

# 将来の21cm線放射の観測による 暗黒エネルギーの性質の解明

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# ◇ Dynamical dark energy

Expansion ratio of the Universe (Hubble parameter)

$$H^2 = H_0 \left[ \Omega_m a^{-3} + \Omega_r a^{-4} + \Omega_k a^{-2} + \Omega_X \exp \left[ 3 \int_a^1 \frac{da'}{a'} (1 + w_X(a')) \right] \right]$$

Contribution of the dark energy

Comoving distance :

$$d(z) = \int_0^z \frac{dz'}{H(z')}$$

Equation of state (EOS)

$$w_X(a') \equiv \frac{p_X(a')}{\rho_X(a')}$$

We investigate sensitivities of 21 cm line to the dynamical dark energy model ( $w_X \neq \text{constant}$ ).

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1. Parametrizations of dynamical dark energy
2. 21cm line observations
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# Dynamical dark energy model

一般にdark energyのEOS  $w_X(a')$ は時間依存

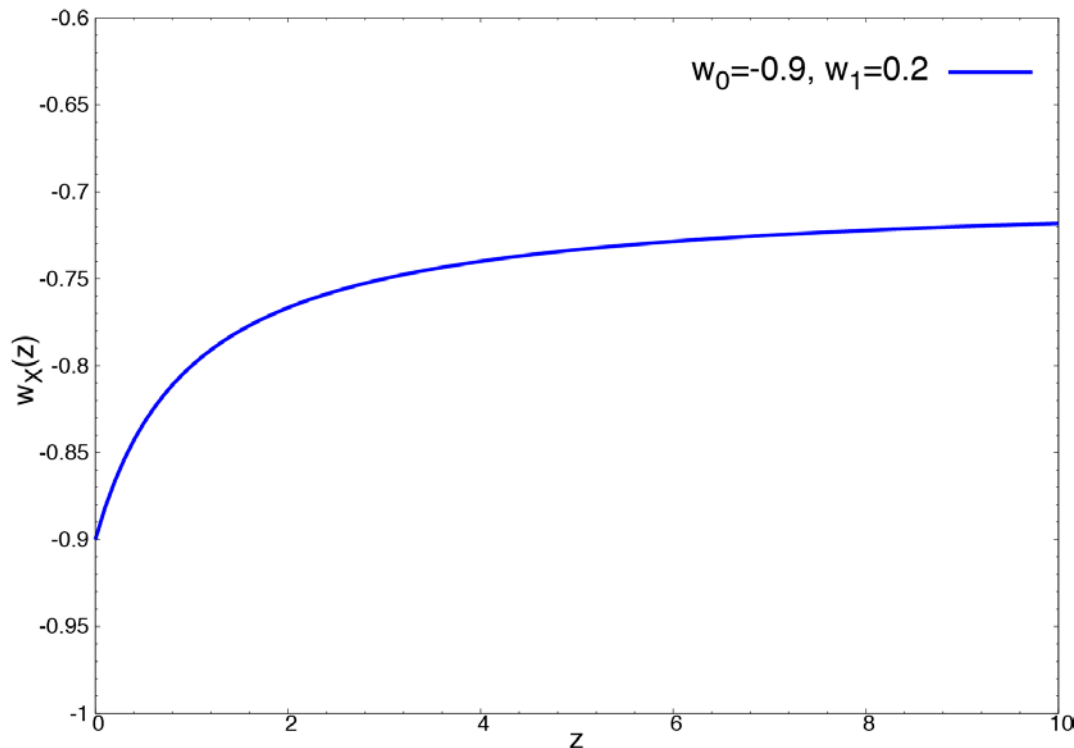
$w_X(a')$ の時間依存性はモデルによって異なる

ここでは特定のmodelを用いず、  
現象論的に数個のパラメータを導入することで  
EOSの時間依存性を表しそれらのパラメータを  
制限するという方法を考慮する

# ◆ Parametrization 1 [Chevallier, Polarski 2001; Linder 2003]

$$w_X(z) = w_0 + (1 - a)w_1 = w_0 + \frac{z}{1 + z}w_1$$

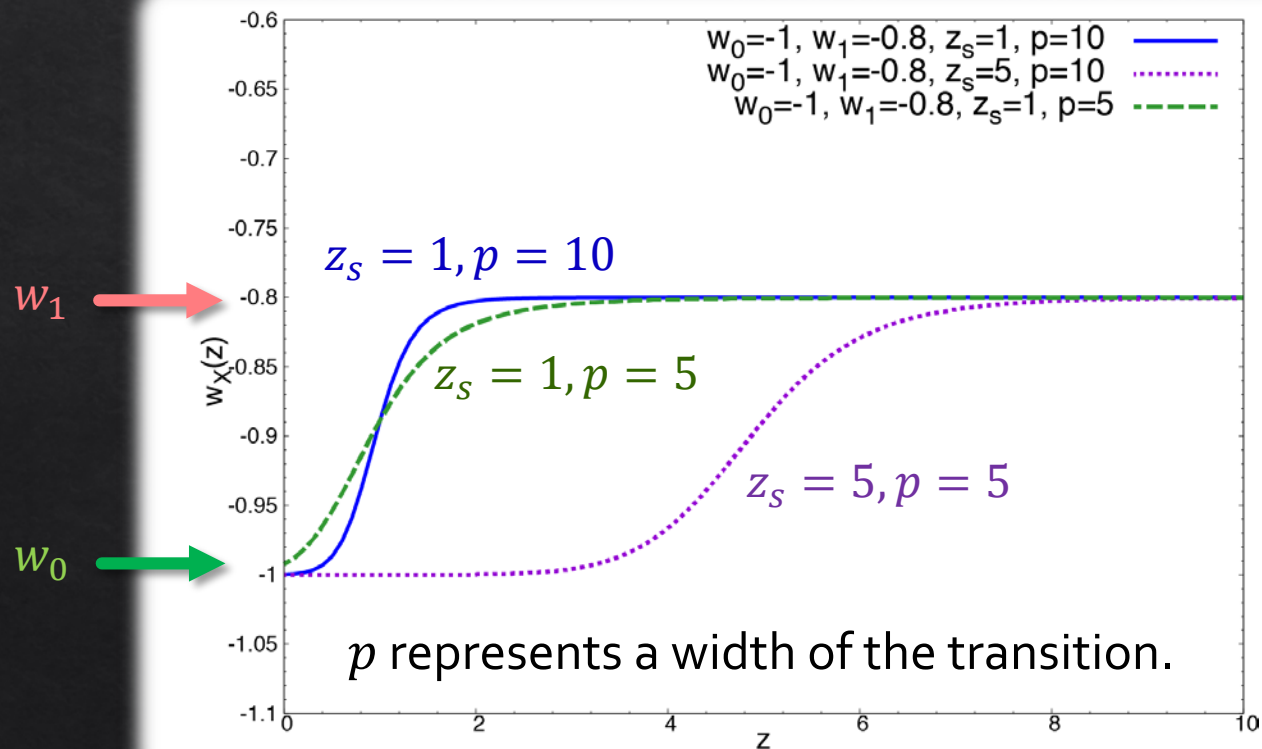
A very simple model :  $w_X$  is proportional to the scale factor  $a$ .



