

素粒子物理学の進展、基礎物理学研究所、2012年7月

# Compact Supersymmetry

飛岡 幸作

Kavli IPMU, 東京大学

arXiv:1206.4993[hep-ph]

Collaboration with

Hitoshi Murayama, Yasunori Nomura and Satoshi Shirai

# Table of my talk

## 1. Introduction

I. Current Situation

II. Weaker constraint by LHC on Compressed Scenario

## 2. Compact Supersymmetry model

I. Scherk-Schwarz mechanism ~ Radion Mediation

II. Model setup

III. Phenomenology

## 3. Summary

# Current Situation

1. Higgs like particle is discovered around 125 GeV

## *Implication to Supersymmetry*

Tightly constraints MSSM, need to boost the Higgs mass by radiative corrections:

- High scale of  $M_{\text{SUSY}}$  (Split) and/or
- Large  $A$  term


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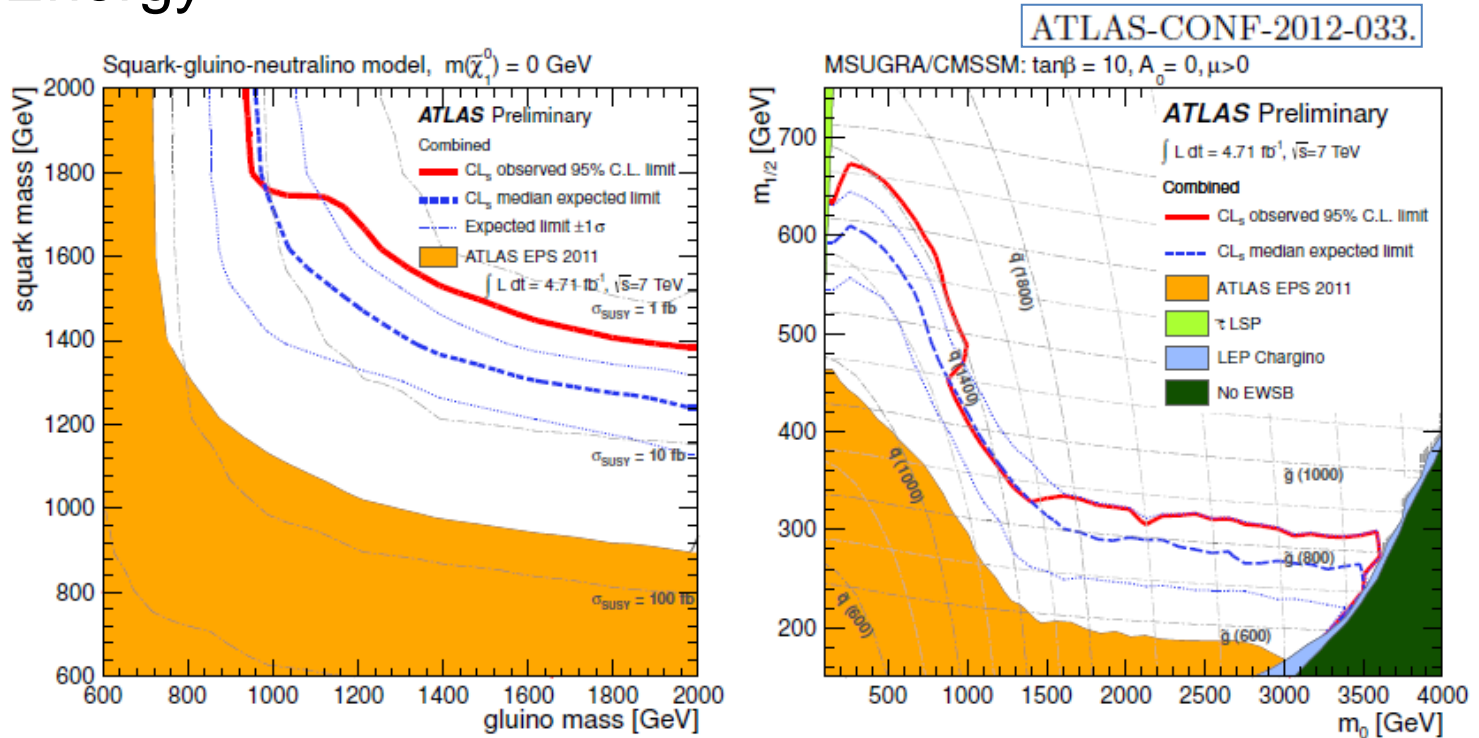
- High scale of  $M_{\text{SUSY}}$  (Split) and/or
- Large  $A$  term



Measure  $M_{\text{SUSY}}$   
arXiv:1207.3608  
Sato, Shirai, Tobioka

# Current Situation

## 2. No signal of Beyond Standard Model with Missing Energy



$$M_{\text{gluino}} = M_{\text{squark}} > 1.4 \text{ TeV for } 5 \text{ fb}^{-1}$$

$$(> 1 \text{ TeV for } 1 \text{ fb}^{-1})$$

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### *Implication to Supersymmetry*

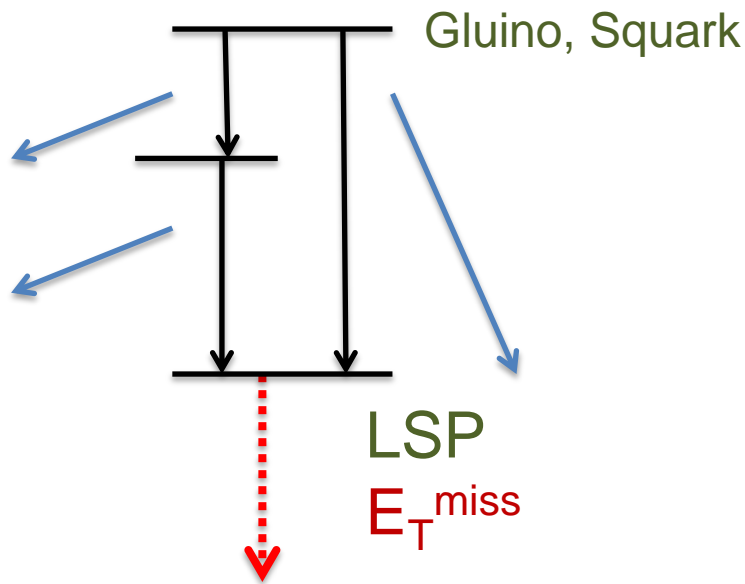
- Split -> Small cross section
- R-parity violation
- Compressed Spectrum -> Small q-value (this talk)

