous dimerization + disorder



Domain wall carries single $S=1/2$

→ RG flow to 1D random-singlet phase (Fisher '94)

2D vortices with spin-1/2 modes



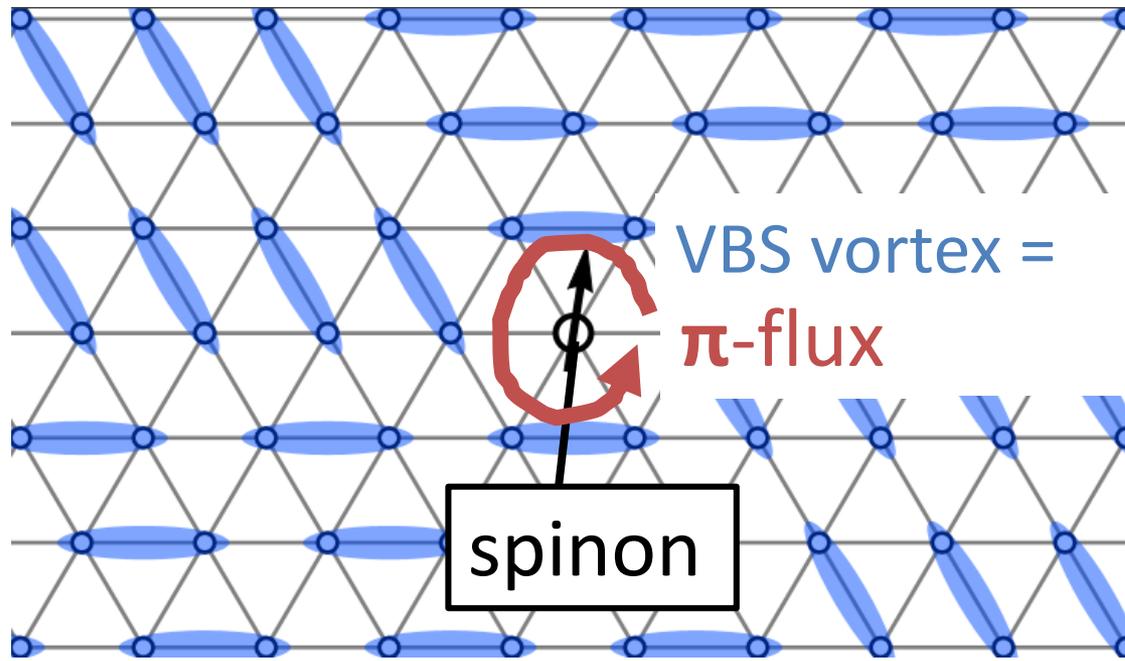
Interpret via Z_2 spin liquid: condense vison vector \mathbf{v}