## ◆ Parametrization 2 [Hannestad, Mortsell 2004]

$$w_X(z) = w_0 w_1 \frac{a^p + a_s^p}{w_1 a^p + w_0 a_s^p} = w_0 w_1 \frac{1 + \left(\frac{1+z}{1+z_s}\right)^p}{w_1 + w_0 \left(\frac{1+z}{1+z_s}\right)^p}$$

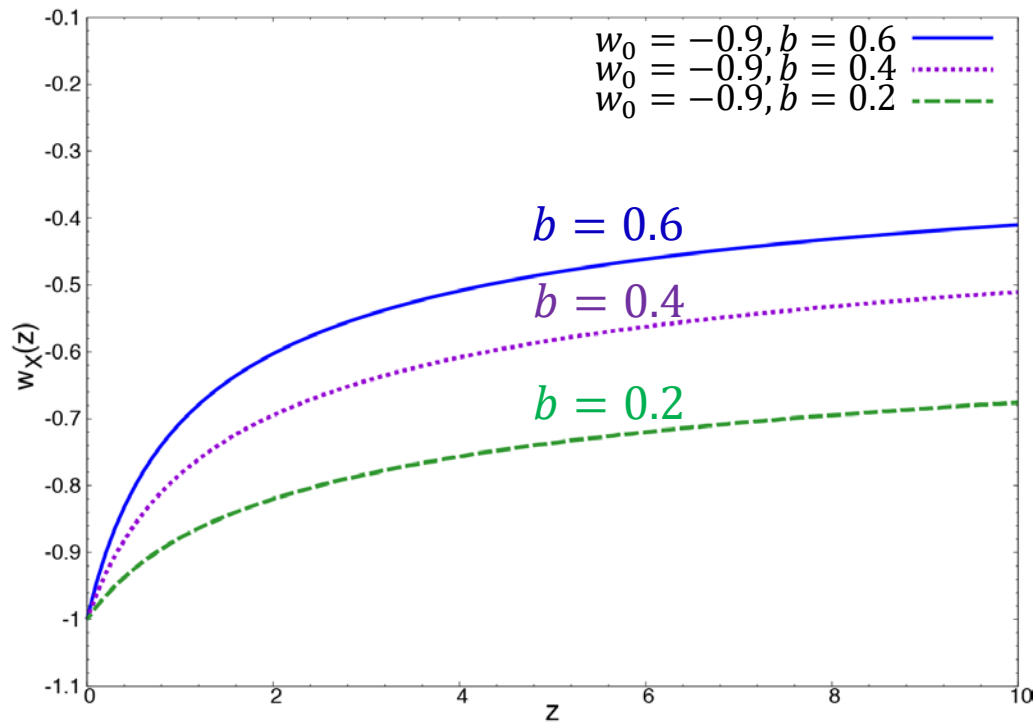
$w_X$  suddenly changes at a transition redshift  $z_s = \frac{1}{a_s} - 1$ .



