

Search for Lepton Flavor Violating Decay at FASER

浅井健人 (埼玉大学)

素粒子物理学の進展2021

09/07/2021 @ オンライン



Based on

T. Araki, **KA**, H. Otono, T. Shimomura, Y. Takubo (ongoing work)

Summary

- FASER can search for the regions where interactions with SM particles are small, and in such regions, new physics may be hidden
- In the region where FASER has a sensitivity, charged lepton flavor violating (CLFV) couplings can be as large as charged lepton flavor conserving (CLFC) couplings
- FASER can explore unsearched region of LFV coupling by search for LFV decays

Introduction

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Charged Lepton Flavor Violation (cLFV)

In the SM

Neutrino Oscillation

➡ cLFV is violating

Prediction : $\sim 10^{-54}$

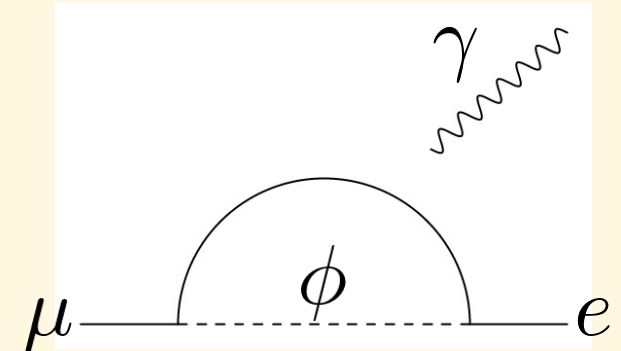
⋈

Experiments : $< 10^{-12}$

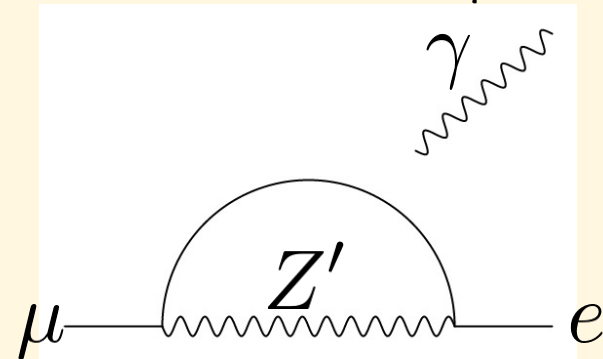
➡ It is impossible to detect cLFV process

Beyond the SM

- Leptophilic scalar



- Extra gauge boson (ex: $U(1)_{L_\mu - L_\tau}$)



- Dark Photon w/ dipole LFV coupling

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Charged Lepton Flavor Violation (cLFV)

In the SM

Neutrino Oscillation

→ cLFV is violating

Prediction : $\sim 10^{-50}$

⌗

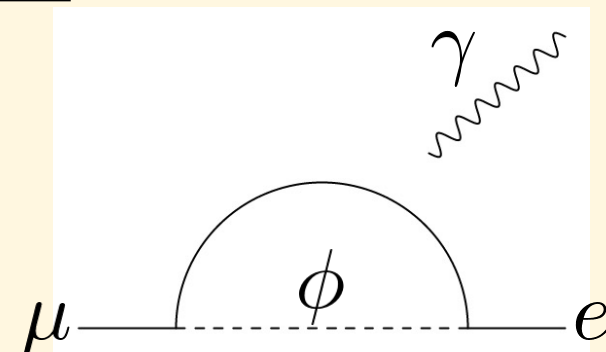
Expe

→ It is
cLFV process

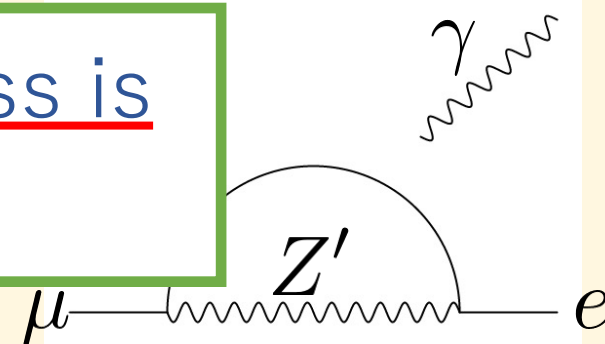
Charged lepton flavor violation process is
a smoking gun signal of new physics

Beyond the SM

- Leptophilic scalar



- Extra gauge boson (ex: $U(1)_{L_\mu - L_\tau}$)



- Dark Photon w/ dipole LFV coupling

Introduction

Origin of cLFV interaction

Ex) Leptophilic scalar boson

$$y_{\alpha\beta} \bar{\ell}_{L,\alpha} \phi \ell_{R,\beta} + \text{h.c.} \quad [\text{flavor basis}]$$



Diagonalization of charged lepton mass matrix

$$\ell_{\alpha} = V_{\alpha i} \ell_i$$

$$\underline{V_{i\alpha}^{\dagger} y_{\alpha\beta} V_{\beta j}} \bar{\ell}_{L,i} \phi \ell_{R,j} + \text{h.c.} \quad [\text{mass basis}]$$



cLFV interaction

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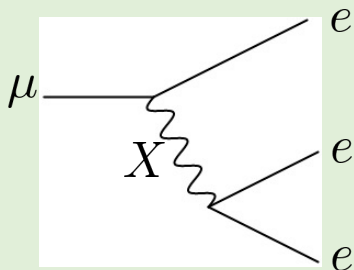
Constraints on cLFV

cLFV process

Exp. limit on BR

Future prospect

$$\mu \rightarrow eee$$



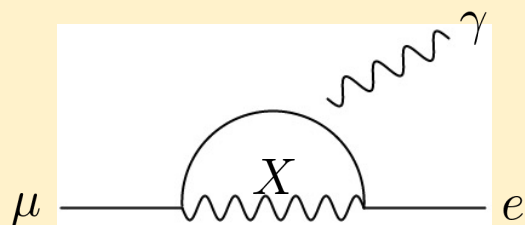
$$1.0 \times 10^{-12}$$

SINDRUM Collaboration
(1988)

$$\approx 10^{-16}$$

PDG (2020)

$$\mu \rightarrow e\gamma$$



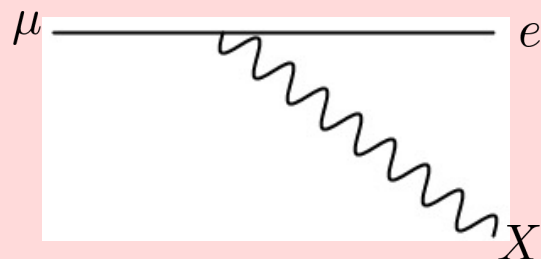
$$4.2 \times 10^{-13}$$

MEG Collaboration
(2016)

$$\approx 6 \times 10^{-14}$$

PDG (2020)

$$\mu \rightarrow eX$$



$$\approx 10^{-5}$$

TWIST Collaboration
(2015)

—————

Introduction

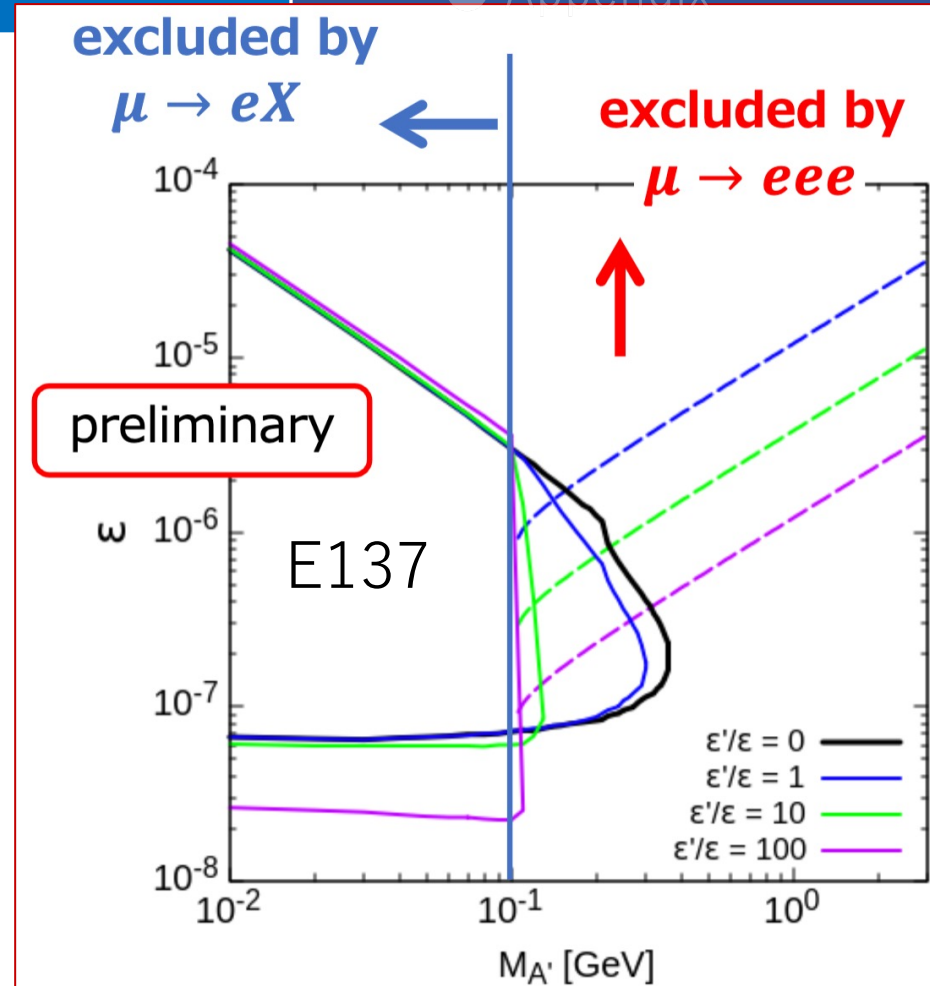
Constraints on cLFV

Ex) Dark photon model

In light mass region ($m_{A'} \sim 0.01 - 1$ GeV),
cLFV decay of muon gives upper bound
as shown in right figure



CLFV coupling can be as large as CLFC
one in such a region ($m_\phi \sim 0.01 - 1$ GeV &
 $y_e \sim 10^{-8} - 10^{-5}$)



From Araki san's slide

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Constraints on cLFV

Ex) Dark photon model

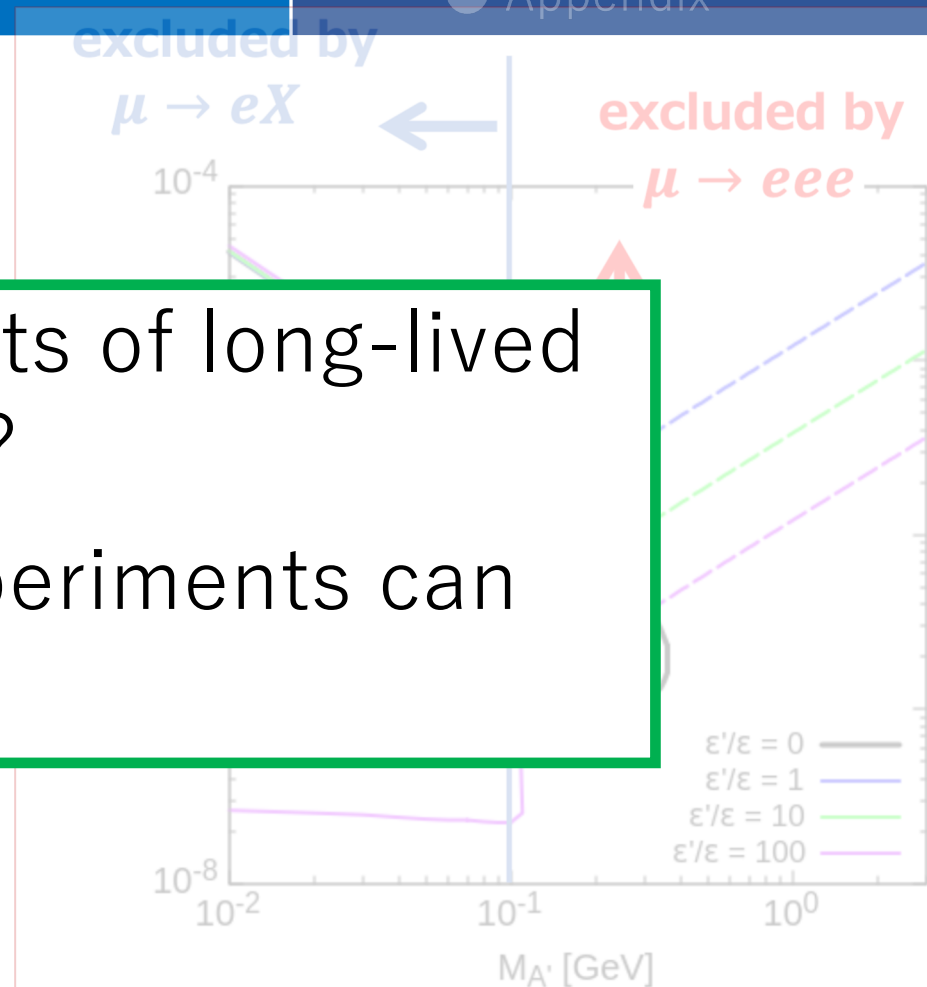
In light mass region ($m_{A'} \sim 0.01 - 1$ GeV).

cLFV
as sh

- CLFV interactions affect results of long-lived particle search experiments?
- Long-lived particle search experiments can detect CLFV decays ?

CLFV

one in such a region ($m_{\phi} \sim 0.01 - 1$ GeV & $y_e \sim 10^{-8} - 10^{-5}$)

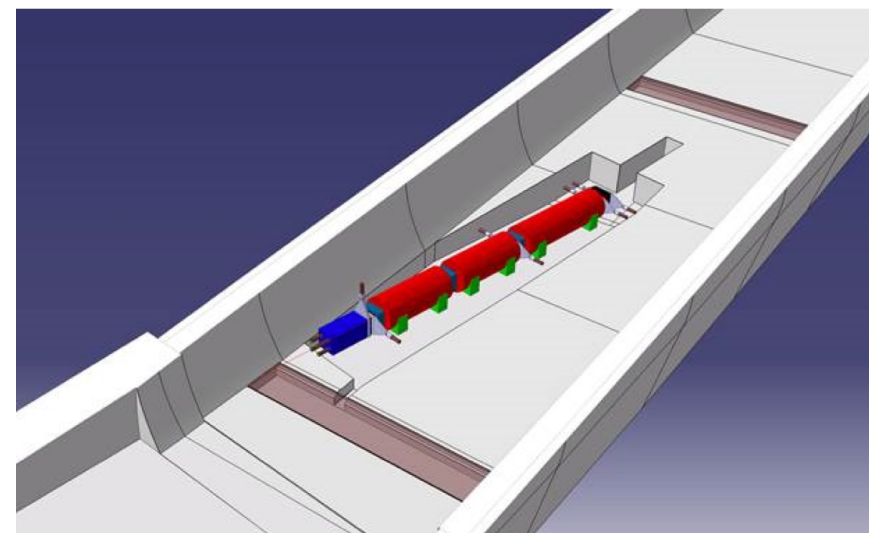
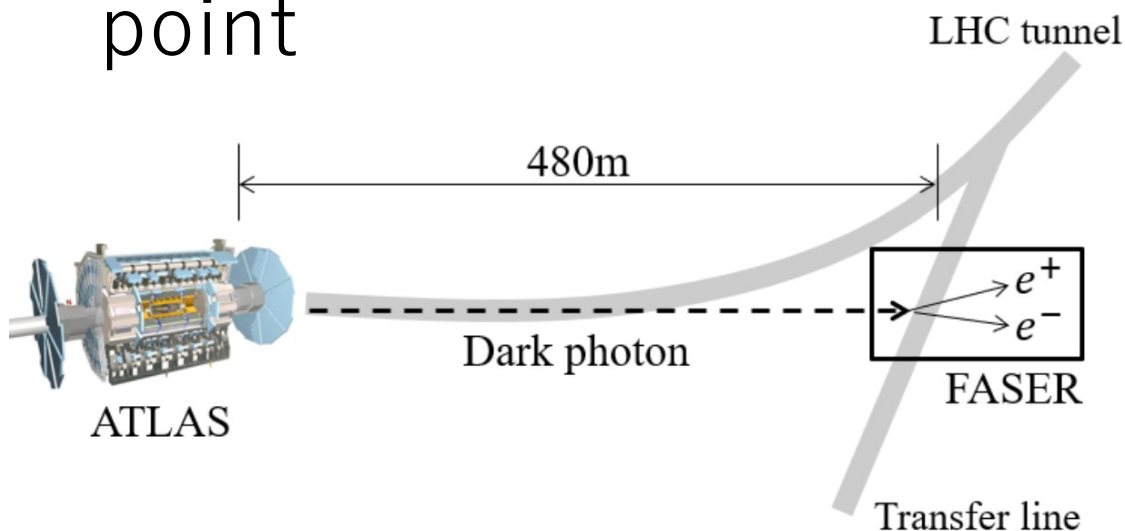


From Araki san's slide

FASER

FASER (ForwArd Search ExpeRiment at the LHC)

- Experiment to search for new light particles (dark photon, dark Higgs, ALP, sterile neutrino...), starting from 2022
- Detector will be placed at 480m from ATLAS collision point



Advantage

1) pp -reaction cross section @ LHC is very large in the direction of beam axis

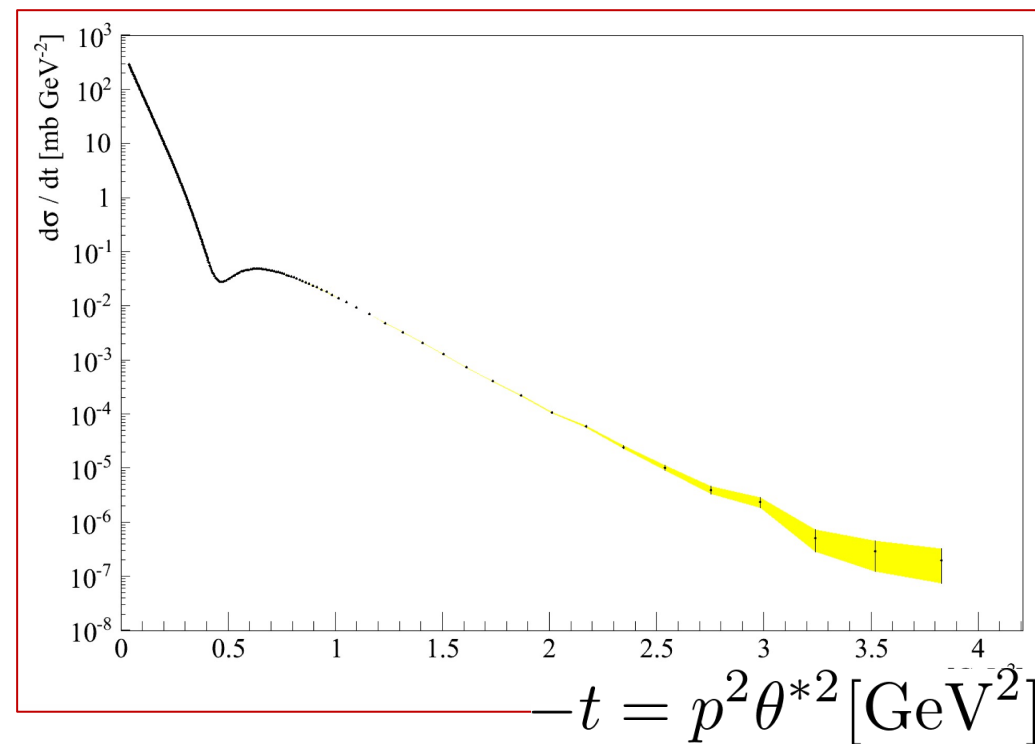
➡ Large production number of BSM particles

2) 480m from ATLAS

→ Long decay width

➡ FASER can search small interaction region

Inelastic scattering cross section of pp collision @ 13TeV LHC



[TOTEM Collaboration, EPJC 79 \(2019\) 10, 861](#)