VBS order parameter = \mathbf{v}^2

\mathbf{v}^2 headless vector $\rightarrow Z_2$ vortices

Z_2 vortex = vison Z_2 gauge field π -flux

Vortices carry spin-1/2 modes



Interpret via Z_2 spin liquid: condense *vison* vector \mathbf{v}

VBS order parameter = \mathbf{v}^2

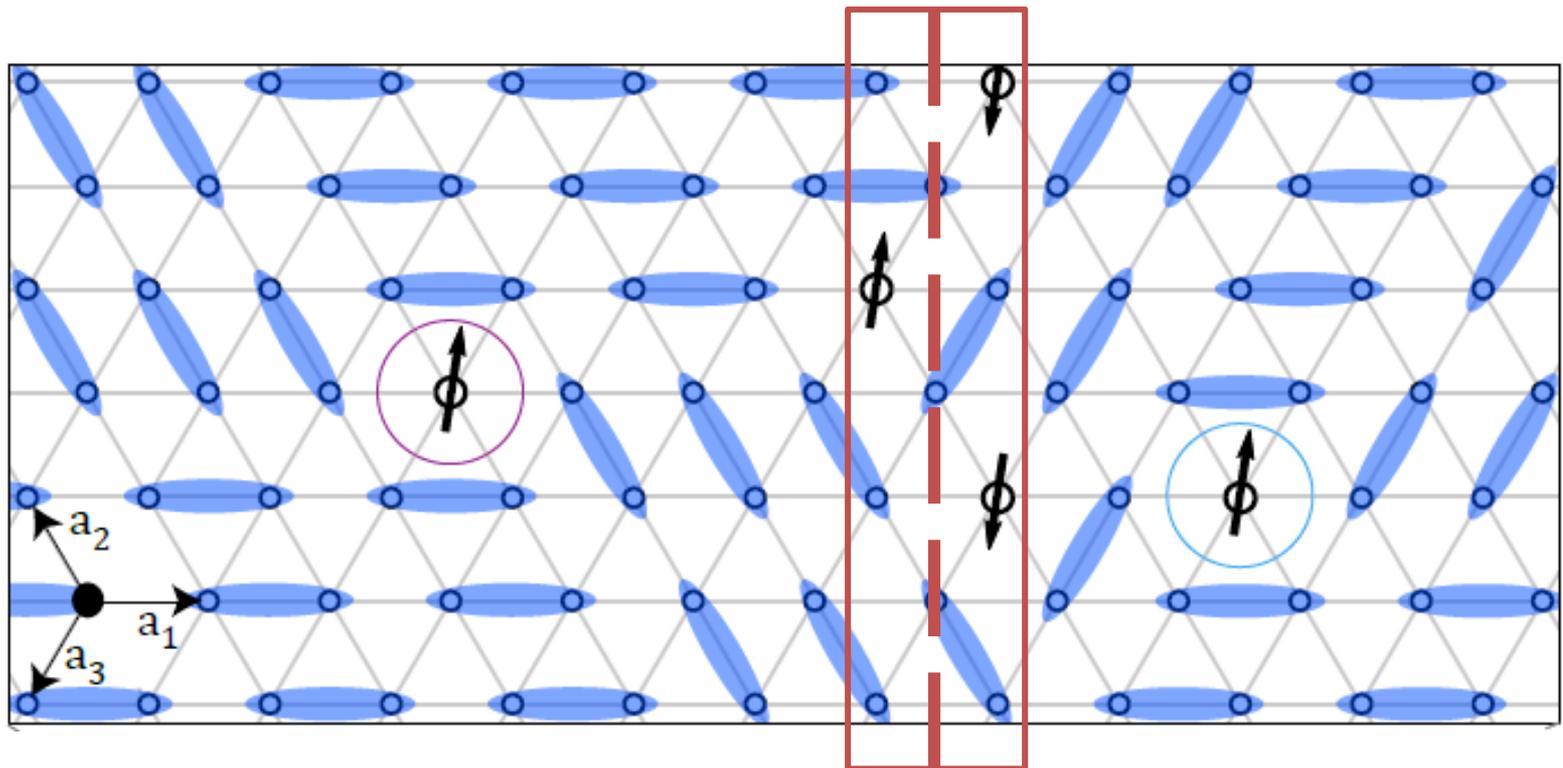
\mathbf{v}^2 headless vector $\rightarrow Z_2$ vortices

Z_2 vortex = vison π -flux = spin-1/2 *spinon*

Details for triangular lattice columnar-VBS:

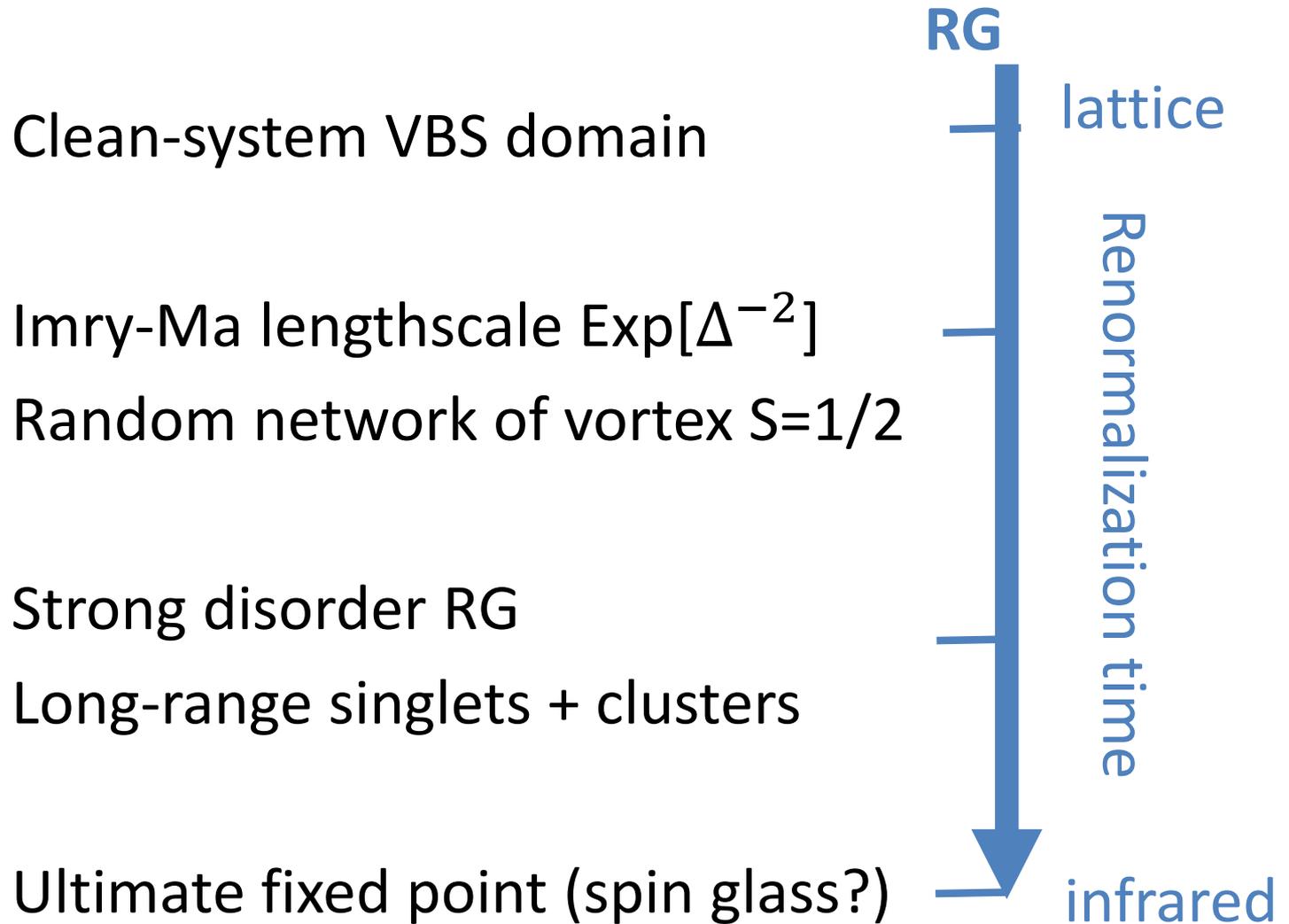
domains cluster into “superdomains”