## ◆ Parametrization 3 [Wetterich 2004]

$$w_X(z) = \frac{w_0}{[1 + b \log(1 + z)]^2}$$

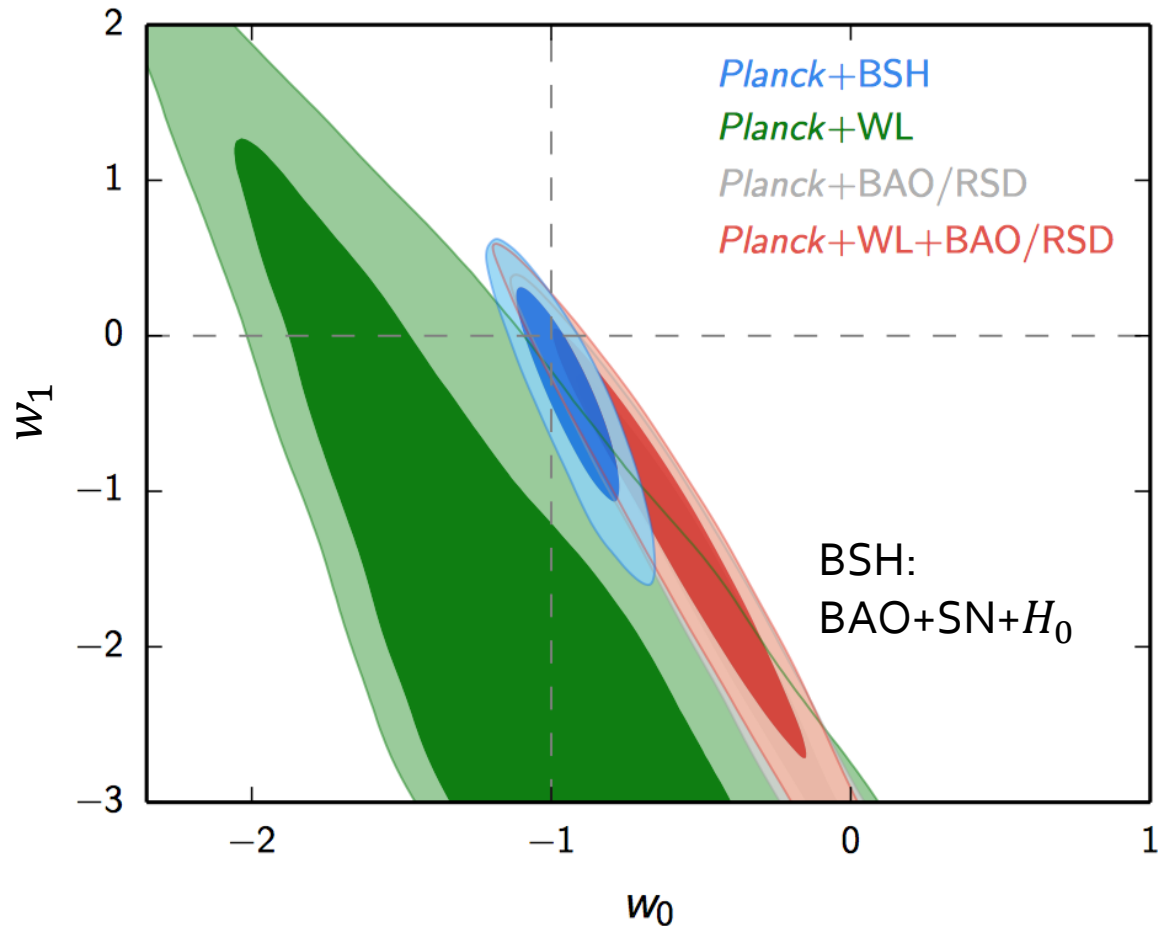
**Early dark energy** : the contribution of dark energy has a relatively large impact in the early epoch.





# ◆ Current constraint of dynamical dark energy

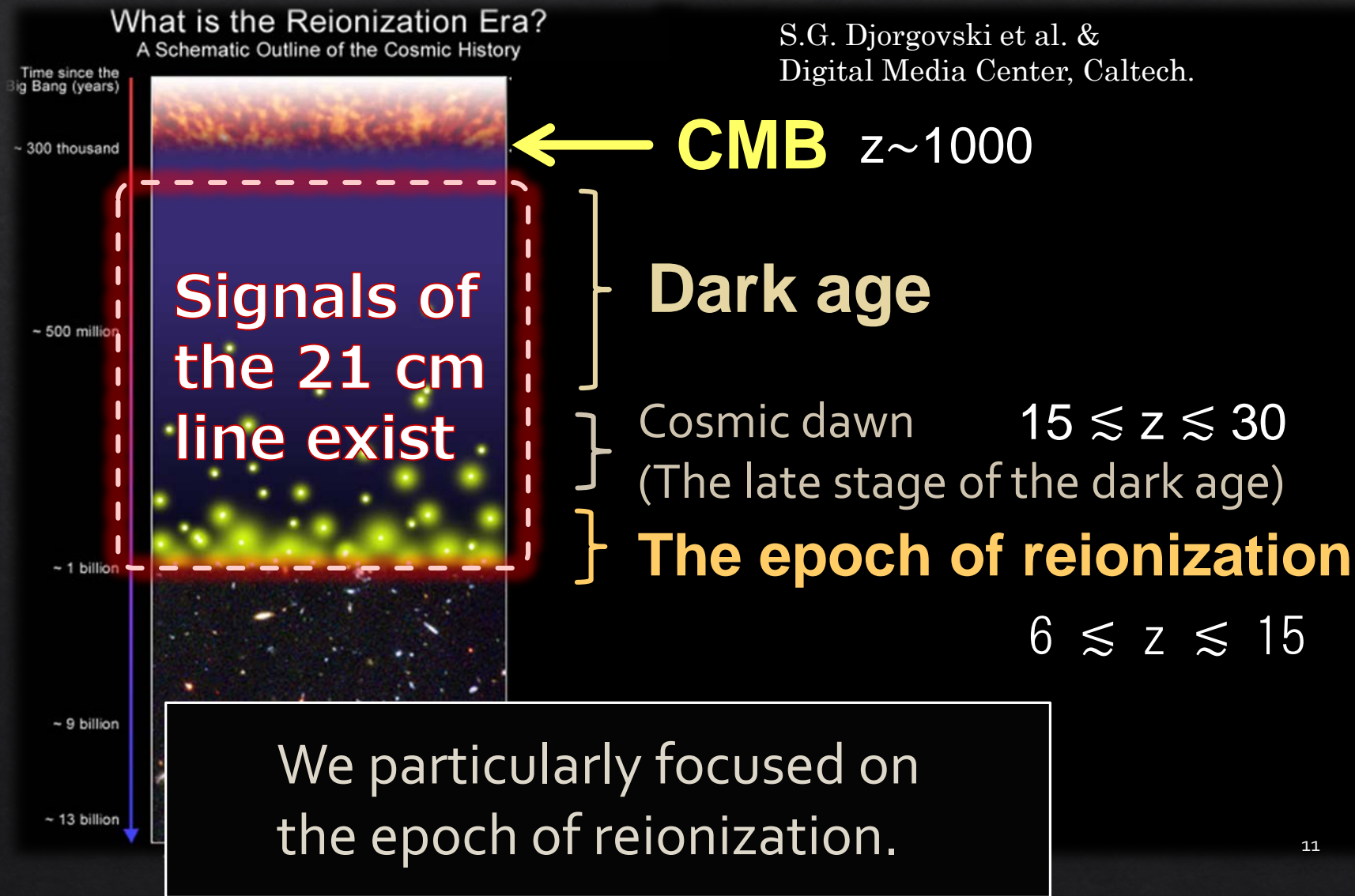
Planck 2015 (Parametrization 1)