# Current Situation

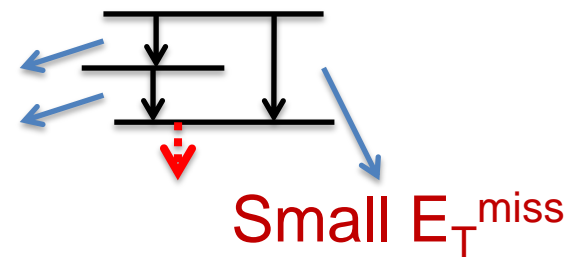
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### MSUGRA like



### Compressed

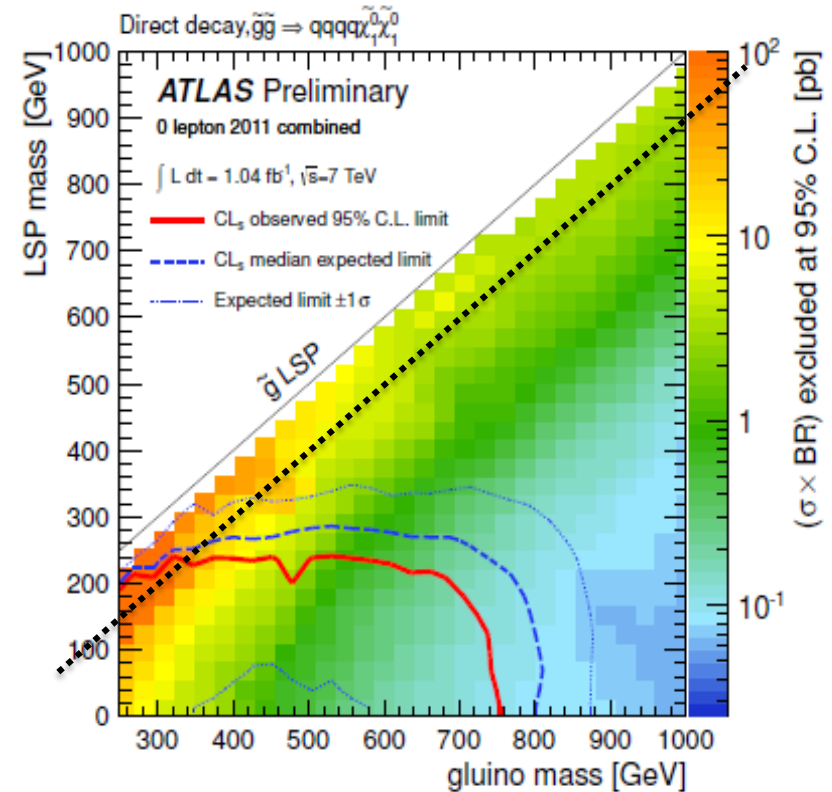


Difficult to extract the signal from the background (ttbar, W+jets)

# Weaker constraint by LHC in Compressed scenario

Experimentalist's Analysis  
ATLAS-CONF-2011-155

- Gluino+LSP model  
 $M_{\text{gluino}} \sim 350 \text{ GeV}$
- Squark+LSP model  
 $M_{\text{squark}} \sim 250 \text{ GeV}$   
( $\Delta m > 100 \text{ GeV}$ )
- Gluino+Squark+LSP model  
 $M_{\text{gluino}} = M_{\text{squark}} \sim 400 \text{ GeV}$   
( $\Delta m = 5 \text{ GeV}$ !)



Actually much weaker constraint,  $400 \text{ GeV} \ll 1 \text{ TeV}$



# Weaker constraint by LHC in Compressed scenario

## Theorist's Analysis

LeCompte, S. P. Martin [arXiv:1111.6897]

Phenomenological model + ATLAS Multijet search

$M_{\text{gluino}} \simeq M_{\text{squark}} \sim 650 \text{ GeV}$  when  $\Delta m \geq 100 \text{ GeV}$

$$M_1 = \left( \frac{1 + 5c}{6} \right) M_{\tilde{g}}$$
$$M_2 = \left( \frac{1 + 2c}{3} \right) M_{\tilde{g}}$$

## Recent analysis

H. K. Dreiner et al. [arXiv: 1207.1613]

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Mostly based on Phenomenological models,  
Scherk-Schwarz SUSY breaking generates **universal soft masses**

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# Scherk-Schwarz mechanism

[Scherk and Schwarz (1979)]

- 5D Minimal SUSY (corresponding to  $\mathcal{N}=2$  in 4D)
- Geometry:  $S^1/Z_2$  (chiral for zero mode,  $\mathcal{N}=1$  in 4D)

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- Non-trivial boundary condition on  $SU(2)_R$  space  
breaks supersymmetry =Scherk-Schwarz mechanism

Non-trivial B.C. 
$$\begin{pmatrix} \lambda_1 \\ \lambda_2 \end{pmatrix} (x_\mu, y + 2\pi R) = e^{-2\pi i \alpha \sigma_2} \begin{pmatrix} \lambda_1 \\ \lambda_2 \end{pmatrix} (x_\mu, y)$$

$y$ : 5<sup>th</sup> dimensional coordinate  
 $R$ : radius of extra dimension

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Resulting KK decomposition 
$$\begin{pmatrix} \lambda_1 \\ \lambda_2 \end{pmatrix} (x_\mu, y) = \sum_{n=0}^{\infty} e^{-i\alpha\sigma_2 y/R} \begin{pmatrix} \lambda_1^{(n)}(x_\mu) \cos[ny/R] \\ \lambda_2^{(n)}(x_\mu) \sin[ny/R] \end{pmatrix}$$

$$\supset \begin{pmatrix} \lambda_1^{(0)}(x_\mu) \cos[\alpha y/R] \\ \lambda_1^{(0)}(x_\mu) \sin[\alpha y/R] \end{pmatrix}$$

5D derivative generates  
(soft) masses in 4D

$$m_n = \begin{cases} \alpha/R & \text{zero mode} \\ (\alpha \pm n)/R & \text{non-zero modes} \end{cases}$$