superdomain-walls carry $S=1/2$ chains



a superdomain maps to square lattice VBS: Z_4 vortices

RG flow arises from weak disorder



Are the $S=1/2$ vortices always natural?
Can stronger disorder pin singlets into a
short-ranged “valence bond glass”?

- (1) Weakly disordered VBS: vortices carrying spin-1/2
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Enforced nucleation of $S=1/2$ defects

Spin-1/2 defects cost energy. Are they necessarily nucleated by disorder?

Limit #1, Imry-Ma weak disorder: topological defects appear between large domains

Limit #2, regime of intermediate disorder:

VBS pattern selection scale \ll Disorder \ll Clean-system spin gap

Map to random-energy dimer model
but now allow monomers/defects

Enforced nucleation of $S=1/2$ defects

Classical dimer model w/ random energies on bonds
with allowed monomers/defects

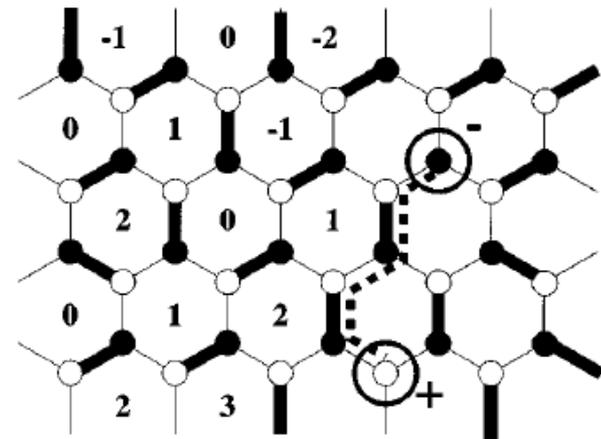
Bipartite lattices: any disorder will nucleate defects

Zeng-Leath-Fisher (PRL 1999)

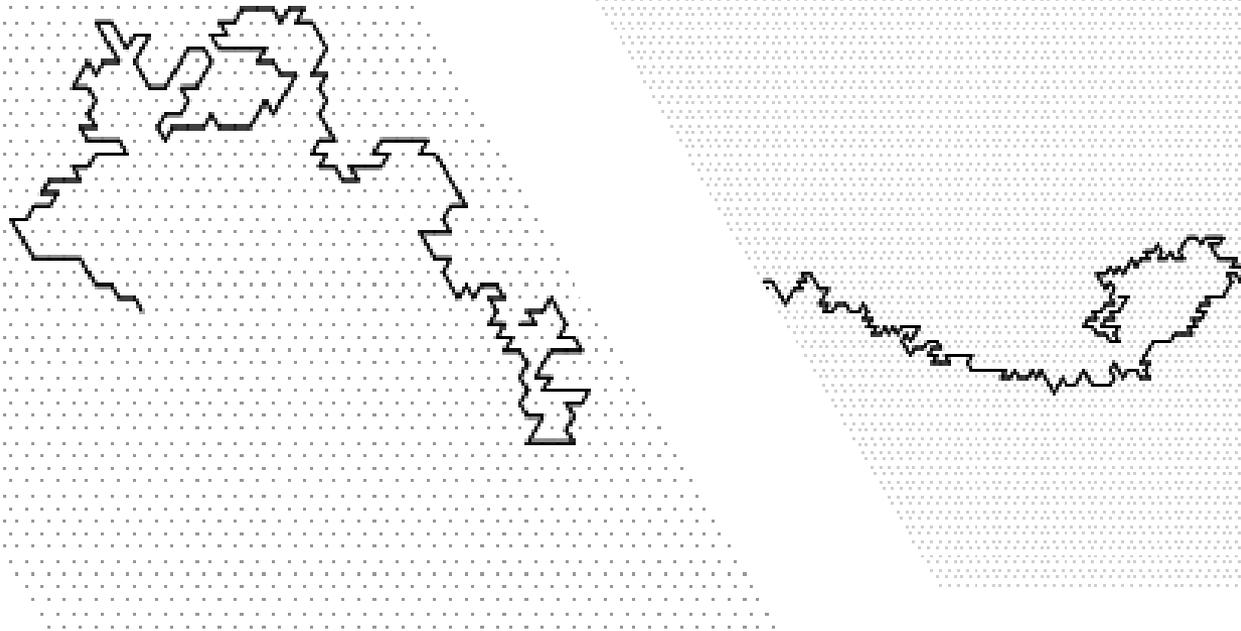
Middleton (PRB 2000)