3) FASER is placed in far-forward region

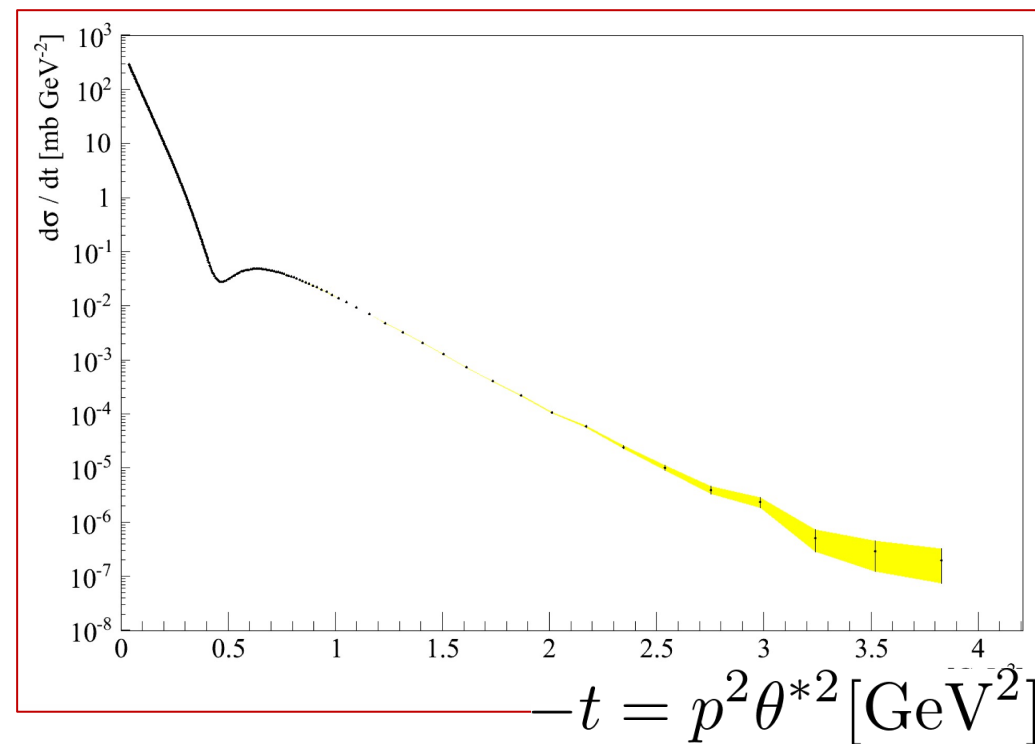
➡ FASER can search light BSM particles

cf.) LHC

New physics searches @ LHC focus on high p_T

➡ LHC can search heavy BSM particles

Inelastic scattering cross section of pp collision @ 13TeV LHC



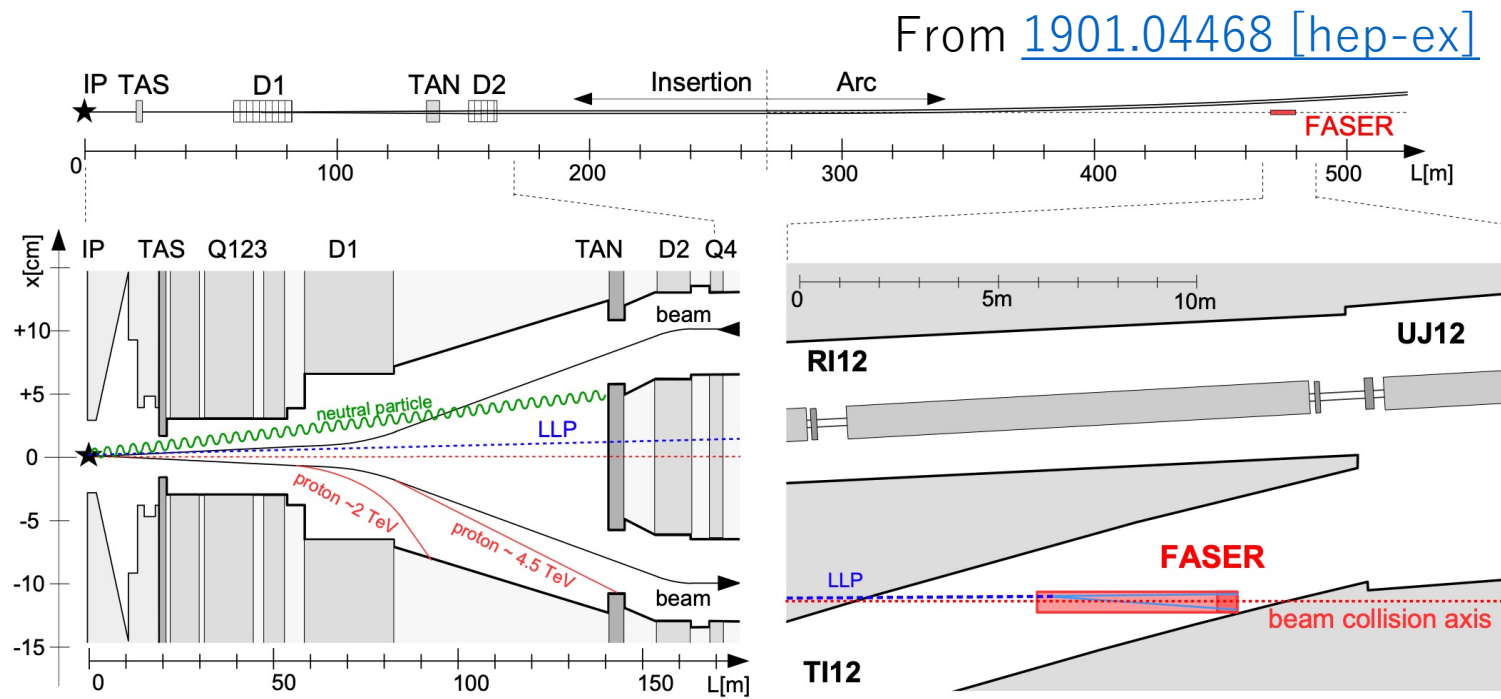
[TOTEM Collaboration, EPJC 79 \(2019\) 10, 861](#)

FASER detector will be situated along the beam collision axis line of sight

Between ATLAS IP and FASER detector, there are LHC materials and natural rock



They eliminate most potential backgrounds



Detector setup

	Depth (m)	Radius (m)
FASER	1.5	0.2
FASER2	5	1

FASER

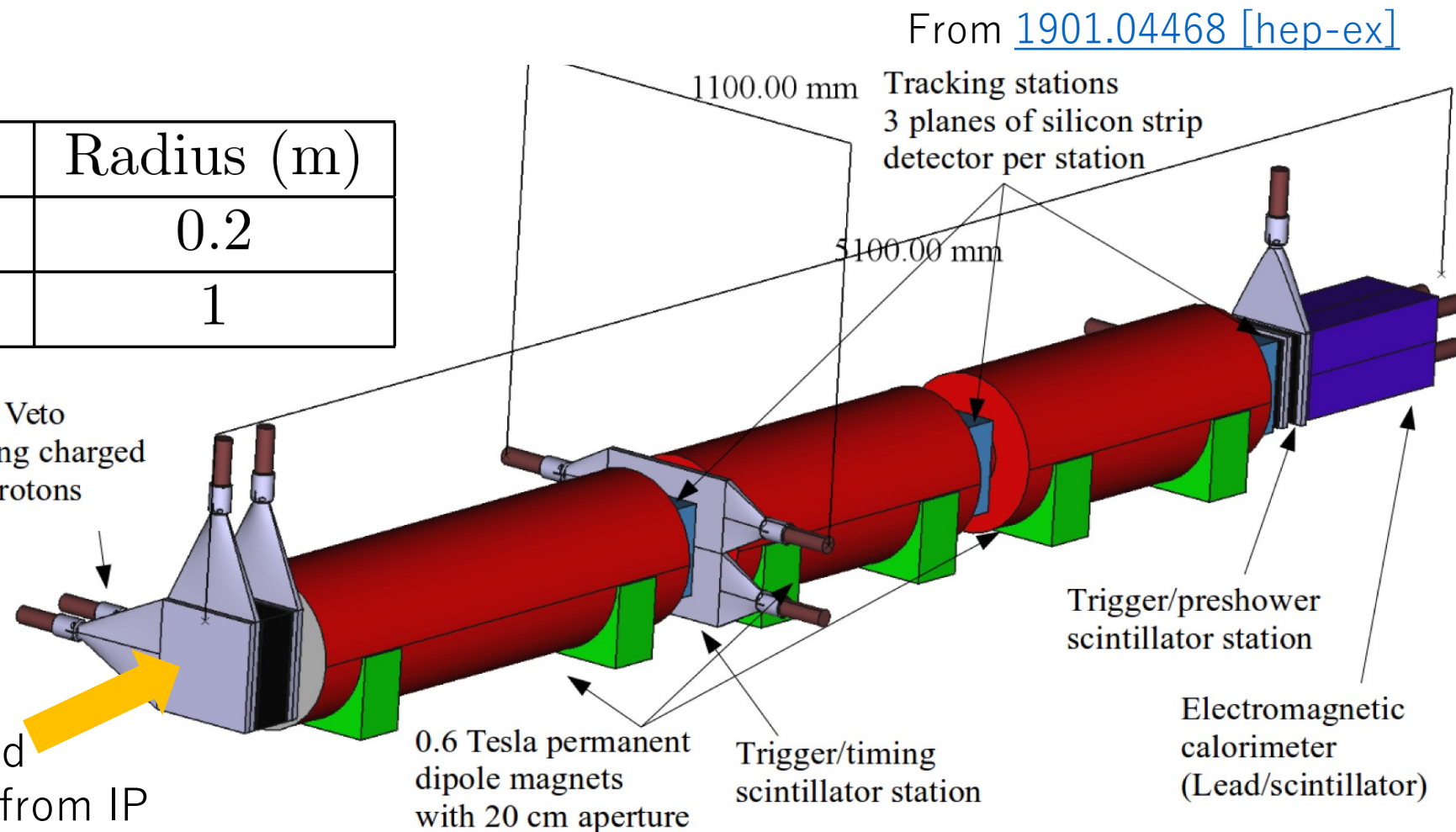
→ 150 fb^{-1}
@ LHC Run3

FASER2

→ 3 ab^{-1}
@ HL-LHC

Scintillator/Pb Veto
to veto incoming charged
particles and protons

Long-lived
particles from IP



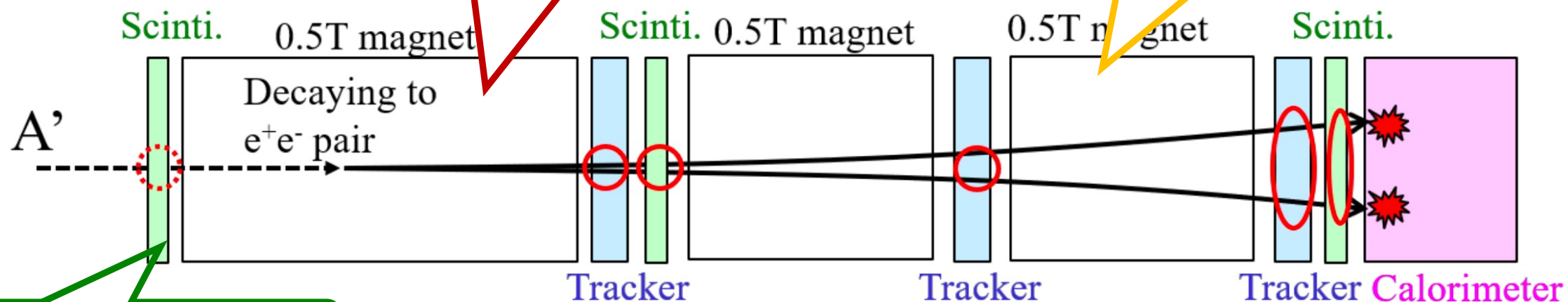
From [1901.04468 \[hep-ex\]](#)

Detector

Detector setup

Spectrometer : separating daughter particles by magnets

1.5m decay volume

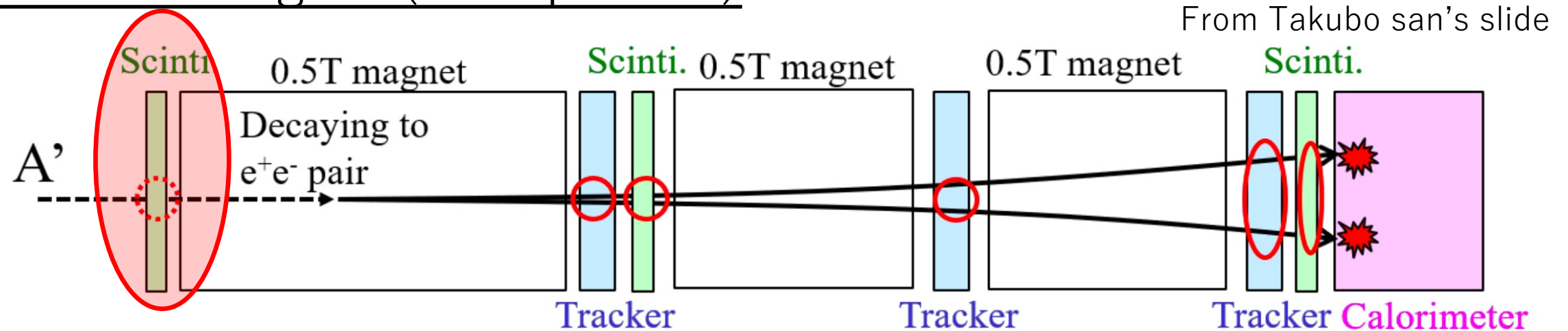


Scintillator veto

Calorimeter : measurement of energies of daughter particles

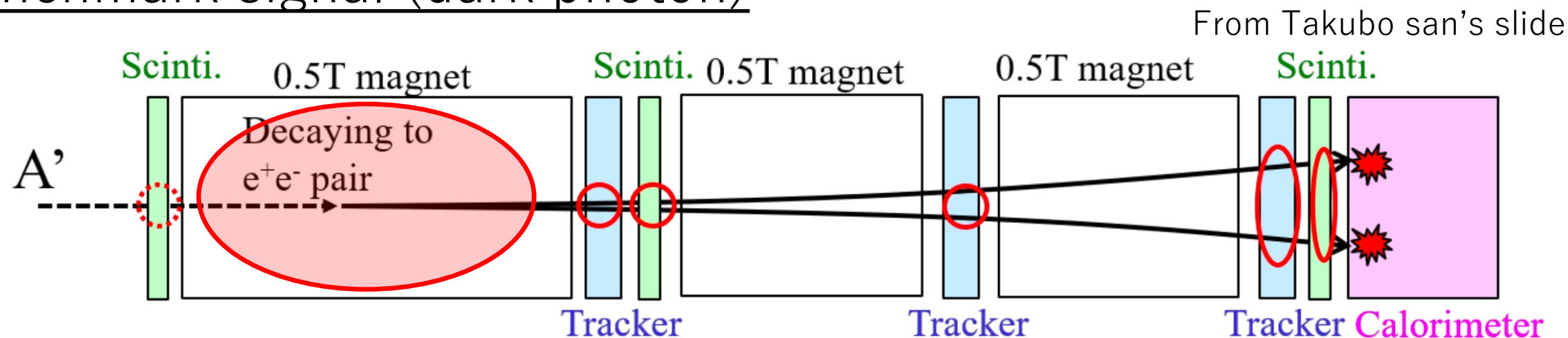
Tracker : Measurement of two oppositely charged particle and their momenta

Benchmark signal (dark photon)



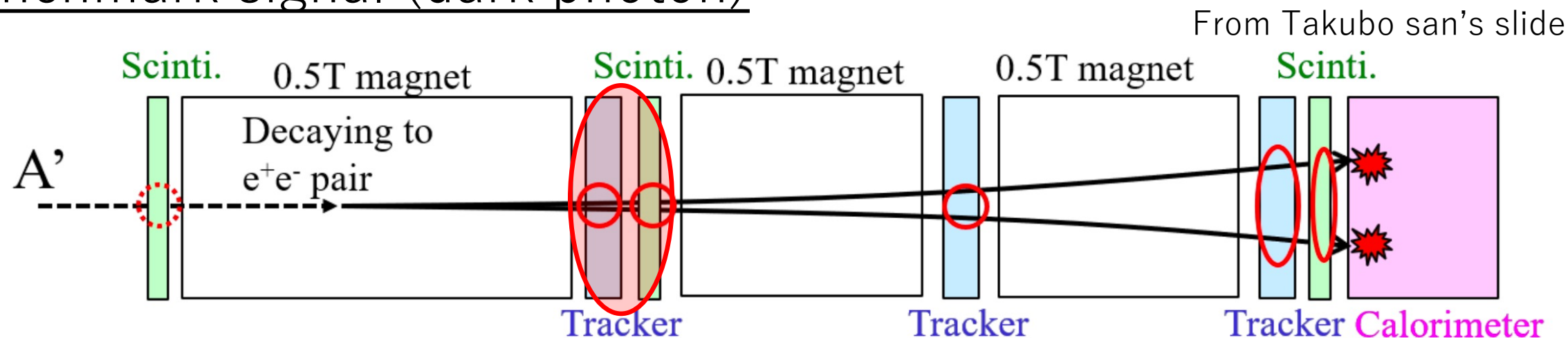
1, No signal in veto scintillator

Benchmark signal (dark photon)



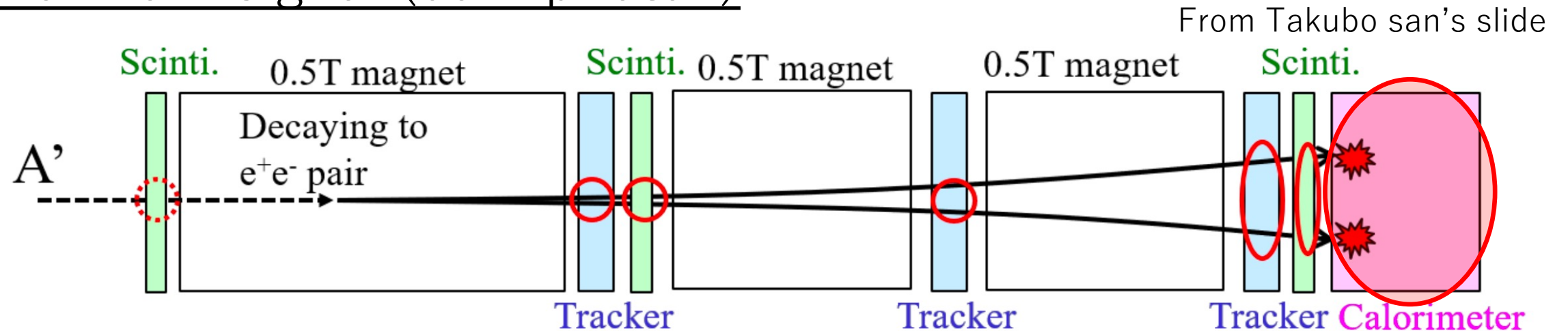
- 1, No signal in veto scintillator
- 2, Decay into e^+ and e^-

Benchmark signal (dark photon)



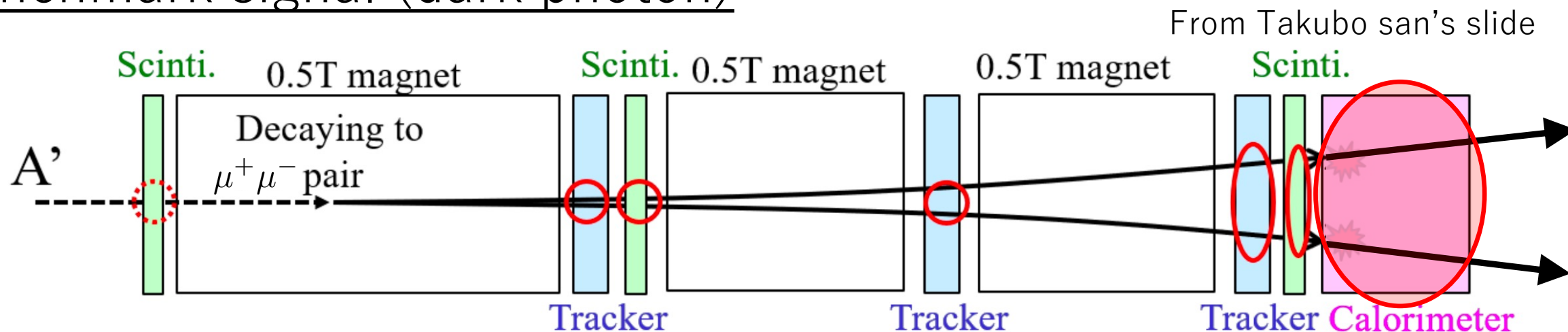
- 1, No signal in veto scintillator
- 2, Decay into e^+ and e^-
- 3, Detection of two high energy charged tracks that emanate from a common vertex