# Field Properties in 5th dimension

Fields:  $V, \chi, \Phi, \Phi^c$

Higgs localized at  $y=0$ :  $H_u(x), H_d(x)$

$V$  : Vector superfield

$\chi$  : Adjoint chiral superfield

$\Phi^{(c)}$ : Hypermultiplet of matter fields

$A_\mu, \lambda_1$

$\lambda_2, A_5, \Sigma$

$\phi^{(c)}, \psi^{(c)}$

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Inversion

$$\begin{pmatrix} V(x, -y) \\ \chi(x, -y) \end{pmatrix} = \begin{pmatrix} V(x, y) \\ -\chi(x, y) \end{pmatrix}$$

$$\begin{pmatrix} \Phi(x, -y) \\ \Phi^c(x, -y) \end{pmatrix} = \begin{pmatrix} \Phi(x, y) \\ -\Phi^c(x, y) \end{pmatrix}$$

Translation (SS mechanism)

For  $SU(2)_R$  doublets, common twist

$$\begin{pmatrix} \lambda_1(x, y + 2\pi R) \\ \lambda_2(x, y + 2\pi R) \end{pmatrix} = e^{-2\pi\alpha\sigma_2} \begin{pmatrix} \lambda_1(x, y) \\ \lambda_2(x, y) \end{pmatrix}$$

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same for gravitinos

For others,  $X(x, y + 2\pi R) = X(x, y)$



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For others,  $X(x, y + 2\pi R) = X(x, y)$

$$m_{1/2, \text{squark, slepton}} = \frac{\alpha}{R}$$

Common soft mass

# Radion Mediation ~ SS mechanism

Radion mediation: SUSY breaking by the Radion superfield vev

$$T = R + iB_5 + \theta\Psi_R^5 + \theta^2 F_T$$

~Dynamical realization of Scherk-Schwarz mechanism

[D.Marti and A.Pomarol(2001), D.Kaplan and N. Weiner(2001)]

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◆ Gauge sector

$$S_5 = \int dx^4 dy \left[ \frac{1}{4g_5^2} \int d^2\theta \left( \frac{T}{R} \right) W^\alpha W_\alpha + \text{h.c.} + \frac{1}{g_5^2} \int d^4\theta \frac{2R}{T + T^\dagger} \left( \partial_5 V - \frac{\chi + \chi^\dagger}{\sqrt{2}} \right)^2 \right]$$

◆ Matter sector

$$S_5 = \int dx^4 dy \left[ \frac{1}{4g_5^2} \int d^4\theta \frac{T + T^\dagger}{2R} \left( \Phi^\dagger e^{-V} \Phi + \Phi^c e^V \Phi^{c\dagger} \right) + \int d^2\theta \Phi^c \left( \partial_5 - \frac{\chi}{\sqrt{2}} \right) \Phi + \text{h.c.} \right]$$

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Radion vev:  $\langle T \rangle = R + F_T \theta^2$   $\frac{F_T}{2} = -\alpha$

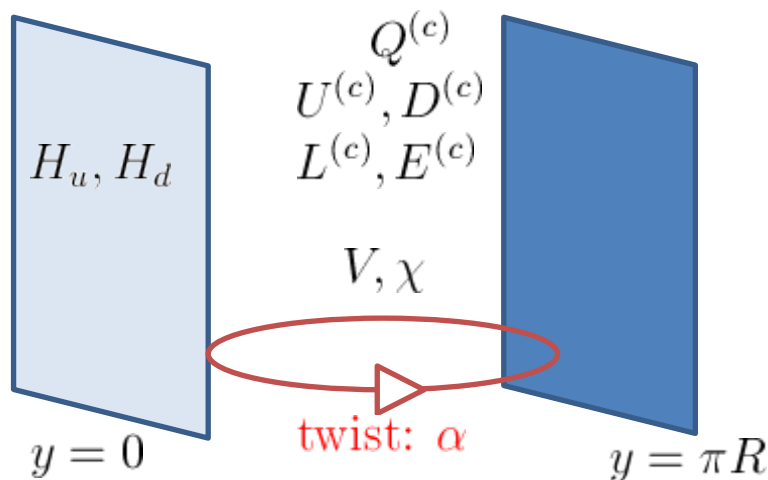
Canonically normalize:  $\Phi^{(c)} \rightarrow \left( 1 - \frac{F_T}{2R} \theta^2 \right) \Phi^{(c)}$ ,  $\chi \rightarrow \left( 1 + \frac{F_T}{2R} \theta^2 \right) \chi$

Actually this theory is identified as SSSB

# Compact Supersymmetry model

Higgs fields and Yukawa interactions are localized on the brane at  $y=0$

$$\mathcal{L}_{\text{brane}} = \delta(y) \int d^2\theta (y_U^{ij} Q_i U_j H_u + y_D^{ij} Q_i D_j H_d + y_E^{ij} L_i E_j H_d + \mu H_u H_d).$$



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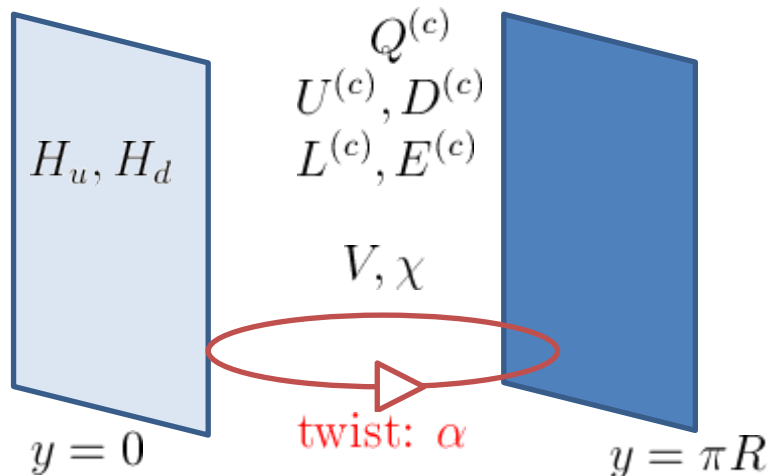
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Large A term is generated by the field redefinition

$$\Phi^{(c)} \rightarrow \left(1 - \frac{F_T}{2R} \theta^2\right) \Phi^{(c)}, \quad y_U^{ij} \underline{Q}_i U_j H_u \rightarrow \left(1 - \frac{F_T}{R} \theta^2\right) y_U^{ij} Q_i U_j H_u \quad \boxed{\frac{F_T}{2} = -\alpha}$$

$$= \left(1 + \frac{2\alpha}{R} \theta^2\right) y_U^{ij} Q_i U_j H_u$$



$$\boxed{A_0 = -\frac{2\alpha}{R}}$$

# Compact Supersymmetry model

- Take  $\alpha \ll 1$
- KK states ( $\sim n/R$ ) are decoupled  $\rightarrow$  MSSM at low energy
- Compact parameter set rather than CMSSM:

At tree level and  
at scale  $\sim 1/R$ ,

$$M_{1/2} = \frac{\alpha}{R}, \quad m_{\tilde{Q}, \tilde{U}, \tilde{D}, \tilde{L}, \tilde{E}}^2 = \left(\frac{\alpha}{R}\right)^2, \quad m_{H_u, H_d}^2 = 0,$$
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- Radiative corrections from at and above  $1/R$  are under control because of symmetries of higher dimensions
- Calculated threshold corrections to the Higgs mass parameters

$$\delta m_{H_u}^2 = \left( -\frac{33y_t^2}{8\pi^2} + \frac{9(g_2^2 + g_1^2/5)}{16\pi^2} \right) \left(\frac{\alpha}{R}\right)^2,$$
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Only three parameters!



$$\frac{1}{R}, \quad \frac{\alpha}{R}, \quad \mu.$$



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Only three parameters!



$$\frac{1}{R}, \quad \frac{\alpha}{R}, \quad \mu.$$

- No physical phase  $\rightarrow$  ✓ CP
- Geometry is universal  $\rightarrow$  ✓ Flavor

## Brane-localized kinetic terms and cutoff

- Radiative corrections from above  $1/R$  generates boundary kinetic terms from dimensional analysis

$$\frac{\delta M_{1/2}}{M_{1/2}}, \frac{\delta m_{\tilde{f}}^2}{m_{\tilde{f}}^2}, \frac{\delta A_0}{A_0} \approx O\left(\frac{1}{16\pi^2} \ln(\Lambda R)\right).$$

- Assume the tree level contributions are same size of radiative ones

$$\sim \frac{y_t^2}{16\pi^2} O(1)$$

Effective theory with tree level estimation of soft parameters is valid for  $\Lambda R \ll 16\pi^2$

## Power of $\mathcal{N}=2$

- $S^1/Z_2$  orbifolding makes zero modes chiral, but higher KK modes consists  $\mathcal{N}=2$  multiplets
- No wavefunction renormalization of hypermultiplet in  $\mathcal{N}=2$  SUSY

$$S_5 = \int dx^4 dy \left[ \frac{1}{4g_5^2} \int d^4\theta \frac{T + T^\dagger}{2R} \left( \Phi^\dagger e^{-V} \Phi + \Phi^c e^V \Phi^{c\dagger} \right) + \int d^2\theta \Phi^c \left( \partial_5 - \frac{\chi}{\sqrt{2}} \right) \Phi + \text{h.c.} \right]$$

- Even log divergences are cancelled out for each KK mode ( $n>0$ )
- Only MSSM( $n=0$ ) particles give log divergences

## Gravitino mass

□ Obviously the SU(2)R doublets should have same soft mass from their 5d derivatives

□ SUSY breaking is from Radion

•GR action

$$M_{pl}^2 \mathcal{R} \rightarrow M_{pl}^2 \left( \frac{T + T^\dagger}{R} \right)^2$$
$$\left( g_{55} \rightarrow \frac{T + T^\dagger}{R} \right)$$

•Gravitino mass

Radion should be canonically normalized

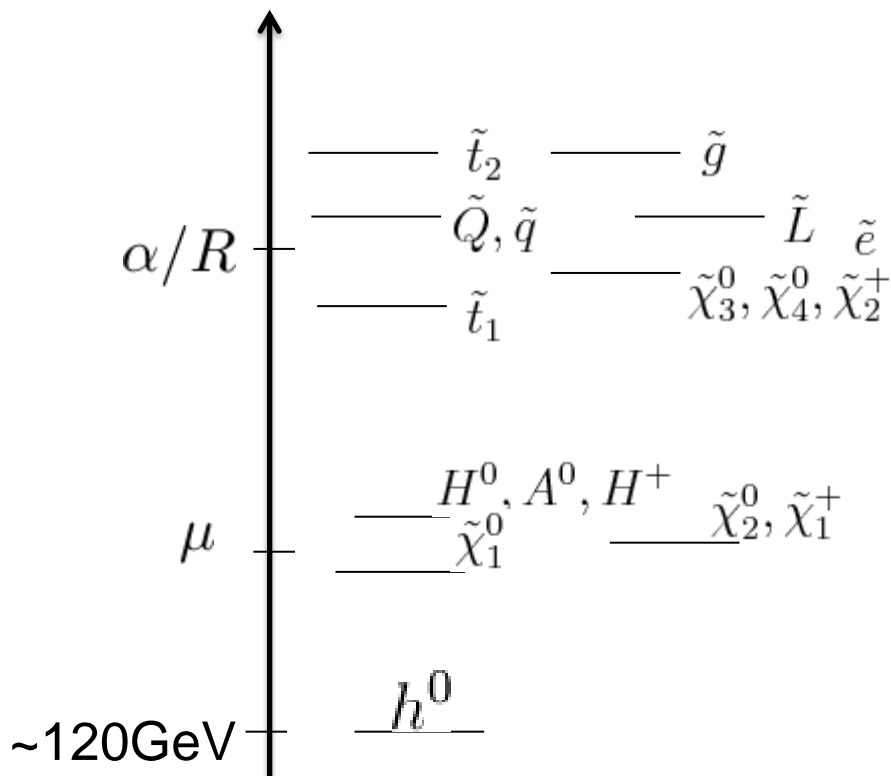
$$M_{3/2} \sim \frac{\langle \mathcal{F} \rangle}{M_{pl}} \sim \frac{(F_T/R)M_{pl}}{M_{pl}}$$