Non-bipartite case unknown;

Study on triangular lattice

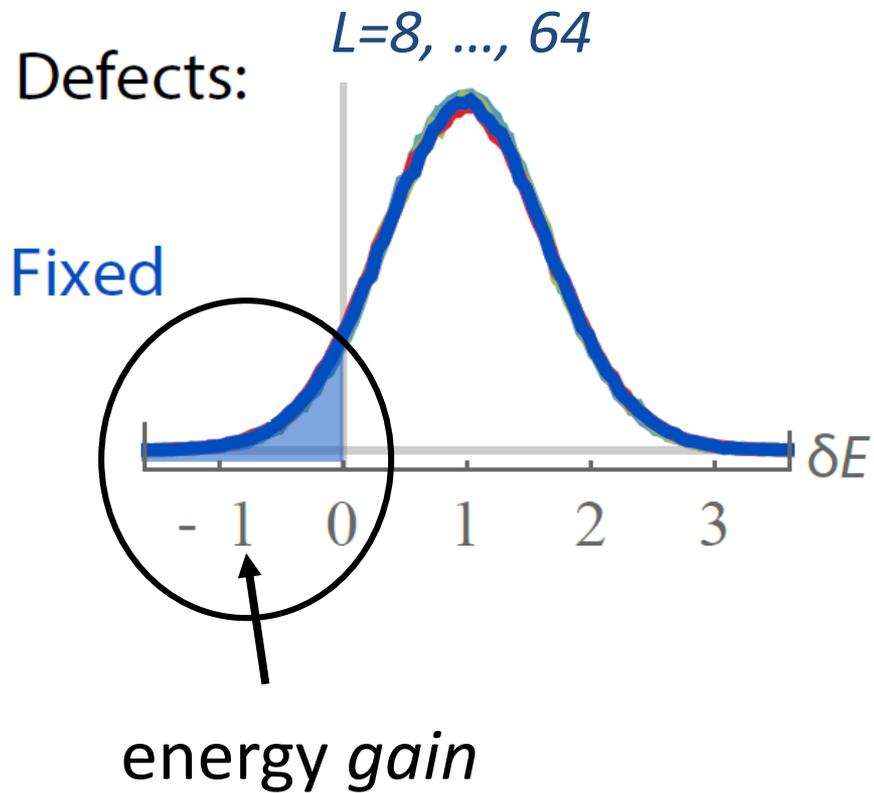


(Fractal defect strings confirm mapping to Ising spin glass)

