



Higgs mechanism without Higgs potential in an extra dimension

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Mysteries of the Standard Model ← → 2

> **Symmetry breaking**

Why does nature take M_H^2 negative?

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> **Chiral theory**

Why chiral theory the Standard Model is?

> **Symmetry breaking**

Why does nature take M_H^2 negative?

> **Chiral theory**

Why chiral theory the Standard Model is?

> **Mass hierarchy**

Why so different the masses of the fermions are?

Purpose



Purpose

In the context of **5d gauge theories on an interval**, we want to figure out the mysteries of the standard model

- > Symmetry breaking
- > Chiral theory
- > Mass hierarchy

naturally.

Setting



> 5d U(1) gauge theory on an interval

- with
- 5d fermions
 - 5d complex scalar

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- > 5d U(1) gauge theory on an interval
 - with
 - 5d fermions
 - 5d complex scalar
- > All fields live in the bulk.
- > No negative square mass for the scalar.
- > General boundary conditions
 - compatible with { 5d gauge invariance
the action principle

Results



Results

5d gauge theories on an interval



The low energy effective theories

4d gauge theories

- + Symmetry breaking
- + Chiral theories
- + Mass hierarchy

5d U(1) gauge field

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The boundary condition for the U(1) gauge field is restricted uniquely by the 5d gauge invariance.

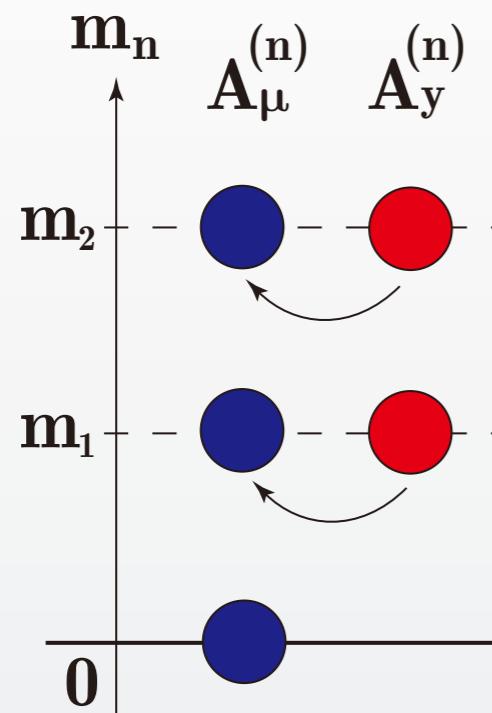
$$\left\{ \begin{array}{l} \partial_y A_\mu(x, 0) = 0 \\ A_y(x, 0) = 0 \end{array} \right. \quad \text{y=0} \qquad \qquad \qquad \left\{ \begin{array}{l} \partial_y A_\mu(x, L) = 0 \\ A_y(x, L) = 0 \end{array} \right. \quad \text{y=L}$$

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> 4d mass spectrum

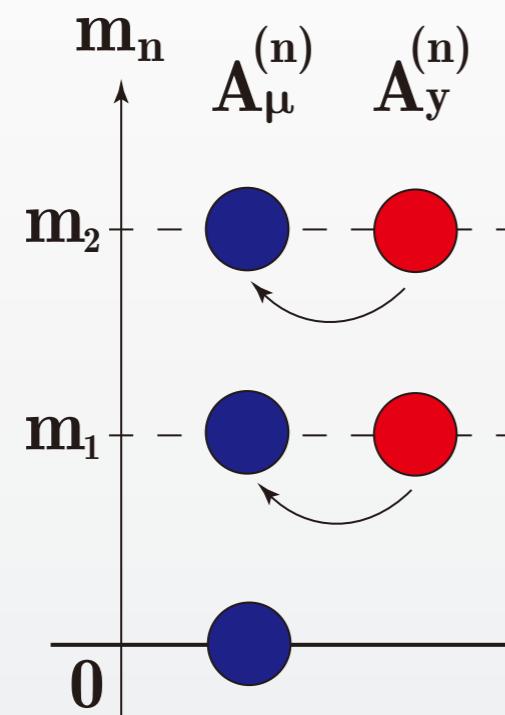


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> 4d mass spectrum



!

A wider class of boundary conditions are allowed to non-abelian gauge fields.

Higgs mechanism with $M^2 > 0$



5d complex scalar field

$$S = \int d^4x \int_0^L dy \left[\Phi^*(x, y) (\partial^\mu \partial_\mu + \partial_y^2) \Phi(x, y) - V(\Phi) \right]$$

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 **>Positive square mass**

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+ **General boundary conditions**
 (2-parameter family)

$$\begin{cases} \Phi(x, 0) + L_+ \partial_y \Phi(x, 0) = 0 \\ \Phi(x, L) - L_- \partial_y \Phi(x, L) = 0 \end{cases} \quad \begin{array}{c} \text{y=0} \quad \text{y=L} \\ \hline \text{No outflow of the probability current.} \end{array}$$

$(-\infty \leq L_{\pm} \leq +\infty)$

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$$V_{4d}(\Phi) = \int_0^L dy \left[\Phi^*(x, y) (-\partial_y^2 + M^2) \Phi(x, y) + (\text{ quartic term }) \right]$$

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> The vacuum of the theory is given by the configuration which minimizes $V_{4d}(\Phi)$.

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- > The vacuum of the theory is given by the configuration which minimizes $V_{4d}(\Phi)$.
- > What $V_{4d}(\Phi)$ looks like when we see it from the point of view of 4d mass eigenstates?

Higgs mechanism with $M^2 > 0$

$$V_{4d}(\Phi) = \int_0^L dy \left[\Phi^*(x, y) (-\partial_y^2 + M^2) \Phi(x, y) + (\text{ quartic term }) \right]$$

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$$\Phi(x, y) = \sum_n \varphi_n(x) f_n(y)$$

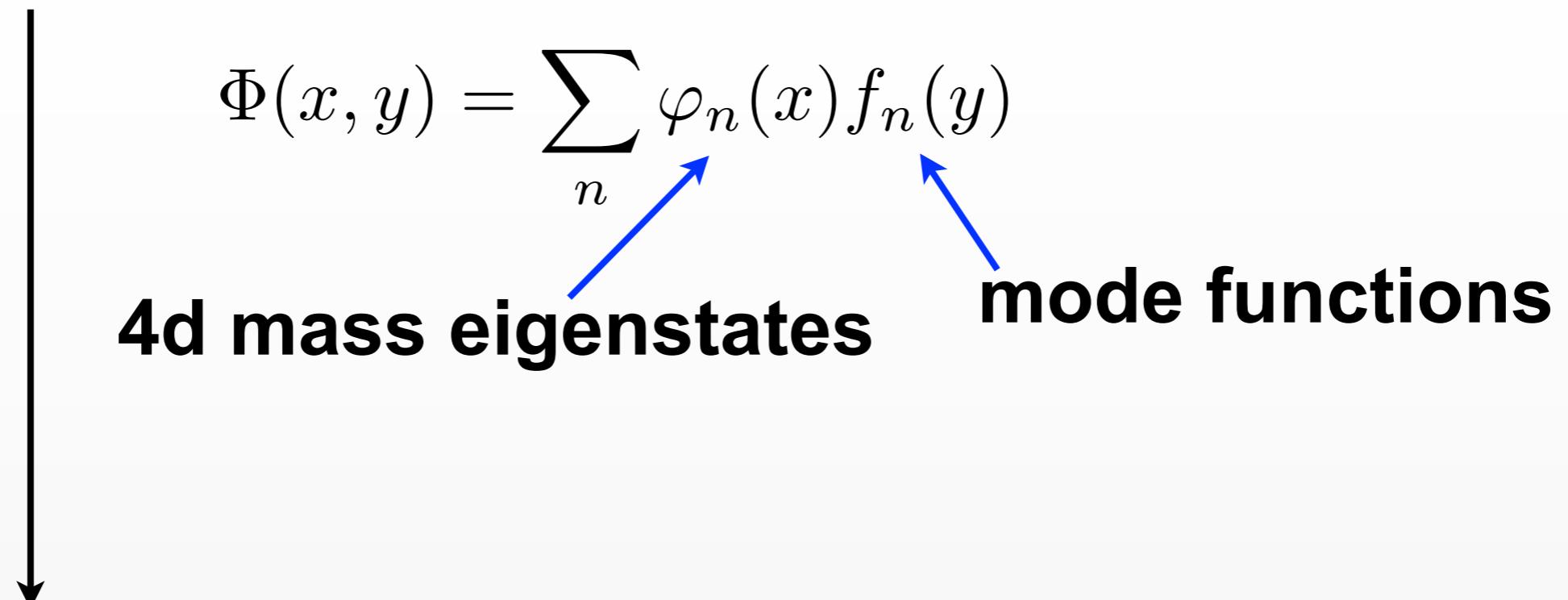


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4d mass eigenstates **mode functions**



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$$\begin{aligned} \Phi(x, y) &= \sum_n \varphi_n(x) f_n(y) \\ -\partial_y^2 f_n(y) &= E_n f_n(y) \\ + \left\{ \begin{array}{l} f_n(0) + L_+ \partial_y f_n(0) = 0 \\ f_n(L) - L_- \partial_y f_n(L) = 0 \end{array} \right. \end{aligned}$$