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# ◇ Cosmological 21 cm radiation



# ◇ Brightness temp $T_b$ of 21 cm line

$$\Delta T_b \equiv T_b - T_{CMB}$$

$$\Delta T_b \left( \frac{\nu_{21}}{1+z}, \mathbf{r}, z \right) \approx 27 x_{\text{HI}} (1 + \delta_b) \left( \frac{\Omega_b h^2}{0.023} \right) \left( \frac{0.15}{\Omega_m h^2} \frac{1+z}{10} \right)^{\frac{1}{2}}$$

Fluctuation of baryon  $\times \left[ 1 - \frac{T_{CMB}}{T_S} \right] \left[ \frac{H(z)/(1+z)}{dv_{||}/dr_{||}} \right] \text{mk}$

$T_b$  : Brightness temp  
of 21 cm line

$T_S$  : Spin temp

$T_{CMB}$  : CMB temp

$x_{\text{HI}}$ : Neutral fraction

$T_S > T_{CMB}$  emission ( $6 \lesssim z \lesssim 15$ )

$T_S < T_{CMB}$  absorption ( $15 \lesssim z$ )

# ◇ スピン温度 $T_S$

$$T_S \text{ の定義 : } \frac{n_1}{n_0} \equiv \frac{g_1}{g_0} \exp\left(-\frac{h\nu_{21}}{k_B T_S}\right)$$

$n_1, n_0$ : スピン1, 0 状態の数密度

## $T_S$ に影響する効果

- (1) 水素原子, 電子, 陽子との衝突
- (2) CMB光子の吸収, 放射
- (3) Ly $\alpha$ 光子の吸収, 放射
- (4) 中性水素比  $x_{HI}$  の変化

$$x_{HI} \equiv n_{HI}/n_H$$

# ◇ 21cm line fluctuation $\delta_{21}$

$$\delta_{21} \equiv \frac{\Delta T_b - \Delta \bar{T}_b}{\Delta \bar{T}_b}$$

$$T_S \gg T_\gamma : z \sim 10$$

Baryon Neutral fraction Peculiar velocity  $\equiv \frac{1 + z}{H(z)} \frac{dv_{p||}}{dr}$



$$\delta_{21} \approx \delta_b + \delta_{x_{\text{HI}}} - \delta_{\partial v}$$

Fourier component (in matter dominated era)

$$\tilde{\delta}_{21} \approx \tilde{\delta}_b + \tilde{\delta}_{\text{HI}} + \mu^2 \tilde{\delta}_{b, \parallel} \quad \mu \equiv \frac{k_{\parallel}}{|k|}$$

# ◇ Power spectrum of 21cm $P_{21}(k, \mu)$

$$\langle \tilde{\delta}_{21}(\mathbf{k}) \tilde{\delta}_{21}^*(\mathbf{k}') \rangle \equiv (2\pi)^3 \delta^D(\mathbf{k} - \mathbf{k}') P_{21}(k, \mu)$$

$$\begin{aligned} P_{T_b} &\equiv (\bar{T}_b)^2 P_{21} \\ &= (\bar{T}_b / \bar{x}_{\text{HI}})^2 \left\{ [\bar{x}_{\text{HI}}^2 P_{\delta\delta} - 2\bar{x}_{\text{HI}} P_{x\delta} + P_{xx}] \right. \\ &\quad \left. + 2\mu^2 [\bar{x}_{\text{HI}}^2 P_{\delta\delta} - \bar{x}_{\text{HI}} P_{x\delta}] + \mu^4 \bar{x}_{\text{HI}}^2 P_{\delta\delta} \right\} \end{aligned}$$

$$\left\{ \begin{array}{l} P_{\delta\delta} : \text{Matter power spectrum} \\ P_{x\delta} = \bar{x}_i P_{\delta_x\delta} : \text{Density-ionization power spectrum} \\ P_{xx} = \bar{x}_i^2 P_{\delta_x\delta_x} : \text{Ionization power spectrum} \end{array} \right.$$

Ionization fraction :  $x_i = 1 - x_{\text{HI}}$

## ◆ u-space vector : $\mathbf{u}$

Fluctuation of 21 cm line  $\delta_{21} \equiv \frac{\Delta T_b - \Delta \bar{T}_b}{\Delta \bar{T}_b}$

Its power spectrum is

$$P_{21}(\mathbf{u}) = \frac{1}{d_A(z)^2 y(z)} P_{21}(\mathbf{k}) \quad \left\{ \begin{array}{l} d_A(z) : \text{commoving} \\ \text{angular diameter distance} \\ y(z) = \frac{\lambda_{21}(1+z)^2}{H(z)} \end{array} \right.$$

### $\mathbf{u}$ : u-space vector

Fourier dual of angle  $\theta_x, \theta_y$  and frequency (line of sight direction) space:  $\Theta \equiv \theta_x \mathbf{e}_x + \theta_y \mathbf{e}_y + \Delta f \mathbf{e}_z$

$$\mathbf{u} = (u_1, u_2, u_3) = \left( k_1 \frac{d_A(z)}{2\pi}, k_2 \frac{d_A(z)}{2\pi}, k_3 \frac{y(z)}{2\pi} \right)$$

We can get information about  $d_A(z)$  and  $H(z)$  from this relation between  $\mathbf{k}$  and  $\mathbf{u}$  space.



# ◇ 21cm line observations

LOFAR



MWA



PAPER



21CMA



## Next generation ~2020

### ◆ SKA (Square Kilometer Array)



in Australia

2018:

Construction starts.

