

Cosmological search of axion dark matter

Yuko Urakawa (KEK, iPNS)

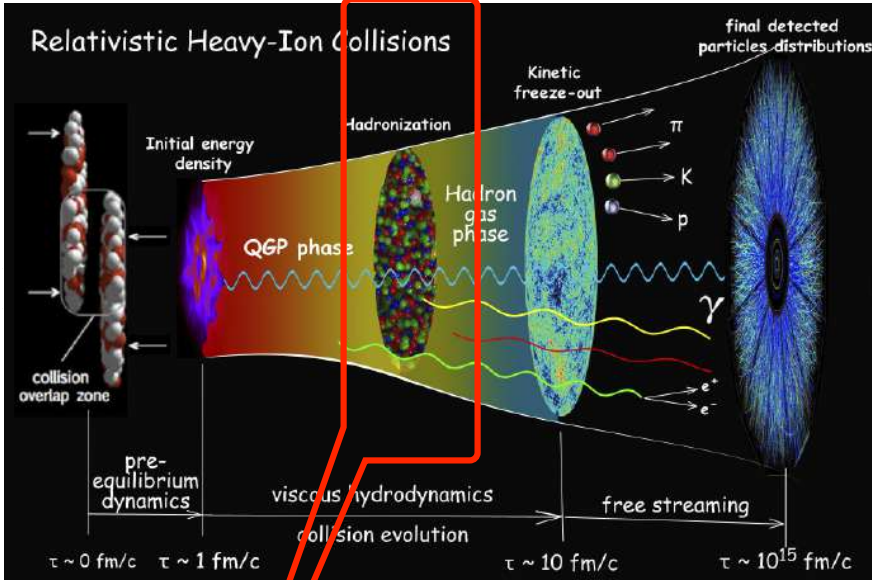


w/ Hayato Fukunaga, Kazuhiro Kogai (Nagoya U.), Naoya Kitajima (Tohoku U.)

Can we learn little bang from big bang?

Little bang

Heinz (13)



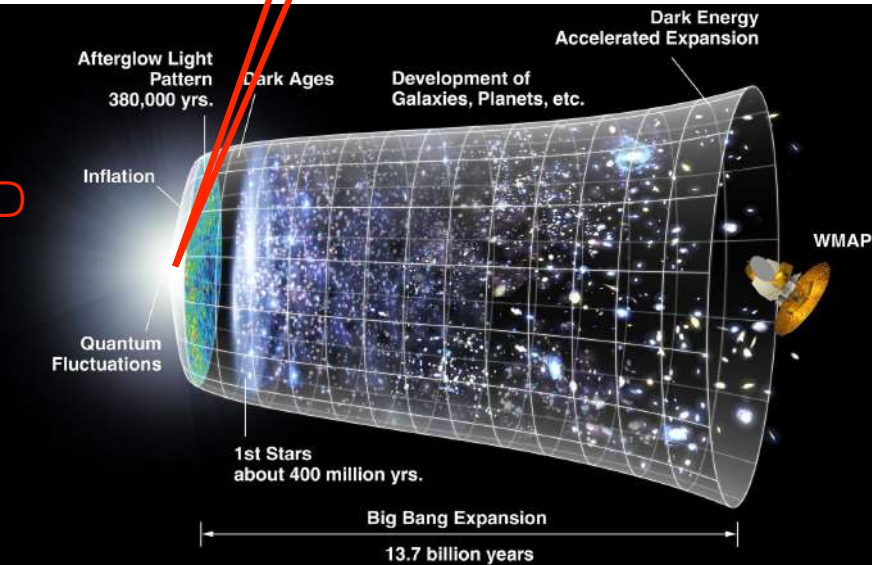
J-PARC



LHC/RHIC

$T \sim \Lambda_{QCD}$

Big bang

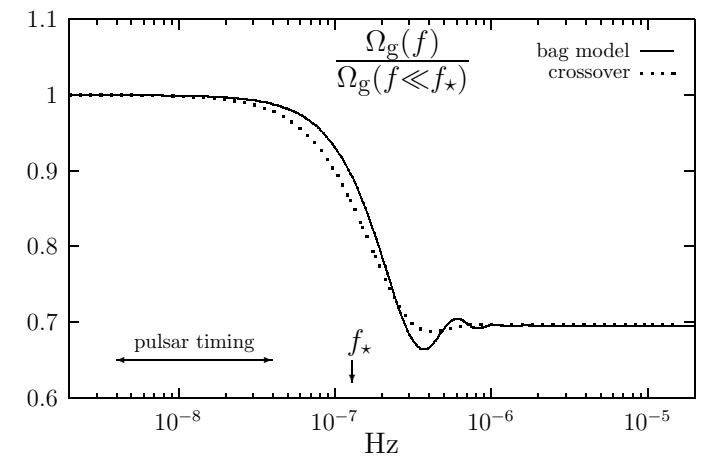
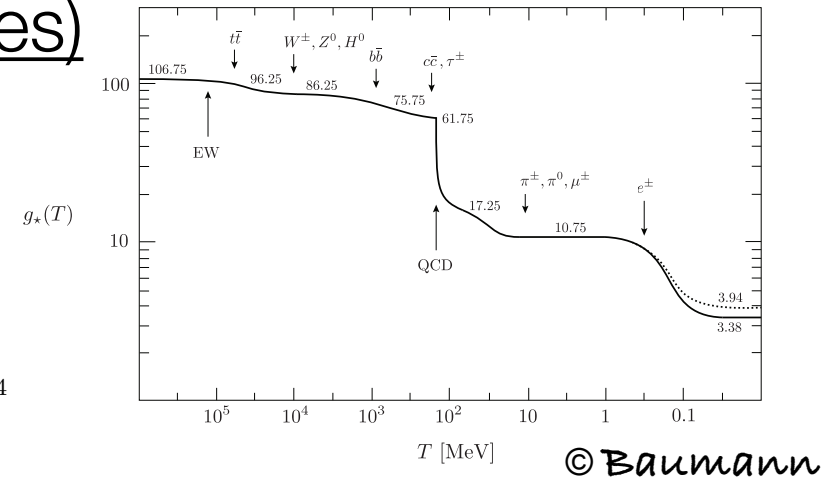


Learning QCD from cosmology

Through radiation (= relativistic species)

- Equation of state
- # of relativistic species

$$\rho_r = \sum_i \rho_i = \frac{\pi^2}{30} g_*(T) T^4$$

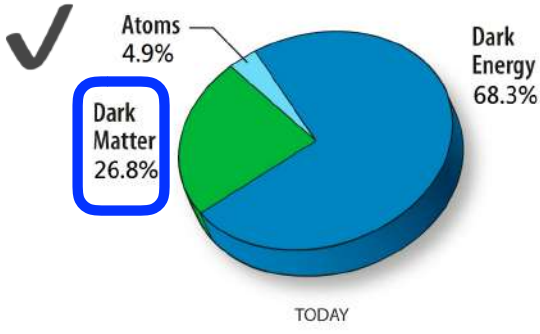
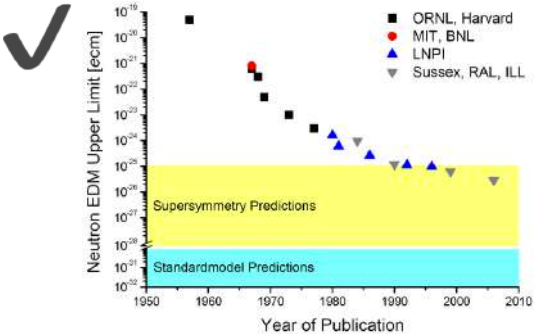
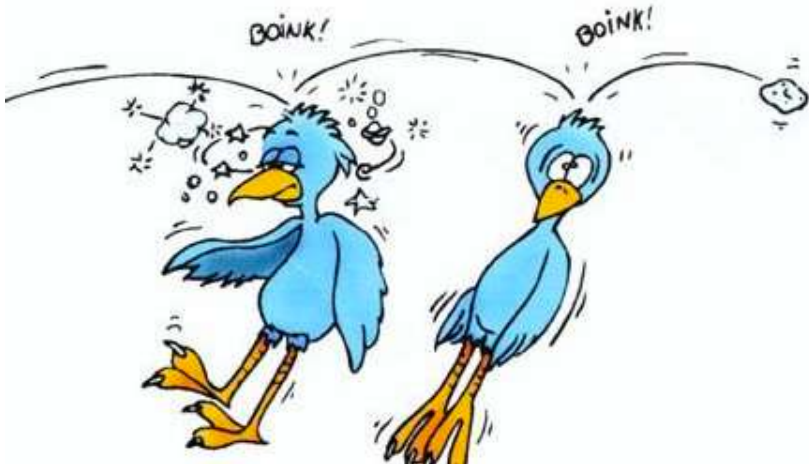


- Expansion history of Universe
- Drop of Ω_{gw} at $f \sim 10^{-7} \text{Hz}$
- PBH abundance w/ M_{\odot}

Through dark matter ?

if 1st order PT : Production of GWs, Magnetic fields ...

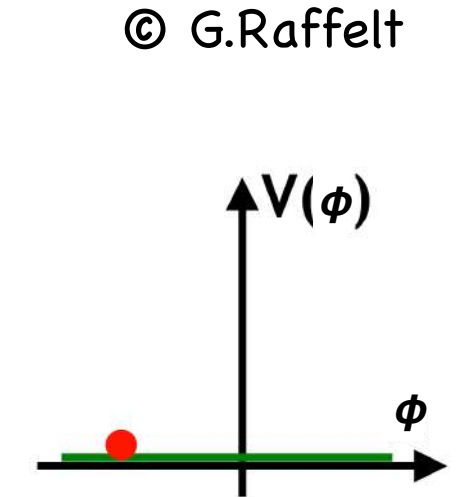
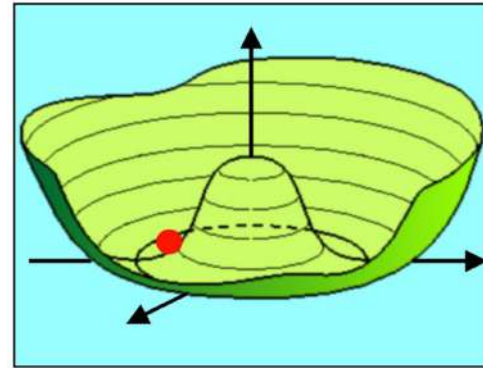
Axion dark matter



Axion as dark matter

- $T \sim f$ Symmetry breaking

azimuthal symmetry
axion is massless



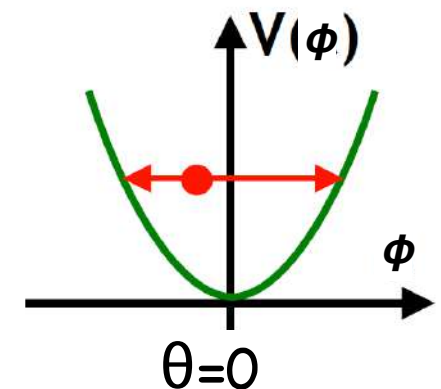
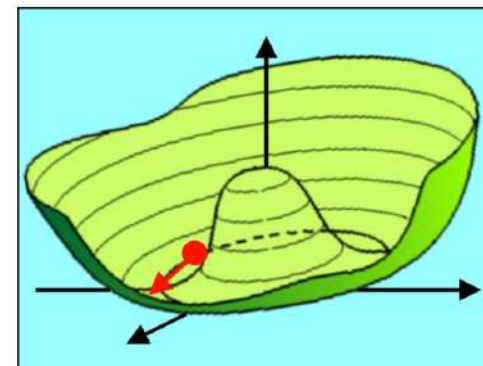
- $T \sim 0.1 \text{ GeV}$

Axion acquires potential through the non-perturbative effect of QCD matters.