Boundary conditions

Higgs mechanism with $M^2 > 0$

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Higgs mechanism with $M^2 > 0$

$$V_{4d}(\Phi) = \int_0^L dy \left[\Phi^*(x, y) (-\partial_y^2 + M^2) \Phi(x, y) + (\text{quartic term}) \right]$$

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$$-\partial_y^2 f_n(y) = E_n f_n(y)$$

$$\int_0^L f_m^*(y) f_n(y) dy = \delta_{m,n}$$

+ $\begin{cases} f_n(0) + L_+ \partial_y f_n(0) = 0 \\ f_n(L) - L_- \partial_y f_n(L) = 0 \end{cases}$

$$V_{4d}(\varphi_n) = \sum_n (\underline{E_n + M^2}) |\varphi_n|^2 + (\text{quartic term})$$

> **Negative square mass if $E_B + M^2 < 0$ due to the eigenvalue of the bound state !!**

Non-trivial phase structures

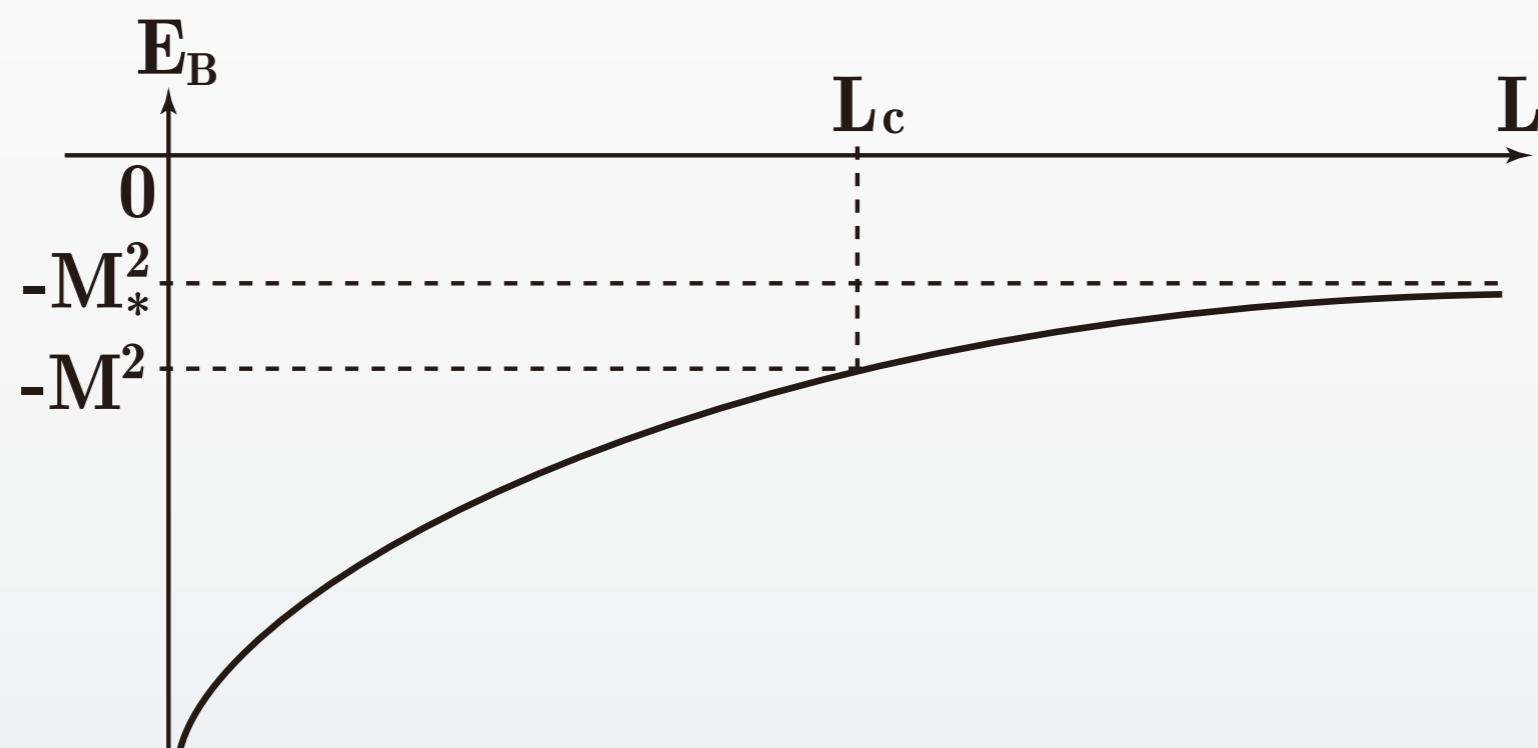
(e.g.) Energy spectrum of the bound state