Benchmark signal (dark photon)



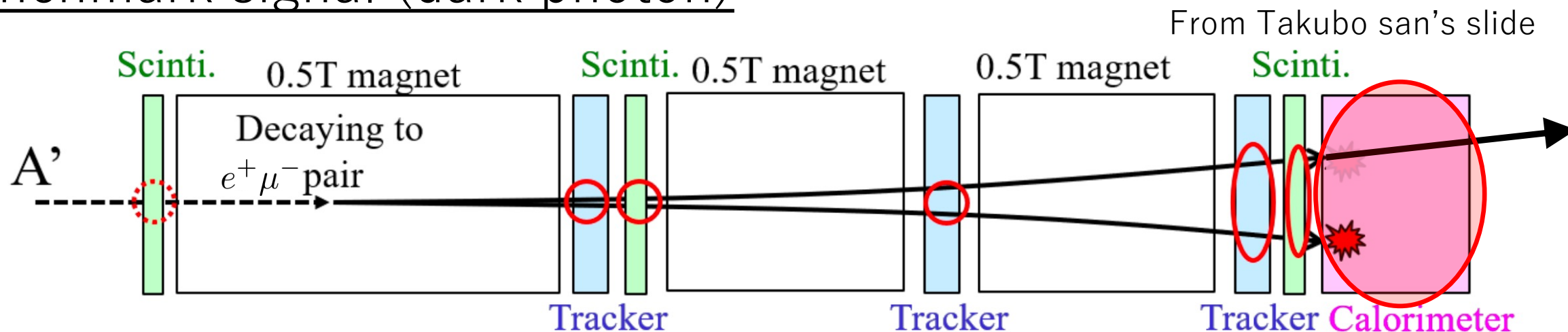
- 1, No signal in veto scintillator
- 2, Decay into e^+ and e^-
- 3, Detection of two high energy charged tracks that emanate from a common vertex
- 4, Measurement of large EM energy in calorimeter

Benchmark signal (dark photon)



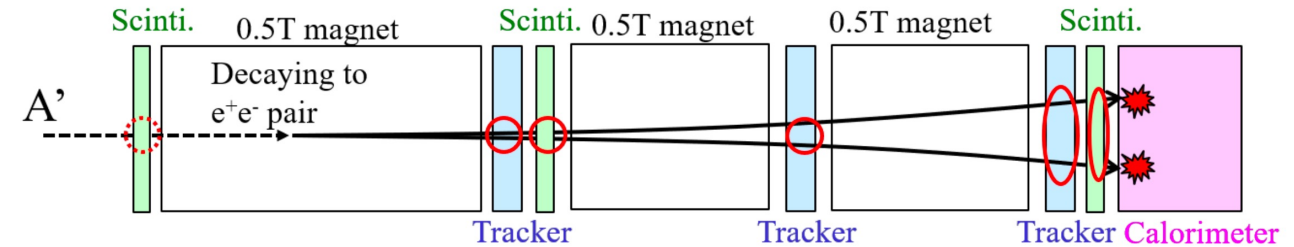
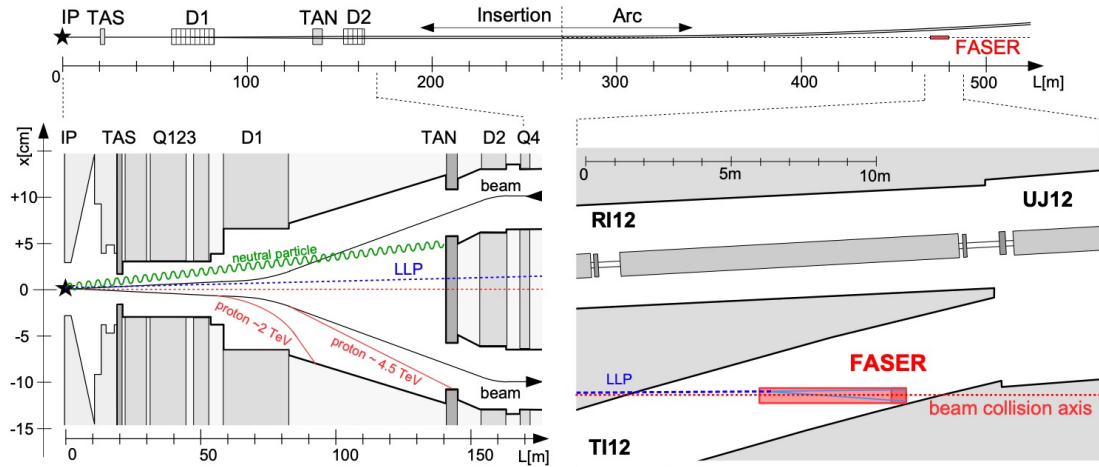
- 1, No signal in veto scintillator
- 2, Decay into μ and μ^-
- 3, Detection of two high energy charged tracks that emanate from a common vertex
- 4, Muons pass through calorimeter as minimum ionizing particles (MIP)

Benchmark signal (dark photon)



- 1, No signal in veto scintillator
- 2, Decay into e^+ and μ^-
- 3, Detection of two high energy charged tracks that emanate from a common vertex
- 4, Muon passes through calorimeter as minimum ionizing particle (MIP), and electron leaves a large EM energy in calorimeter

Background @ FASER



Almost all potential backgrounds are eliminated by rock and LHC materials

Main background

- EM/HD shower from muon bremsstrahlung : 8×10^3 events @ LHC Run3
- CC/NC neutrino events : A few (~ 100 GeV) events @ LHC Run3

➡ Can be eliminated by charged particle veto with efficiency of 99.99%

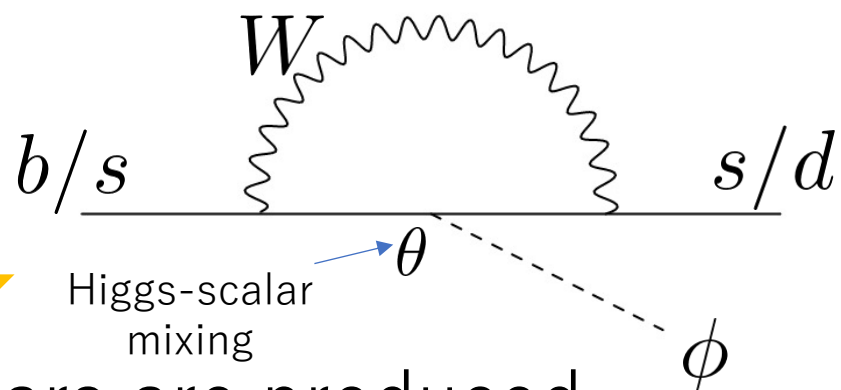
Production of BSM particles

Production @ ATLAS

B and K mesons are produced at ATLAS IP

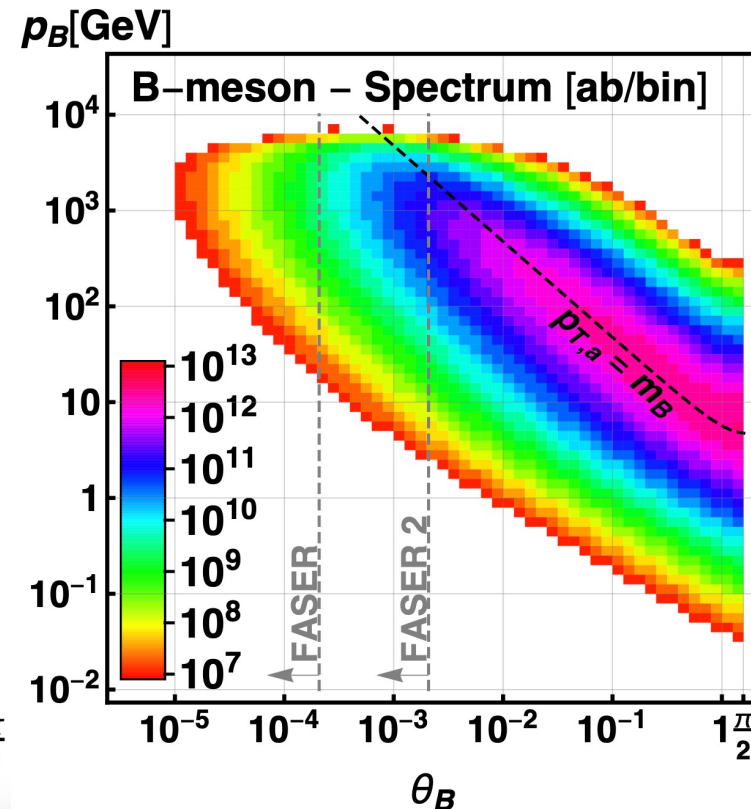
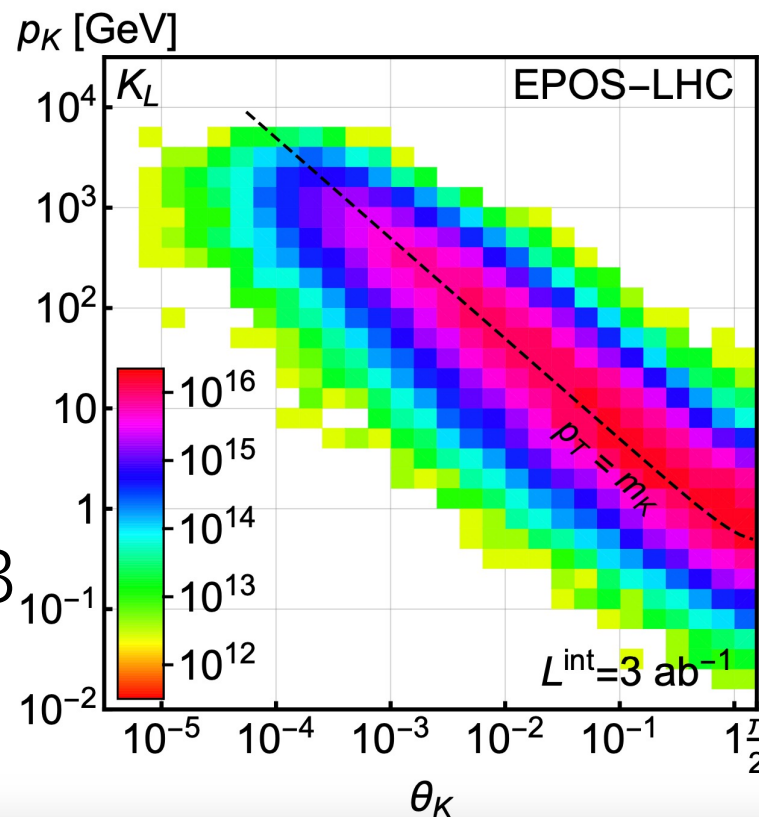
[J. L. Feng, I. Galon, F. Kling, S. Trojanowski, PRD 97 \(2018\) 5, 055034](#) ;

[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)

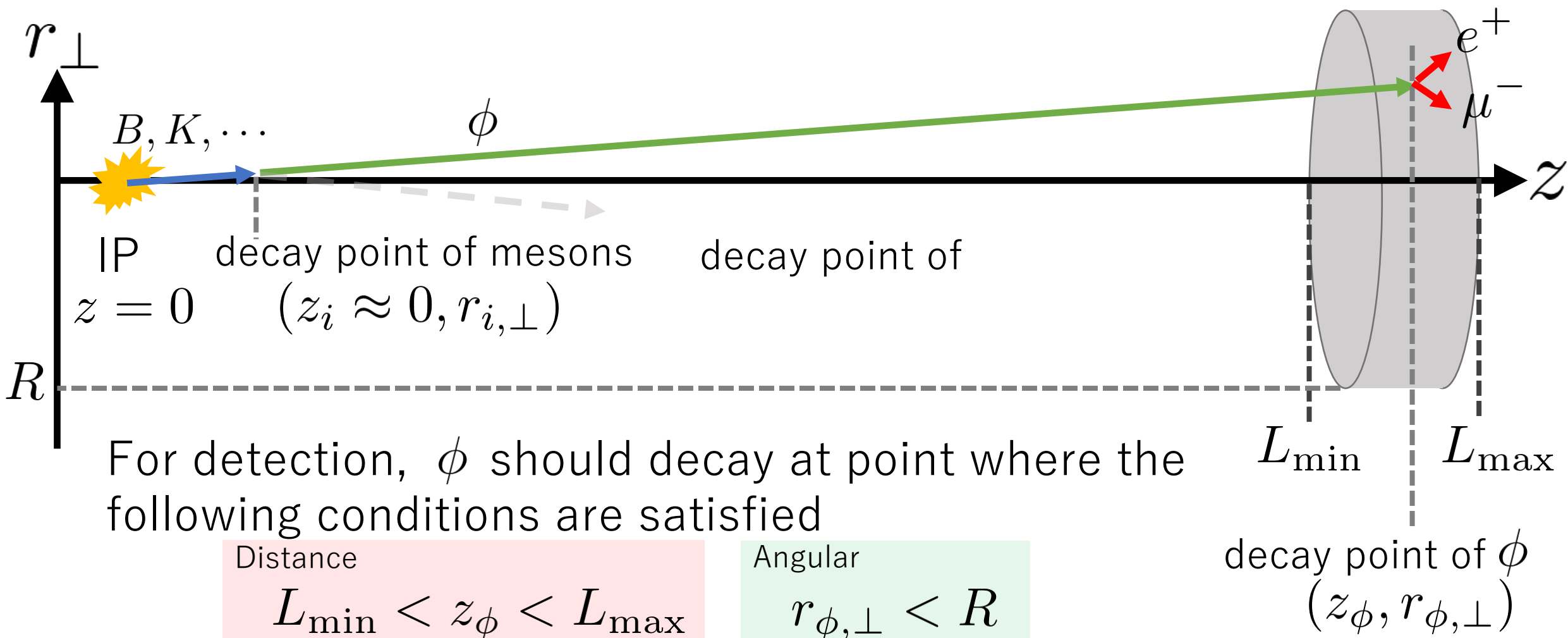


Scalars are produced by rare decays of K and B

$$\text{Br}(B \rightarrow X_s \phi) \simeq 5.7 \left(1 - \frac{m_\phi^2}{m_b^2}\right)^2 \theta^2$$



Geometrical acceptance



Geometrical acceptance

Probability of dark photon decay in detector

$$\mathcal{P}_{\phi}^{\text{det}}(\mathbf{p}_{\phi}) = \underbrace{\frac{1}{\bar{d}_{\phi} \cos \theta_{\phi}} \int_{z_{\phi, \min}}^{z_{\phi, \max}} dz_{\phi} e^{-\frac{z_{\phi}}{\bar{d}_{\phi} \cos \theta_{\phi}}}}_{\text{Distance}} \times \underbrace{\Theta(R - r_{\phi, R}) \Theta(R - r_{\phi, F})}_{\text{Angular}}$$

$\bar{d}_{\phi}$: decay length of ϕ in lab frame $\Theta(x)$: Heaviside step function

θ_{ϕ} : angular of ϕ with respect of beam axis (z-axis)

$r_{\phi, F}(R)$: distance of ϕ from beam axis at $z_{\phi} = L_{\min} (L_{\max})$

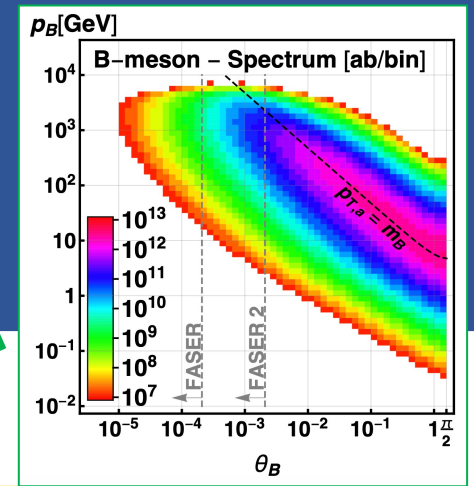
FASER

Number of events

Number of dark photons which decay in detector

$$\begin{aligned} N_{\text{cLFV}} &= \mathcal{L} \sum_{i:\text{meson}} \int dp_i d\theta_i \int d\mathbf{p}_\phi \frac{d\sigma_{pp \rightarrow iX}}{dp_i d\theta_i} \text{Br}(i \rightarrow X\phi) \text{Br}(\phi \rightarrow e\mu) \\ &\quad \times \mathcal{P}_\phi^{\text{det}}(\mathbf{p}_\phi) \\ &\sim 150 \text{ fb}^{-1} \times 10^{11} \text{ ab} \times 10^{-7} \times \mathcal{O}(0.1) \times 10^{-2} \\ &\sim \mathcal{O}(1) \end{aligned}$$

for FASER (LHC Run3)



Result

Result

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Origin of cLFV interaction

Ex) Leptophilic scalar

$$\phi \bar{e}_{L,i} \begin{pmatrix} 0 & \lambda & 0 \\ \lambda & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} e_{R,j} + h.c. \quad [\text{mass basis}]$$

cLFV interaction

SM Higgs - ϕ mixing : θ



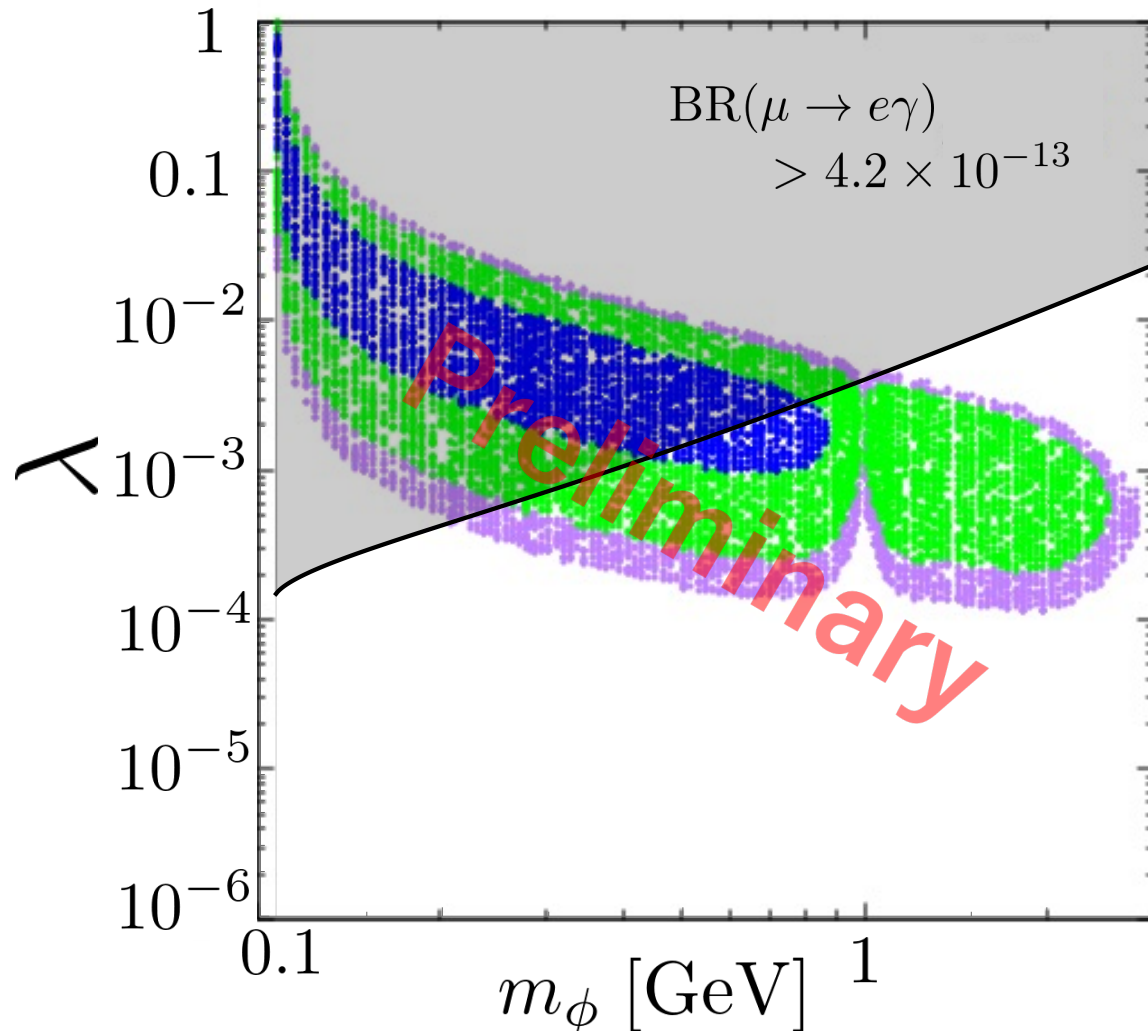
Couplings with the other SM fermions

$$\theta \sum_f \frac{m_f}{v} \phi \bar{f} f$$

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of signal events (Leptophilic scalar)