→ dark matter

No number changing process

$$\rho_a \sim m_a n_a \propto 1/(\text{Volume})$$



Topological susceptibility

Effective field theory treatment

Free energy : $F(\theta, T) = -\ln Z/\mathcal{V}$

$$V_{\text{eff}} \sim F(\theta, T) - F(0, T)$$

$$\theta = \frac{\phi}{f}$$

$$\frac{\partial^2}{\partial \theta^2}$$

$$F(\theta, T) - F(0, T) \sim \frac{1}{2} \chi(T) \theta^2 \quad |\theta| \ll 1$$

$$\mathcal{L} \ni \theta q(x) \quad q(x) \equiv \frac{\alpha_s}{8\pi^2} \theta F_{\mu\nu a} F^{\mu\nu a}$$

$$\chi = \int d^4x \langle q(x) q(0) \rangle_{\theta=0}$$

$$m_a^2(T) f^2 = \chi(T)$$

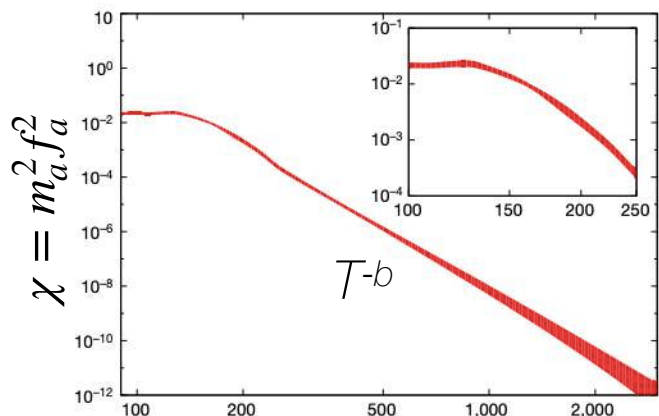
QCD meets cosmology

$$T \propto 1/a$$

Temperature dep. axion potential from QCD dynamics

$$V(\theta, T) = \frac{\chi(T)}{\chi(T=0)} V_0(\theta)$$

Axion abundance



Borsanyi+ (16) T [MeV]

$b \sim 8.2$

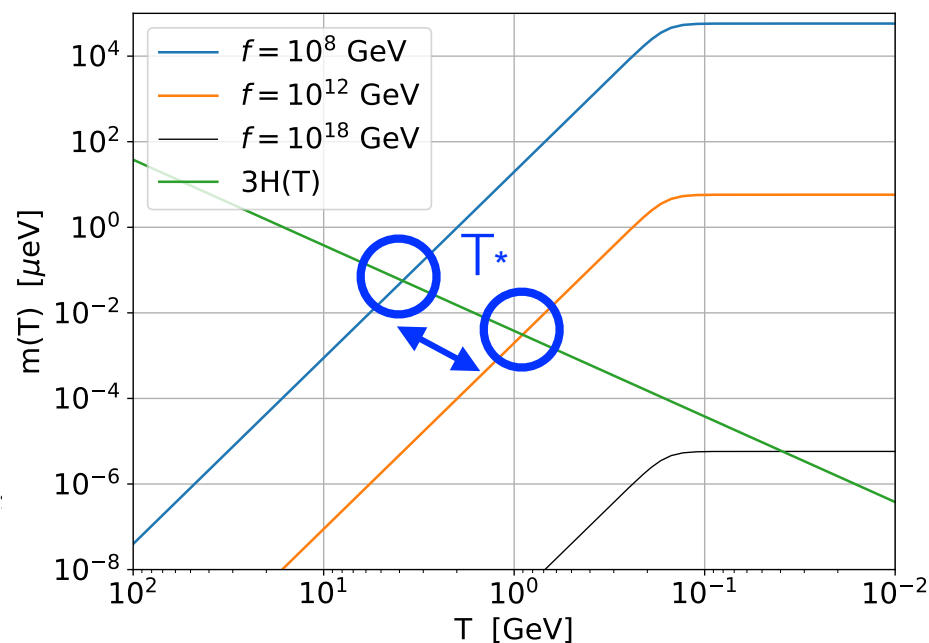
Pheno.constraints

$$10^8 \text{ GeV} < f < 10^{12} \text{ GeV}$$

$$T_* \simeq 1.0 \text{ GeV} \left(\frac{10^{12} \text{ GeV}}{f} \right)^{0.16}$$

(Misalignment) Axion abundance

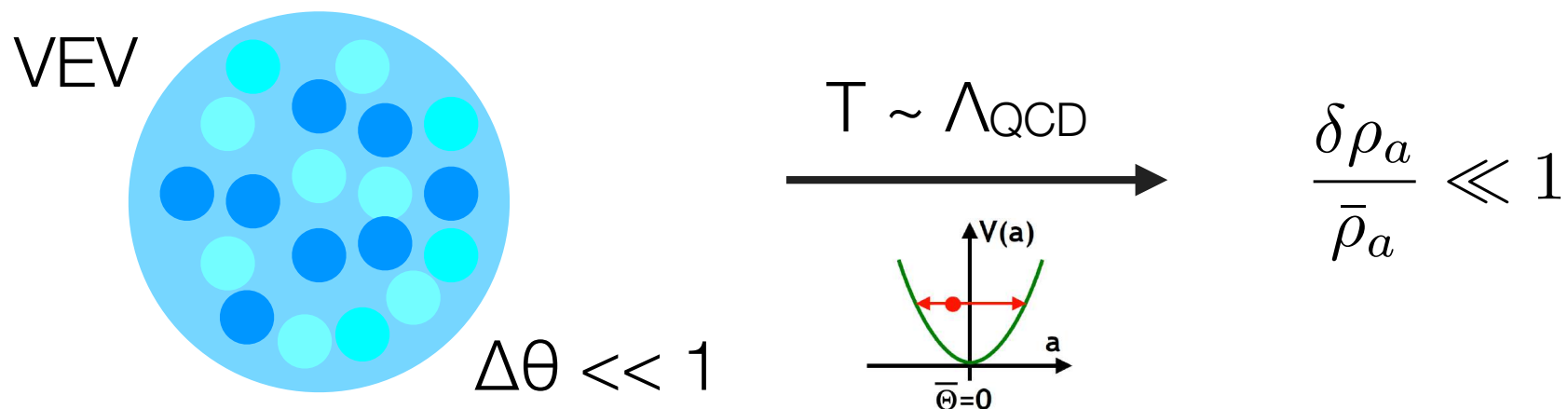
Axion starts oscillation $3H(T^*) \sim m(T^*)$



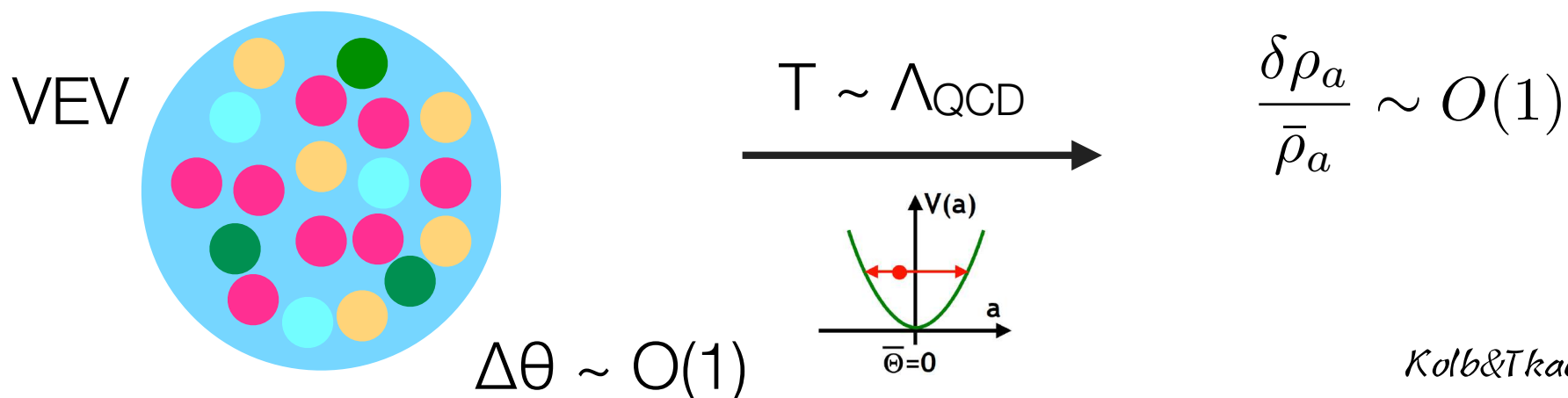
$$\Omega_a h^2 \sim 0.1 \left(\frac{f}{10^{12} \text{ GeV}} \right)^{1.16} \theta_i^2$$

2 axion scenarios

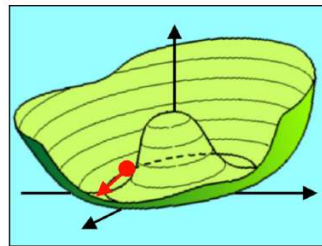
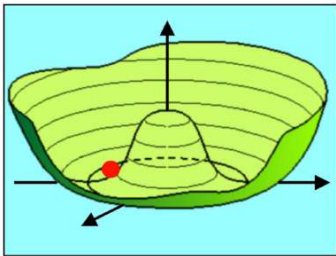
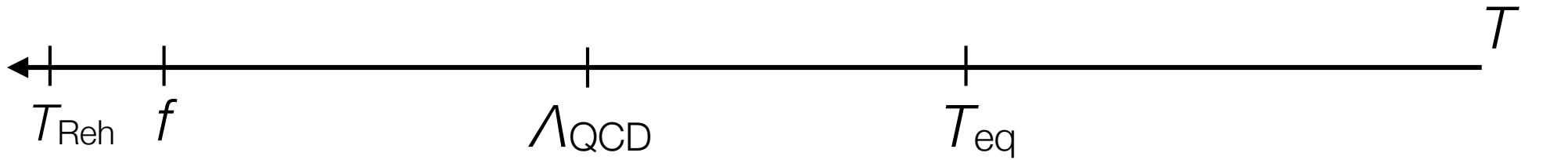
i) PQ symmetry was broken during inflation $f \geq H_{\text{inf}}$



ii) PQ symmetry got broken after inflation $f < H_{\text{inf}}$



AMC formation in post-inflationary scenario



Gravitational collapse

Virialization *Kolb&Tkacev(93, 94)*

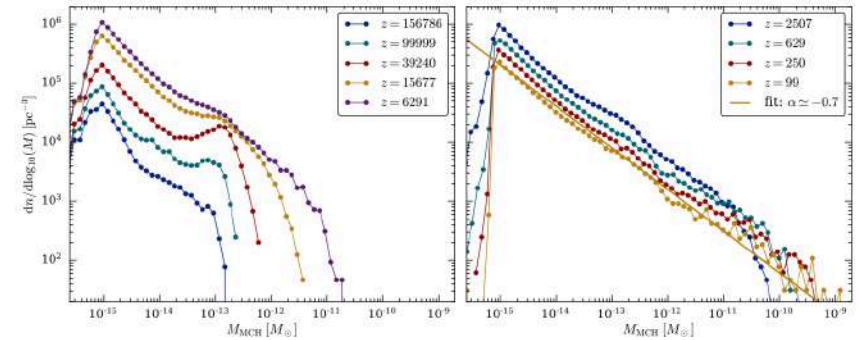
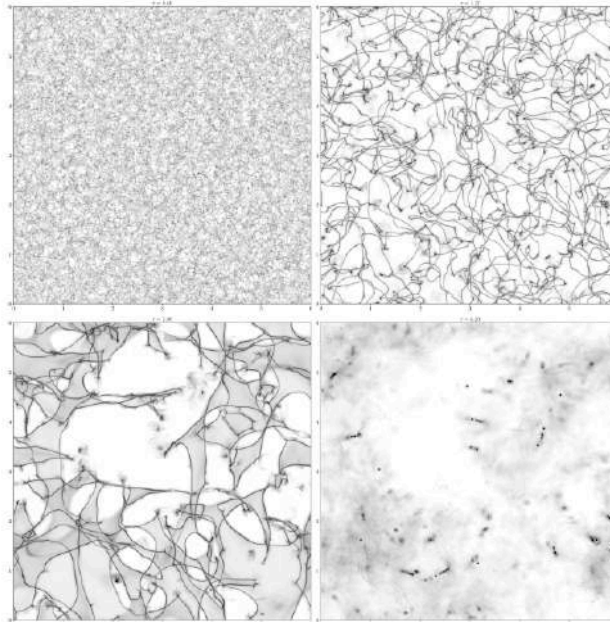


FIG. 2. MC-HMF at different redshifts z separated into times before (left) and after matter-radiation equality (right). The slope of the MC-HMF at $z_f = 99$ is $\alpha \simeq -0.7$.