$$L_{\pm} > 0 > L_{\mp}$$

with

$$L_{\pm} + L_{\mp} < 0$$

L_{\pm}, M^2 : fix



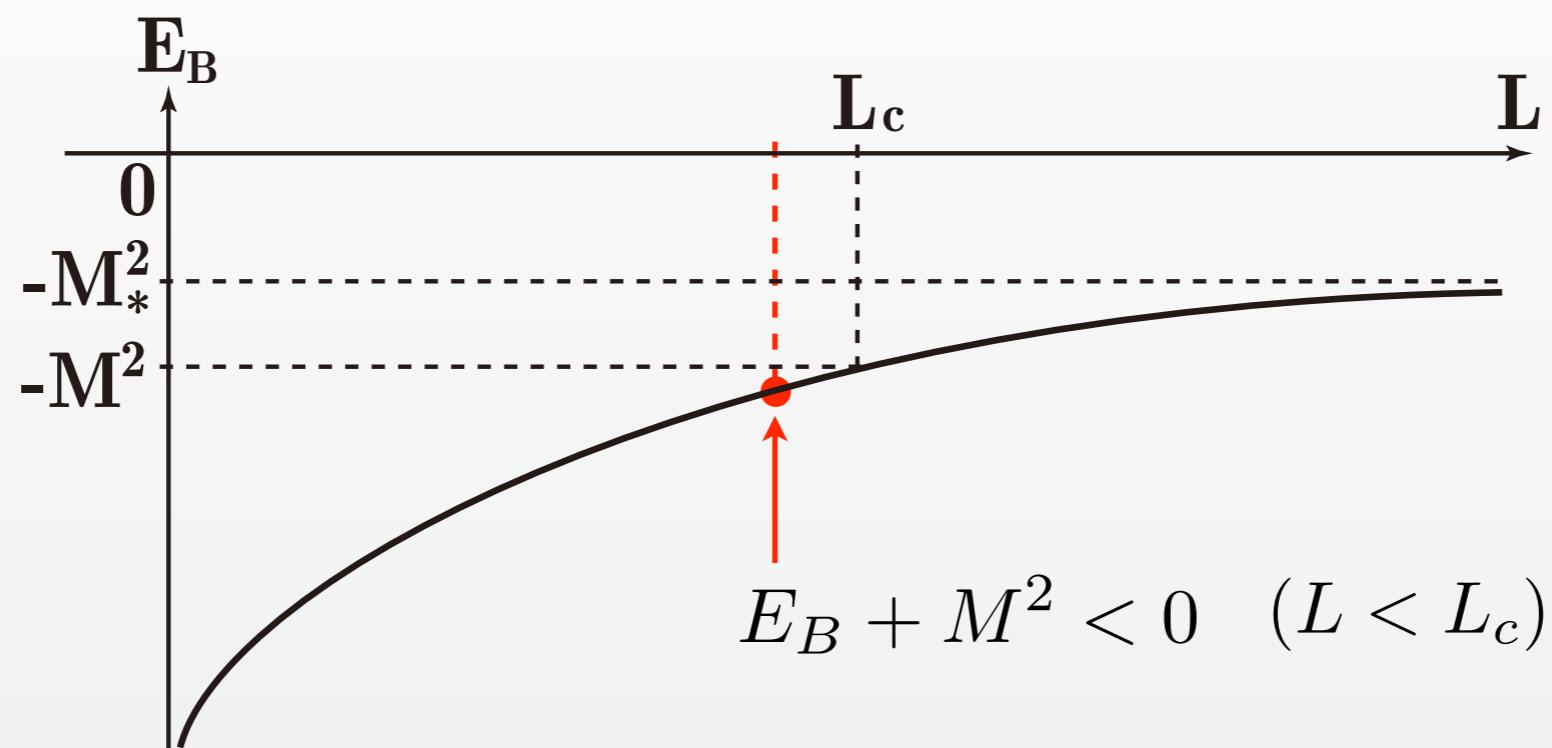
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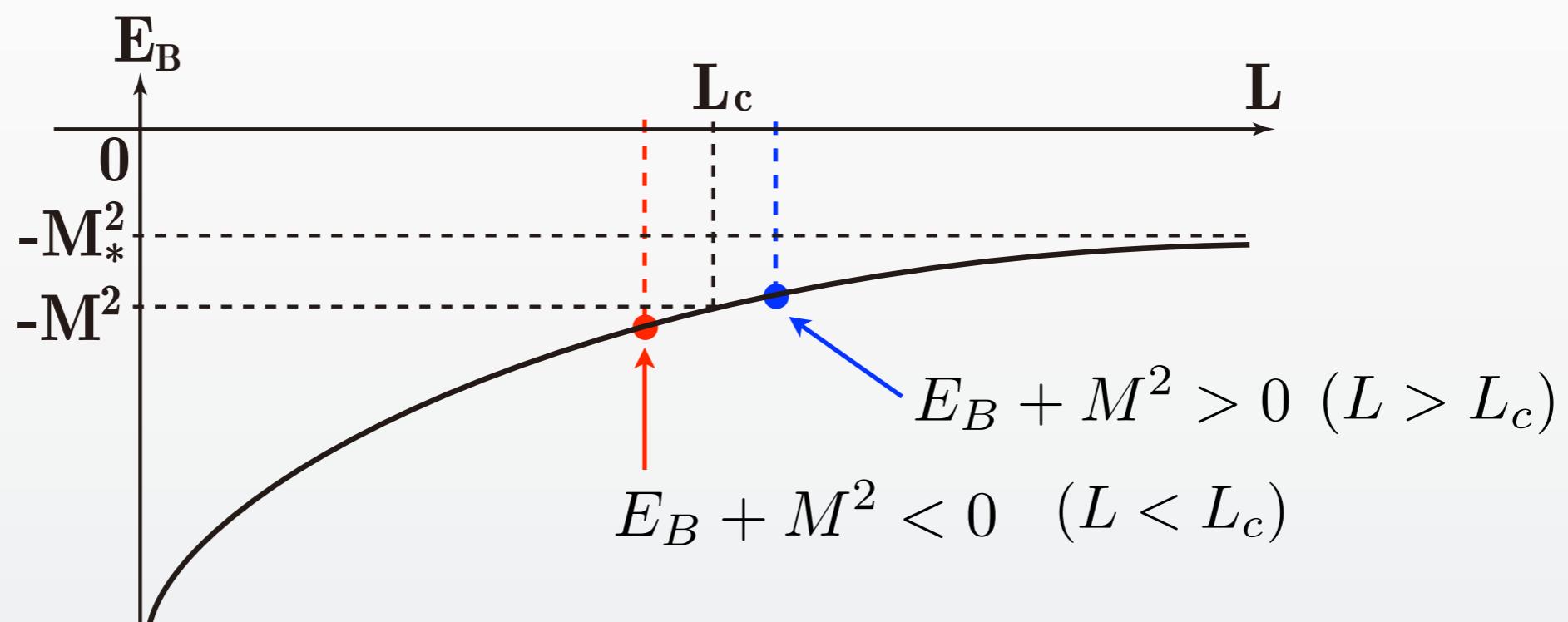
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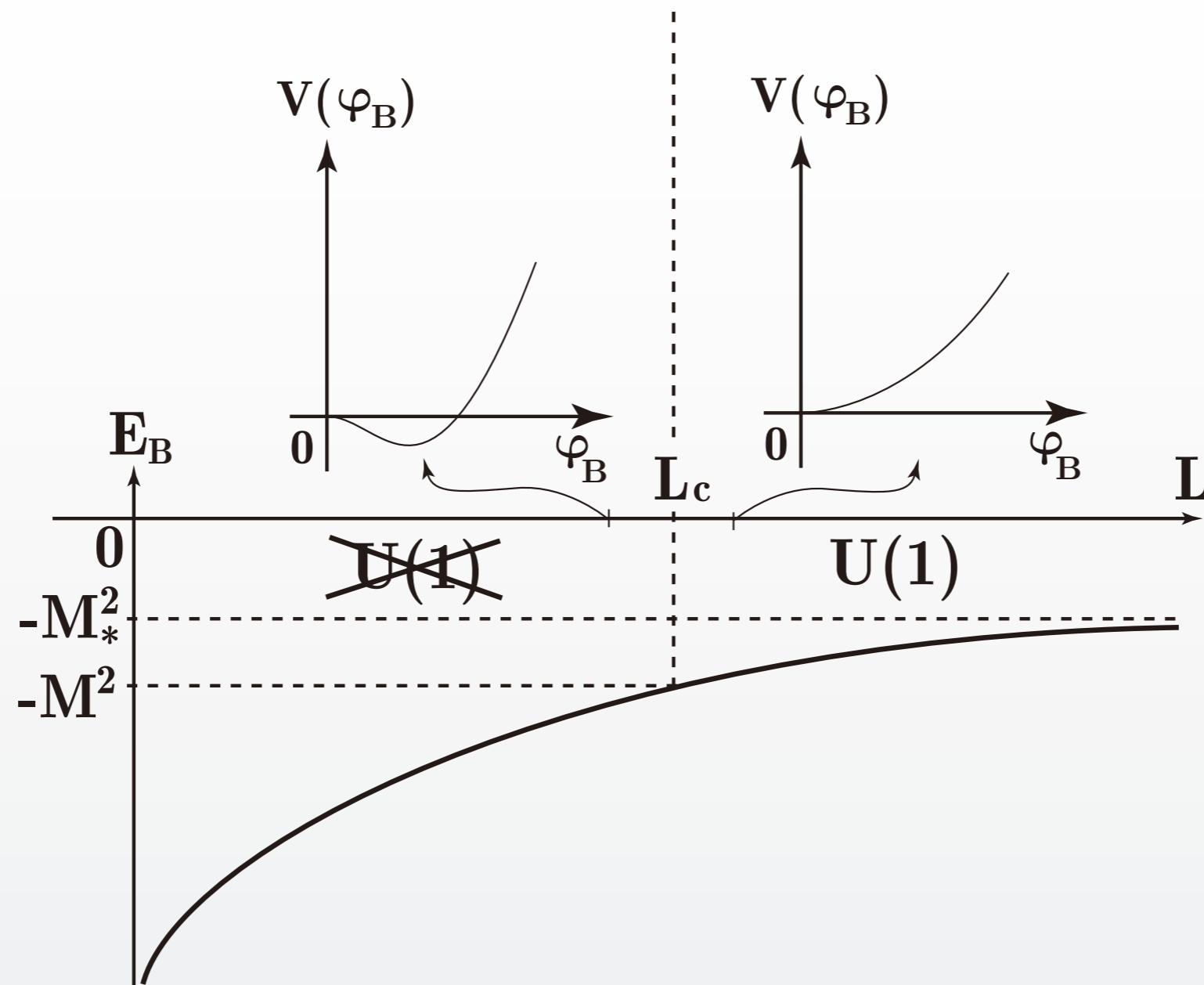
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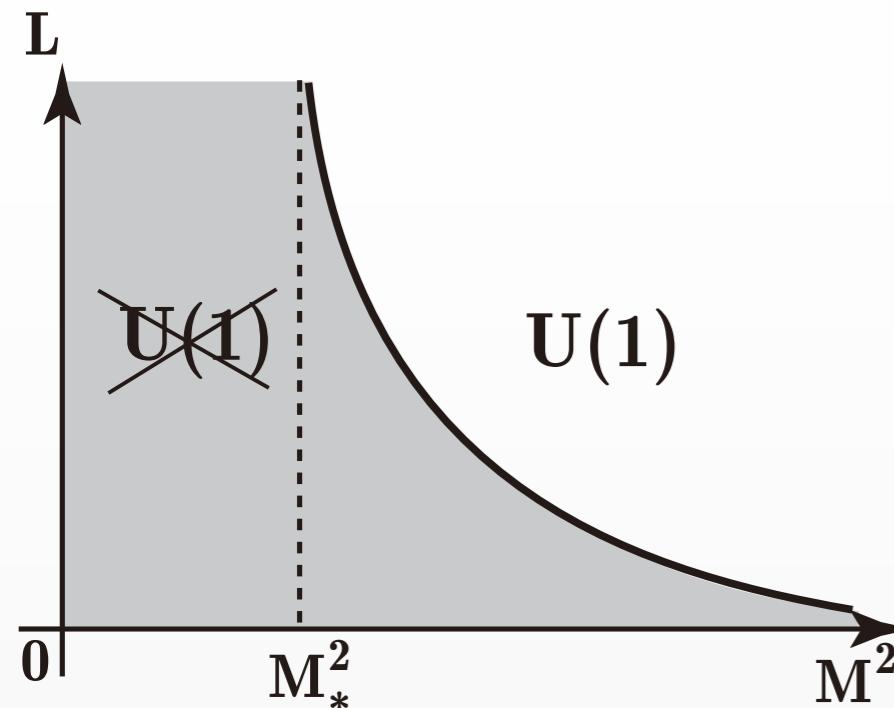
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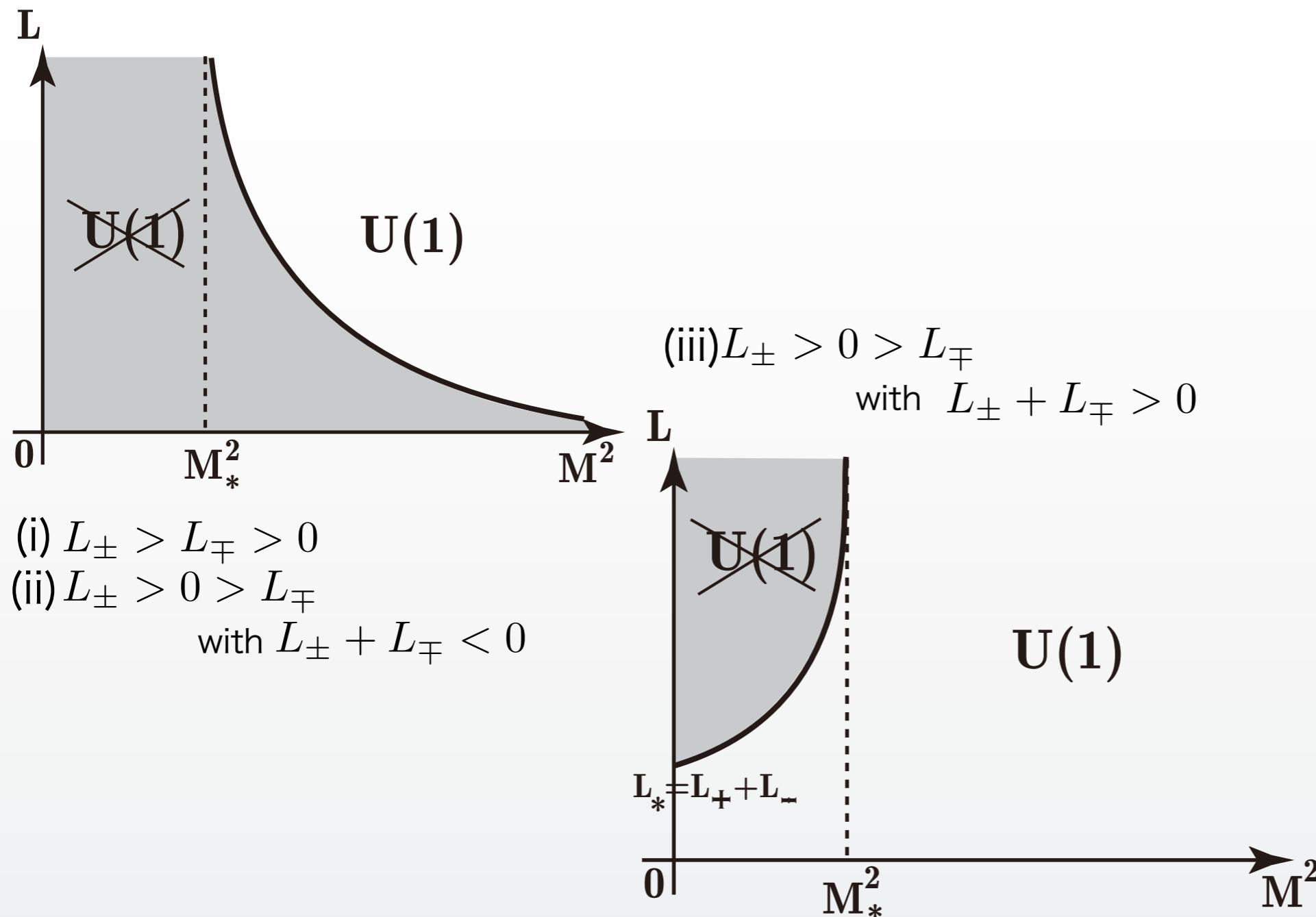