Parameters

$$\theta = 5 \times 10^{-5}$$

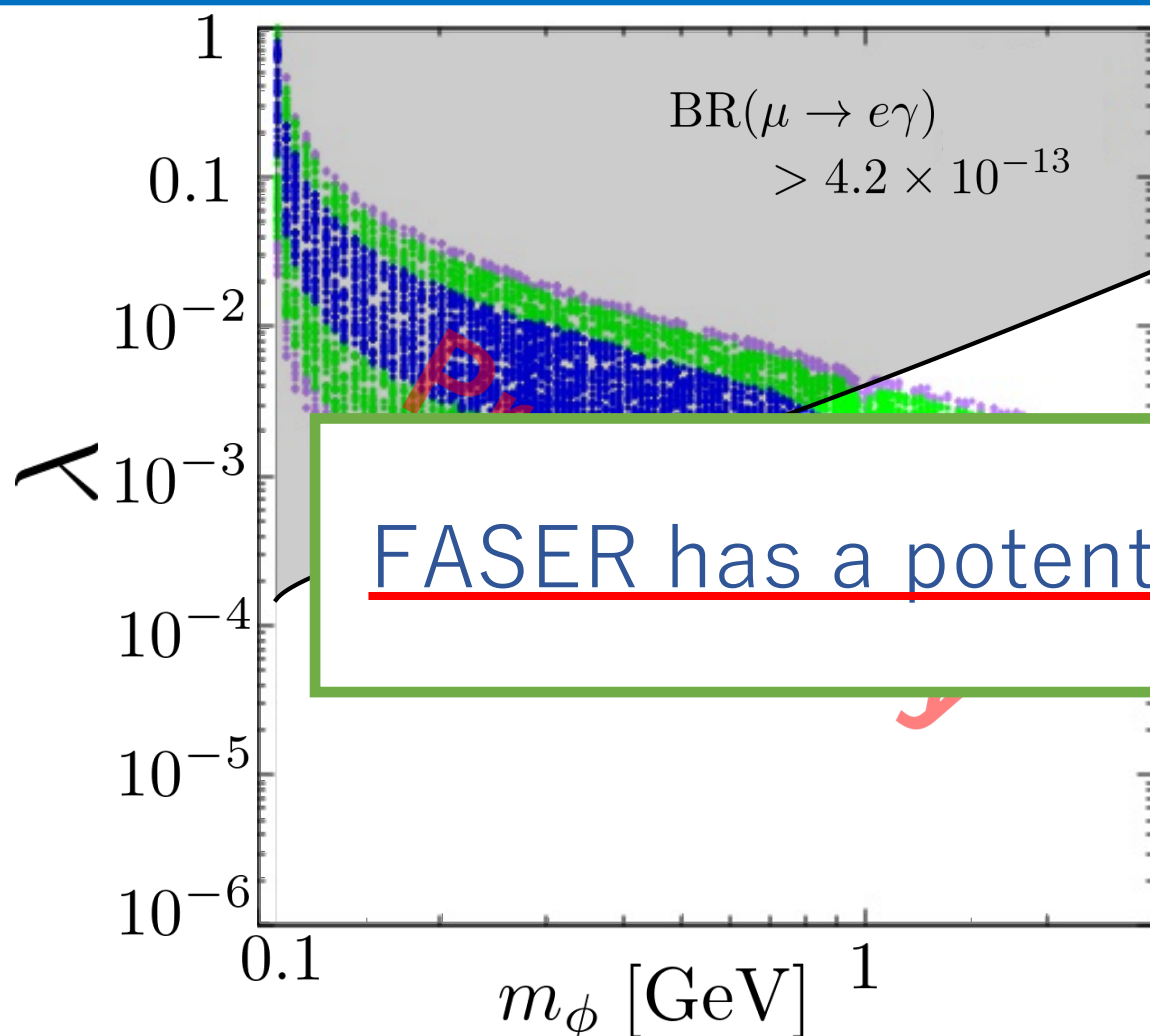
of LFV decay signals

$3 \sim 9$	●
$10 \sim 99$	●
$100 \leq$	●

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of signal events (Leptophilic scalar)



Parameters

$$\theta = 5 \times 10^{-5}$$

of LFV decay signals



Summary

- FASER can search for the regions where interactions with SM particles are small, and in such regions, new physics may be hidden
- In the region where FASER has a sensitivity, charged lepton flavor violating (CLFV) couplings can be as large as charged lepton flavor conserving (CLFC) couplings
- FASER can explore unsearched region of LFV coupling by search for LFV decays

Appendix

Schedule of FASER

FASER実験の歴史と今後の予定

- 2018年8月、LOI (Letter Of Intent)をLHC委員会に提出 ([arXiv:1811.10243](#)).
- 2018年11月、TP (Technical Proposal)をLHC委員会に提出 ([arXiv:1812.09139](#))
- 2019年3月、CERNに公式に承認された
- 検出器の建設と運転費用は、Simons財団とHeising-Simons財団が提供
- 2020年秋に検出器を実験サイトに設置する予定
- 2021年はコミッショニング作業
- データ取得はLHC Run3が始まる2022年に開始

From Takubo san's slide

Appendix

Spectrometer

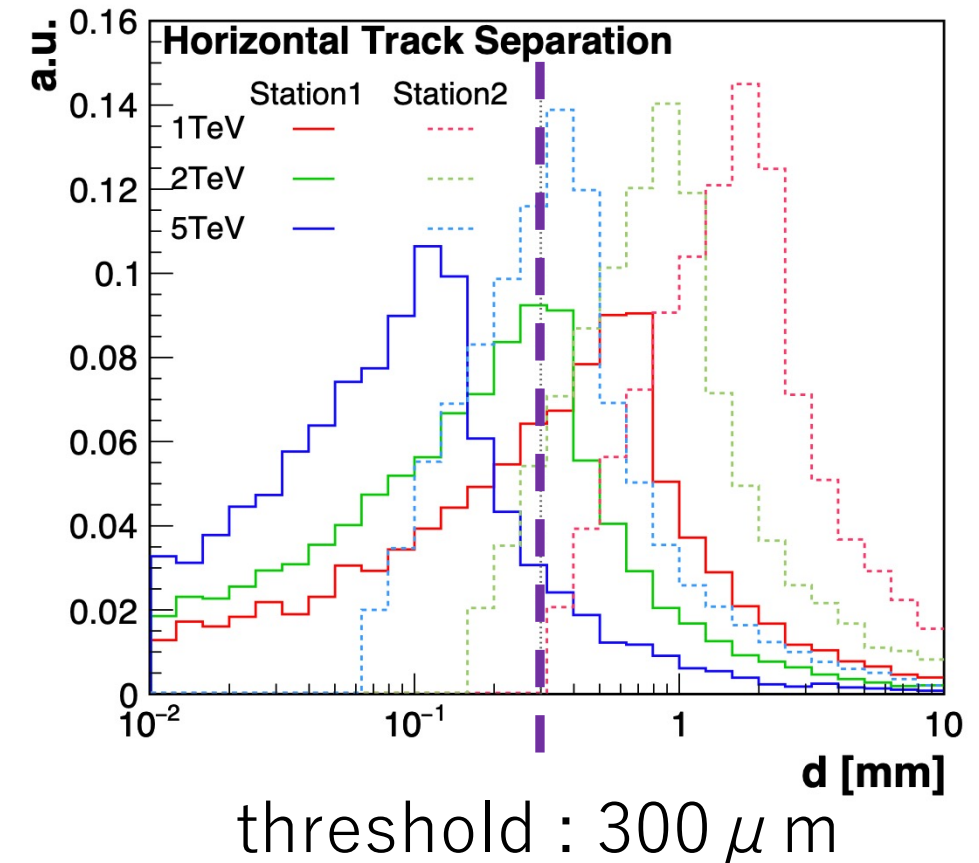
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Ex) e^+e^- pair from decay of 100 MeV dark photon are separated by 0.5T magnet



Can identify two tracks

From [1811.10243 \[hep-ex\]](#)



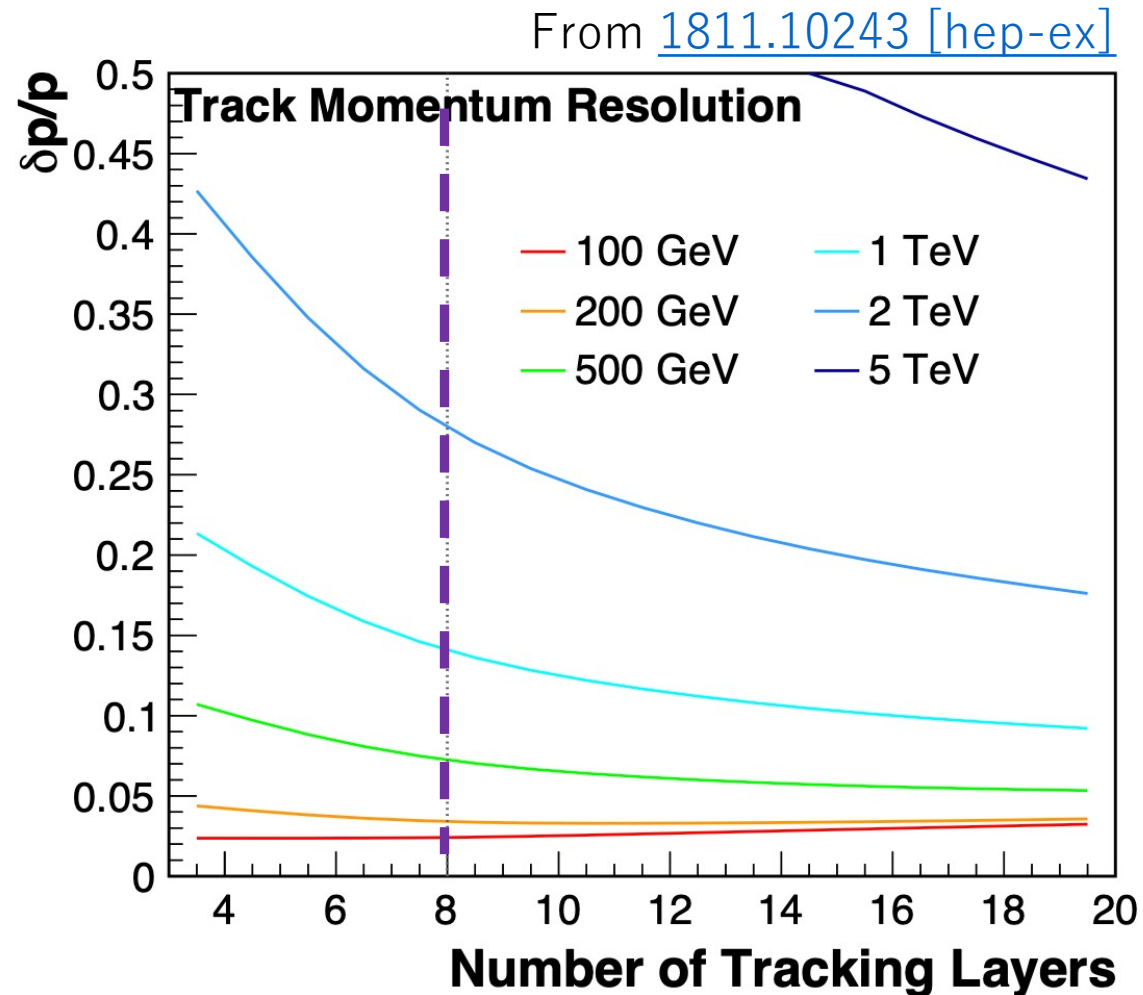
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Spectrometer

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Momentum resolution

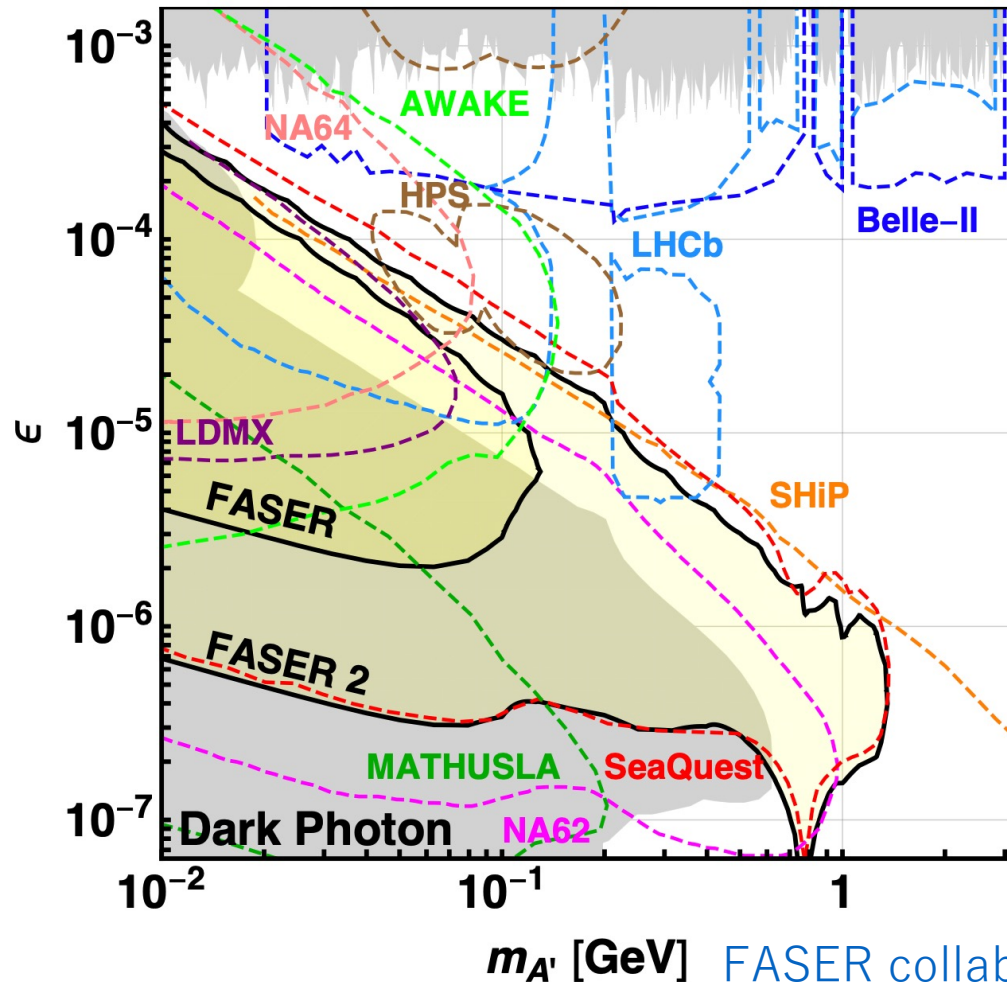
The lower energy e^+e^- have,
the better momentum resolution is.



Appendix

Expected sensitivity

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Dark photon

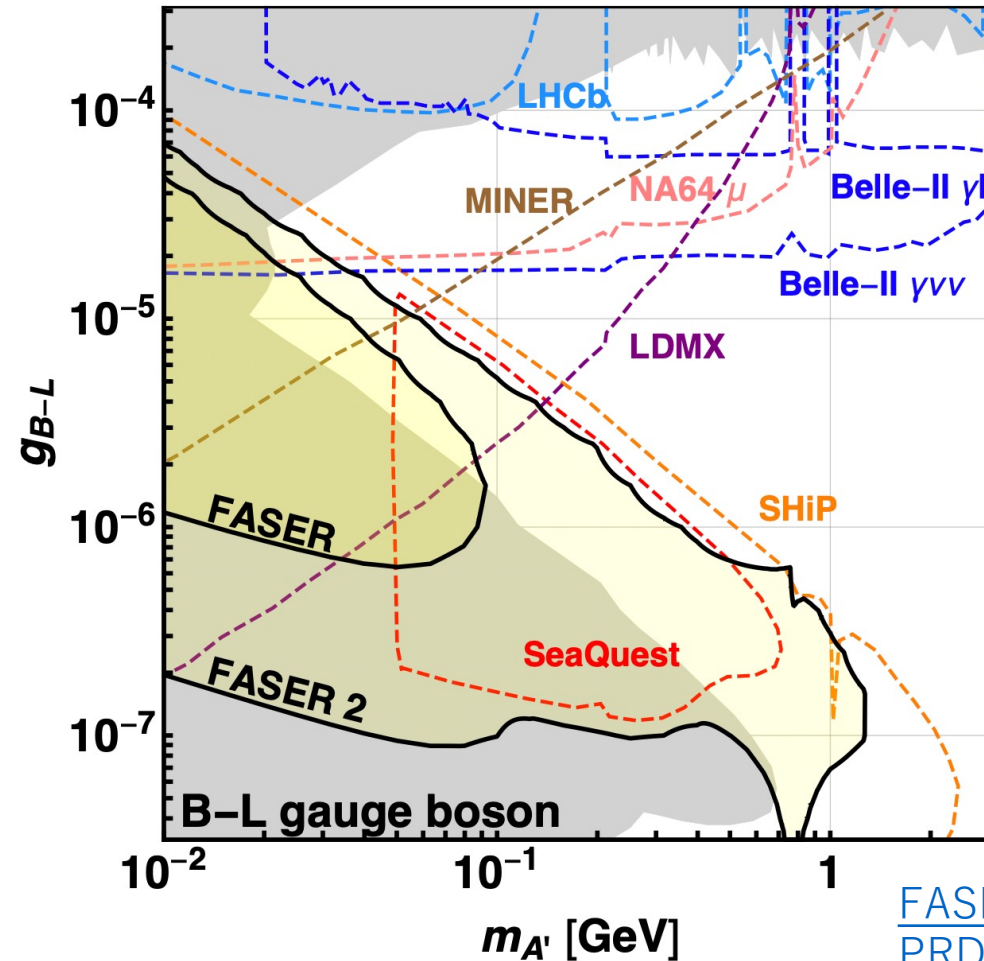
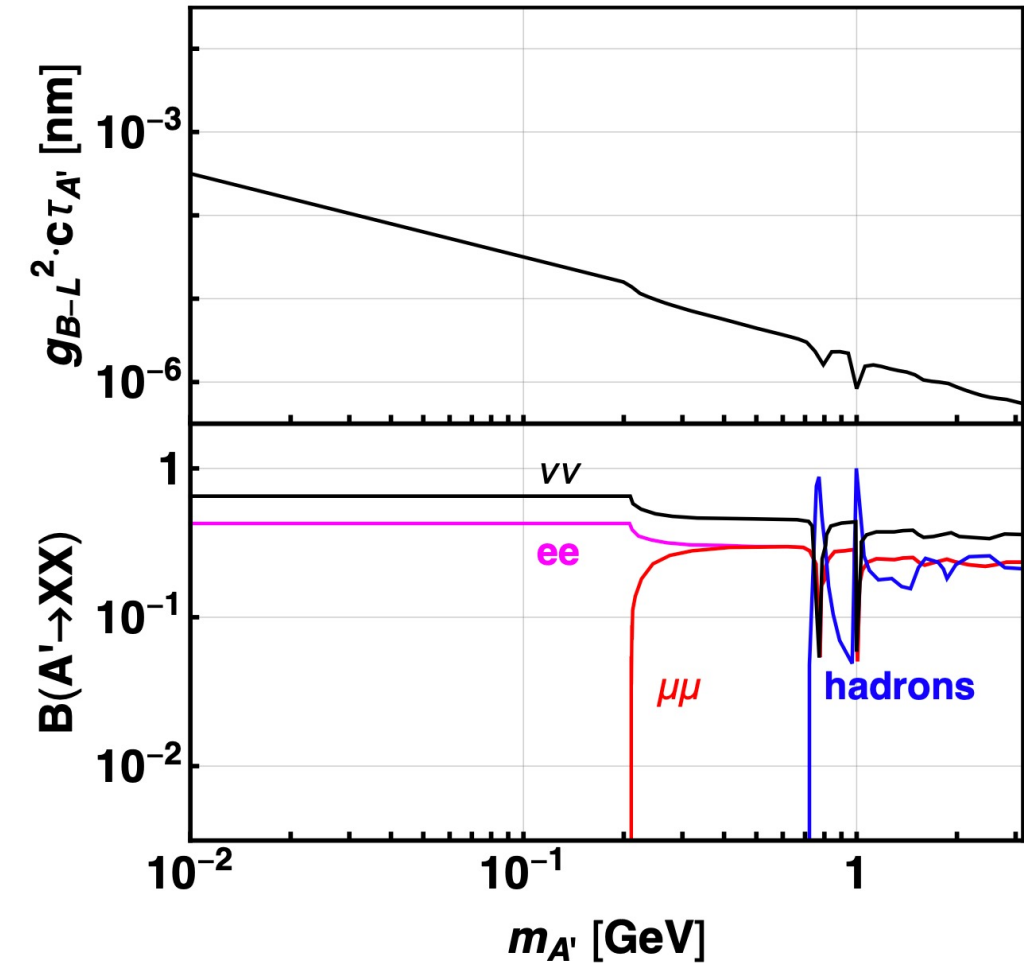
FASER(2) can explore parameter region with $\epsilon \sim 10^{-5} - 10^{-4}$ ($10^{-7} - 10^{-4}$)