Eggmeier et al(19)

See also *Kavanagh et al. (20)*, *Xiao et al. (21)*

Vaquero, Redondo, Stadler (18)

See also *Fluery&Moore(15),...*

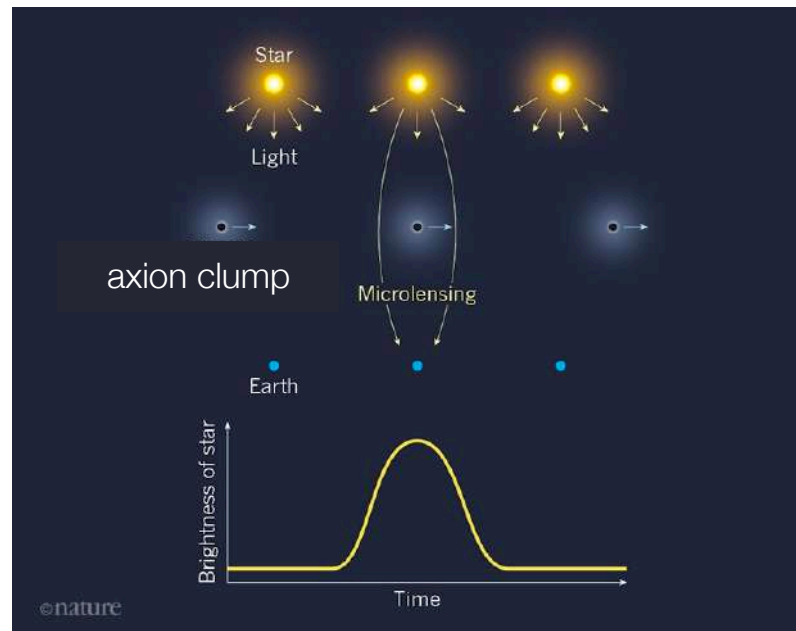
Axion minicluster (AMC) search

- Microlensing
- Pulsar Timing array

.... and so on

typical velocity
of DM

$$v \sim 10^{-3}c$$



*Kolb&Tkacev(95),
Fairbairn et al. (17),
Dai&Miralda-Escude(19)*

Is AMC the smoking gun of post-inflation scenario?



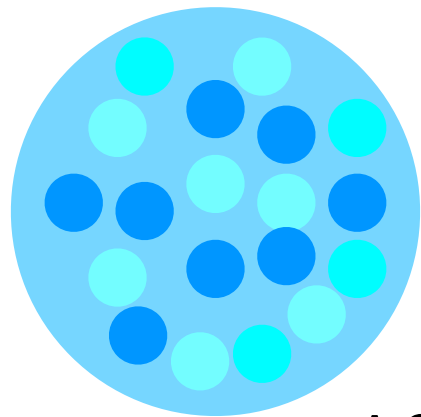
If yes, a detection of AMC gives the upper bound on f .

$$f < H_{\text{inf}}/2\pi < 10^{13} \text{ GeV (PLANCK/BICEP/Keck)}$$

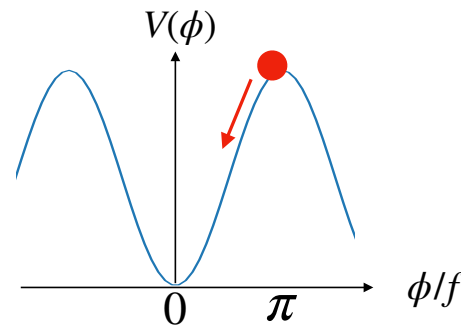
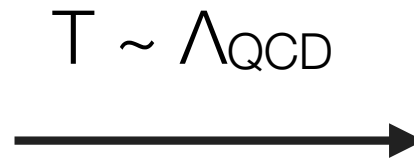


Is that true?

PQ symmetry broken during inflation $f \geq H_{\text{inf}}$



$$\Delta\theta \ll 1$$



$$\frac{\delta\rho_a}{\bar{\rho}_a} \sim O(1) ?$$

$$V(\phi) = m(T)^2 f^2 (1 - \cos(\phi/f))$$

1) Interaction btwn QCD matters and axion

✓ - Direct int. **NEW!!**

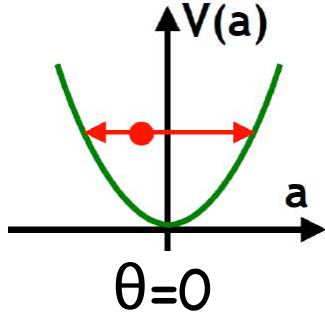
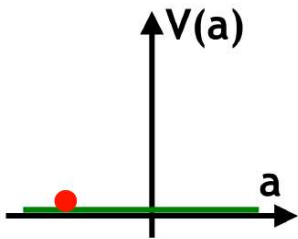
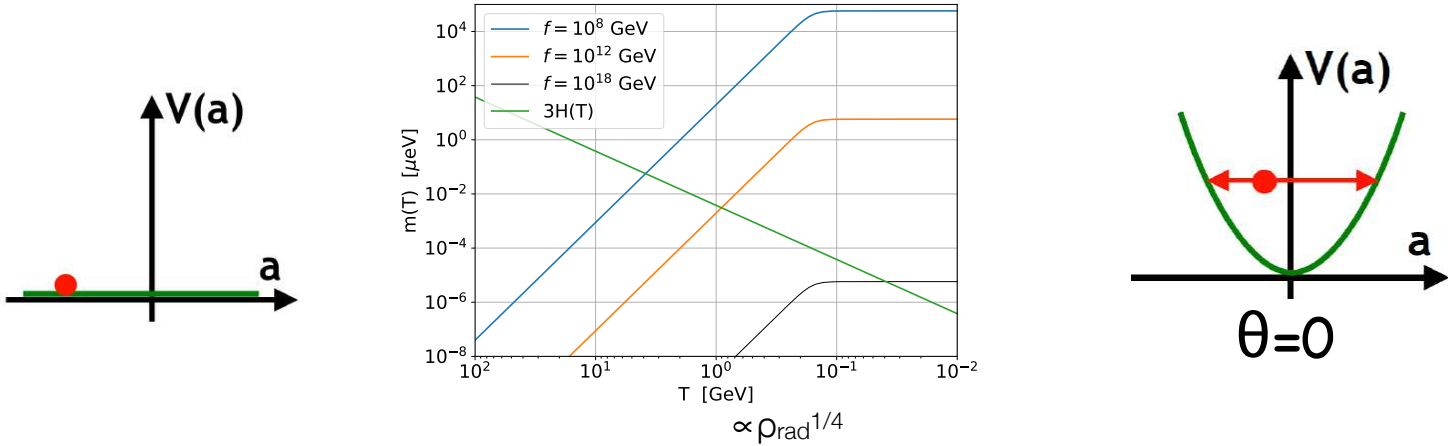
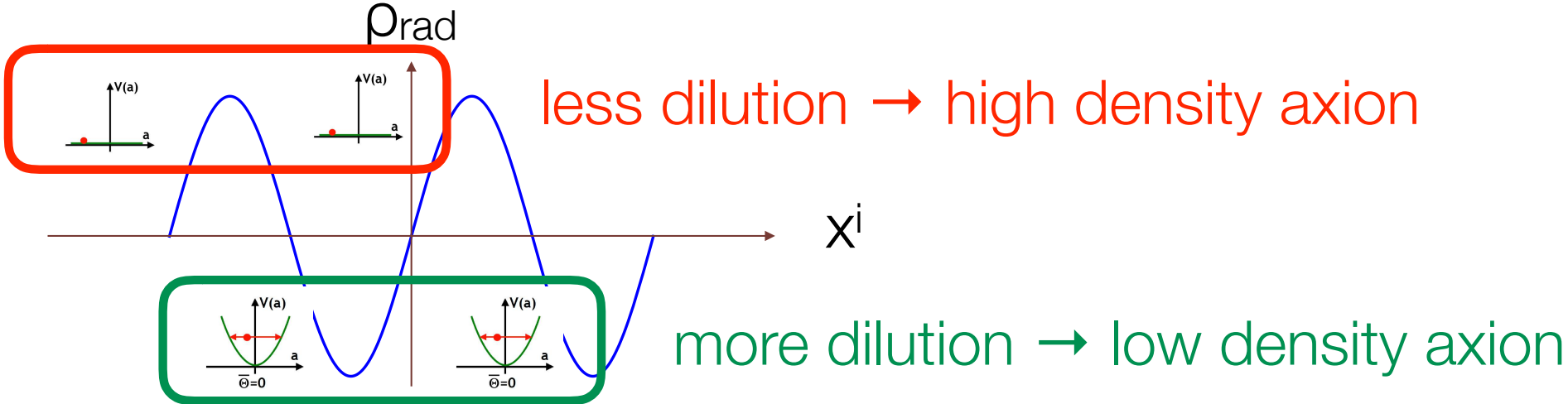
- Indirect int. via gravity (clustering) *Arvanitaki et al. (19)*

2) Self interaction $\lambda\phi^4$ ($\lambda < 0$) attractive force
for $\phi/f \sim \pi$

Direct interaction w/QCD matters (~ radiation)

Inhomogeneous Universe

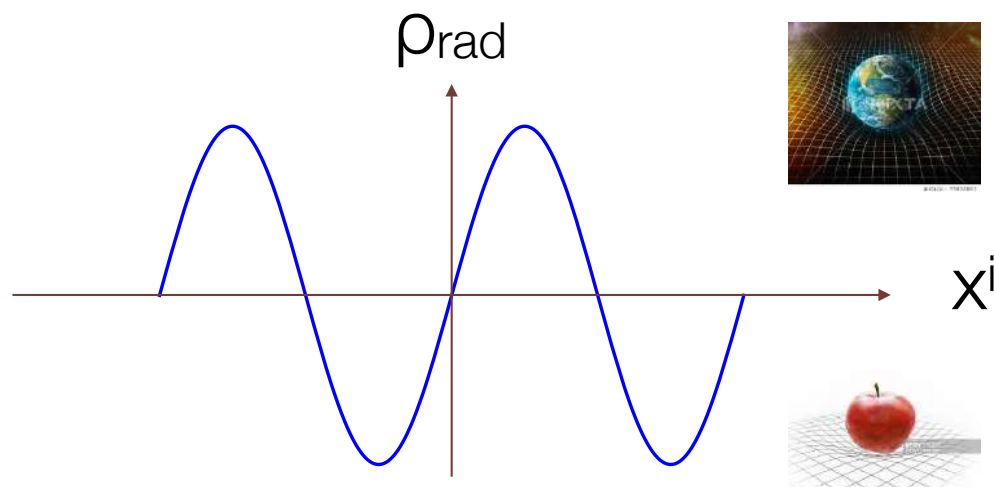
Kitajima, Kogai, Y.U. (in prep.)