Non-trivial phase structures

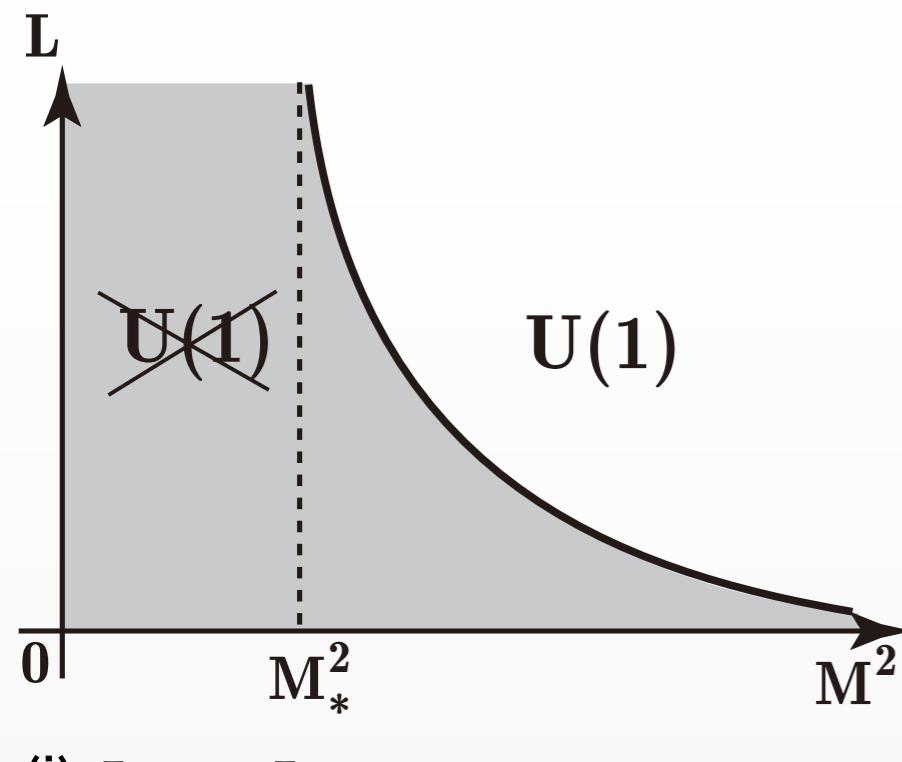


- (i) $L_{\pm} > L_{\mp} > 0$
- (ii) $L_{\pm} > 0 > L_{\mp}$
with $L_{\pm} + L_{\mp} < 0$

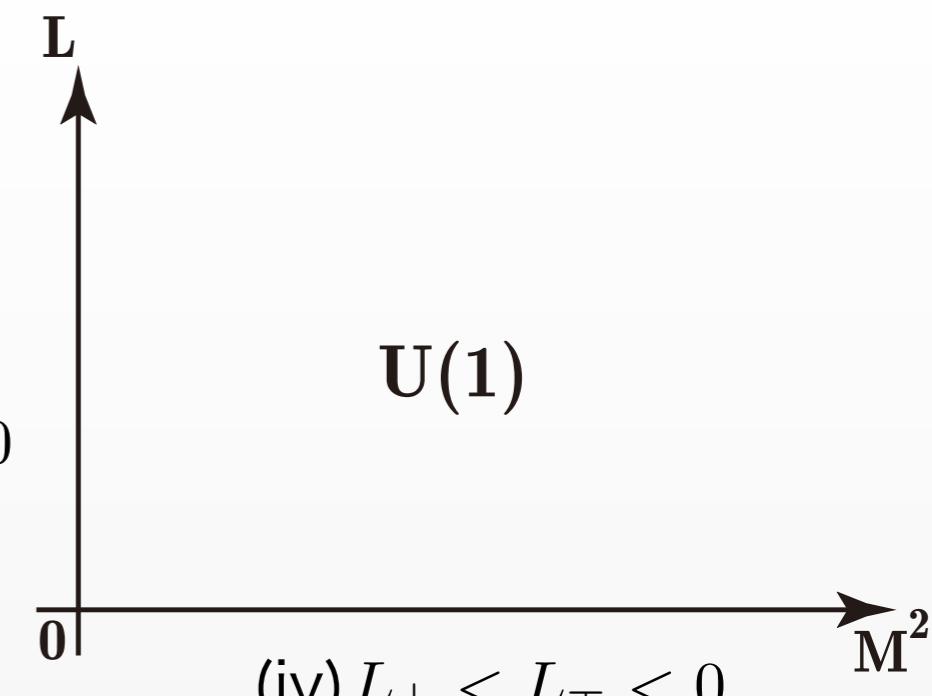
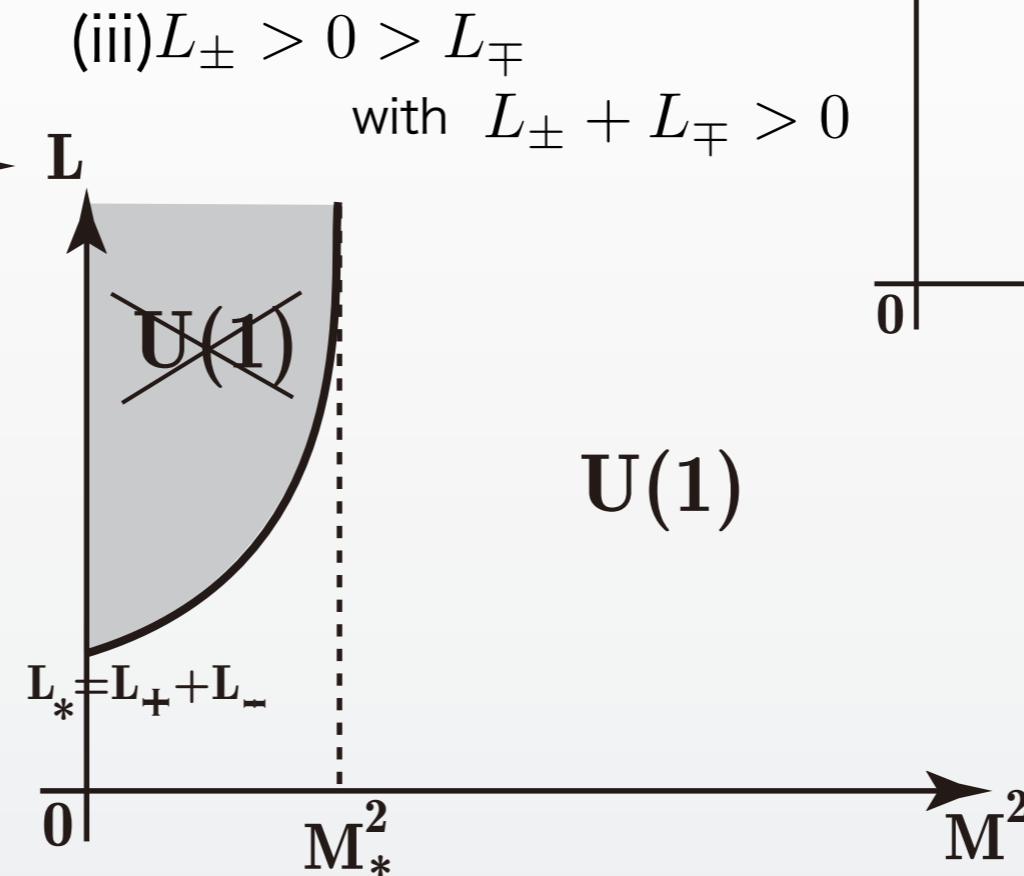
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- (ii) $L_{\pm} > 0 > L_{\mp}$
with $L_{\pm} + L_{\mp} < 0$