<http://www.skatelescope.org/>

### ◆ Omniscope



Max Tegmark,  
Matias Zaldarriaga

Phys. Rev. D 82,  
103501 (2010)

From Max Tegmark's presentation

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# ◆ Fisher Information matrix $F_{ij}$

$$F_{\alpha\beta} \equiv \left\langle \frac{\partial^2 \ln L(\boldsymbol{\theta}|\mathbf{x})}{\partial \theta_\alpha \partial \theta_\beta} \right\rangle$$

$L(\boldsymbol{\theta}|\mathbf{x})$ : Likelihood function

$\theta_{\alpha\beta}$ : theoretical parameters

$\mathbf{x}$ : data vector

## Cramér-Rao bound

$$V_{\alpha\beta}(\hat{\boldsymbol{\theta}}) \geq (F^{-1})_{\alpha\beta} \quad V_{\alpha\beta}(\hat{\boldsymbol{\theta}}) : \text{variance of } \hat{\boldsymbol{\theta}}$$

We can estimate minimum variance of  $\hat{\boldsymbol{\theta}}$ .

# ◆ Cosmological parameter set

## Fiducial parameters

$$(\Omega_m h^2, \Omega_b h^2, \Omega_\Lambda, n_s, A_s, \tau, Y_p)$$

$$= (0.1417, 0.02216, 0.6914, 0.9611, 2.214 \times 10^{-9}, 0.0952, 0.25)$$

### **Parametrization of dark energy models**

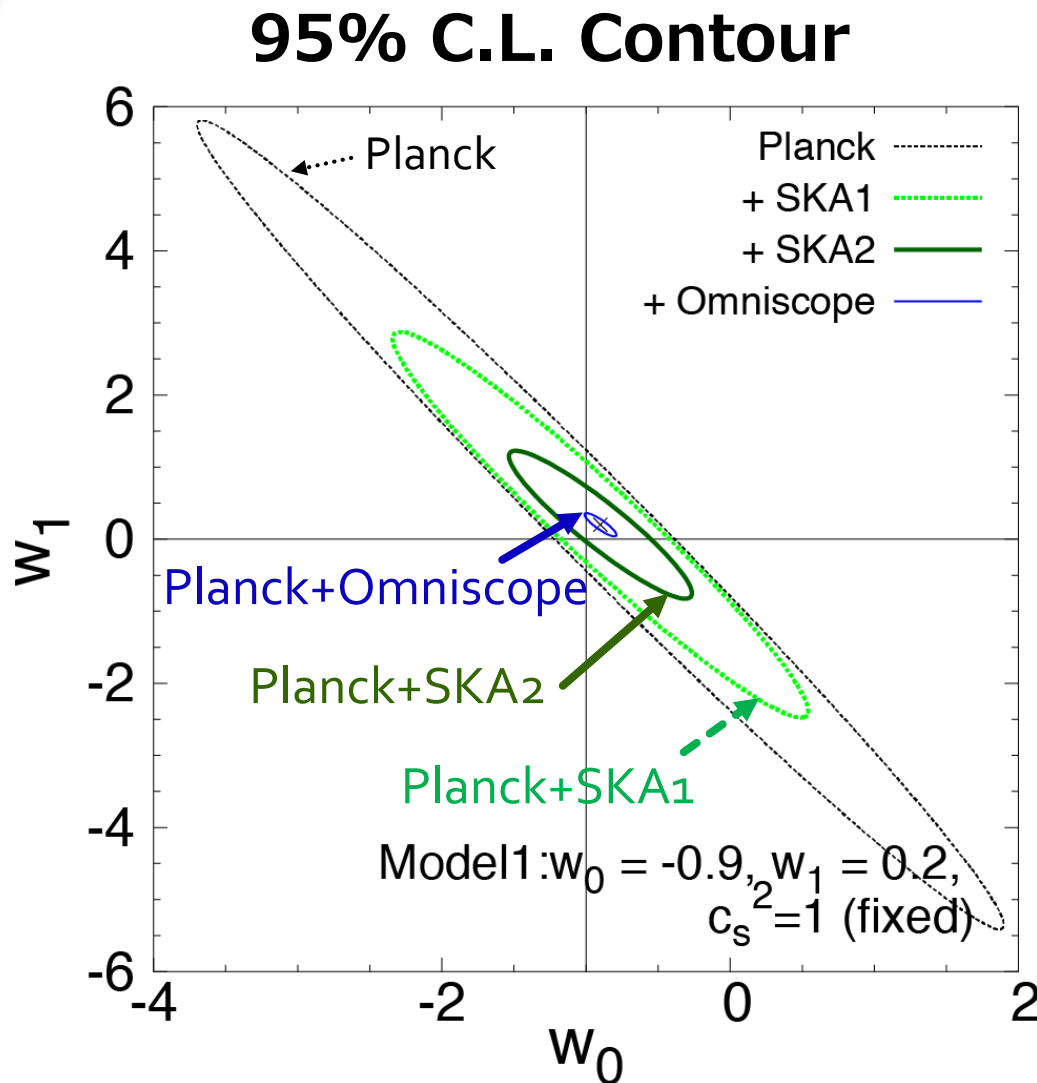
( 1 ) Parametrization 1  $w_0 = -0.9, w_0 = 0.2$