Indirect interaction w/QCD matters (~ radiation)

Inhomogeneous Universe

Kitajima, Kogai, Y.U. (in prep.)



Through eq.

$$\nabla_{\mu} T^{\mu}_{\nu}(\text{axion}) = -\nabla_{\mu} T^{\mu}_{\nu}(\text{rad})$$



indirect



direct

Linear perturbation

Kitajima, Kogai, Y.U. (in prep.)

$\delta\phi$: Inhomogeneity of axion

at longitudinal gauge

Φ : Gravitational potential

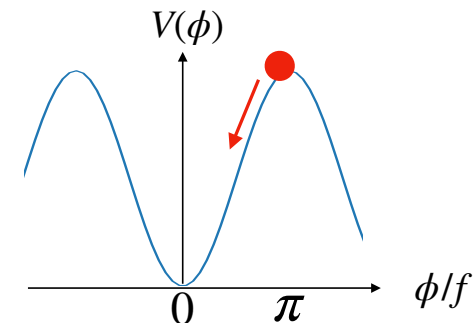
Equations

$$\partial_t^2 \delta\phi + 3H\partial_t \delta\phi + \left(\frac{k}{a}\right)^2 \delta\phi + \underbrace{V_{\phi\phi}\delta\phi}_{\text{incl. self-int.}} + \underbrace{4\dot{\phi}\partial_t\Phi}_{\text{indirect int.}} - 2V_{\phi}\Phi + \underbrace{V_{\phi\rho_r}\delta\rho_r}_{\text{direct int.}} = 0$$

$$\Delta\Phi \sim 4\pi G\delta\rho \quad + \text{GR corrections}$$

IC $\Delta \equiv \pi - \phi_i/f$ (+ slow-roll initial velocity)

For $\Delta \ll 1$, the self-interaction was initially important.



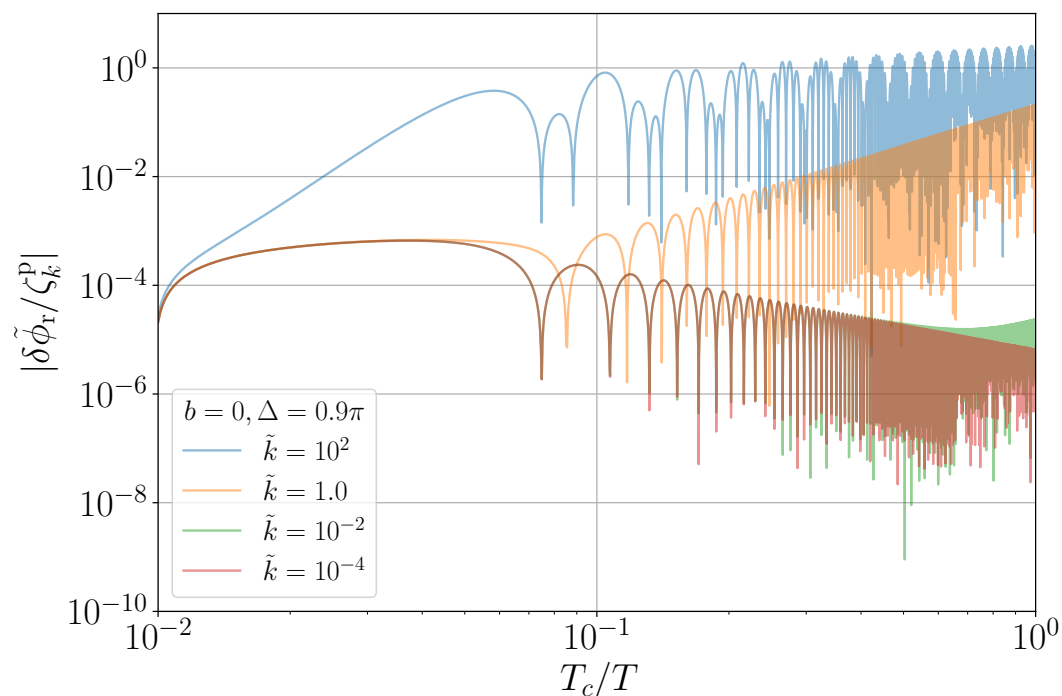
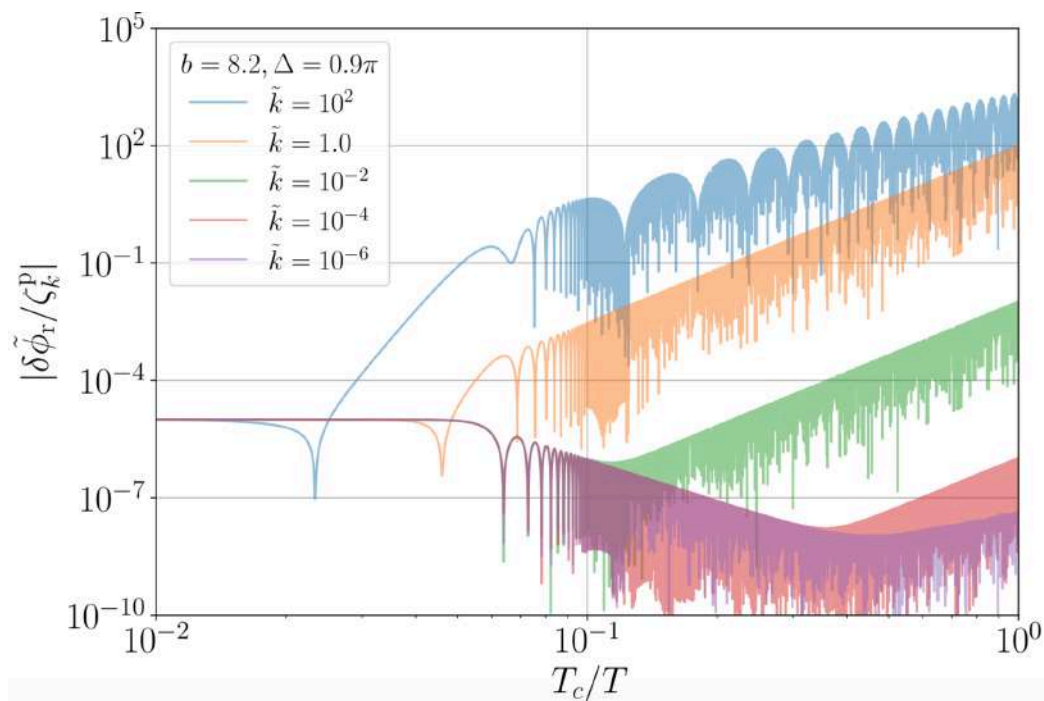
Linear perturbation: $\Delta=O(1)$

Kitajima, Kogai, Y.U. (in prep.)

$$\partial_t^2 \delta\phi + 3H \partial_t \delta\phi + \left(\frac{k}{a}\right)^2 \delta\phi + \underbrace{V_{\phi\phi}}_{\text{no self-int.}} \delta\phi + \underbrace{4\dot{\phi}\partial_t\Phi - 2V_\phi\Phi}_{\text{indirect int.}} + \underbrace{V_{\phi\rho}}_{\text{direct int.}} \delta\rho_r = 0$$

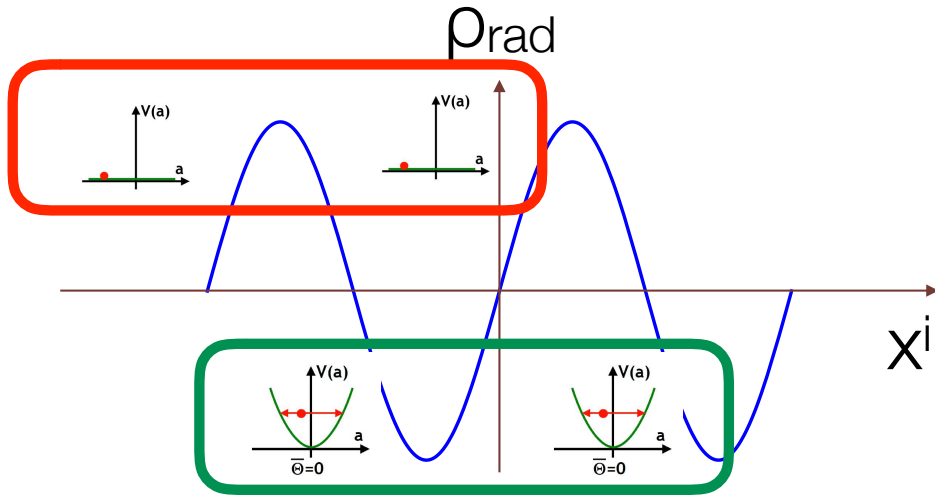
(axion) $V_{\phi\rho_r} \neq 0$, $V_{\phi\phi} \sim m^2$

(ALP like) $V_{\phi\rho_r} = 0$, $V_{\phi\phi} \sim m^2$



$\zeta^p \sim$ Primordial amplitude (of radiation)

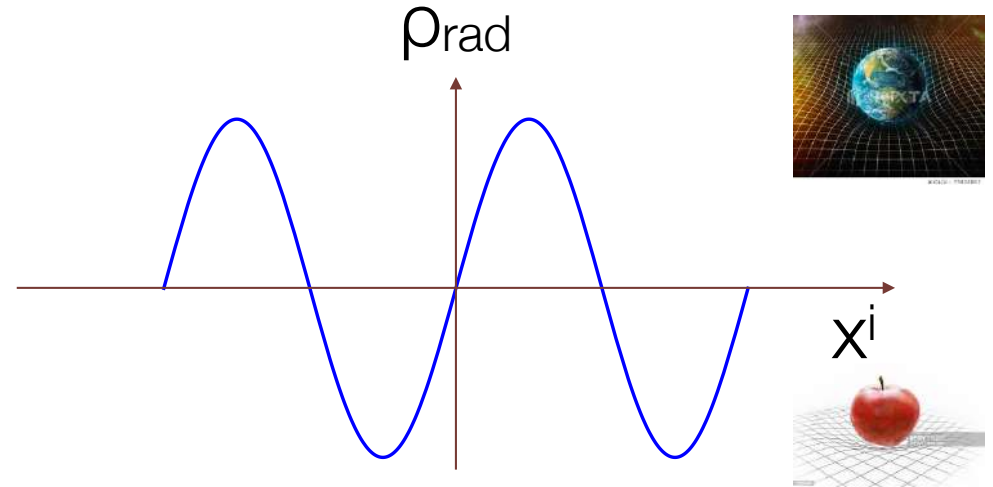
Significance of direct interaction



Kitajima, Kogai, Y.U. (in prep.)