Chiral theory

Chiral theory

Boundary conditions for the fermion are restricted to 3-cases by the action principle.



$$\begin{cases} (\partial_y + M_F) \Psi_R(x, 0) = 0 \\ \Psi_L(x, 0) = 0 \end{cases}$$

or

$$\begin{cases} \Psi_R(x, 0) = 0 \\ (-\partial_y + M_F) \Psi_L(x, 0) = 0 \end{cases}$$

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Bulk mass

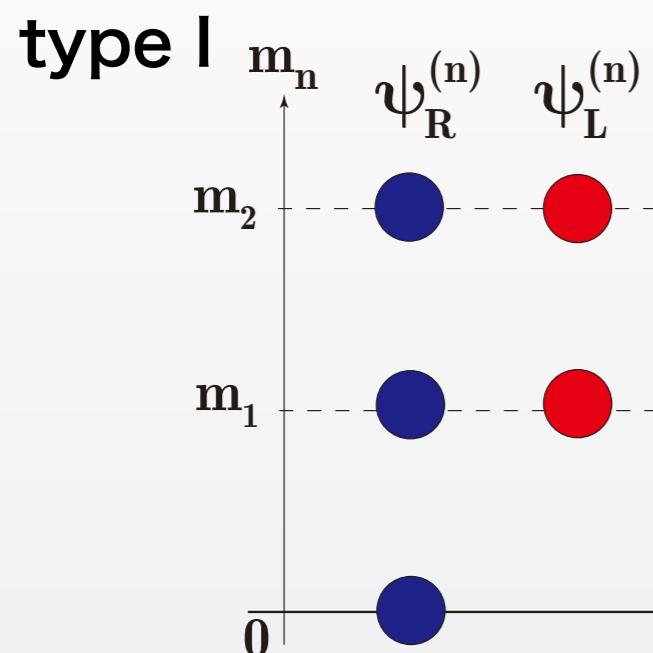
$$\begin{cases} \left(\partial_y + \frac{M_F}{\text{Bulk mass}} \right) \Psi_R(x, 0) = 0 \\ \Psi_L(x, 0) = 0 \end{cases} \quad \text{or} \quad \begin{cases} \Psi_R(x, 0) = 0 \\ \left(-\partial_y + \frac{M_F}{\text{Bulk mass}} \right) \Psi_L(x, 0) = 0 \end{cases}$$

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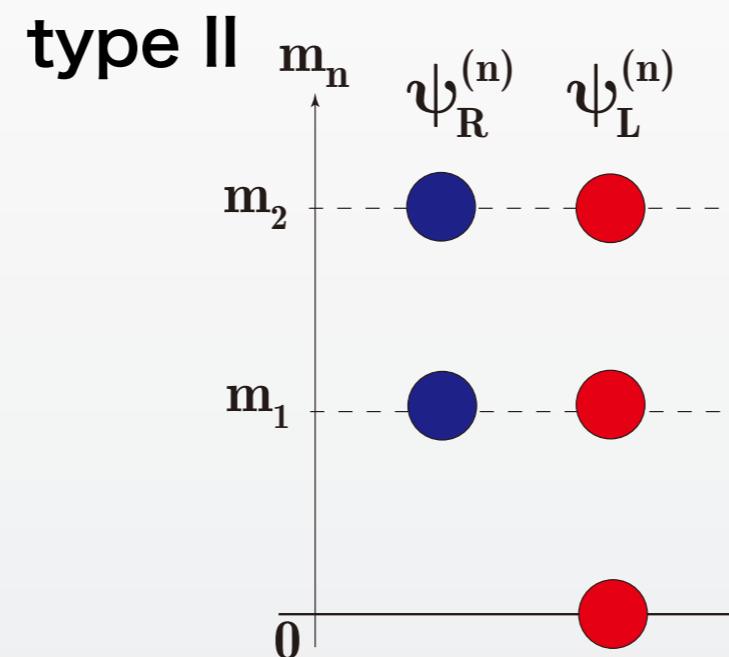
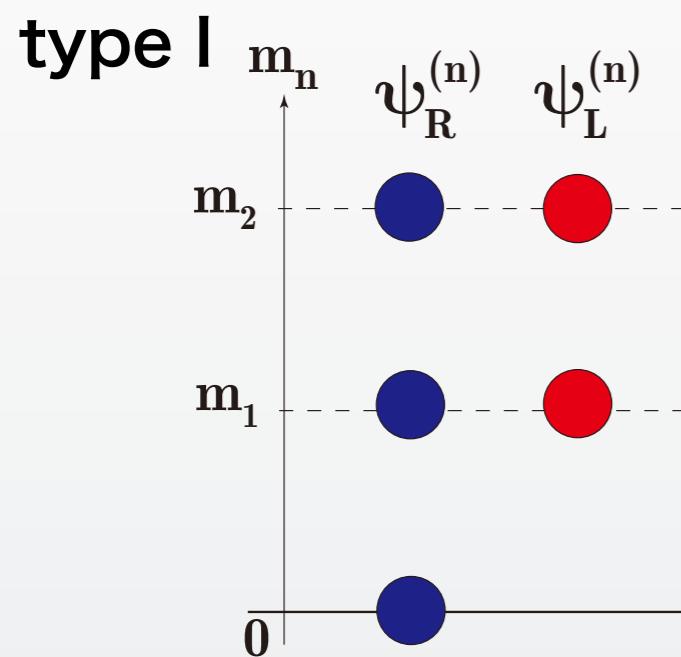
$$\begin{array}{ccc}
 & \xrightarrow{\text{type I}} & \\
 \left\{ \begin{array}{l} (\partial_y + M_F) \Psi_R(x, 0) = 0 \\ \Psi_L(x, 0) = 0 \end{array} \right. & \longleftrightarrow & \left\{ \begin{array}{l} (\partial_y + M_F) \Psi_R(x, L) = 0 \\ \Psi_L(x, L) = 0 \end{array} \right. \\
 \text{or} & & \text{or} \\
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Chiral theory

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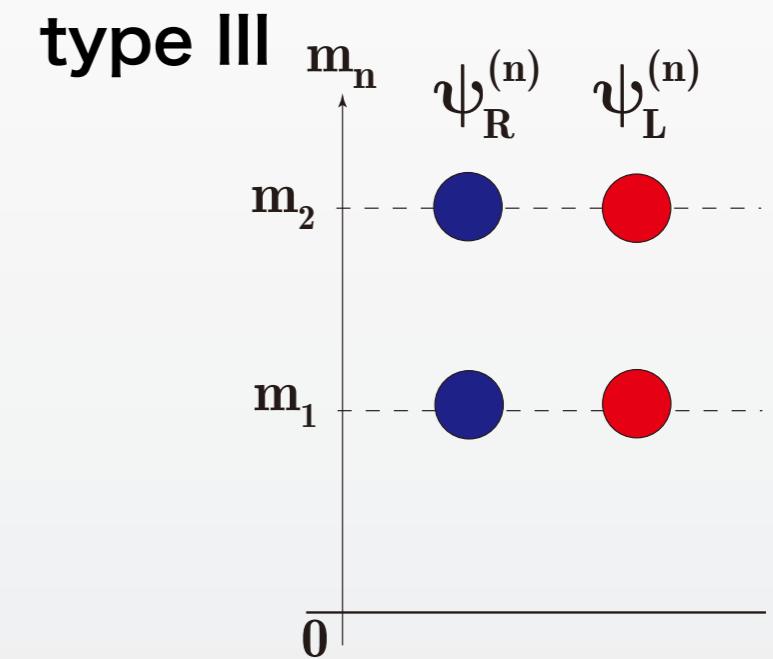
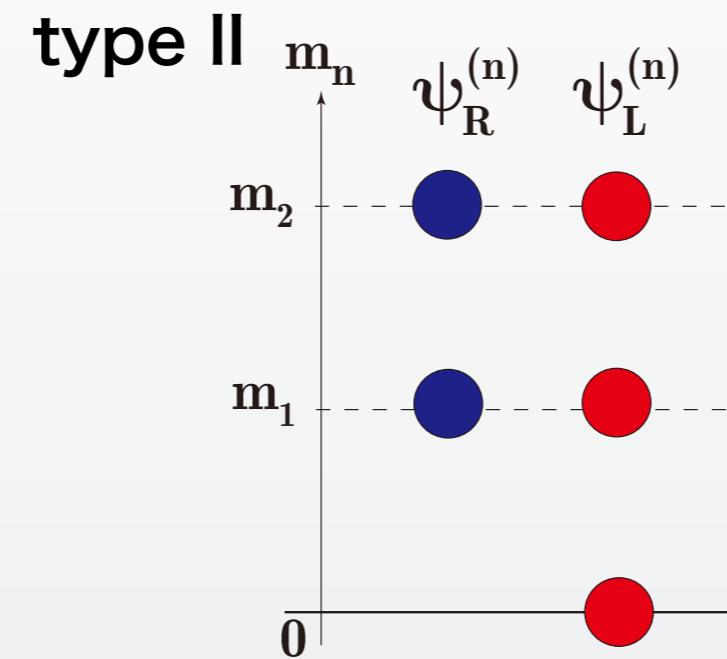
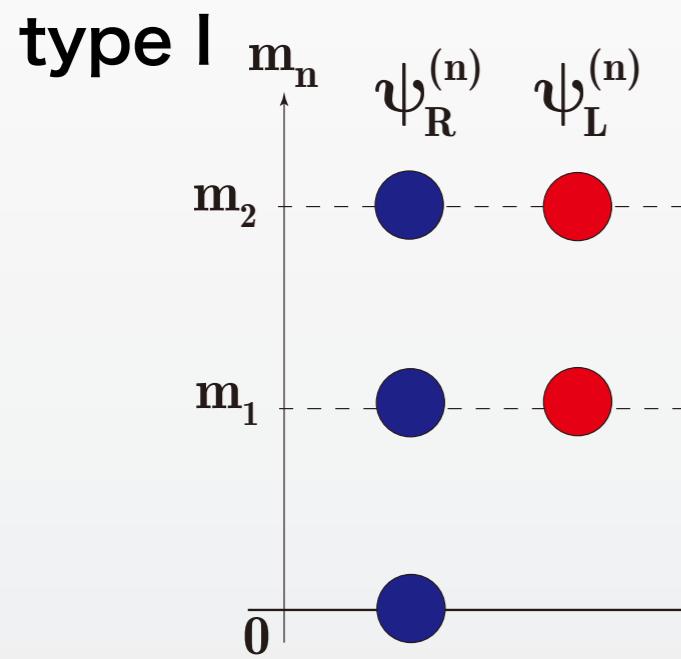
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 \text{or} & & \text{or} \\
 \left\{ \begin{array}{l} \Psi_R(x, 0) = 0 \\ (-\partial_y + M_F) \Psi_L(x, 0) = 0 \end{array} \right. & \xleftrightarrow{\text{type II}} & \left\{ \begin{array}{l} \Psi_R(x, L) = 0 \\ (-\partial_y + M_F) \Psi_L(x, L) = 0 \end{array} \right.
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Chiral theory