[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)

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Expected sensitivity ($U(1)_{B-L}$ gauge boson)

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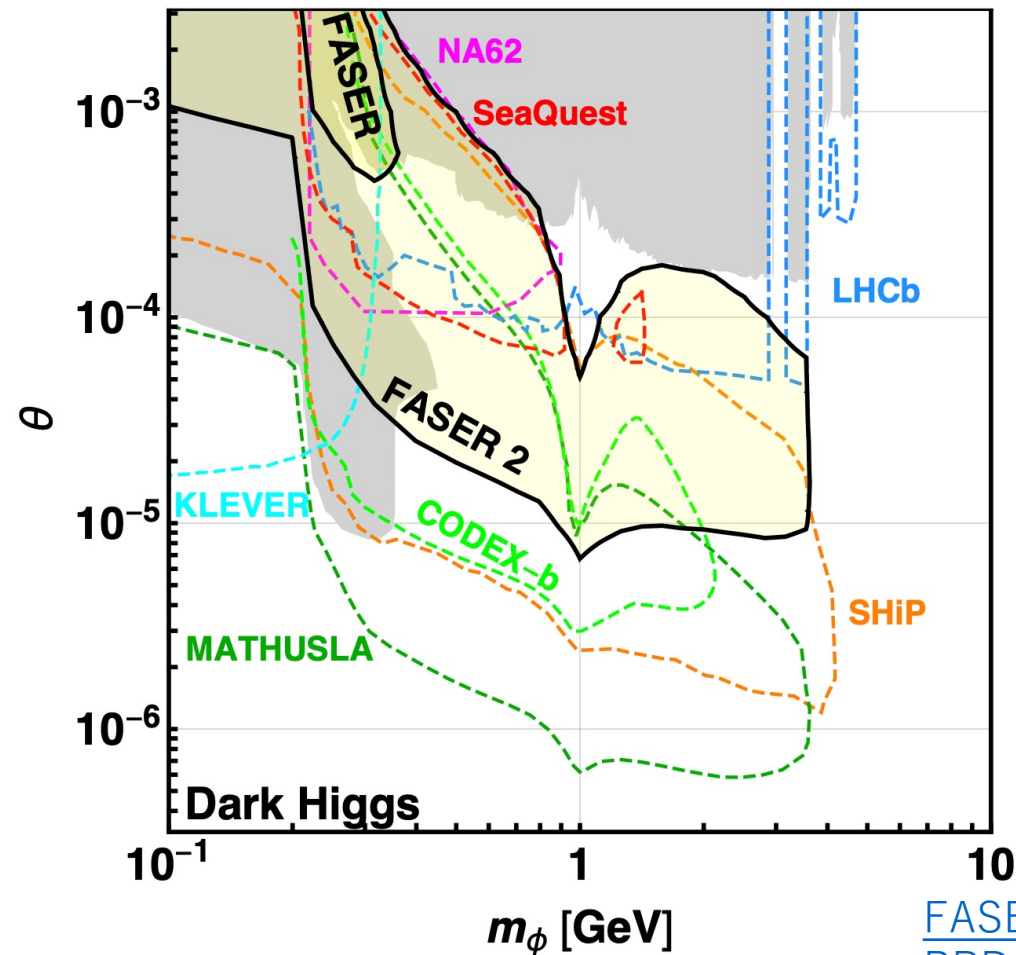
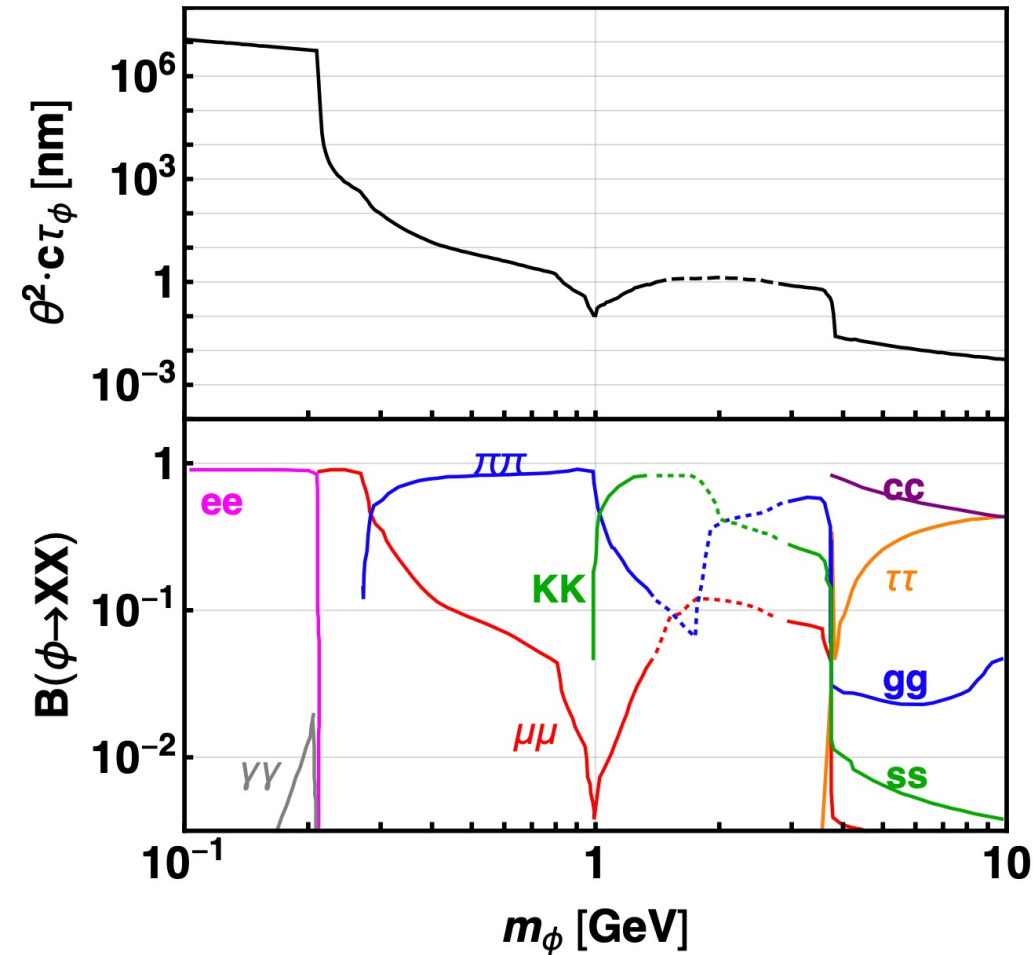


[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)

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Expected sensitivity (Dark Higgs boson)

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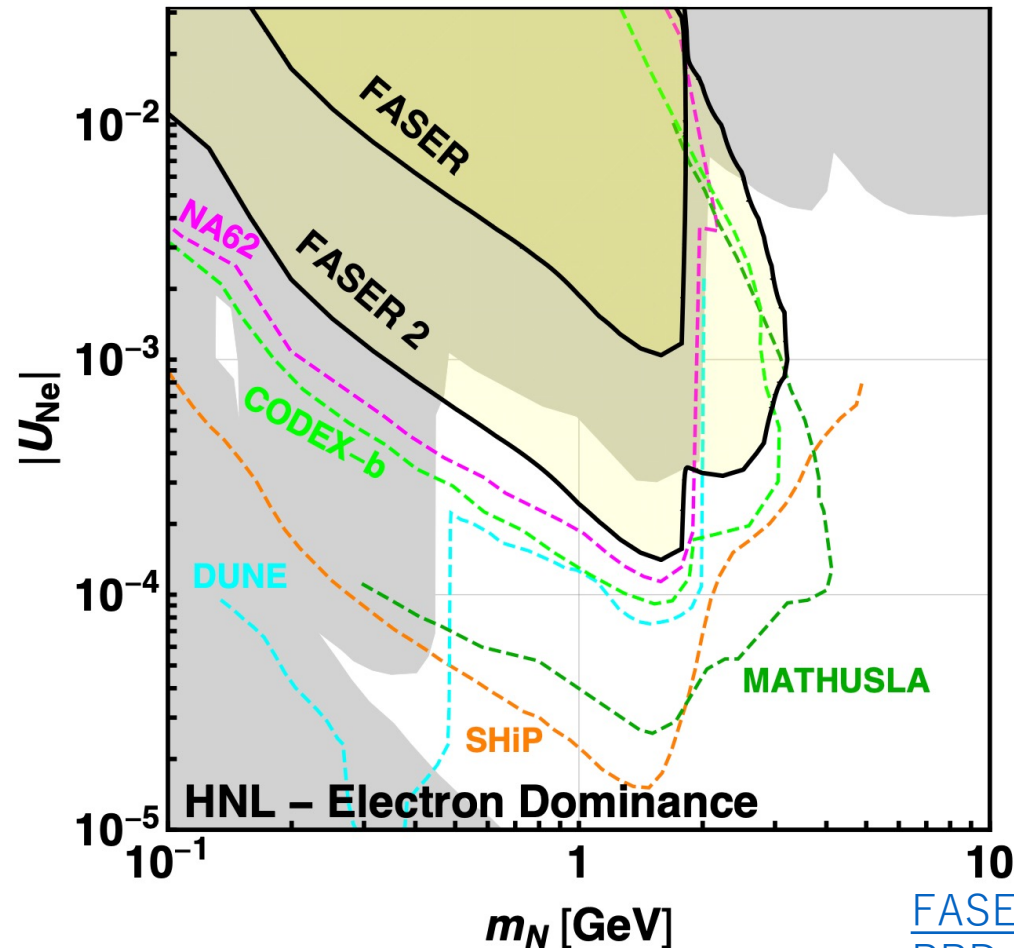
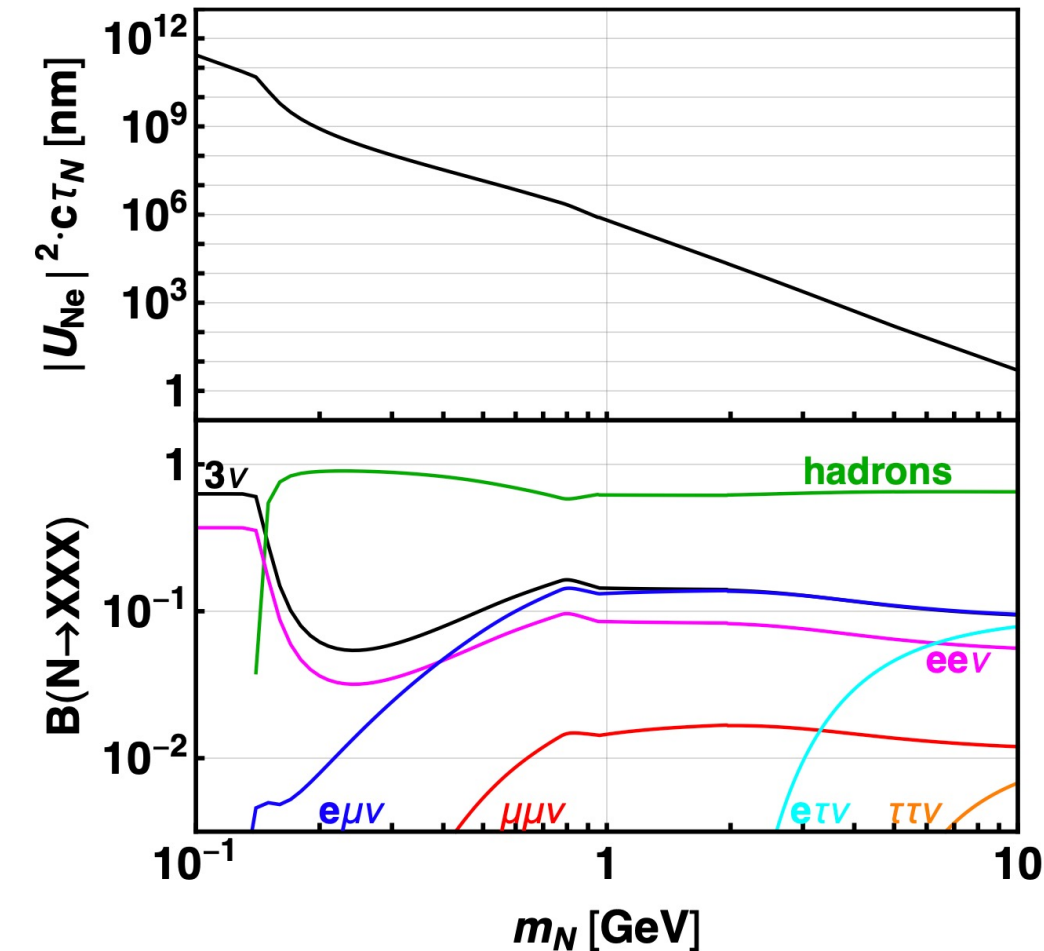


[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)

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Expected sensitivity (Heavy neutral lepton)

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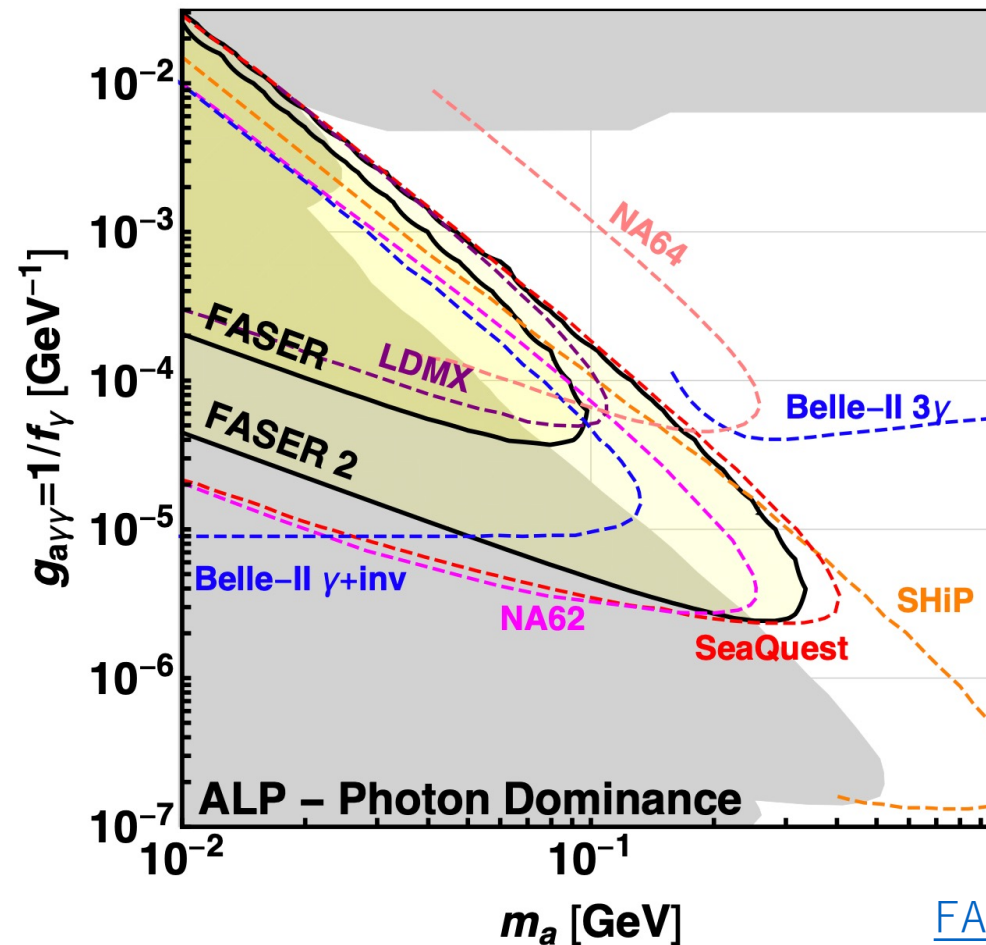
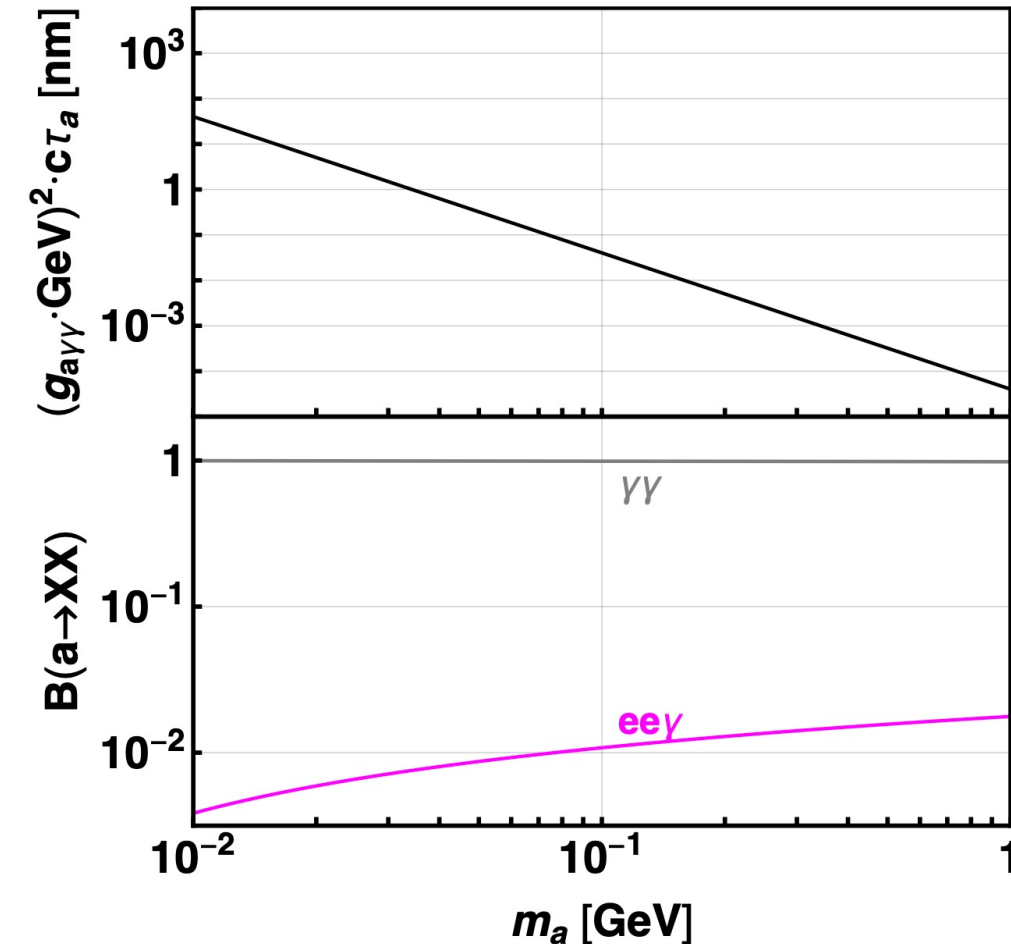