$\delta\rho_{\text{rad}}/\rho_{\text{rad}} \sim \text{const. oscillation}$

\gg



Arvanitaki et al. (19)

$\Phi \sim \text{decaying oscillation}$

$$\Delta\Phi \sim 4\pi G\delta\rho \sim 4\pi G\delta\rho_{\text{rad}}$$

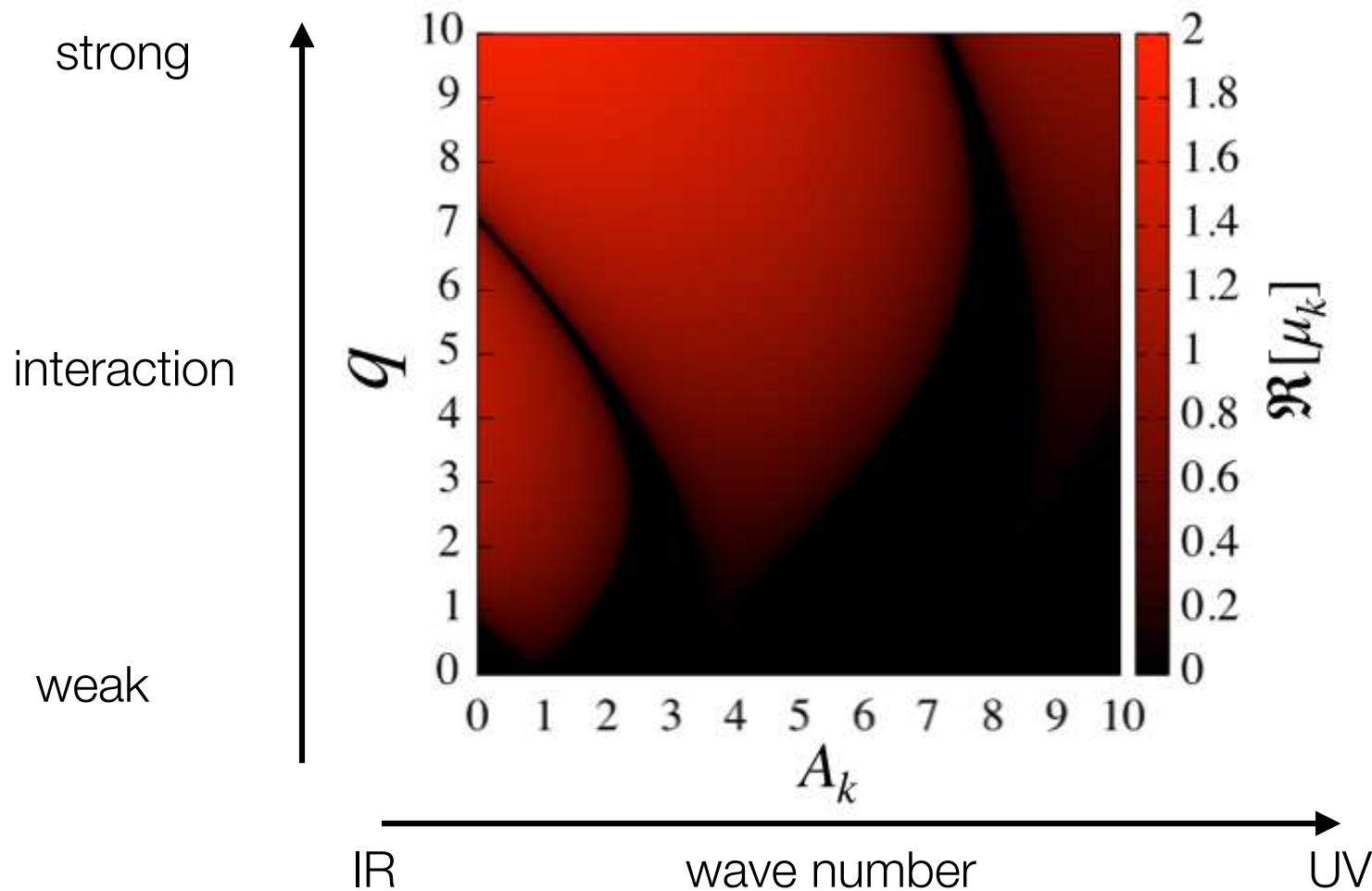
$$\Phi \sim (\delta\rho_{\text{rad}}/\rho_{\text{rad}}) \times (aH/k)^2 \propto (\delta\rho_{\text{rad}}/\rho_{\text{rad}}) \times a^{-2}$$

(in subhorizon)

Impact of self-interaction

Floquet theorem (1883)

axion in coherent oscillation \rightarrow resonant growth $\delta\phi \propto e^{\mu_k m t}$



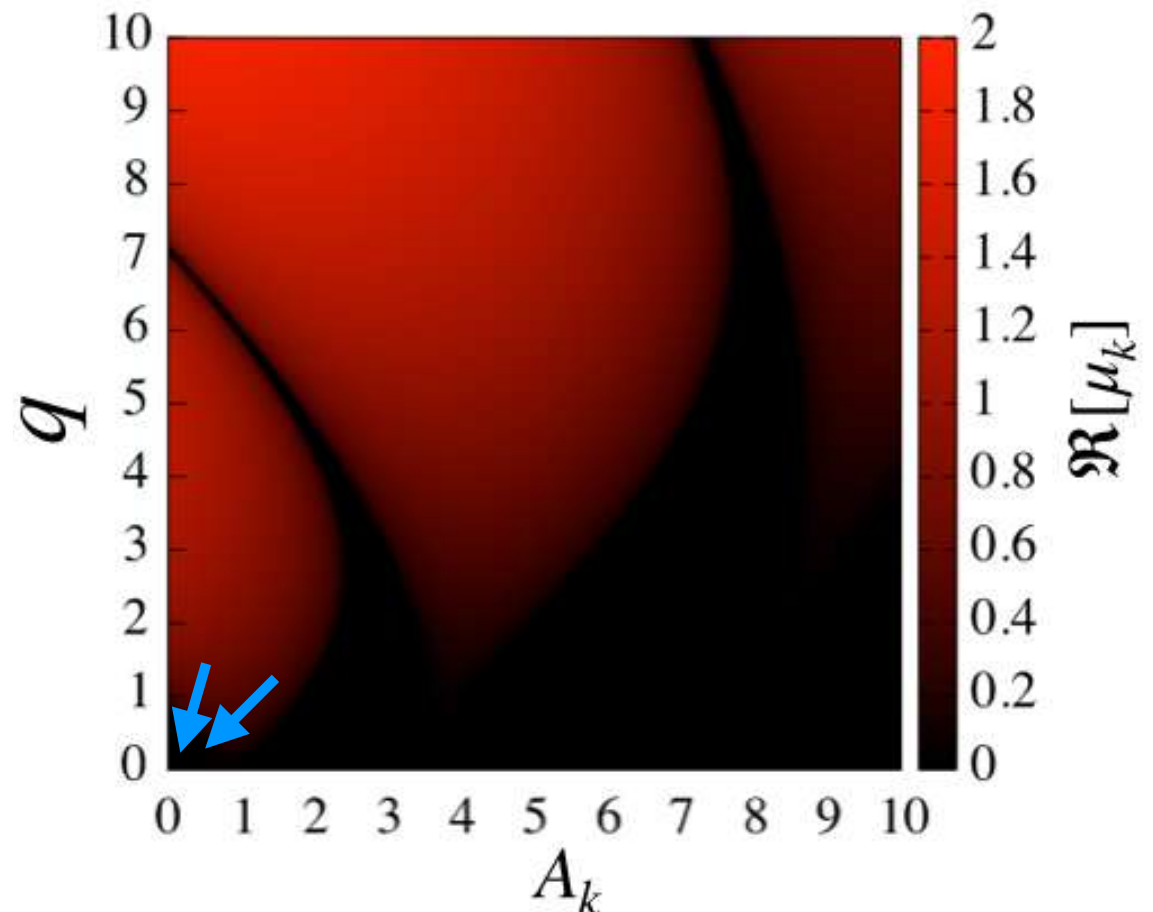
Parametric Resonance in cosmology

Cosmic expansion makes Bose enhancement inefficient.

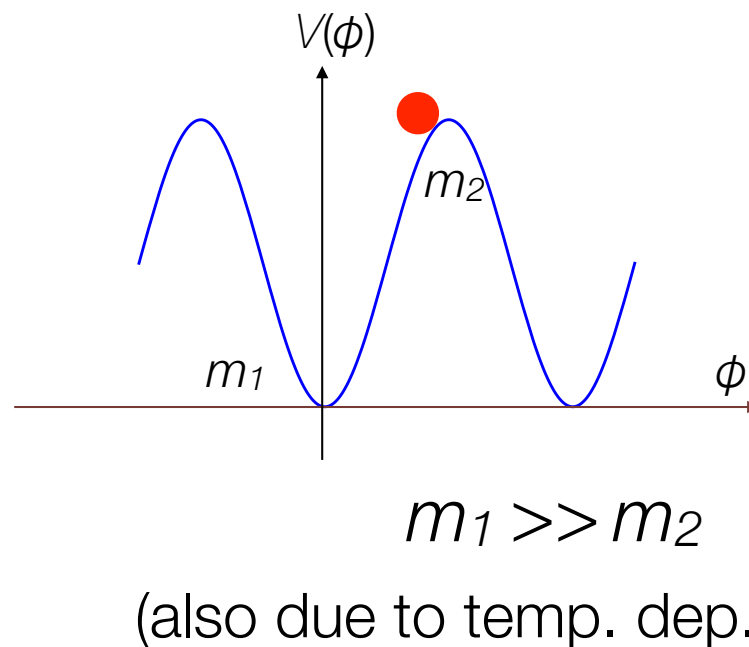
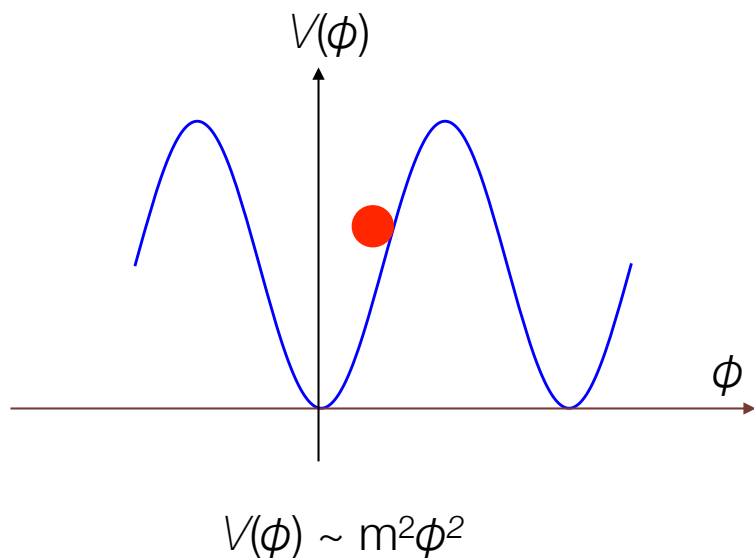
$$q(t) \propto \phi_{\text{bg}}^{\#}(t) \propto \rho_{\text{bg}}^{\#\#}(t) \quad (\#, \#\# > 0) \quad A_k(t) \sim \left(\frac{k}{a(t)m} \right)^2$$