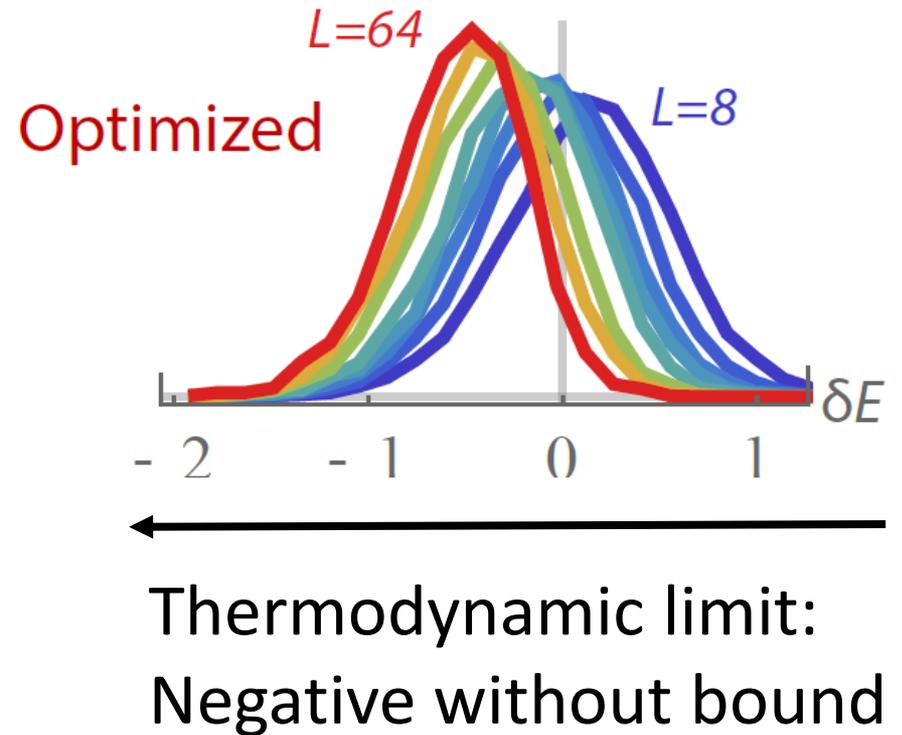
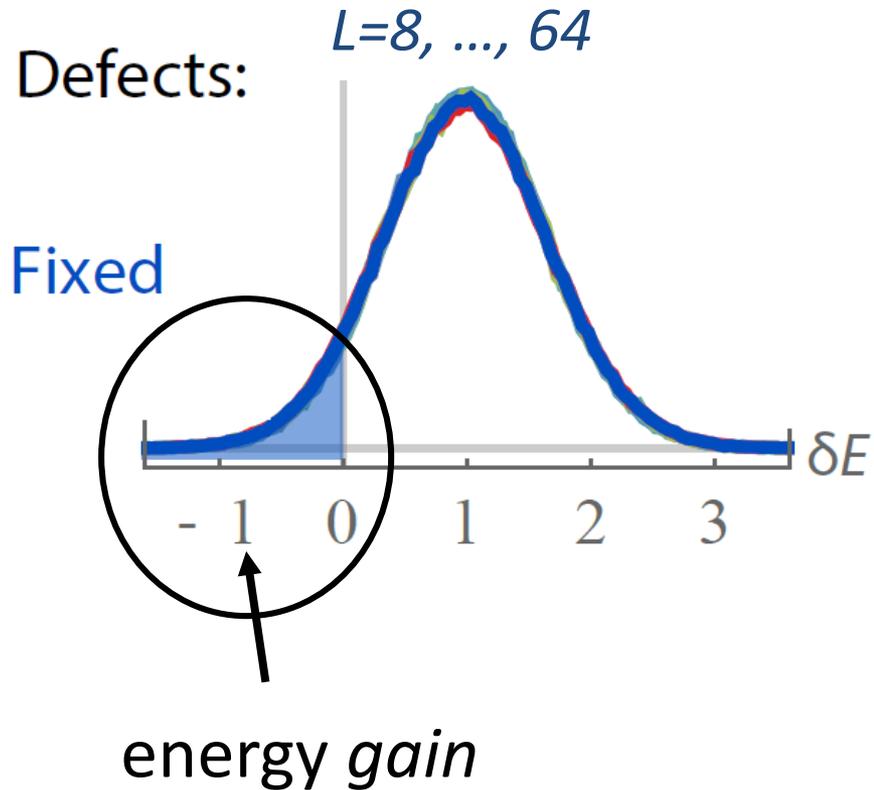


Fractal dimension $d_f = 1.28$

Energy distribution for two fixed defects

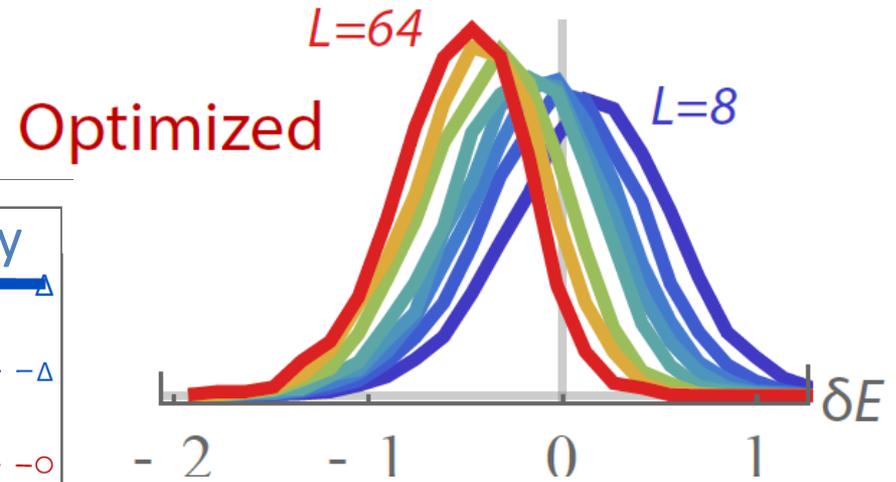
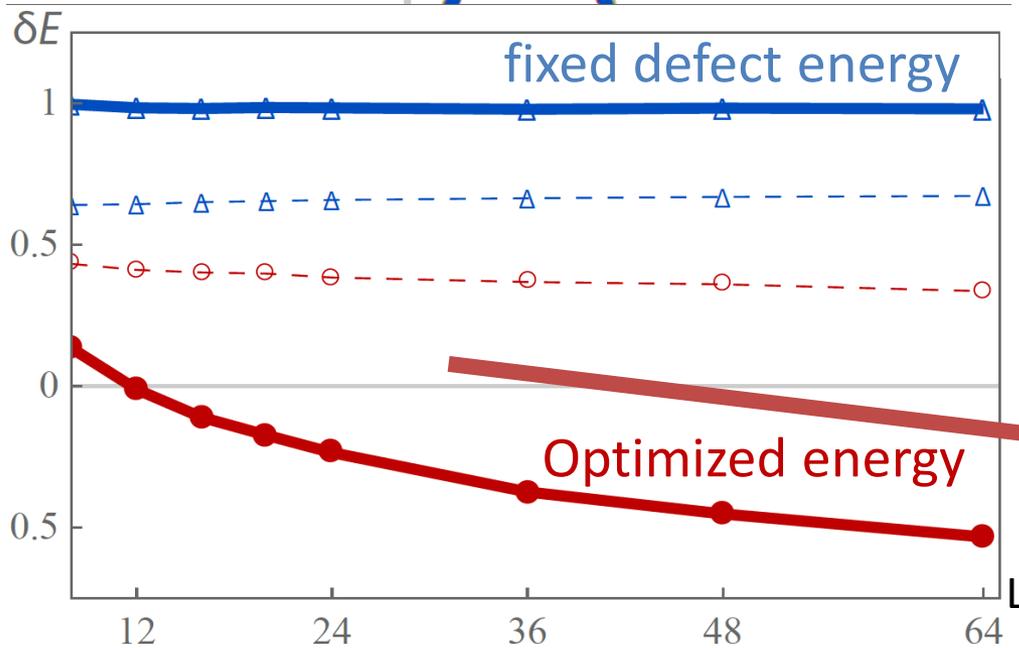