$$M_{1/2, \text{ squark, slepton}} = M_{3/2} = \frac{\alpha}{R}$$

# Spectrum

Point1:  $\alpha/R = 1400$  GeV,  $1/R = 10^4$  GeV

Point2:  $\alpha/R = 800$  GeV,  $1/R = 10^5$  GeV



Higgsino like LSP

$\sim \mu$

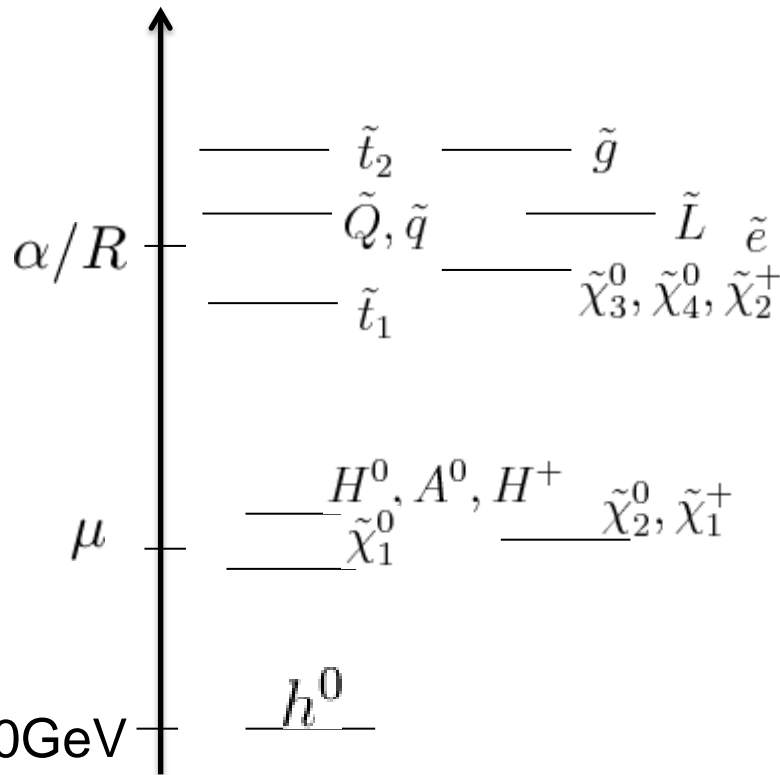
Particle	Point1	Point2	Particle	Point1	Point2
$\tilde{g}$	1494	949	–	–	–
$\tilde{u}_L$	1467	939	$\tilde{u}_R$	1459	925
$\tilde{d}_L$	1469	942	$\tilde{d}_R$	1458	924
$\tilde{b}_2$	1460	924	$\tilde{b}_1$	1430	875
$\tilde{t}_2$	1557	988	$\tilde{t}_1$	1267	681
$\tilde{\nu}$	1411	822	$\tilde{\nu}_\tau$	1410	822
$\tilde{e}_L$	1413	826	$\tilde{e}_R$	1406	812
$\tilde{\tau}_2$	1417	823	$\tilde{\tau}_1$	1402	809
$\tilde{\chi}_1^0$	767	630	$\tilde{\chi}_2^0$	777	671
$\tilde{\chi}_3^0$	1384	755	$\tilde{\chi}_4^0$	1410	821
$\tilde{\chi}_1^\pm$	771	642	$\tilde{\chi}_2^\pm$	1409	817
$h^0$	125	120	$H^0$	819	718
$A^0$	819	717	$H^\pm$	822	722

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Higgsino like LSP

$\sim \mu$

Spectra are available!

<http://www-theory.lbl.gov/~shirai/compactSUSY.php>

# Spectrum and LHC constraint

More compressed as  $\mu \rightarrow \alpha/R$ , i.e. larger  $1/R$  ( $Q_0$ )

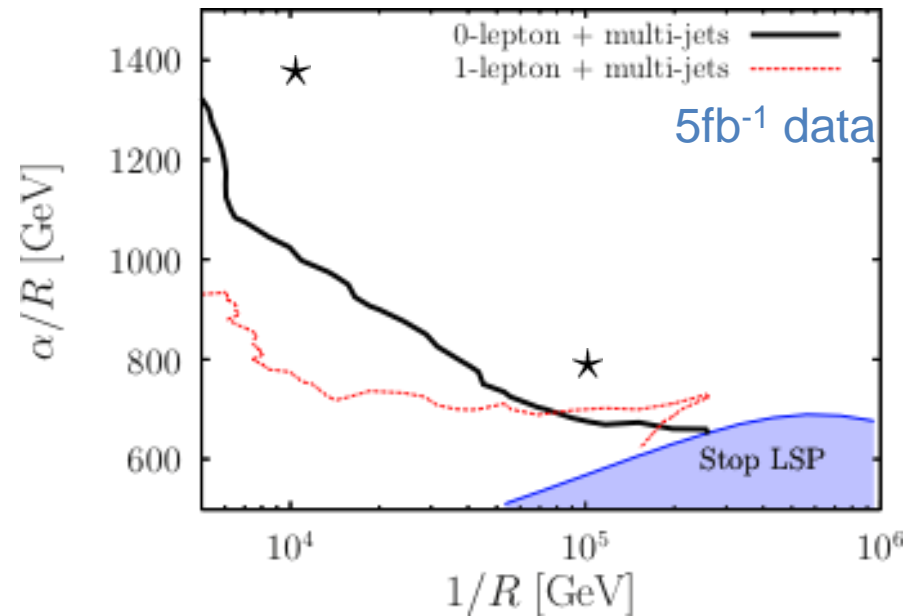
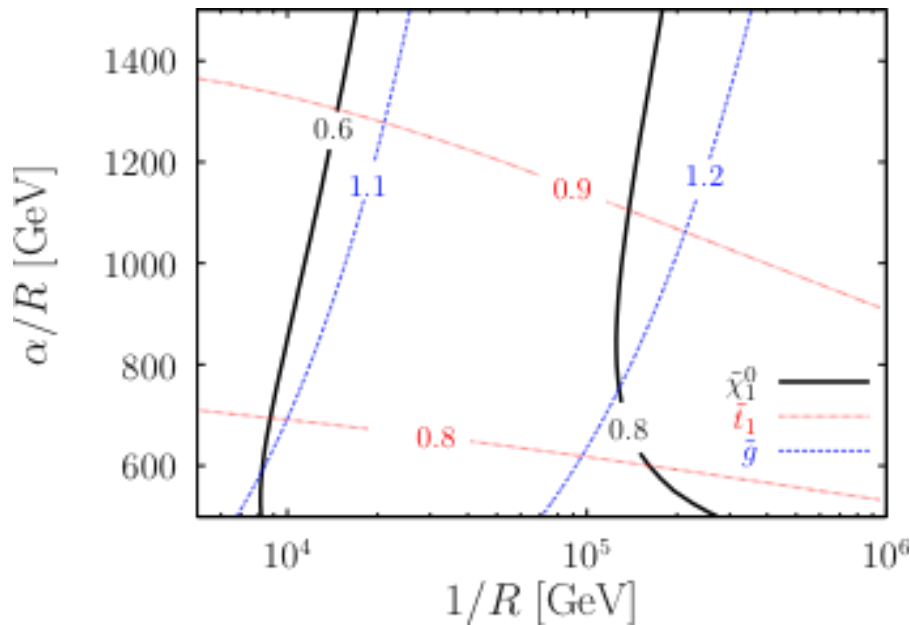
$$m_{H_u}^2 + |\mu|^2 \simeq m_z^2 \cos 2\beta/2 \quad (\text{EWSB})$$