Time scale of \swarrow
= $1/H_{\text{osc}}$

Time scale of exp.growth
= Period of oscillation
 $\mathcal{T} = 2\pi/m$



Parametric Resonance in cosmology



Onset of oscillation

$$H_{\text{osi}} \sim m$$

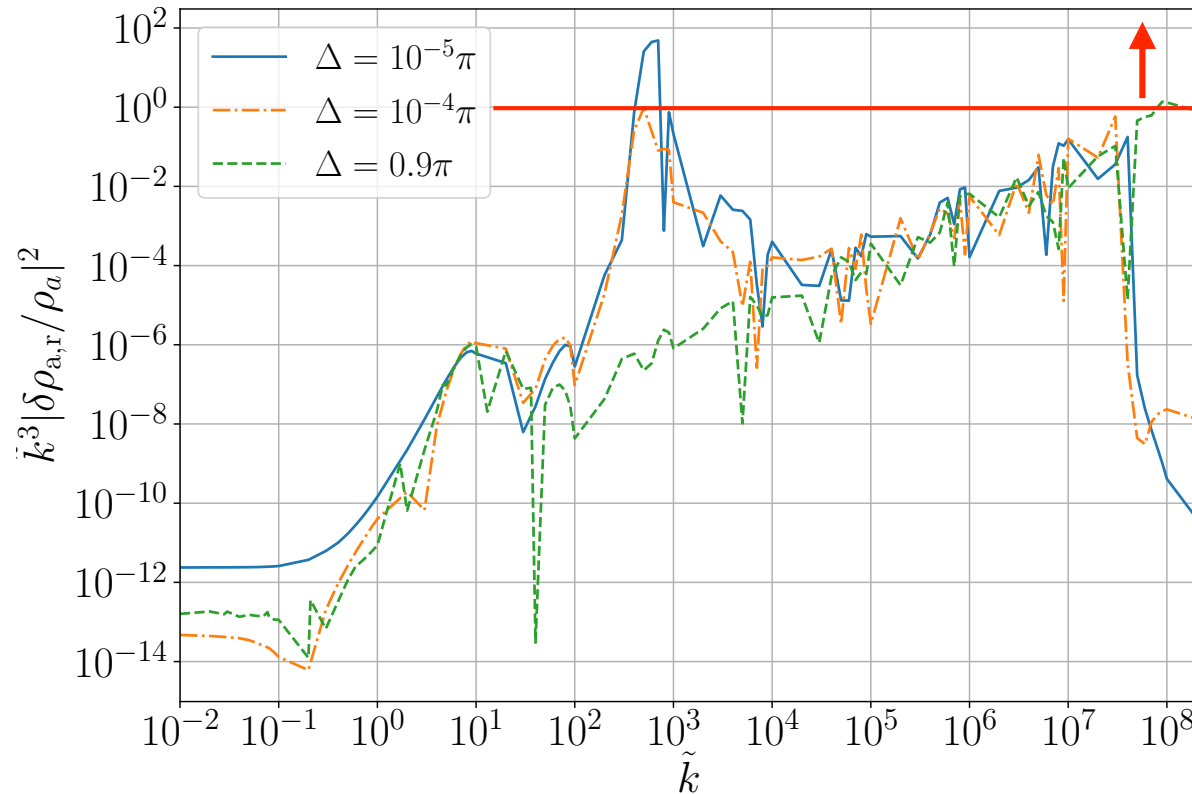
$$\mathcal{T} \sim 1/m \sim 1/H_{\text{osi}}$$

$$H_{\text{osi}} \sim m_2$$

$$\mathcal{T} \sim 1/m_1 \ll 1/H_{\text{osi}} \sim 1/m_2$$

Power spectrum of axion density perturbation

break down of perturbation



$$\tilde{k} \equiv k / (a_c H_c)$$

$$\mathcal{P}_\zeta = 2.1 \times 10^{-9}$$

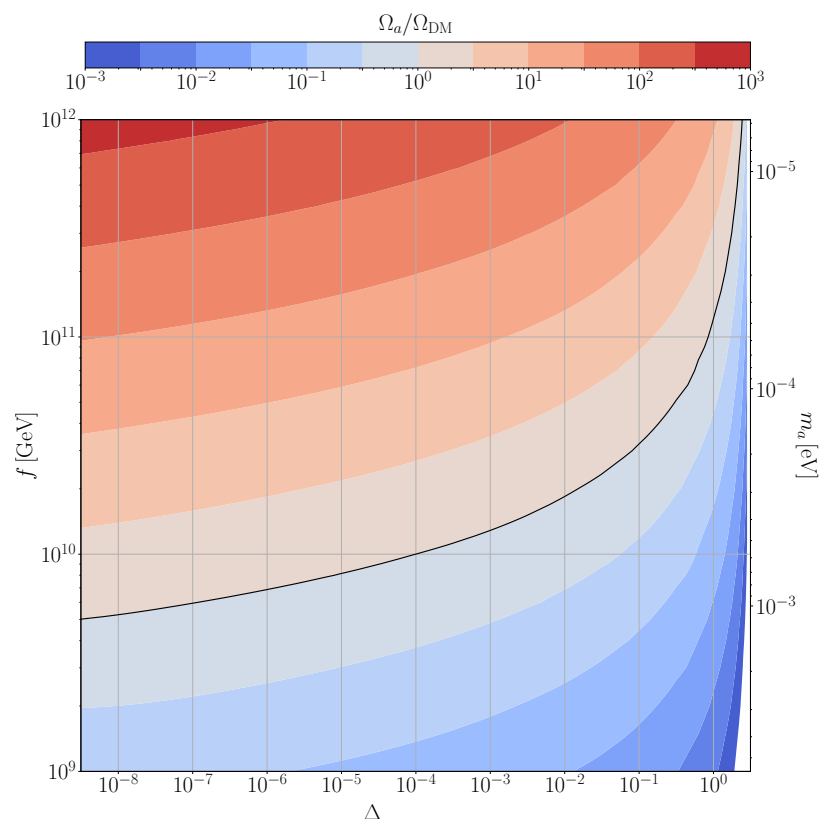
Kitajima, Kogai, Y.U. (in prep.)

- Radiation sources $\delta\rho_a$ only in sub horizon scales
- Resonance peak for fine-tuned IC, $\Delta \ll 1$

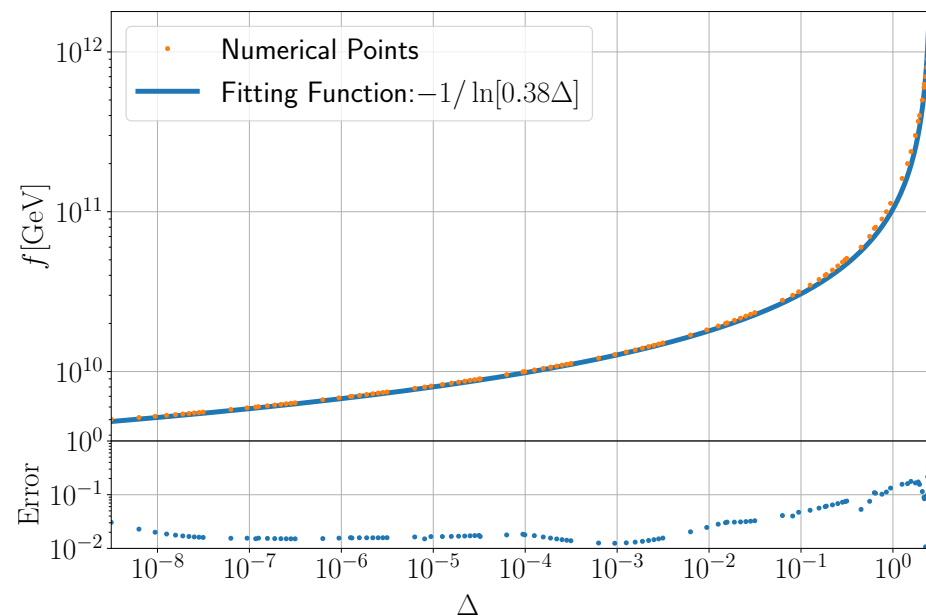
Axion abundance

For $\Delta \ll 1$, Ω_a is enhanced due to the delayed oscillation.

Turner(86), Lyth(92), ...



For $\Omega_a = \Omega_c$, i.e., DM=axion



Oscillation starts earlier as \searrow
(black line: $\Omega_a = \Omega_c$)

fitting

$$\left(\frac{f(\Delta)}{10^{11} [\text{GeV}]} \right) = -\frac{1}{[\log(0.38\Delta)]}$$

Kitajima, Kogai, Y.U. (in prep.)

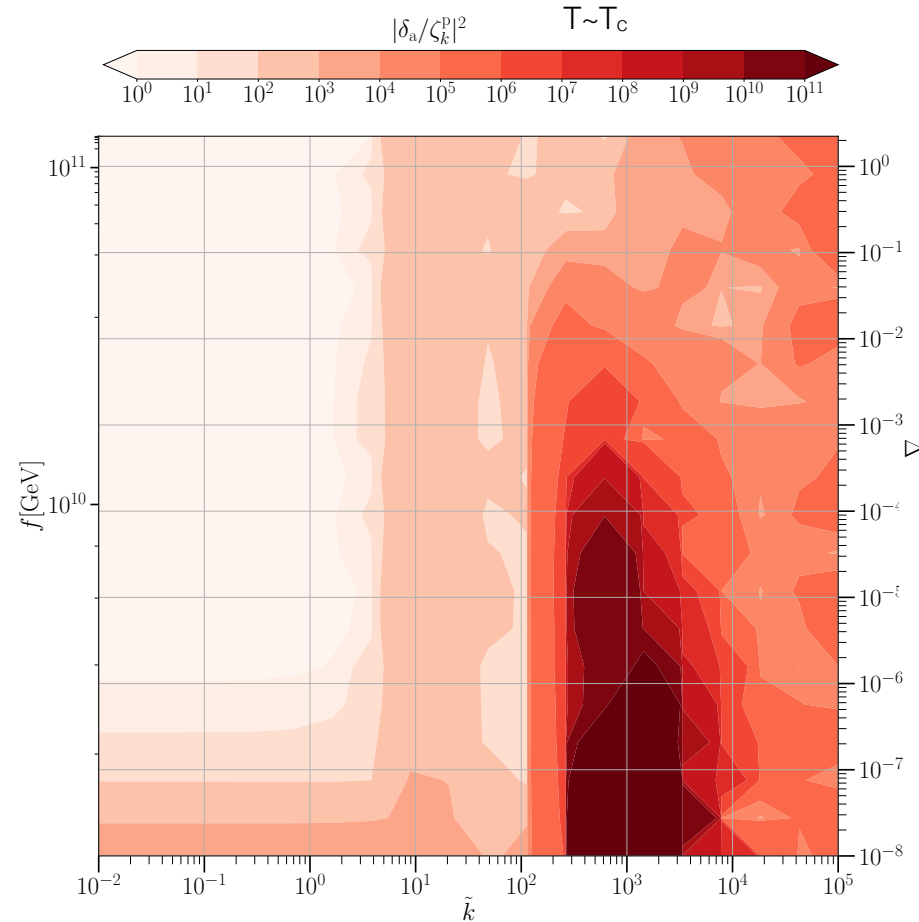
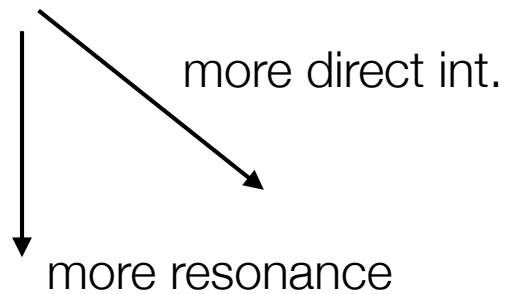
$$\Delta \equiv \pi - \phi_i / f$$

AMC formation

Kitajima, Kogai, Y.U. (in prep.)