Energy distribution for partially-optimized defects



Energy distribution for partially-optimized defects

Defects: $L=8, \dots, 64$



Divergent energy gain from disorder: defects will always nucleate

Adding weak/stronger disorder to destroy
VBS-symmetry-breaking / spin-liquid
necessarily nucleated gapless spin excitations.
Is this a general principle?

- (1) Weakly disordered VBS: vortices carrying spin-1/2
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Is this a general principle?

Naively expect disordered state to be featureless, spin-gapped

But vortices/monomers with spin-1/2 appeared!

Recall Lieb-Schultz-Mattis-Hastings-Oshikawa theorem (*LSM*):

$S=1/2$ per unit cell \Rightarrow featureless states must be gapless

Here: **LSM** with **disorder**? gapless **spins**?

Disordered-LSM conjectures

Given spin rotations and *statistical* translations
with $S=1/2$ per unit cell

Conjecture restrictions for featureless ground states
(if no symmetry-breaking/topological order)

2D: must have gapless *spin* excitations
(e.g. long-range singlets)

1D: spin correlations at least algebraically-long-ranged
General argument in 1D (Adam Nahum)

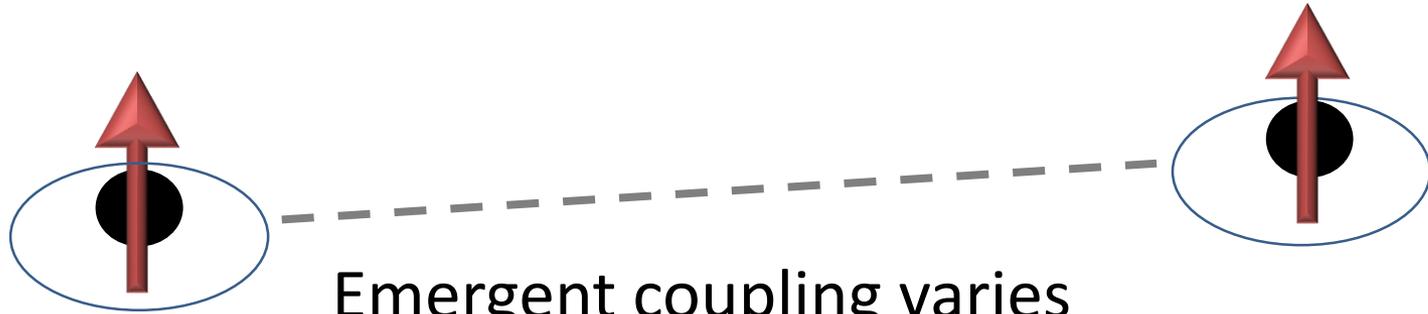
3D: forthcoming

[alternative formulations via quantum information]

What are implications of this enforced RG flow?
Can this physics be observed?