( 2 ) Parametrization 2  $w_0 = -0.9, w_0 = -0.35,$   
 $a_s = 0.5, p = 100$

( 3 ) Parametrization 3  $w_0 = -0.9, b = 1$

# ◆ Expected constraint of Parametrization 1

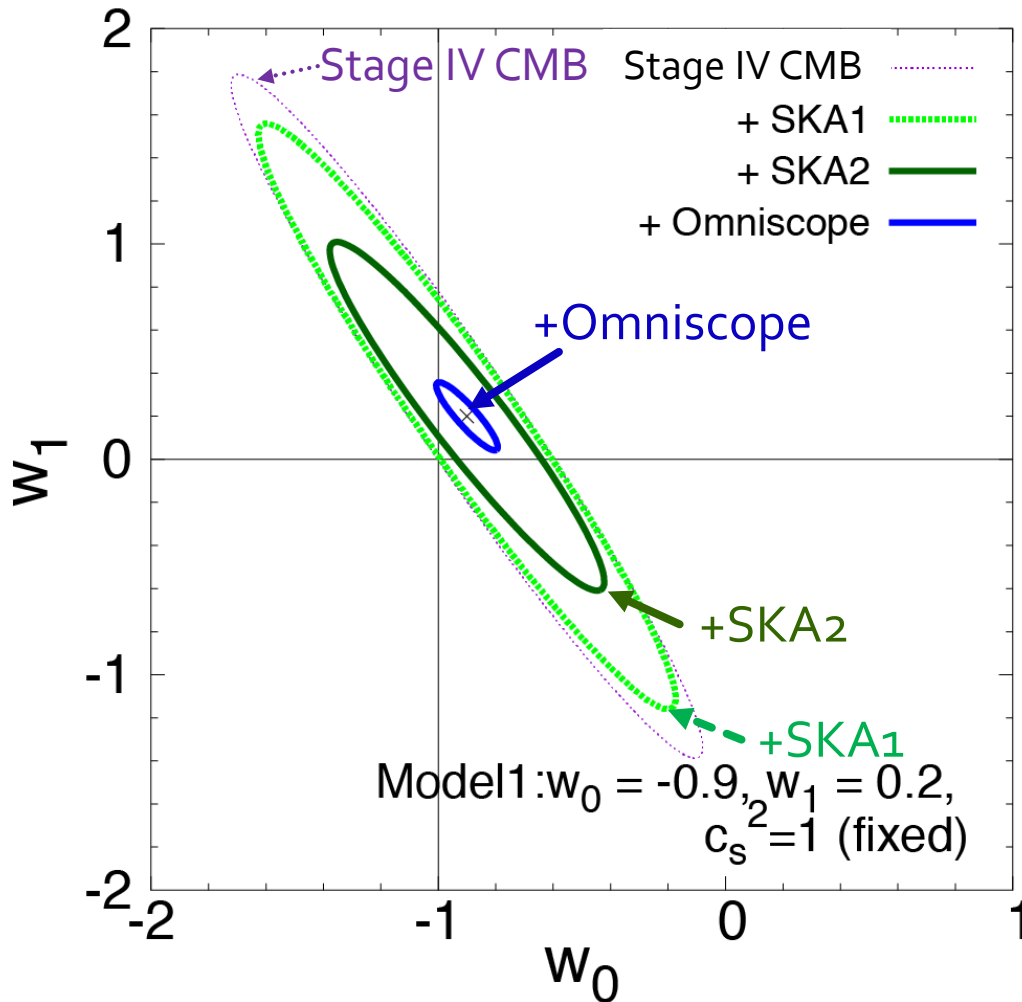
(Preliminary result)



A constraint of **SKA 1 + Planck** is **twice** as good as that of Planck only.

# Expected constraint of Parametrization 1

## 95% C.L. Contour



(Preliminary result)

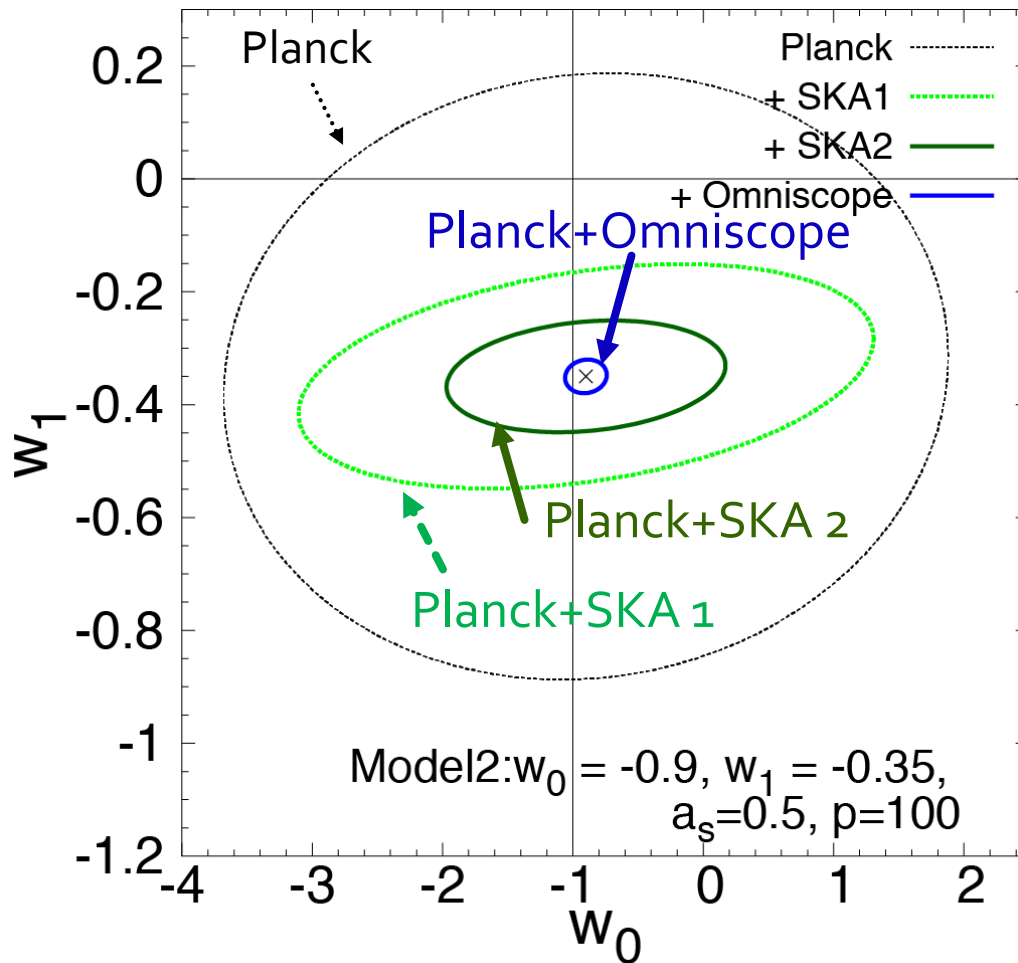
Stage IV CMB:  
(we assume a CMBPol  
like experiment)

A constraint of **SKA 2 +  
stage IV CMB** is twice  
as good as that of  
Stage IV CMB only.