Enhancement at $T=T_c$

$$\times \mathcal{P}_\zeta \rightarrow \sim \langle (\delta\rho_a/\rho_a)^2 \rangle$$



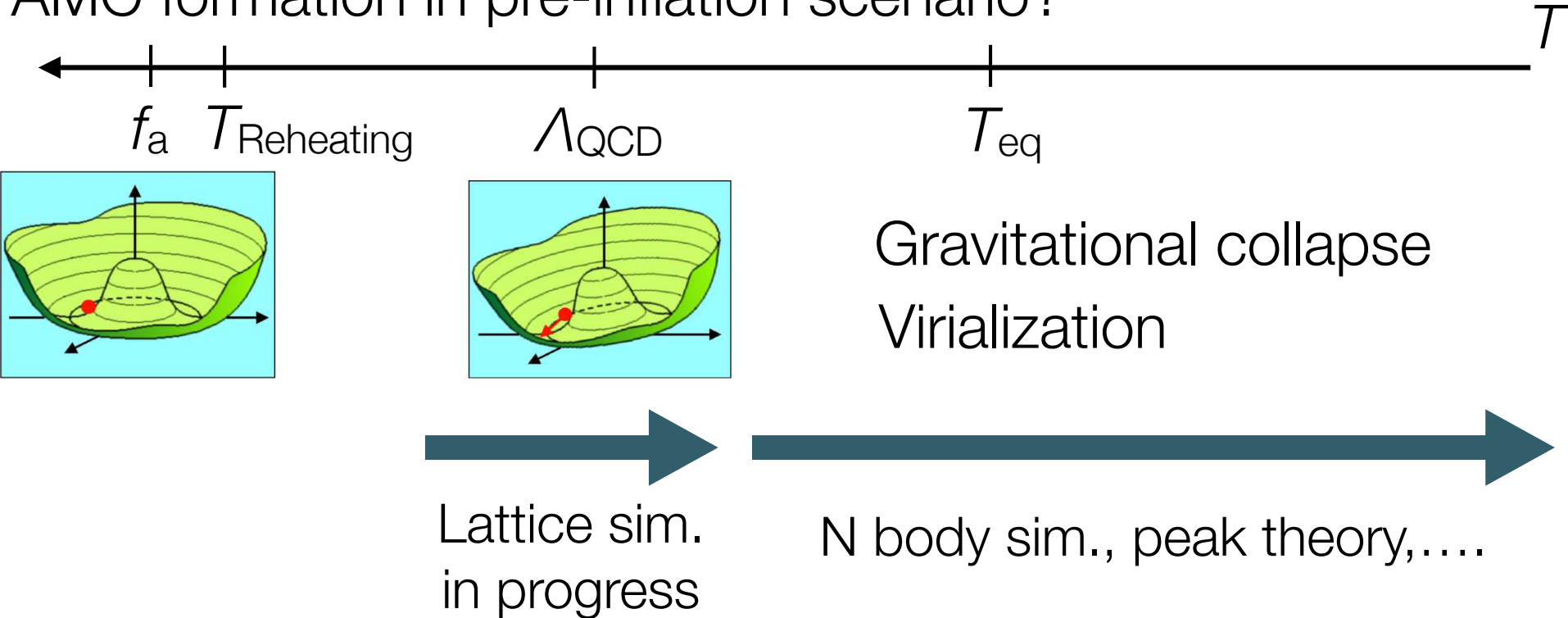
$\delta\rho_a/\rho_a$ reaches $O(1)$ at $T=T_c$,

1) $\mathcal{P}_\zeta > 10^{-3}$ at $k \sim a_c H_c$

2) For $\mathcal{P}_\zeta \sim 10^{-9}$, fine-tuned IC $w/\Delta < 10^{-3}$

Future issues

AMC formation in pre-inflation scenario?



1) $\mathcal{P}_\zeta > 10^{-3}$

if $\mathcal{P}_\zeta > 10^{-2}$



$O(M_\odot)$ PBH

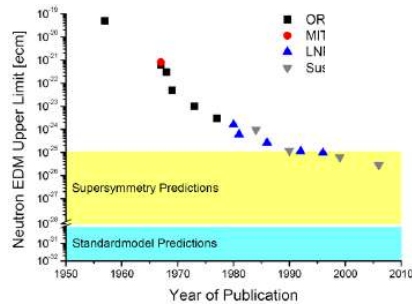
2) Fine-tuned IC

* More realistic input from QCD physics

If time permits....

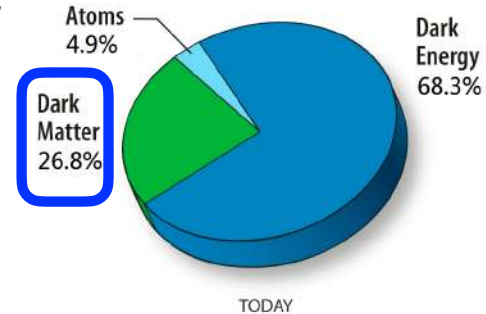
ALP dark matter

×



like
particle

✓



nHz GW signature from ALP dark matter

w/Kitajima (Tohoku), Soda (Kobe)

Phys. Rev. Lett. 126. (2021) 12, 121301
Kitajima, Soda, Y.U.

partly based on JCAP 10 (2018) 014

EPJC 78 (2018) 9, 779

ALP dark matter w/GUT scale f

axion w/ $f \sim 10^{15-16} \text{GeV}$ ($H_{\text{inf}} < f$)

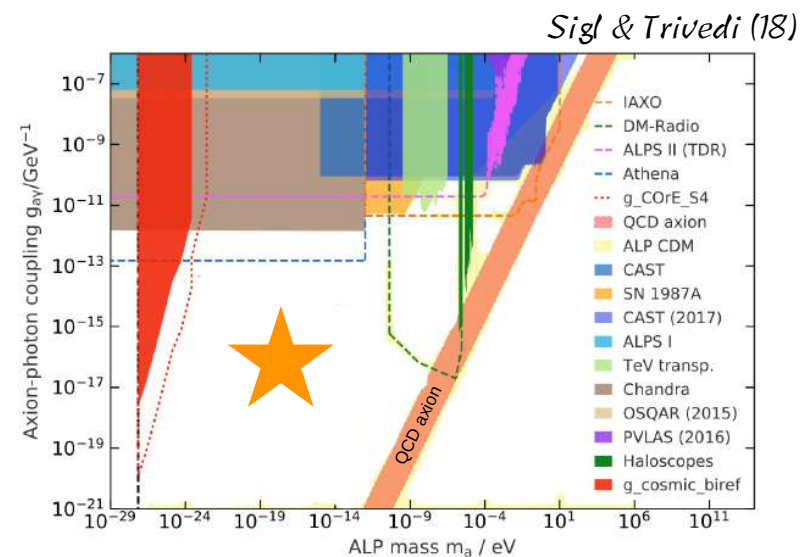
anthropic window



Model building

Surceek & Witten (06)

Weakly coupled $\propto 1/f$



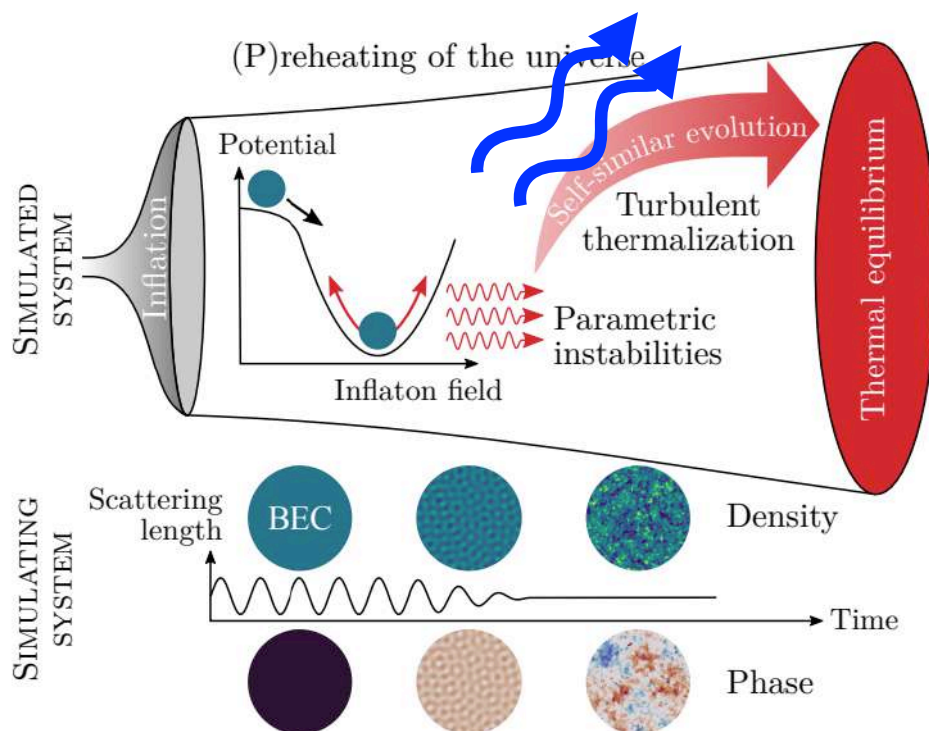
non-thermal production $\Omega_a h^2 \sim 2.8 \left(\frac{g_{*osc}}{10}\right)^{-1/4} \left(\frac{m}{10^{-14} \text{eV}}\right)^{1/2} \left(\frac{f\theta_i}{10^{16} \text{GeV}}\right)^2$

Recall reheating

Transition from inflation to thermalized Universe

Traschen, Brandenberger (90), Kofman, Linde, Starobinsky (97)
Khlebnikov, Tkachev (96),

Gravitational Waves (GWs)



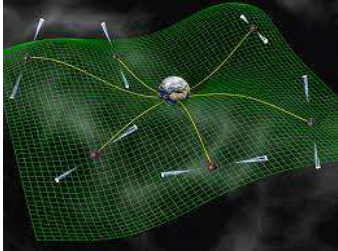
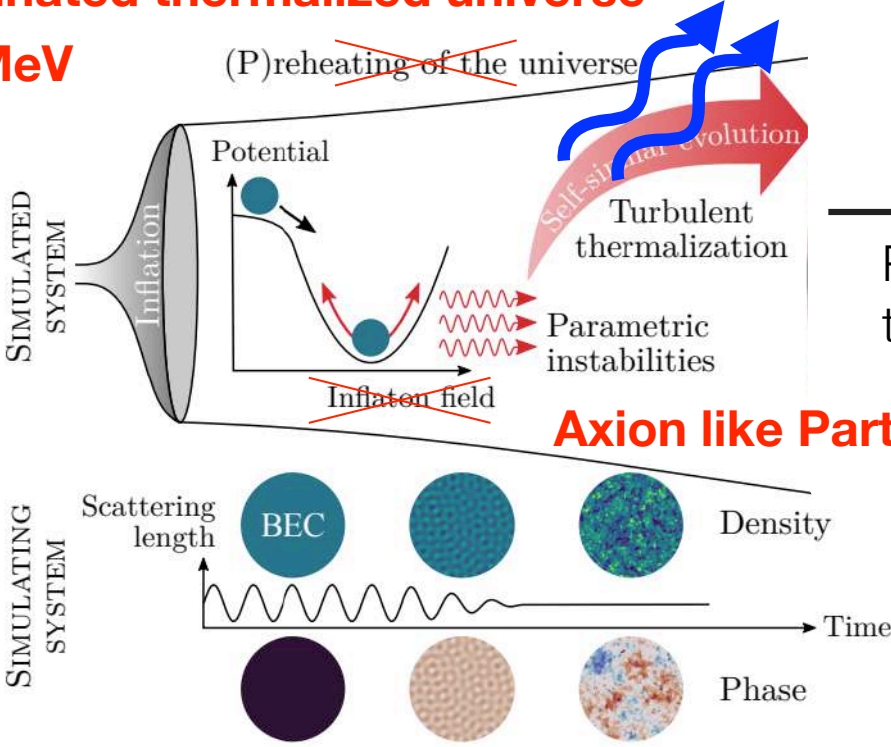
Typically very high frequency

Drawing from Chatrchyan et al. (20)