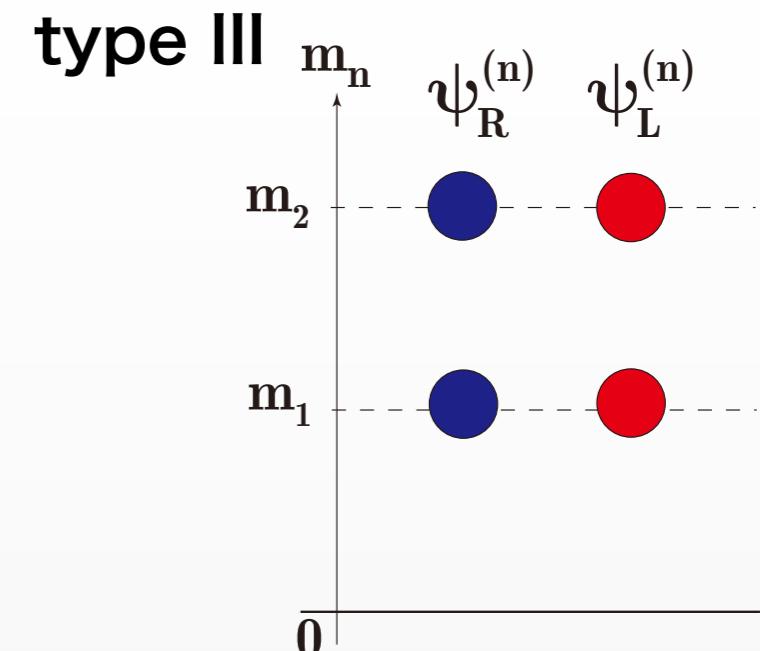
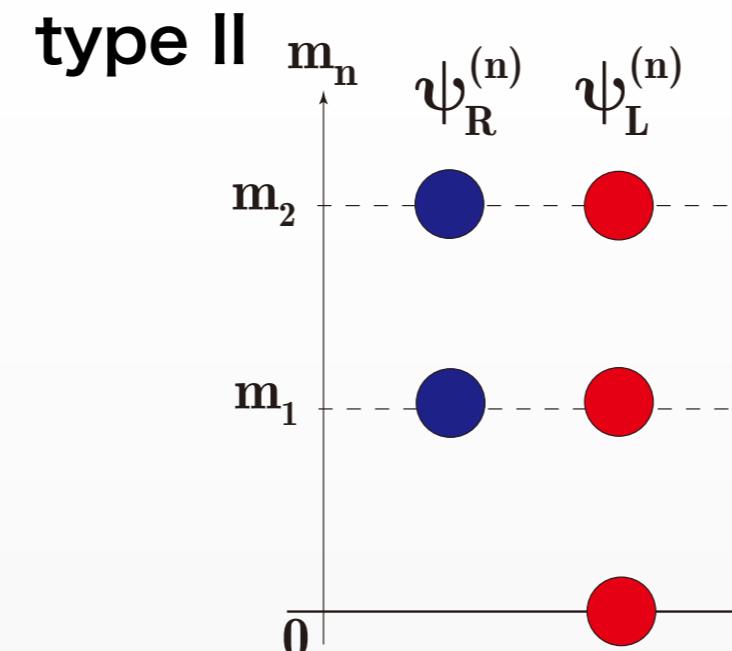
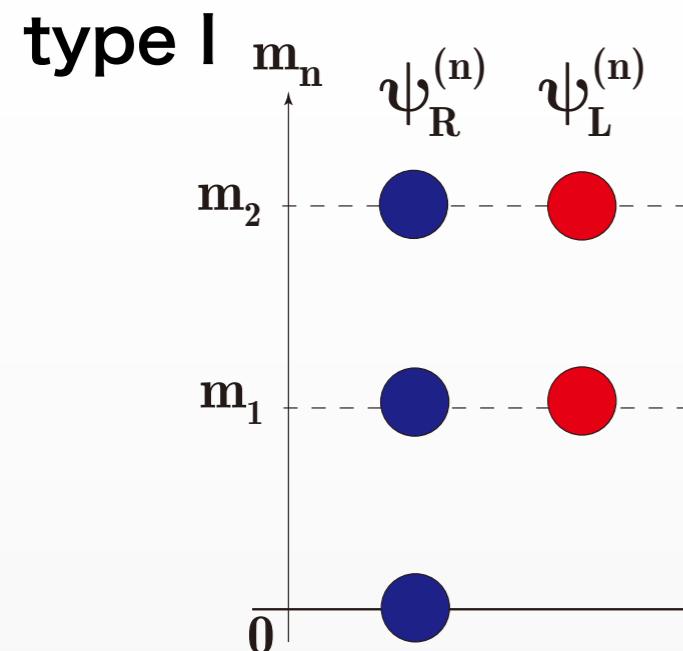
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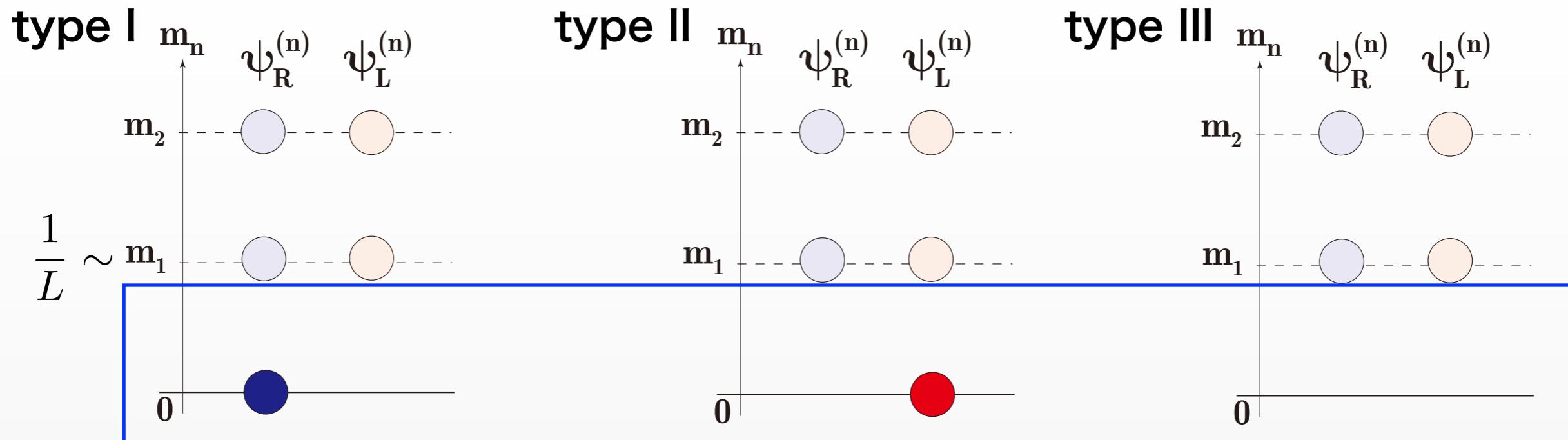
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Chiral theory

Boundary conditions for the fermion are restricted to 3-cases by the action principle.