# ◆ Expected constraint of Parametrization 2

(Preliminary result)

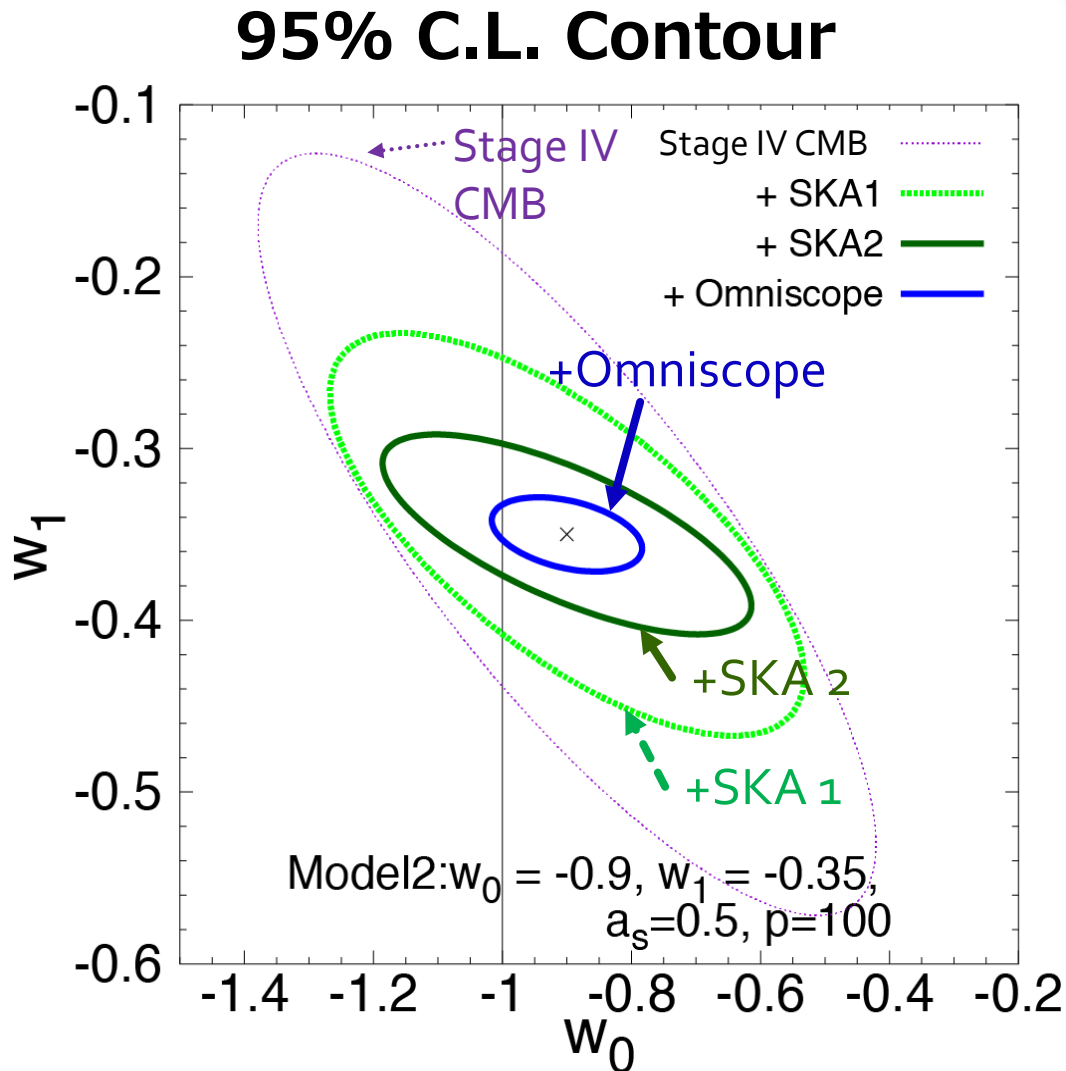
## 95% C.L. Contour



For  $a_s \sim 0.5$ , constraints of 21 cm + Planck are several times as good as that of Planck only.

# Expected constraint of Parametrization 2

(Preliminary result)