$$\delta m_{H_u}^2 = \left( -\frac{33y_t^2}{8\pi^2} + \frac{9(g_2^2 + g_1^2/5)}{16\pi^2} \right) \left( \frac{\alpha}{R} \right)^2$$

$$\Delta m_{H_u}^2 \sim \frac{3y_t^2}{8\pi^2} (m_{Q_3}^2 + m_{u_3}^2 + |A_0|^2) \Delta \ln \left( \frac{Q}{Q_0} \right)$$

• Mass in unit of  $\alpha/R$

• Collider bound



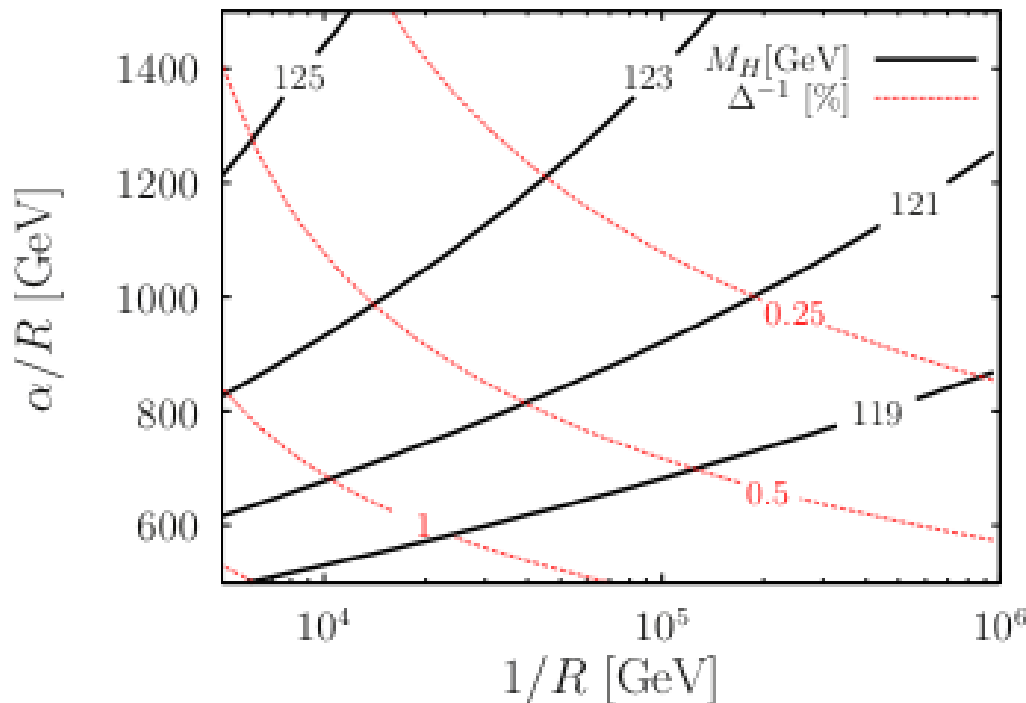
➔ Compressed

Spectra are available!

<http://www-theory.lbl.gov/~shirai/compactSUSY.php>

# Higgs mass and tuning

$m_t = 173.2 \text{ GeV}$



□ Theoretical error of Higgs mass is not small

$$|\Delta M_H| \approx 2 - 3 \text{ GeV}$$

• Also deviation from top mass

$$\Delta m_t = \pm 0.9 \text{ GeV}$$

$$\Rightarrow \Delta M_H \approx \pm 1 \text{ GeV}$$

□ Fine tuning of sub-% level mainly from  $\mu$

➔ better than CMSSM

$$\Delta^{-1} \equiv \min_x |\partial \ln m_Z^2 / \partial \ln x|^{-1} \text{ with } x = \alpha, \mu, 1/R, y_t, g_3, \dots$$

Spectra are available!

<http://www-theory.lbl.gov/~shirai/compactSUSY.php>



# Possible NMSSM extension

Work in progress [Murayama, Nomura, Shirai, KT]

## Singlet Hypermultiplet in the bulk

$$W_{NMSSM} = (\lambda S H_u H_d + \frac{1}{3} \kappa S^3) \delta(y)$$

□ Again, soft parameters of singlet are automatically determined by SS mechanism

$$V_{soft}^{NMSSM} = (a_\lambda S H_u H_d + \frac{1}{3} a_\kappa S^3 + \text{h.c.}) + m_s |S|^2$$

where

$$a_\lambda = -\frac{\alpha}{R}, \quad a_\kappa = -\frac{3\alpha}{R}, \quad m_s^2 = \left(\frac{\alpha}{R}\right)^2,$$