[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)

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Expected sensitivity (Axion-like particle)

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[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)

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Origin of cLFV interaction

Ex) $U(1)_{L_\mu - L_\tau}$ gauge boson

Non-universal charge

$$g_{\mu\tau} Z'_\mu \bar{\ell}_\alpha \gamma^\mu \begin{pmatrix} 0 & 0 & 0 \\ 0 & +1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \ell_\beta \quad [\text{flavor basis}]$$

Diagonalization of charged lepton mass matrix

$$\ell_\alpha = V_{\alpha i} \ell_i$$

$$g_{\mu\tau} Z'_\mu \bar{\ell}_i \gamma^\mu V_{i\alpha}^\dagger \begin{pmatrix} 0 & 0 & 0 \\ 0 & +1 & 0 \\ 0 & 0 & -1 \end{pmatrix} V_{\beta j} \ell_j \quad [\text{mass basis}]$$

cLFV interaction

Dark Photon Model

Extension of the Standard Model by a dark U(1) gauge symmetry



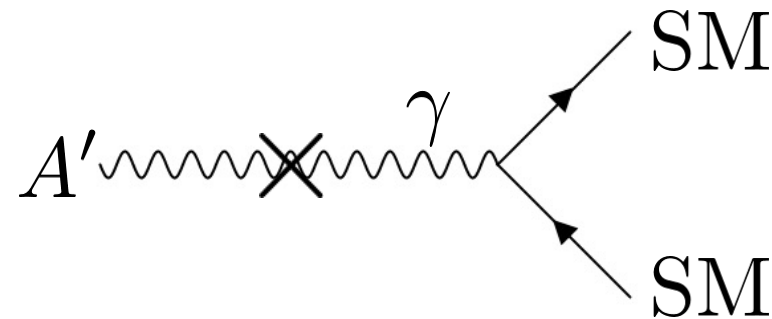
$$SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_{\text{dark}}$$

Mixing between SM photon and dark photon appears through gauge kinetic mixing

$$\mathcal{L}_{\text{mix}} = -\frac{\epsilon}{2} B_{\mu\nu} X^{\mu\nu}$$



$$\mathcal{L}_{\text{int}} = \epsilon e A'_\mu J_{\text{EM}}^\mu$$



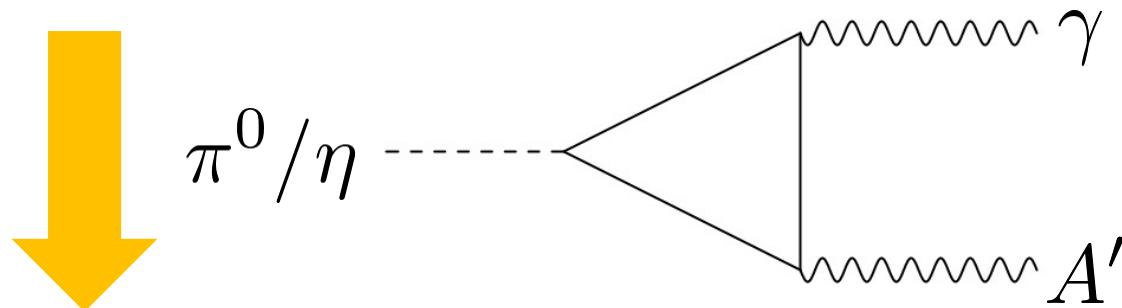
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Production of BSM particles

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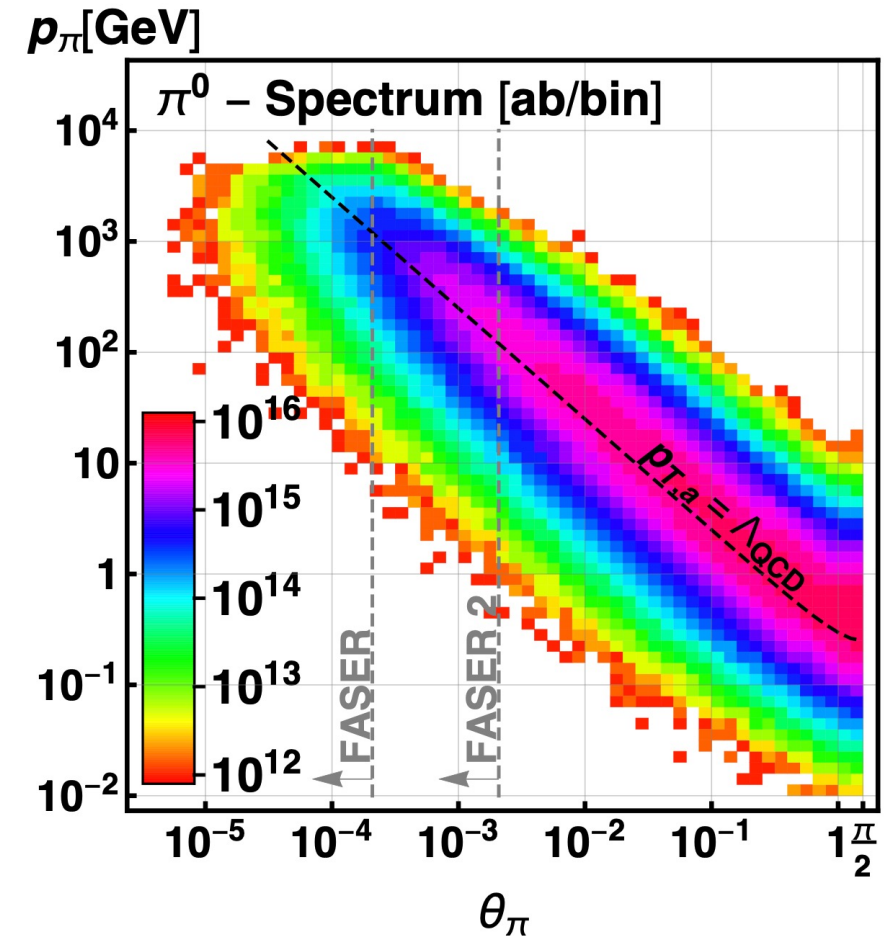
Production @ ATLAS

π^0 and η are produced
at ATLAS interaction point



Dark photons are produced
by rare decays of π^0 and η
$$\text{BR}(\pi^0 \rightarrow \gamma A') \sim 2\epsilon^2 \times \text{BR}(\pi^0 \rightarrow \gamma\gamma) \approx 1$$

[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)



Appendix

Decay length

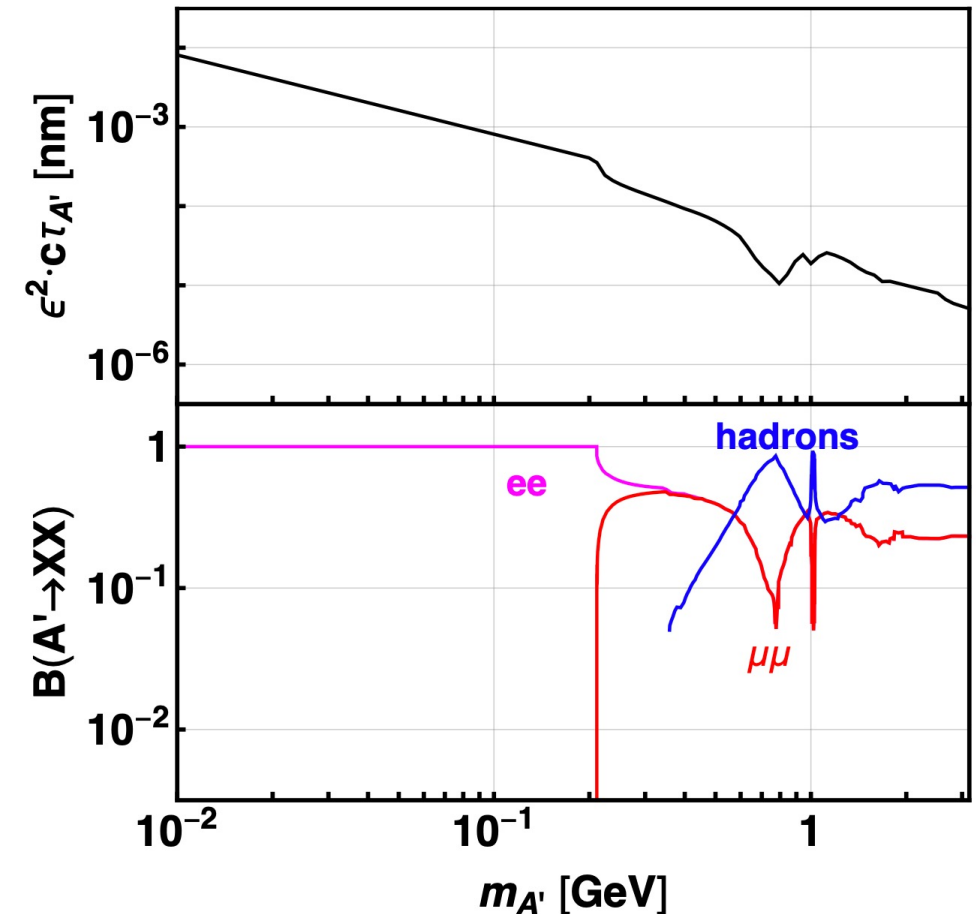
Dark photon decays into SM particles through $\gamma - A'$ mixing

Decay length

$$d = c\tau_{A'}^{\text{rest}} \beta\gamma = c\tau_{A'}^{\text{rest}} \frac{p_{A'}}{m_{A'}} \\ \sim 10^{-2} \text{ m} \cdot \frac{10^3 \text{ GeV}}{0.1 \text{ GeV}} = 100 \text{ m}$$

for $m_{A'} = 0.1 \text{ GeV}$, $\epsilon = 10^{-5}$

[FASER collaboration, PRD 99 \(2019\) 9, 095011](#)



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Dark photon mass

Origin of dark photon mass

→ Spontaneous symmetry breaking by VEV of $U(1)_{\text{dark}}$ -charged scalar

Interaction between $U(1)_{\text{dark}}$ -charged scalar and gauge boson

$$\mathcal{L} \supset g'^2 \phi^\dagger \phi A'_\mu A'^\mu \xrightarrow[\text{SSB}]{\langle \phi \rangle = v_\phi / \sqrt{2}} \underbrace{\frac{1}{2} m_{A'}^2 A'_\mu A'^\mu}_{A' \text{ mass term}} + \underbrace{g' m_{A'} \phi A'_\mu A'^\mu}_{\phi A' A' \text{ coupling}}$$

→ $\phi Z' Z'$ coupling

Dark photon mass

Interaction between $U(1)_{\text{dark}}$ -charged scalar and dark photon

$$\mathcal{L} \supset g' m_{A'} \phi A'_\mu A'^\mu \quad \longrightarrow \quad \phi \text{ decay into } A' \text{ pair}$$

We assume that $U(1)_{\text{dark}}$ -charged scalar is much heavier than dark photon ($m_{A'} \ll m_\phi$)

$\longrightarrow$ Almost all ϕ decay into dark photon

$$\Gamma(\phi \rightarrow A' A') = \frac{g'}{8\pi} \frac{m_{A'}^2}{m_\phi} \sqrt{1 - \frac{4m_{A'}^2}{m_\phi^2}} \left[2 + \frac{m_\phi^4}{4m_{A'}^4} \left(1 - \frac{2m_{A'}^2}{m_\phi^2} \right)^2 \right]$$

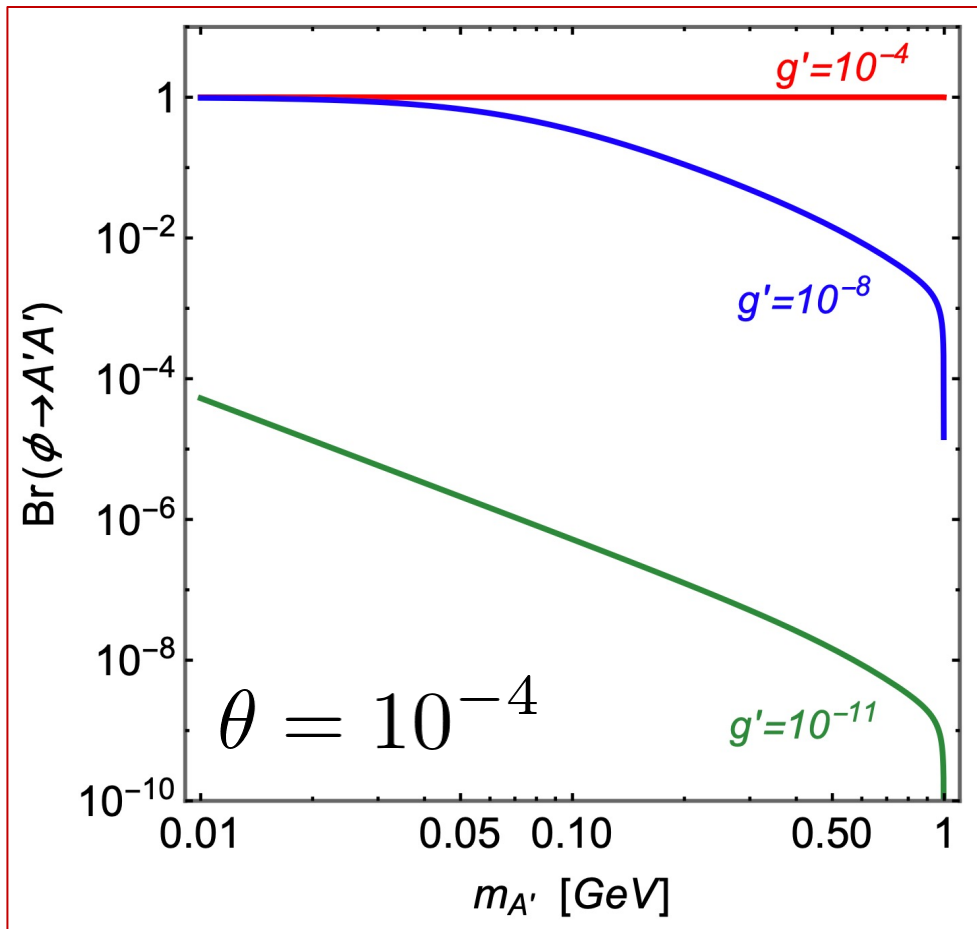
enhancement factor $\gg 1$

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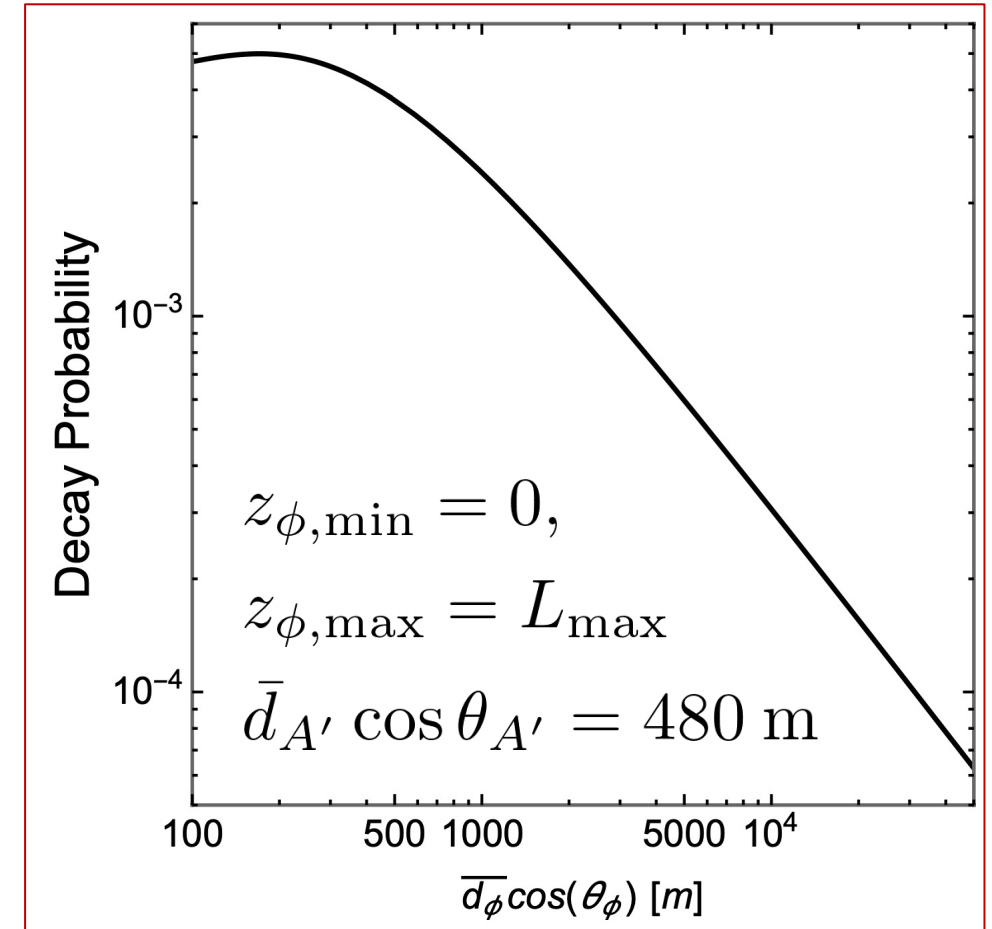
Dark photon

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Branching fraction of $\phi \rightarrow A'A'$