> 4d low energy effective theory at $E < 1/L$ is chiral theory, irrespective of the bulk mass.

Mass hierarchy

Mass hierarchy

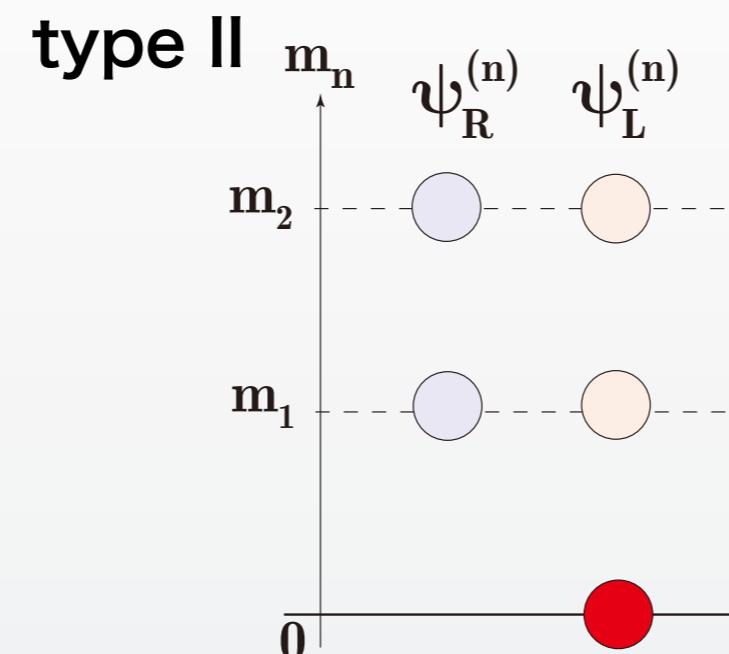
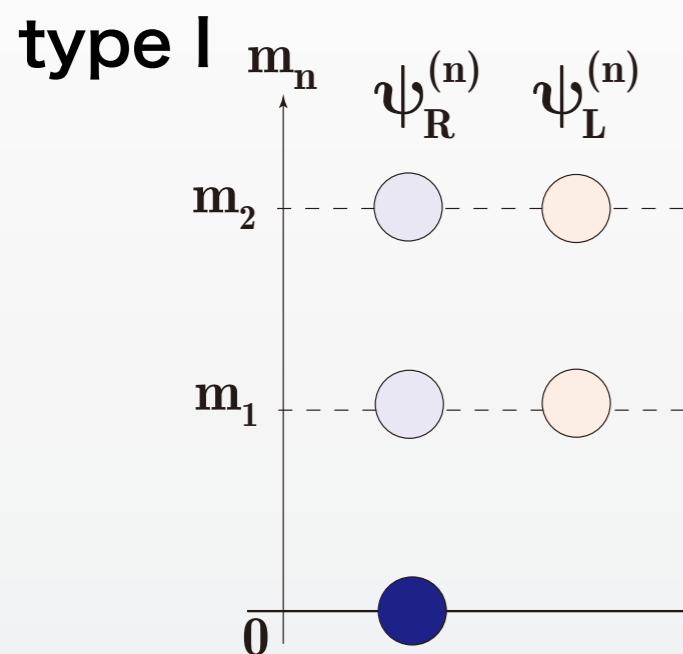
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Mass hierarchy

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$$\Psi_i(x, y) = \psi_{i_R}(x)\mathcal{F}_i(y) + \dots$$

$$\Psi_j(x, y) = \psi_{i_L}(x)\mathcal{G}_j(y) + \dots$$



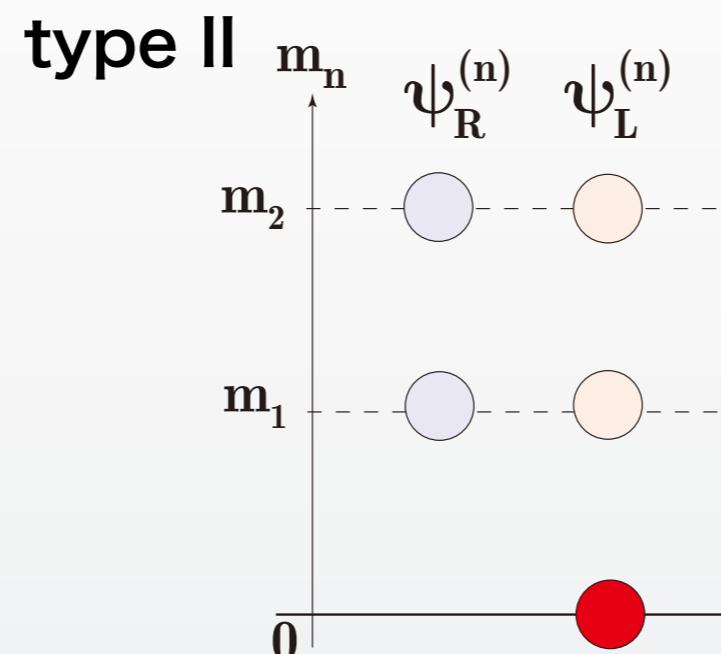
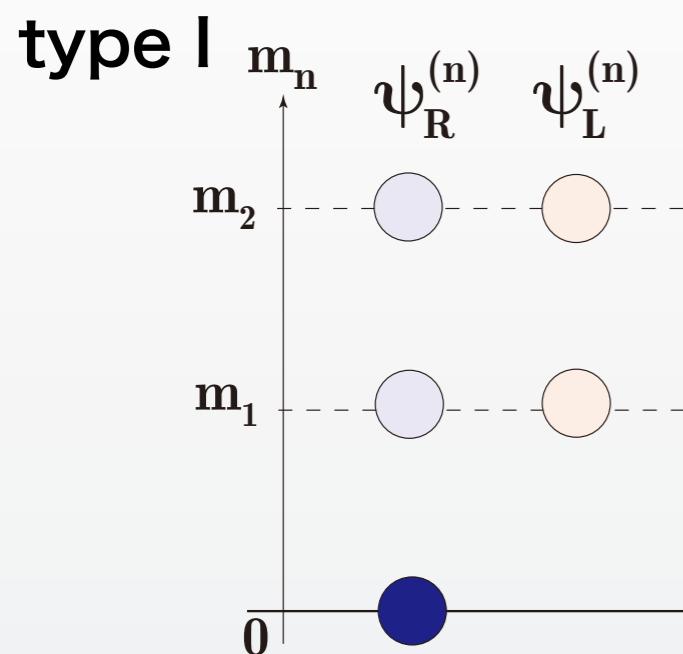
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**mode functions
of the zero modes.**



Mass hierarchy

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$$\Psi_j(x, y) = \psi_{i_L}(x)\mathcal{G}_j(y) + \dots$$

$$\longrightarrow \begin{cases} (\partial_y + M_F)\mathcal{F}_i(y) = 0 \\ (\partial_y + M_F)\mathcal{G}_j(y) = 0 \end{cases}$$

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$$\rightarrow \left\{ \begin{array}{l} (\partial_y + \underline{M_F})\mathcal{F}_i(y) = 0 \\ (\partial_y + \underline{M_F})\mathcal{G}_j(y) = 0 \end{array} \right.$$

Bulk mass

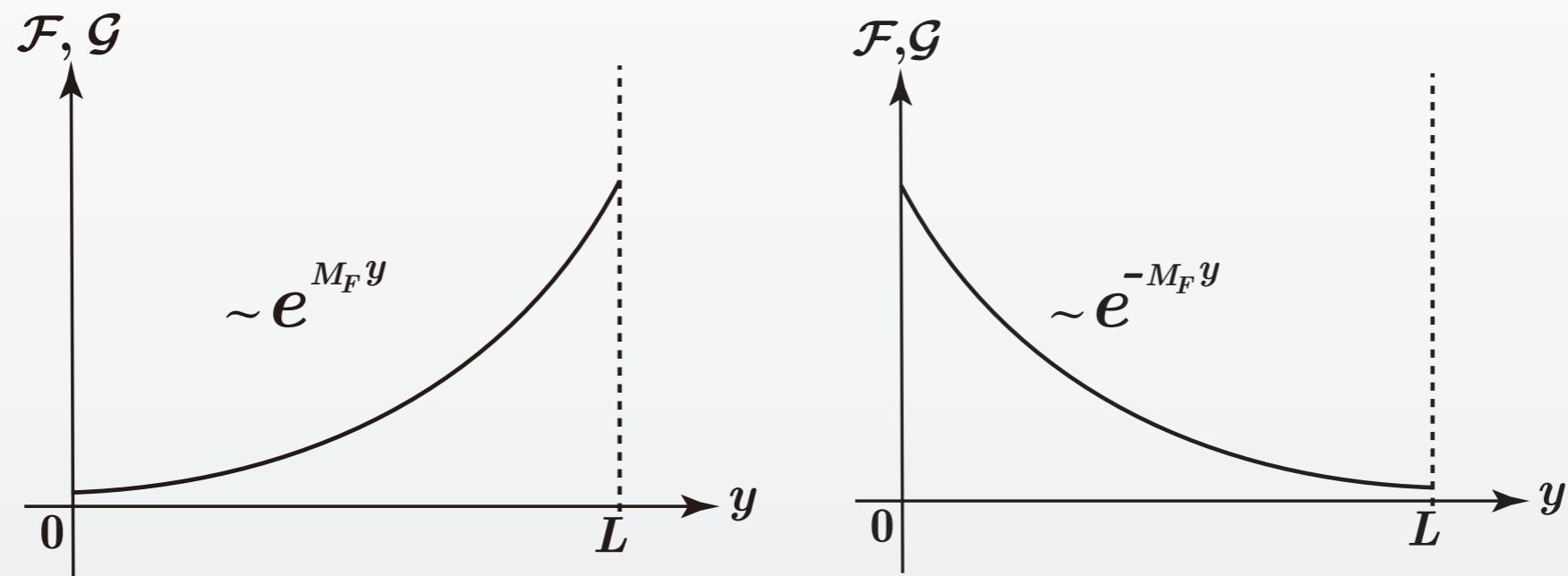
Mass hierarchy

> Chiral zero mode of the fermion is localized at one of the boundaries.