- No CP violation source
- Cubic term of  $S$  is important for vacuum

□ Not necessarily consider the Landau pole

- Relatively free  $\lambda, \kappa$  realize various Higgs mass

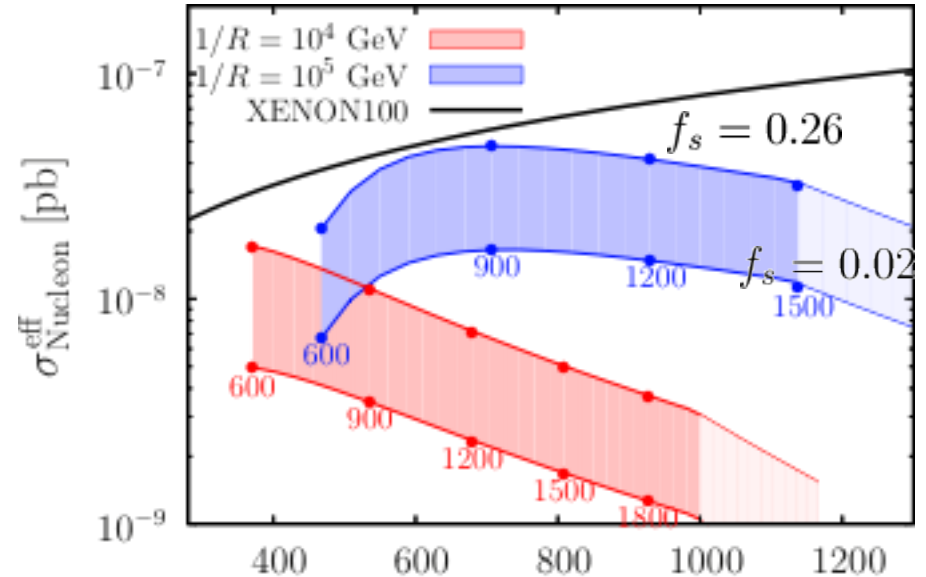
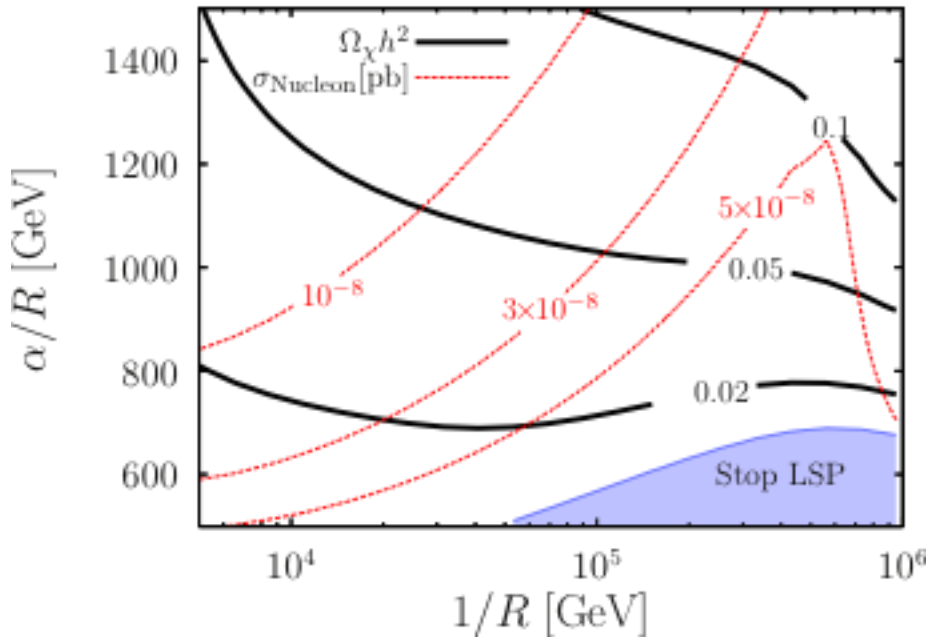
# Dark Matter Nature

❑ Thermal relic of LSP is not enough for observed DM density unless  $LSP \gtrsim TeV$

$$\Omega_{DM} h^2 \simeq 0.1$$

- Relic abundance
- Spin-Indep. cross section with a nucleon

• Effective DM-nucleon cross section



$$\sigma_{Nucleon}^{eff} \equiv \sigma_{Nucleon} \frac{\min\{\Omega_\chi, \Omega_{DM}\}}{\Omega_{DM}}$$

❑ Direct detection of DM does not exclude this scenario, and the future update will be interesting

# Summary

- ❑ Compressed scenario is rather difficult to test at the LHC
- ❑ This scenario is realized by **Compact SUSY model**
- ❑ This model has only 3 parameters (No Flavor and CP problem)
- ❑ Less fine-tuned than CMSSM, and alive in sub-TeV
- ❑ Large radiative corrections boost Higgs mass (up to 125GeV?)
- ❑ LSP is sub-dominant component of DM

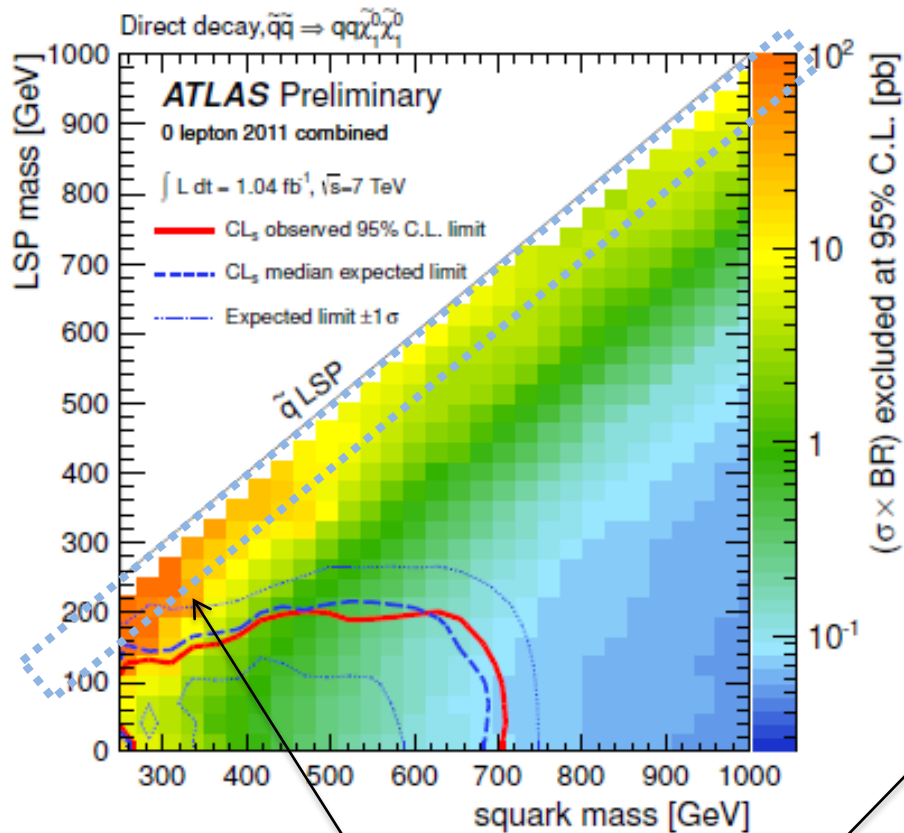
## Future work

- Higgs sector and DM in NMSSM
- Non-thermal production of DM
- Non-trivial radiative corrections from KK modes
- Correspondence of SS mechanism and Radion mediation