Decay probability of ϕ

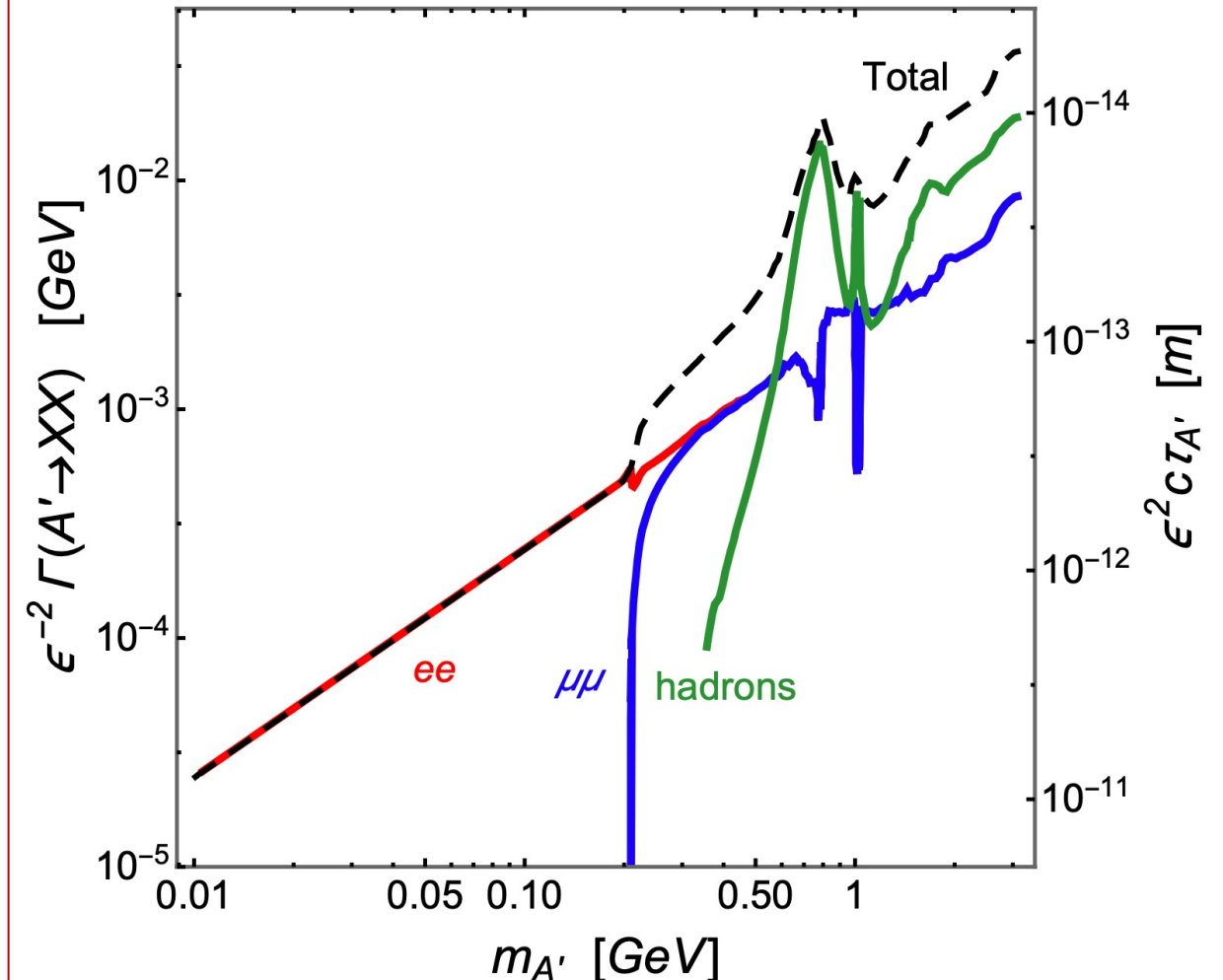


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Dark photon

Decay length of A'

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Decay of scalar boson

Almost all ϕ decay into dark photon

Nonzero ϕ -Higgs mixing : θ

➡ A small number of ϕ can decay into SM particles

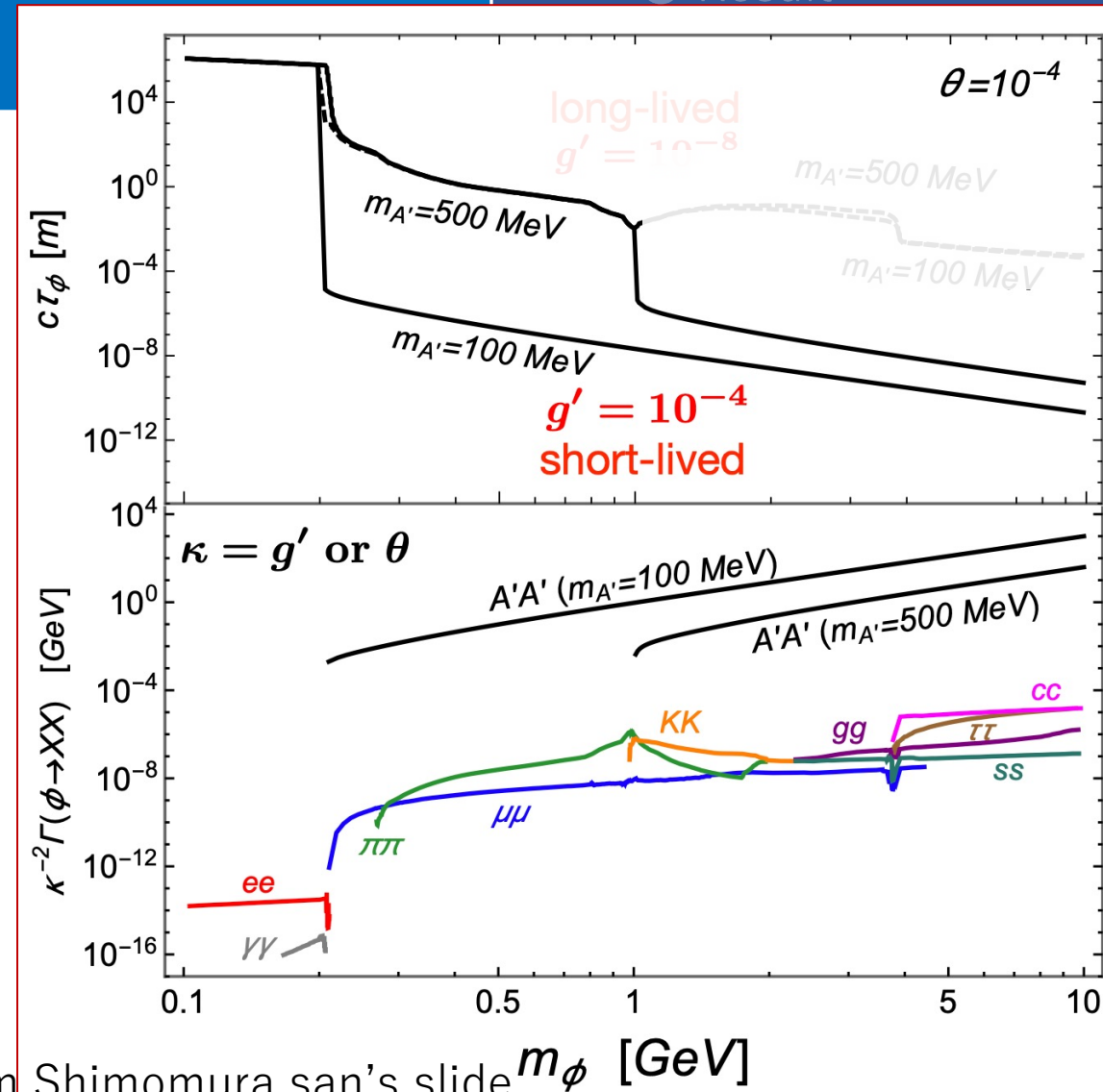
for $g' > \theta$

ϕ : short-lived, A' : long-lived

for $g' \ll \theta$

ϕ : long-lived

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From Shimomura san's slide

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Decay of scalar boson

Almost all ϕ decay into dark photon

Nonzero ϕ -Higgs mixing : θ

➡ A small number of ϕ can decay into SM particles

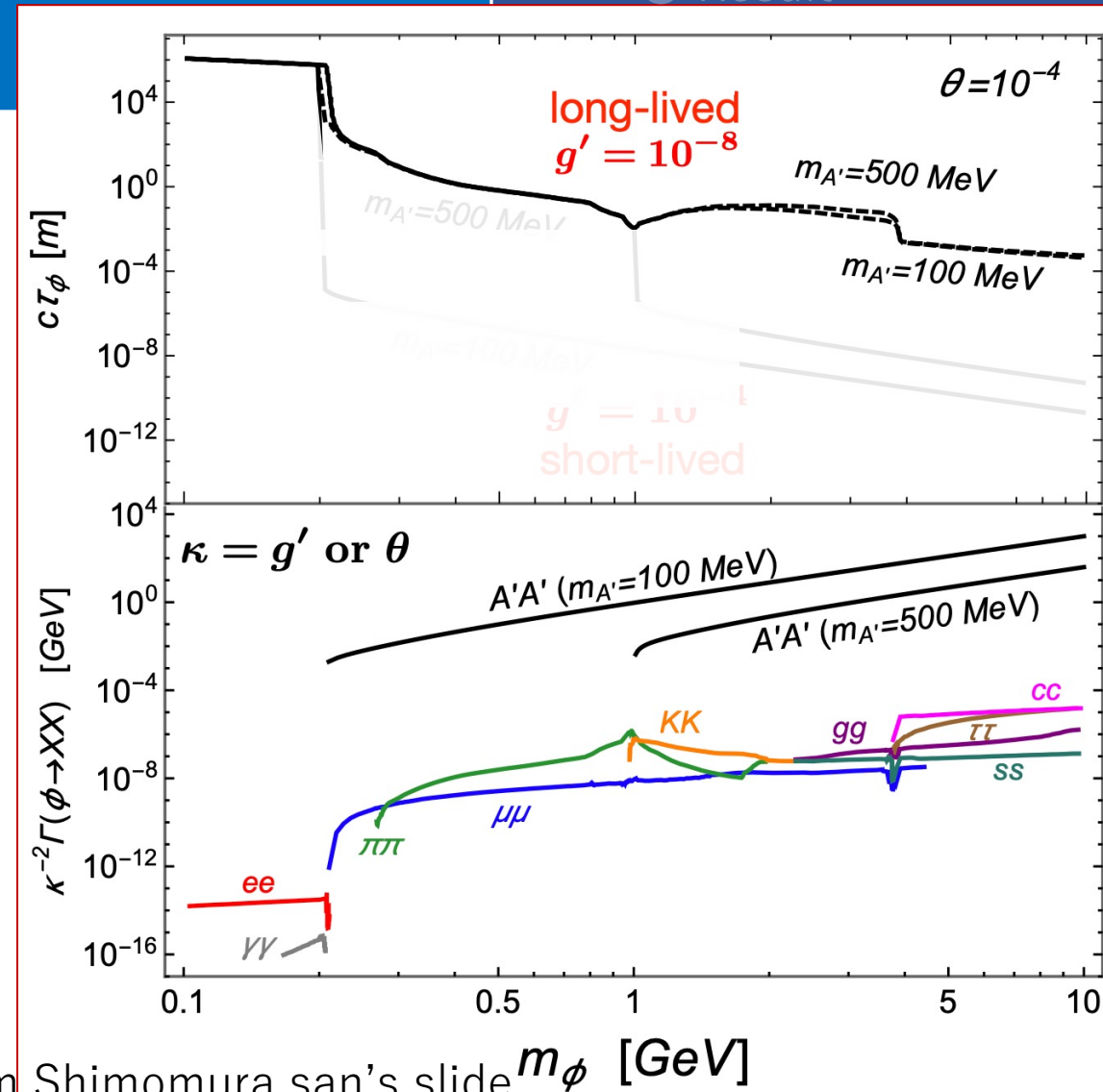
for $g' > \theta$

ϕ : short-lived, A' : long-lived

for $g' \ll \theta$

ϕ : long-lived

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From Shimomura san's slide

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Decay of scalar boson

Almost all ϕ decay into dark photon

Nonzero ϕ -Higgs mixing : θ

➡ A small number of ϕ can decay into SM particles

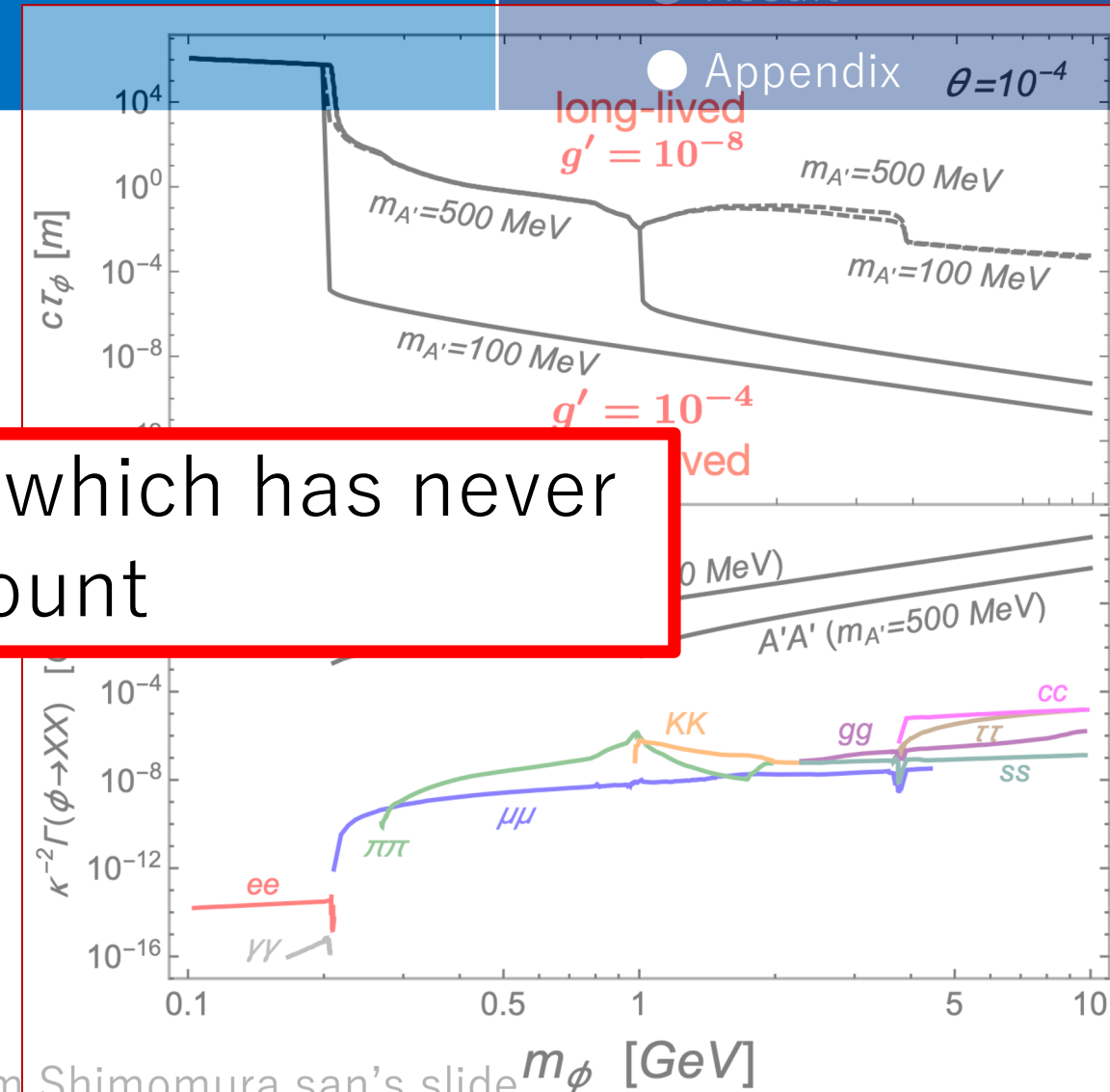
Dark photon source which has never been taken into account

for $g' > \theta$

ϕ : short-lived, A' : long-lived

for $g' \ll \theta$

ϕ : long-lived



From Shimomura san's slide

Scalar production

Production processes

$$B \rightarrow X_s \phi, \quad K^\pm \rightarrow \pi^\pm \phi,$$

$$K_{L(S)} \rightarrow \pi^0 \phi, \quad \eta' \rightarrow \eta \phi$$

$$\text{Br}(B \rightarrow X_s \phi) \simeq 5.7 \left(1 - \frac{m_\phi^2}{m_b^2}\right)^2 \theta^2,$$

$$\text{Br}(K^\pm \rightarrow \pi^\pm \phi) = 2.0 \times 10^{-3} \frac{2p_\phi^0}{m_K} \theta^2,$$

$$\text{Br}(K_L \rightarrow \pi^0 \phi) = 7.0 \times 10^{-3} \frac{2p_\phi^0}{m_K} \theta^2,$$

$$\text{Br}(K_S \rightarrow \pi^0 \phi) = 2.2 \times 10^{-6} \frac{2p_\phi^0}{m_K} \theta^2,$$

$$\text{Br}(\eta' \rightarrow \eta \phi) = 7.2 \times 10^{-5} \frac{2p_\phi^0}{m_{\eta'}} \theta^2,$$

Number of scalars

7.1×10^{13} B mesons are produced in LHC Run3 (150 fb^{-1})

➡ Scalars ϕ are produced by rare decays of B mesons

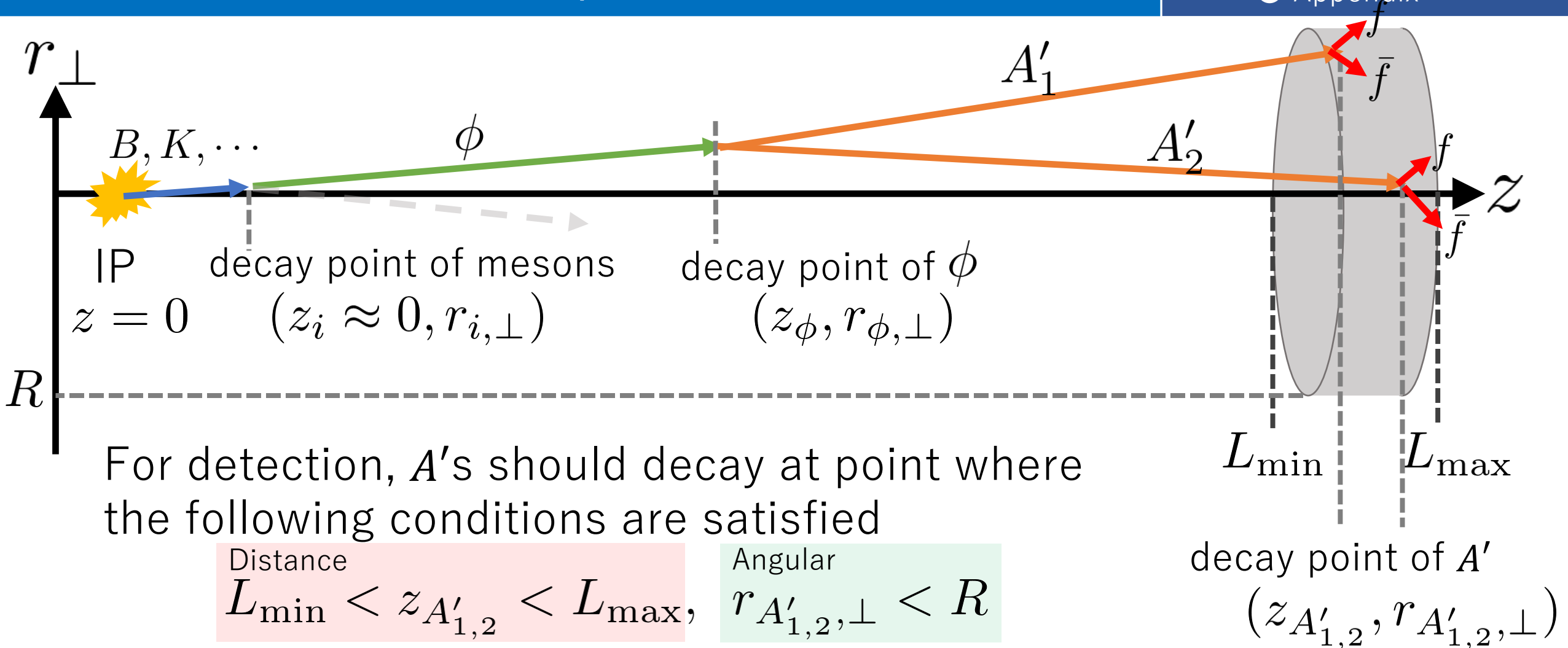
$$\text{BR}(B \rightarrow X_s \phi) \sim \theta^2$$

$$\Rightarrow N_{B \rightarrow \phi} \sim 10^6 \quad (\text{for } \theta = 10^{-4})$$

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Geometrical acceptance

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Geometrical acceptance

Probability of dark photon decay in detector

$$\mathcal{P}_{A'}^{\text{det}}(\mathbf{p}_{A'}, \mathbf{p}_{\phi}) = \frac{1}{\bar{d}_{\phi} \cos \theta_{\phi}} \int_{z_{\phi, \min}}^{z_{\phi, \max}} dz_{\phi} e^{-\frac{z_{\phi}}{\bar{d}_{\phi} \cos \theta_{\phi}}} \frac{1}{\bar{d}_{A'} \cos \theta_{A'}} \int_{z_{A', \min}}^{L_{\max}} dz_{A'} e^{-\frac{z_{A'} - z_{\phi}}{\bar{d}_{A'} \cos \theta_{A'}}}$$

Distance

$$\times \Theta(R - r_{A', R}) \Theta(R - r_{A', F})$$

Angular

$\bar{d}_{\phi(A')}$: decay length of ϕ (A') in lab frame $\Theta(x)$: Heaviside step function