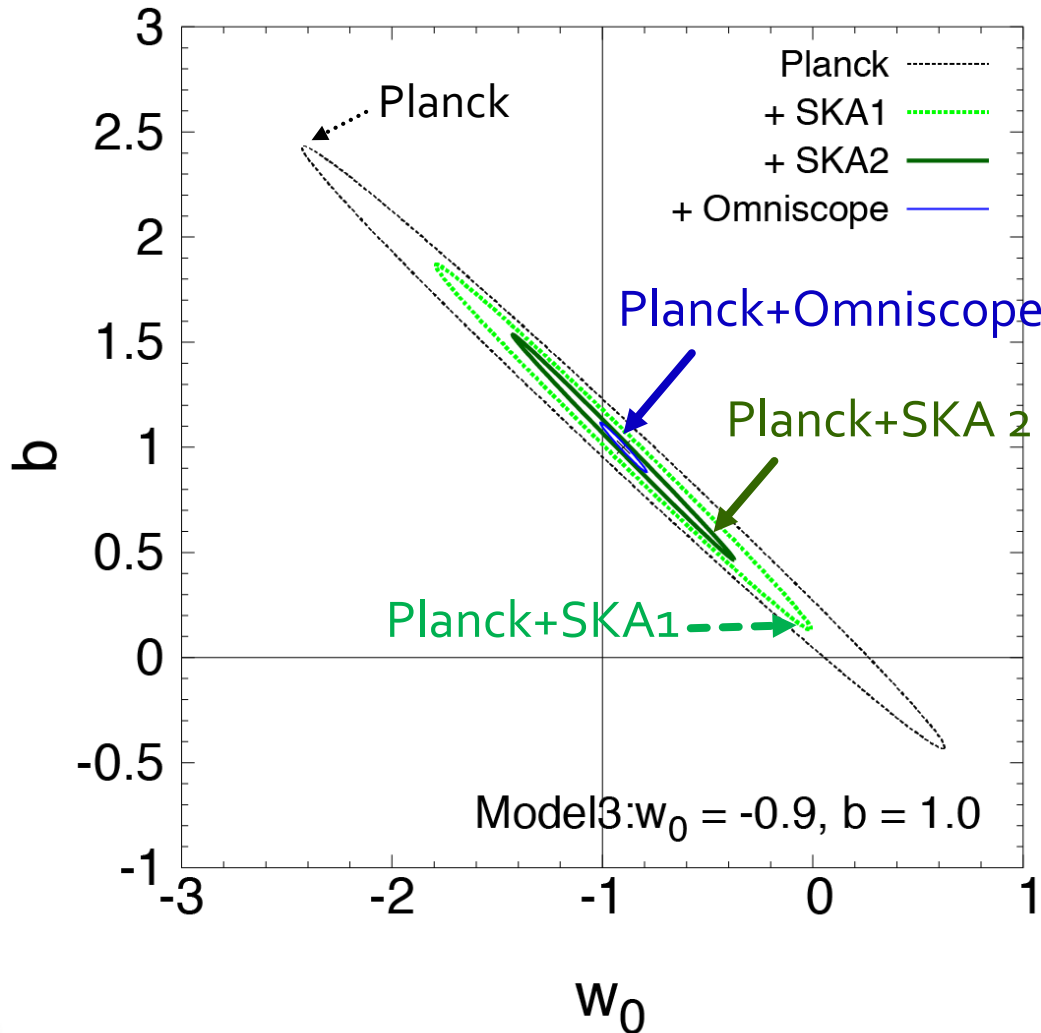
For  $a_s \sim 0.5$ , constraints of 21 cm + stage IV CMB are several times as good as that of stage IV CMB only.



# ◆ Expected constraint of Parametrization 3

(Preliminary result)

## 95% C.L. Contour

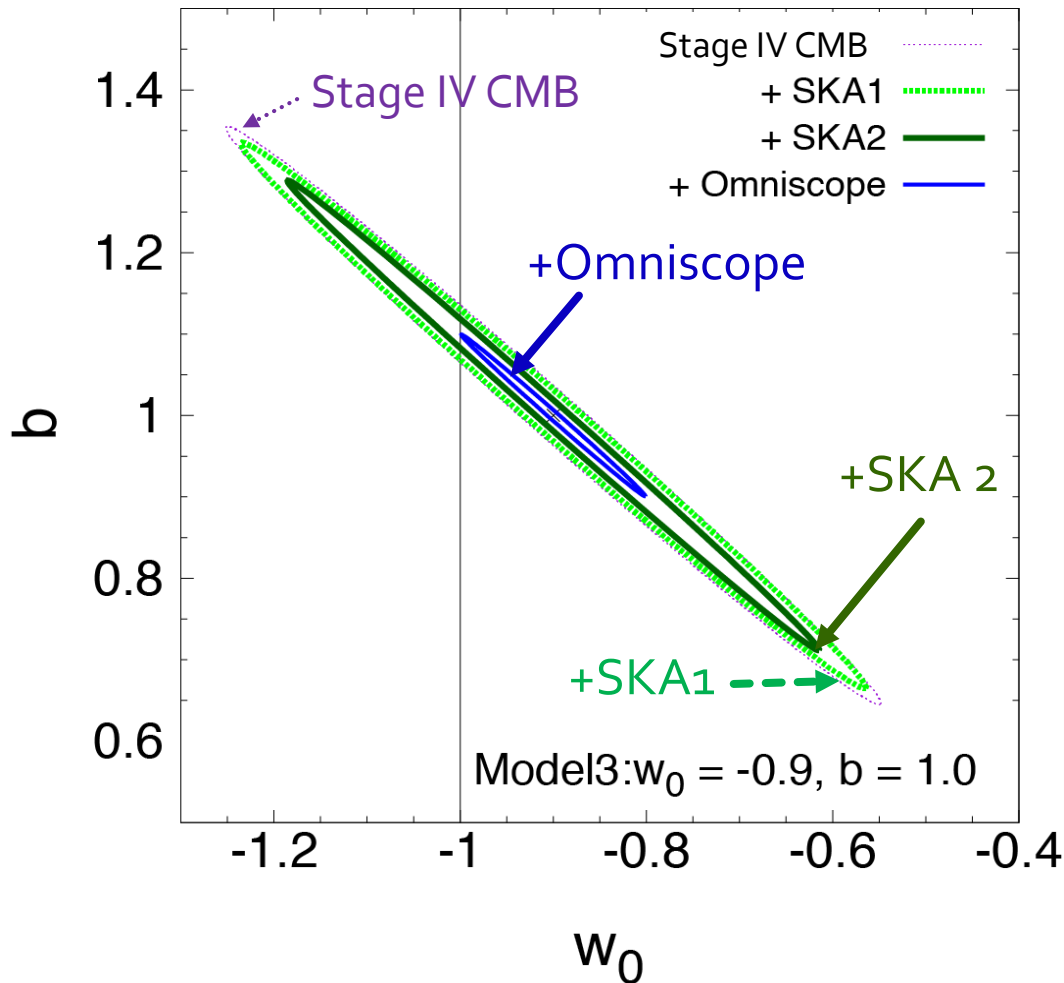


For the early dark energy model, constraints of 21 cm + Planck are several times as good as that of Planck only.

# Expected constraint of Parametrization 3

(Preliminary result)

## 95% C.L. Contour



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## 4. Summary

- We studied sensitivities of **21cm** line to the **dynamical dark energy**.
- **SKA** can strongly improve constraints of **dynamical dark energy**.
- For **suddenly transition model**, constraints of **21 cm + Planck** are **several times** as good as that of **Planck** only.
- For **early dark energy model**, constraints of the combination are **several times** as good as that of **Planck** only.