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Random network of spin-1/2 emerges: Shows random-singlets and some spin freezing



Emergent coupling varies
exponentially with separation

→ Power-law density of states (as in Si:P)

2D random-singlet phase

ultimate fixed point likely has frozen moments

$$\text{YbMgGaO}_4: C(T) \sim T^{0.7}$$

$$\text{YbZnGaO}_4: C(T) \sim T^{0.6}$$

both: anomalous low- T spin freezing

Ma *et al.*

1709.00256

Relevance to YbMgGaO_4 :

Summary, Predictions

Summary, “post-dictions”:

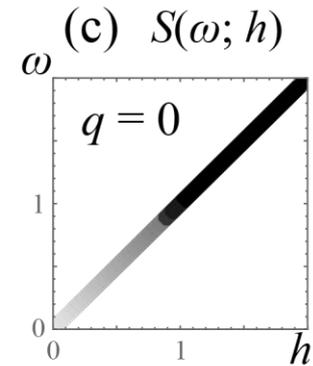
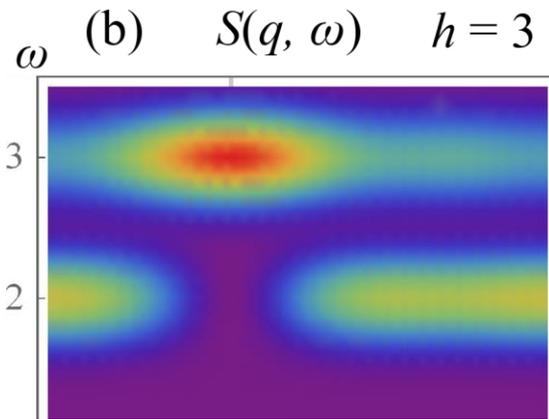
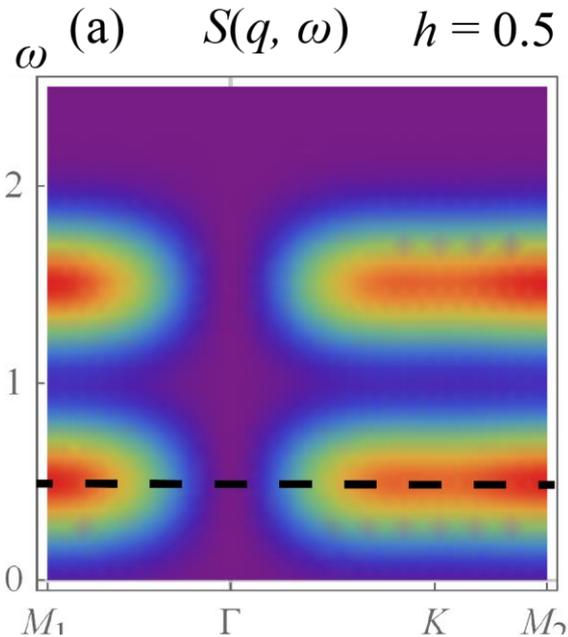
1. No magnetic order
2. Short-ranged singlets at energies of order J
3. Power-law density of states at low energies, $C(T) \sim T^\alpha$

Nontrivial predictions:

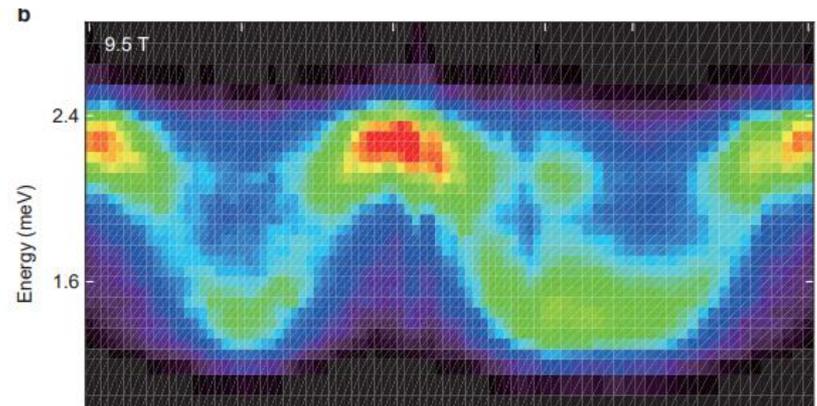
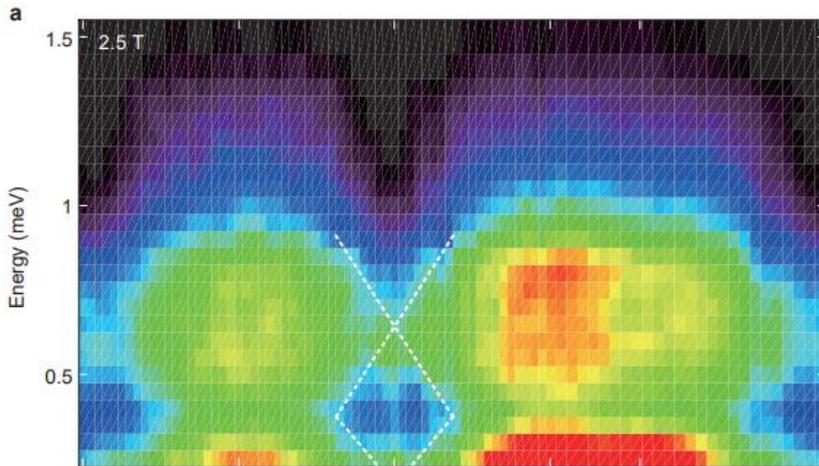
1. Thermal conductivity $\kappa(T) \sim T^{1.9}$ (glassy phonons) ✓
2. Possible short-ranged VBS order [$q=M$]
3. Some glassy freezing at $T=0$ ✓
4. Behavior in a magnetic field: $\kappa(T), C(T)$
and $S(q, \omega)$ ✓