$\theta_{\phi(A')}$: angular of ϕ (A') with respect of beam axis (z-axis)

$r_{A', F}(R)$: distance of A' from beam axis at $z_{A'} = L_{\min} (L_{\max})$

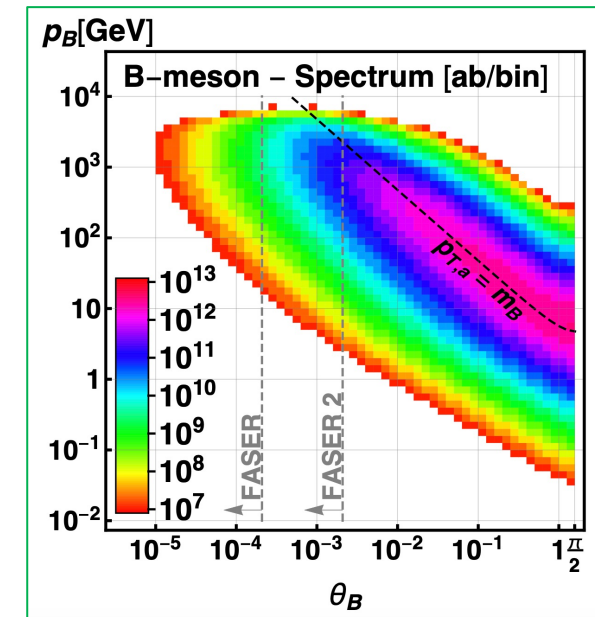
Number of events

Number of LFV decays in FASER detector

$$N_{\text{cLFV}} = \mathcal{L} \sum_{i:\text{meson}} \sum_{j=1,2} \int dp_i d\theta_i \int d\mathbf{p}_{A'} \int d\mathbf{p}_\phi \frac{d\sigma_{pp \rightarrow iX}}{dp_i d\theta_i} \text{Br}(i \rightarrow \tilde{X} \phi) \text{Br}(\phi \rightarrow A'_1 A'_2) \\ \times \mathcal{P}_{A'_j}^{\text{det}}(\mathbf{p}_{A'}, \mathbf{p}_\phi) \times \text{BR}(Z' \rightarrow e\mu)$$

Assumption & Setup

- Calculated for FASER2 case
- Background free



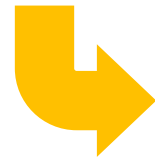
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Ex) $U(1)_{L_\mu-L_\tau}$ gauge boson

$$g_{\mu\tau} Z'_\mu \bar{\ell}_i \gamma^\mu V_{i\alpha}^\dagger \begin{pmatrix} 0 & 0 & 0 \\ 0 & +1 & 0 \\ 0 & 0 & -1 \end{pmatrix} V_{\beta j} \ell_j$$



cLFV interaction

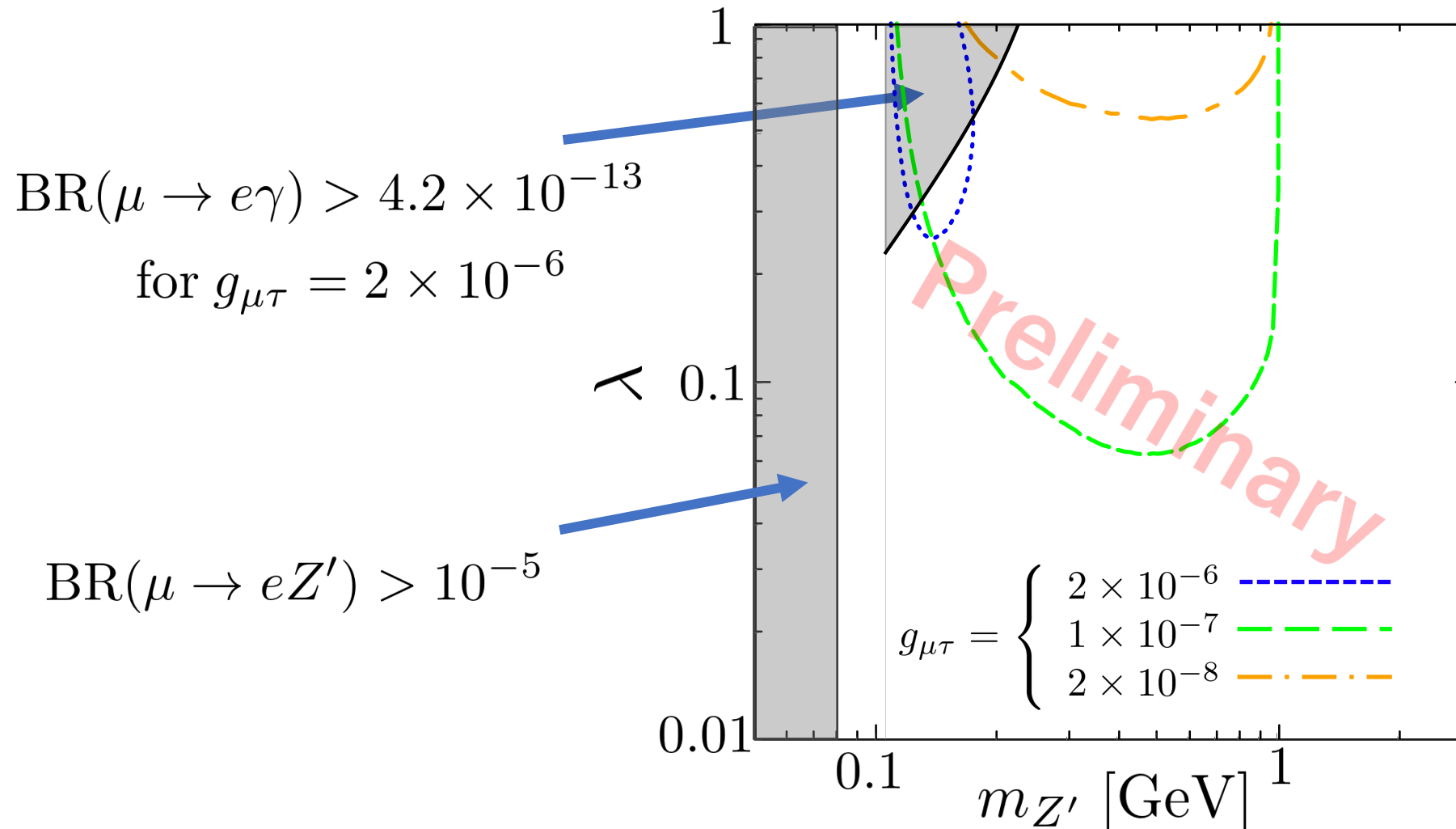
$$\begin{pmatrix} 0 & \lambda & \lambda \\ \lambda & 1 & \lambda \\ \lambda & \lambda & -1 \end{pmatrix}$$

$\lambda(\leq 1)$: Ratio of strength of interaction (LFV / LFC)

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of signal events ($U(1)_{L_\mu-L_\tau}$ gauge boson)



95% C.L. exclusion
contour @ FASER2

Parameters

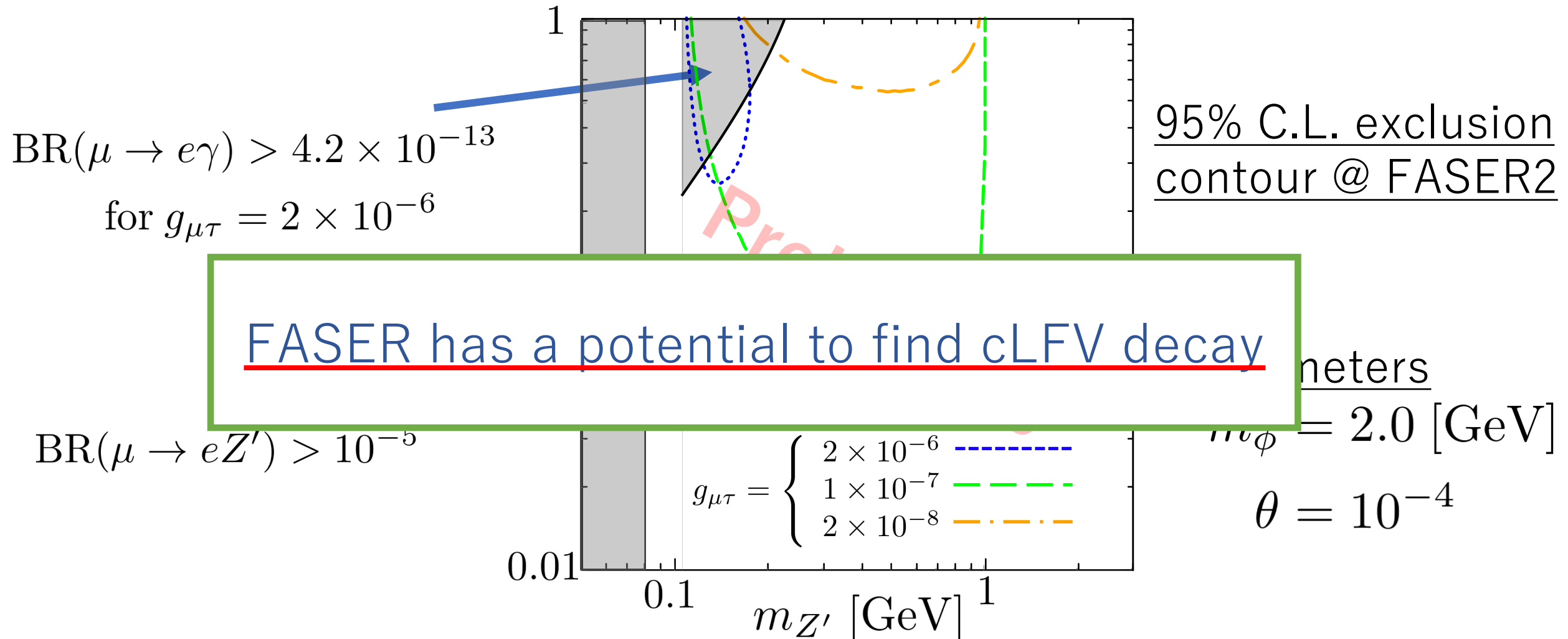
$$m_\phi = 2.0 [\text{GeV}]$$

$$\theta = 10^{-4}$$

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of signal events ($U(1)_{L_\mu-L_\tau}$ gauge boson)



Background

- Rock & LHC infrastructure eliminate most background



Main background [150 fb^{-1} @ LHC Run3]

- Muon brems. $\rightarrow$ photon
: 80000 events
- CC / NC interactions of neutrinos ($E_\nu \gtrsim 100 \text{ [GeV]}$)
: $O(1)$ events



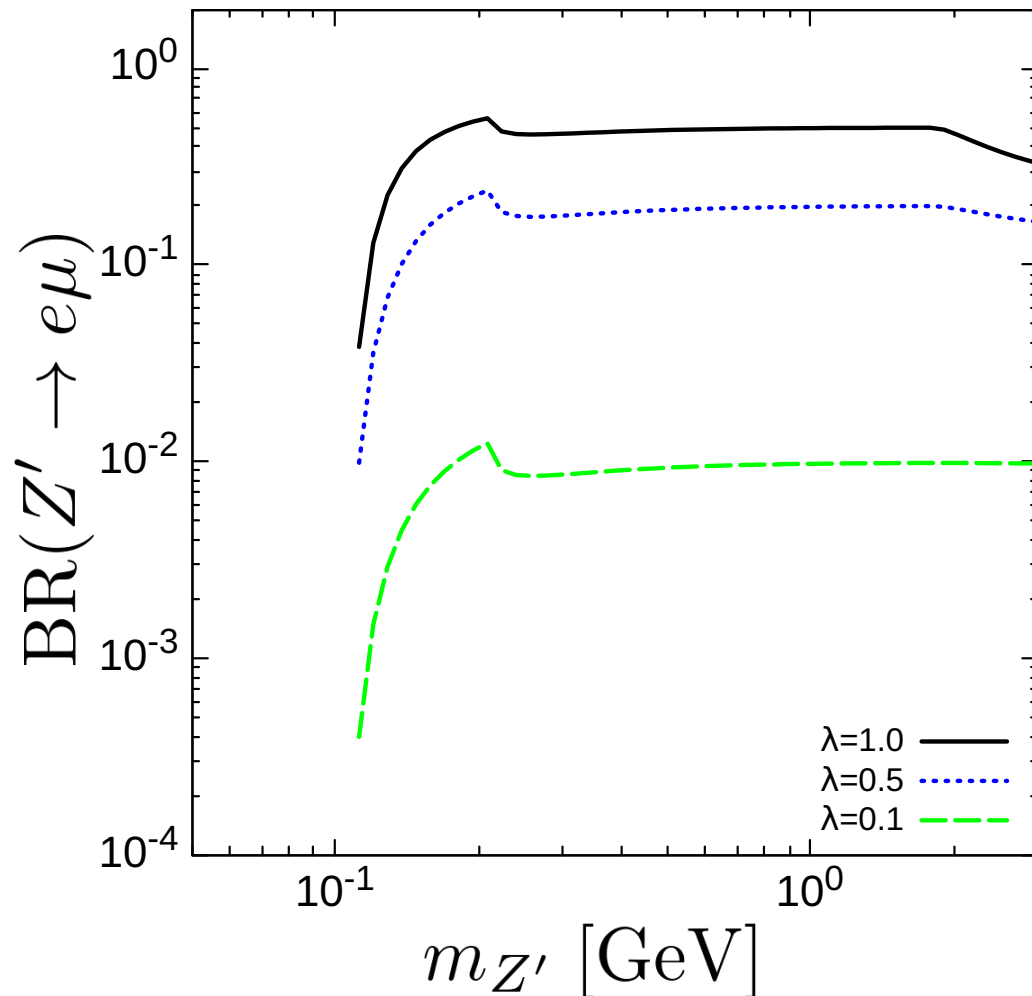
Veting entering charged particles
with an efficiency of 99.99%

Almost background free

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Branching ratio of LFV decay ($U(1)_{L_\mu-L_\tau}$ gauge boson)



Decay width

$$\Gamma(Z' \rightarrow \ell \bar{\ell}') = \frac{V^2}{24\pi} m_{Z'} \lambda \left(1, \frac{m_\ell^2}{m_{Z'}^2}, \frac{m_{\ell'}^2}{m_{Z'}^2} \right) \times \left[2 - \frac{m_\ell^2 - 6m_\ell m_{\ell'} + m_{\ell'}^2}{m_{Z'}^2} - \frac{(m_\ell^2 - m_{\ell'}^2)^2}{m_{Z'}^4} \right]$$

$$\lambda(1, a, b) = \sqrt{1 + a^2 + b^2 - 2a - 2b - 2ab}$$

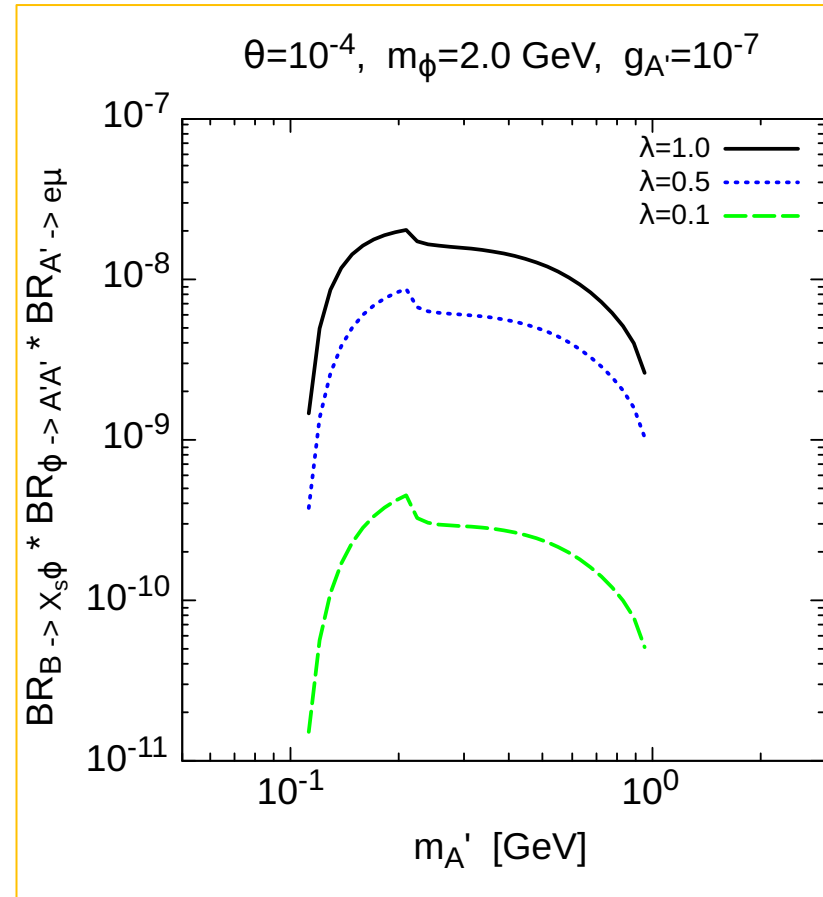
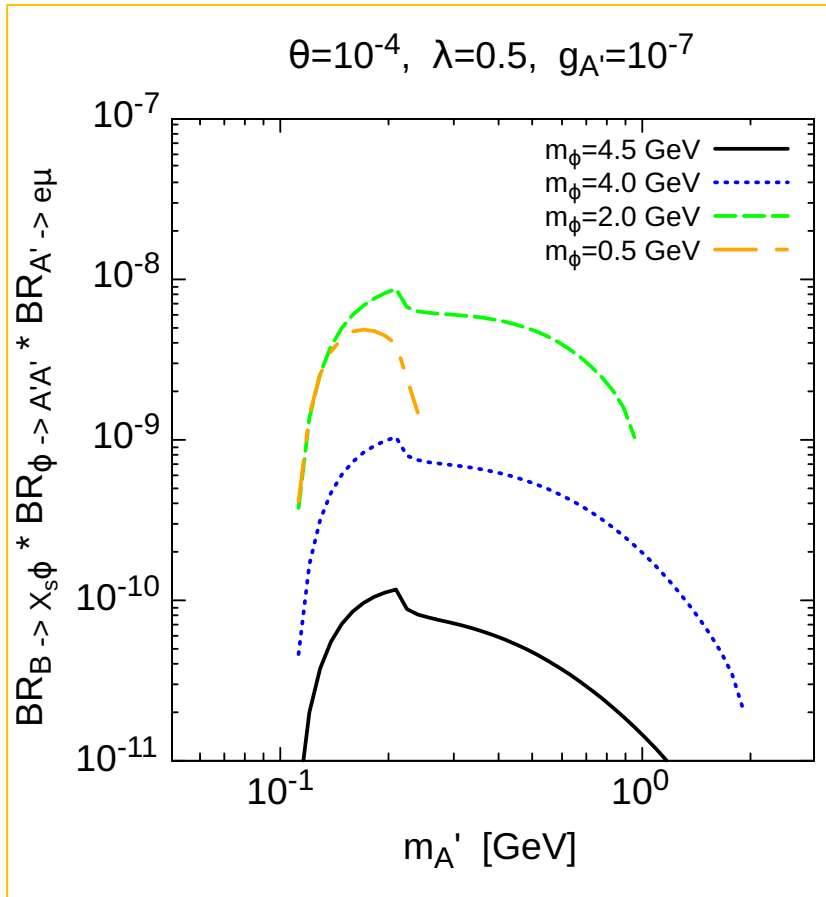
where

$$V = \begin{cases} g_{\mu\tau} & [\text{LFC}] \\ \lambda g_{\mu\tau} & [\text{LFV}] \end{cases}$$

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Branching ratio of LFV decay ($U(1)_{L_\mu-L_\tau}$ gauge boson)



Decay width

$$Br(B \rightarrow X_s \phi) \simeq 5.7 \left(1 - \frac{m_\phi^2}{m_b^2}\right)^2 \theta^2$$

$$Br(\phi \rightarrow A'A') = \frac{\Gamma_{\phi \rightarrow A'A'}}{\Gamma_{\phi \rightarrow A'A'} + \Gamma_{SM}}$$

$$\simeq 1 \text{ (for } g_{A'} > 10^{-6} \text{)}$$