Bottom-line story

Detectable GWs via pulsar observations

Radiation dominated thermalized universe
 e.g. $T = \text{MeV}$

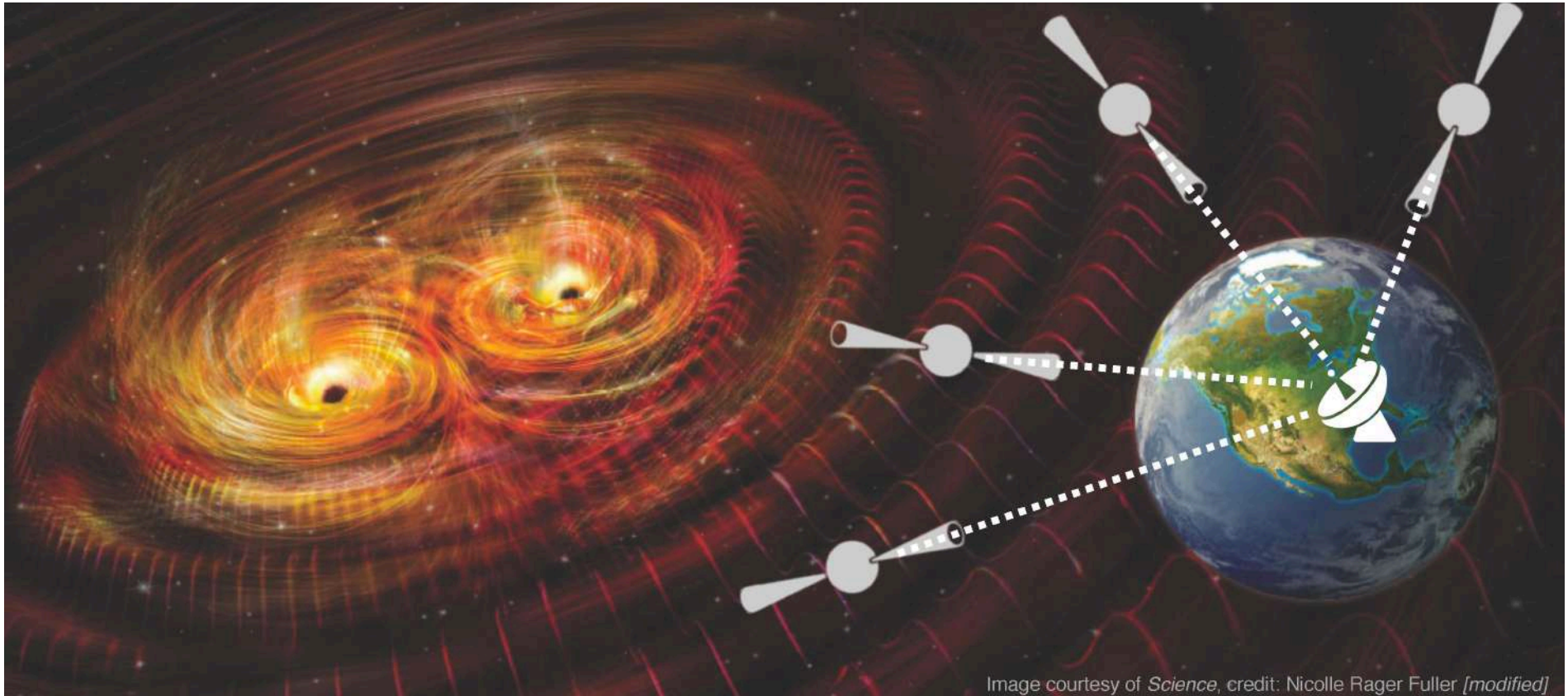


Relic ALP becomes DM, dominating the Universe around $T \sim \text{eV}$

Axion like Particle (ALP) dark matter

Drawing from Chatrchyan et al. (20)

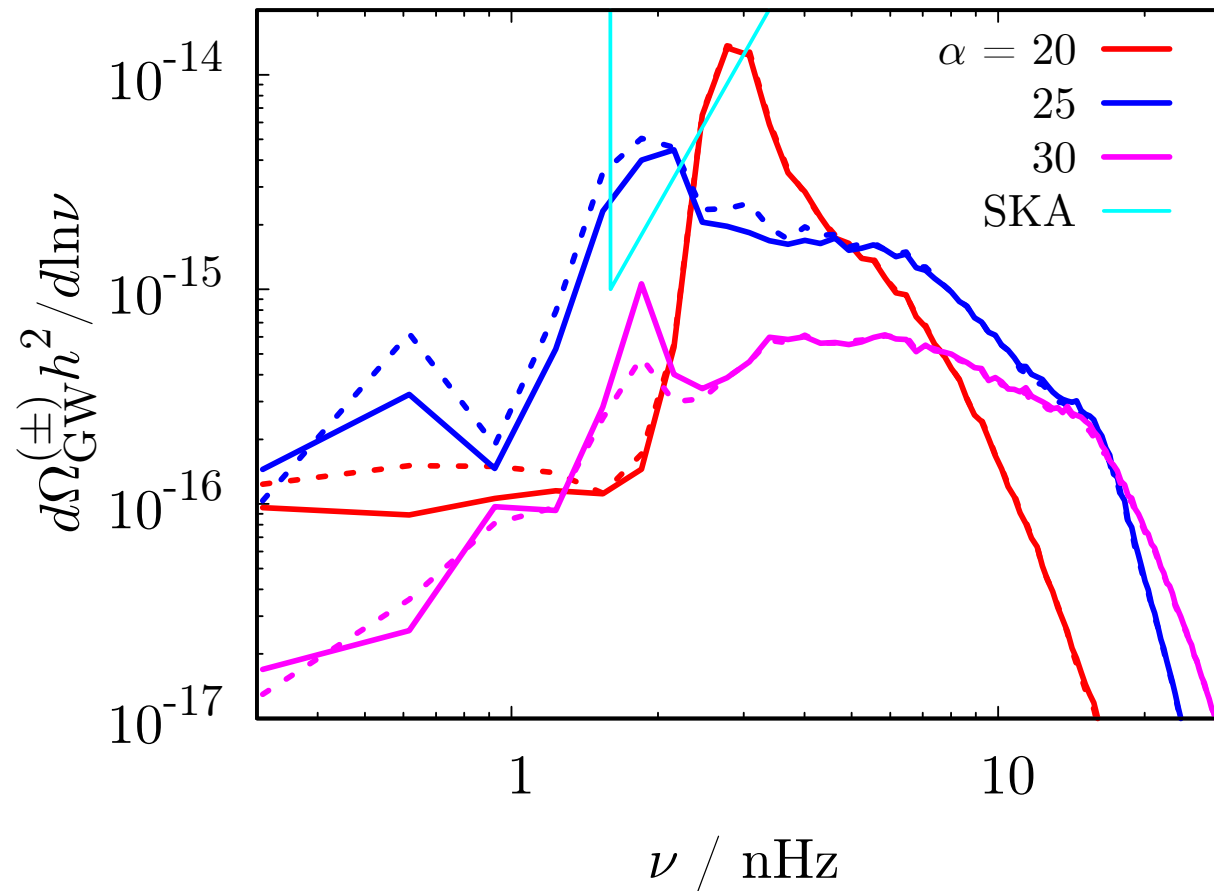
nHz GW from pulsar timing array (PTA)



Time residual of periodic emissions from pulsars
(rapidly rotating NSs)

nHz GW signature from ALP dark matter

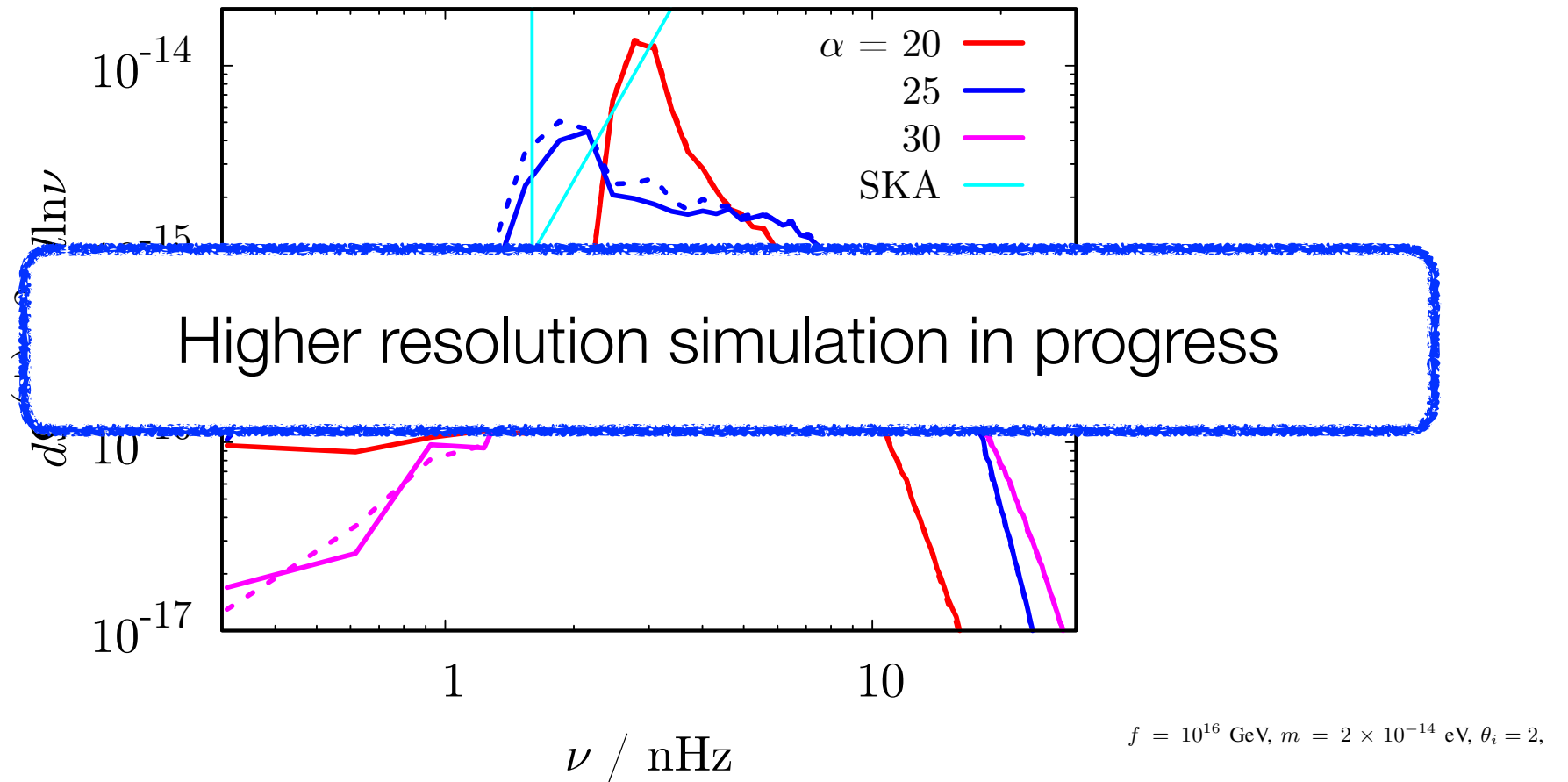
Kitajima, Soda & Y.U. (20)



since $\propto \frac{\phi}{f} F \tilde{F}$, GWs are circularly polarized!

nHz GW signature from ALP dark matter

Kitajima, Soda & Y.U. (20)



since $\propto \frac{\phi}{f} F \tilde{F}$, GWs are circularly polarized!

ALP dark matter w/GUT scale f

axion w/ $f \sim 10^{15-16} \text{GeV}$



Weakly coupled w/SM

$$\propto 1/f$$

Model building *Surcek & Witten (06)*

PTA window *Kitajima, Soda, YU (20)*