$$\Psi_i(x, y) = \psi_{i_R}(x)\mathcal{F}_i(y) + \dots$$

$$\Psi_j(x, y) = \psi_{i_L}(x)\mathcal{G}_j(y) + \dots$$

$$\rightarrow \left\{ \begin{array}{l} (\partial_y + M_F)\mathcal{F}_i(y) = 0 \\ (\partial_y + M_F)\mathcal{G}_j(y) = 0 \end{array} \right.$$



Mass hierarchy

- > Chiral zero mode of the fermion is localized at one of the boundaries.
- > Vacuum Expectation Value of the scalar field necessarily possesses y-dependence.

$$\langle \Phi(x, y) \rangle = \phi(y)$$

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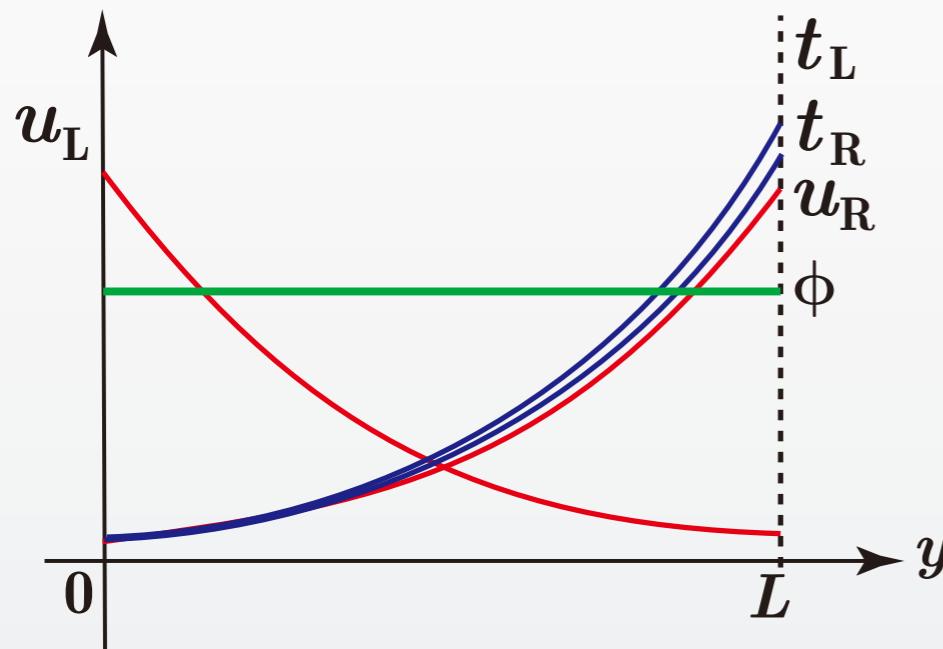
$$m_{ij} = \lambda^{(5)} \int_0^L \mathcal{F}_i^*(y) \phi(y) \mathcal{G}_j(y) dy$$

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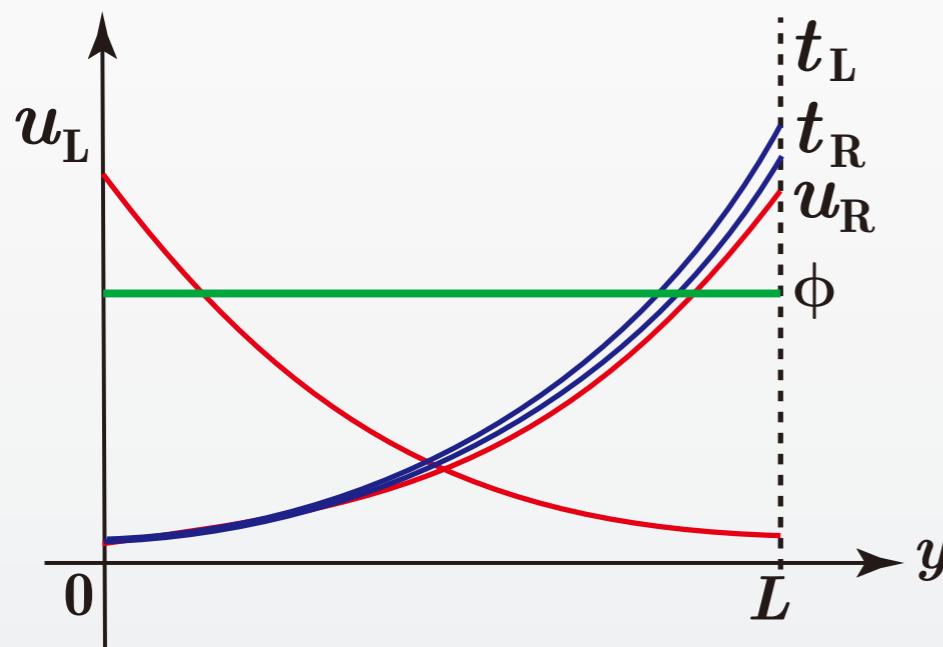


Mass hierarchy

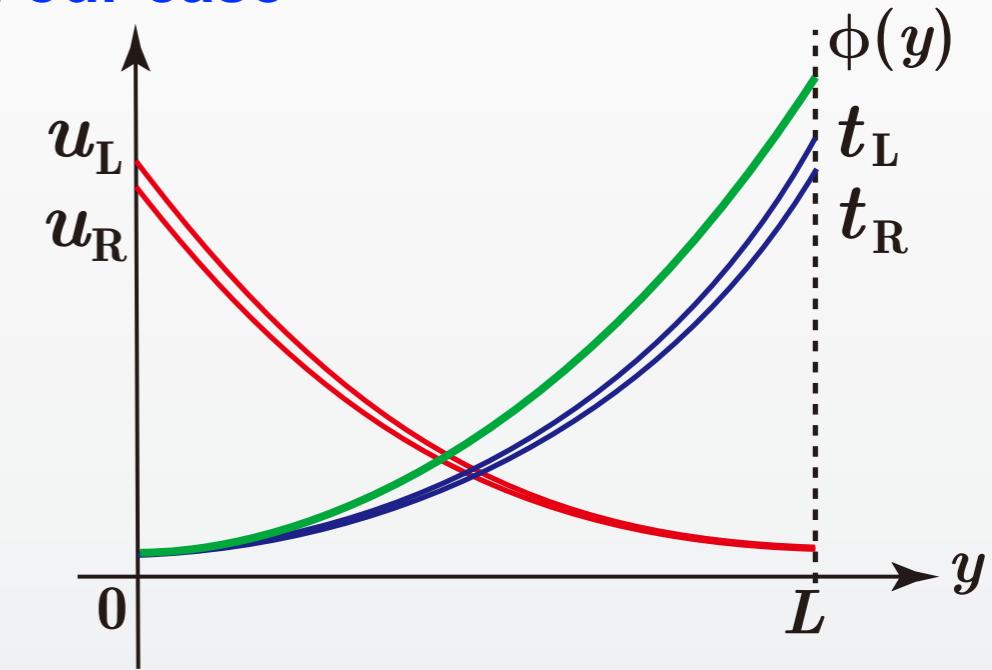
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(i) The case when VEV has no y-dep.



(ii) In our case



This can lead the **hierarchical masses**.

Conclusions and Discussions



12

5d gauge theories on an interval



The low energy effective theories

4d gauge theories

- + Symmetry breaking
- + Chiral theories
- + Mass hierarchy

Conclusions and Discussions



> Challenges for the future

5d gauge theories on an interval

$$\begin{pmatrix} E \\ \nu_E \end{pmatrix}, \quad E', \quad B_M, \quad W_M, \quad \Phi, \quad \dots \quad \dots$$

$$\downarrow \quad \quad \quad \downarrow \quad \quad \quad \downarrow \quad \quad \quad \downarrow \quad \quad \quad \downarrow$$

$$\left(\begin{array}{c} e \\ \nu_e \end{array} \right)_L \quad e_R \quad B_\mu \quad W_\mu \quad h$$

The low energy effective theories

the Standard Model

- > Mass of the Higgs is different from the SM.
- > Mass hierarchy appears naturally.



Appendix

Mass of the Higgs

> Mass of the Higgs is given by the eigenvalue of the following eigenvalue equation.

$$\left(-\frac{d^2}{dy^2} + M^2 + 3\lambda\phi^2(y) \right) J(y) = m^2 J(y)$$
$$+ \begin{cases} J(0) + L_+ \partial_y J(0) = 0 \\ J(L) - L_- \partial_y J(L) = 0 \end{cases}$$