Recently measured $S(q, \omega)$ in magnetic fields consistent with random singlets



random singlets



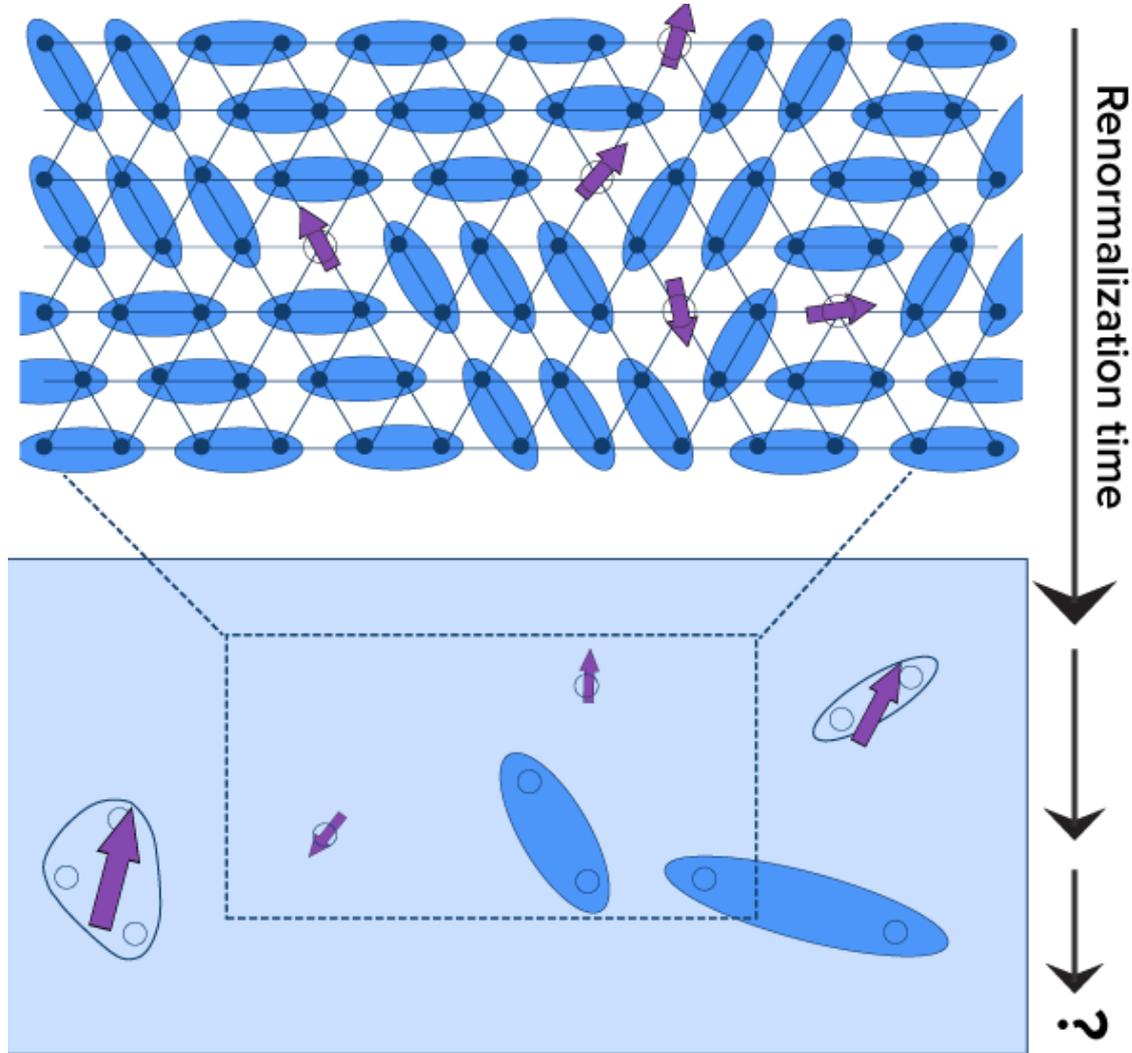
Shen et al, 1708.06655

Conclusions: Frustration + Disorder

RG flow to random network of spin-1/2

Anomalous power-laws

Enforced by disordered-LSM?



Outlook:

Spin liquids + defects

Disordered-LSM proofs

Numerical access

Other materials

For more, see [arXiv:1710.06860](https://arxiv.org/abs/1710.06860)