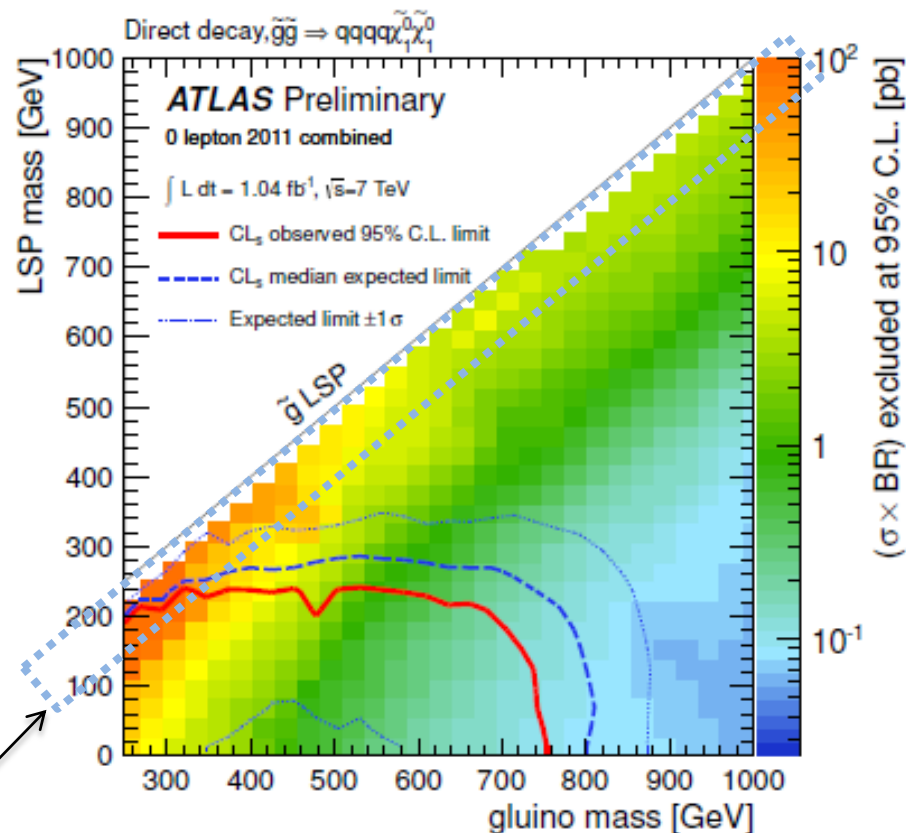
Thank you for your attention

ありがとうございました

Squarks+LSP model



Glauino+LSP model



100 GeV  
splitting

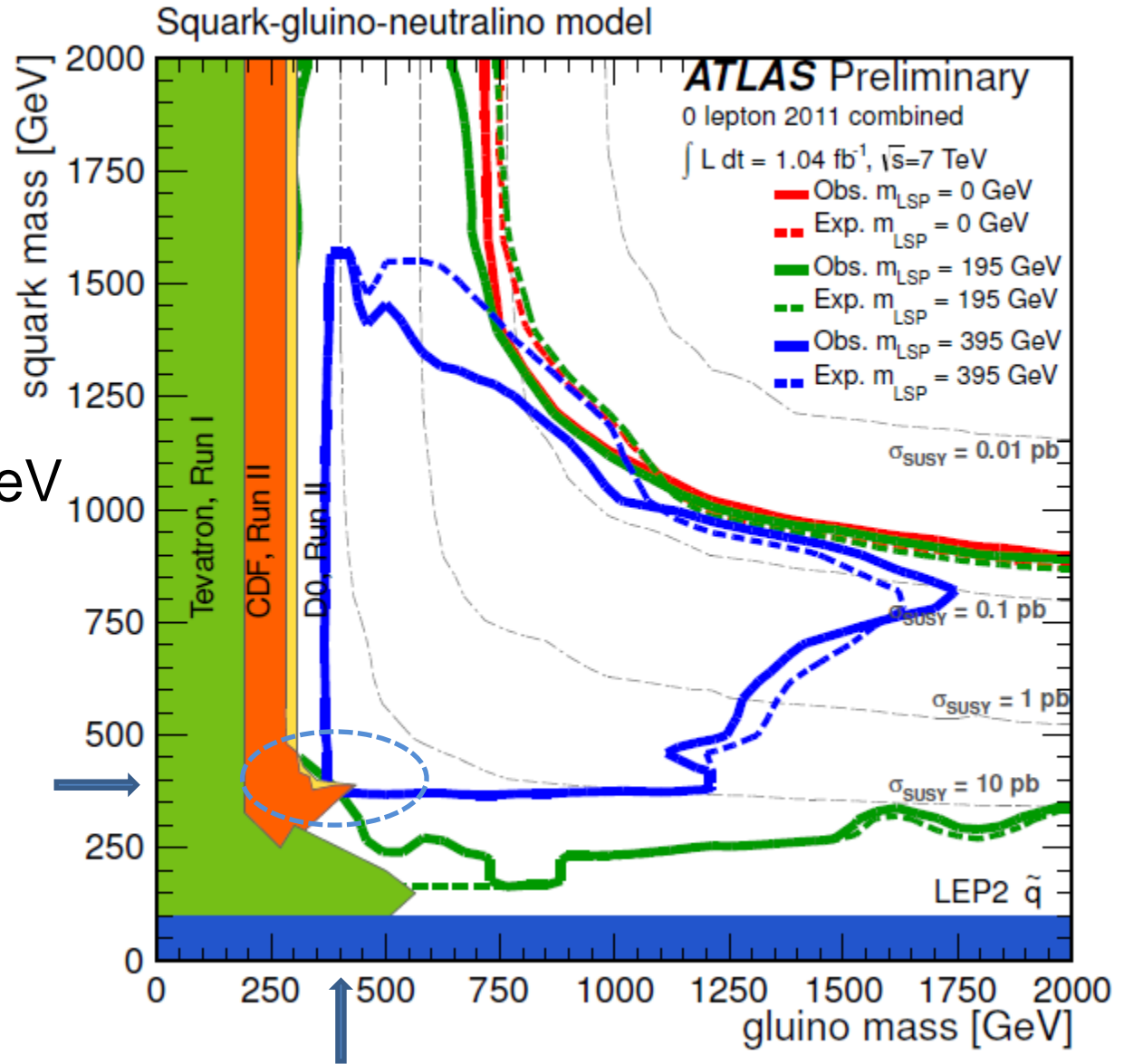
$M_{\text{squark}} \sim 250 \text{ GeV},$

$M_{\text{gluino}} \sim 350 \text{ GeV}$

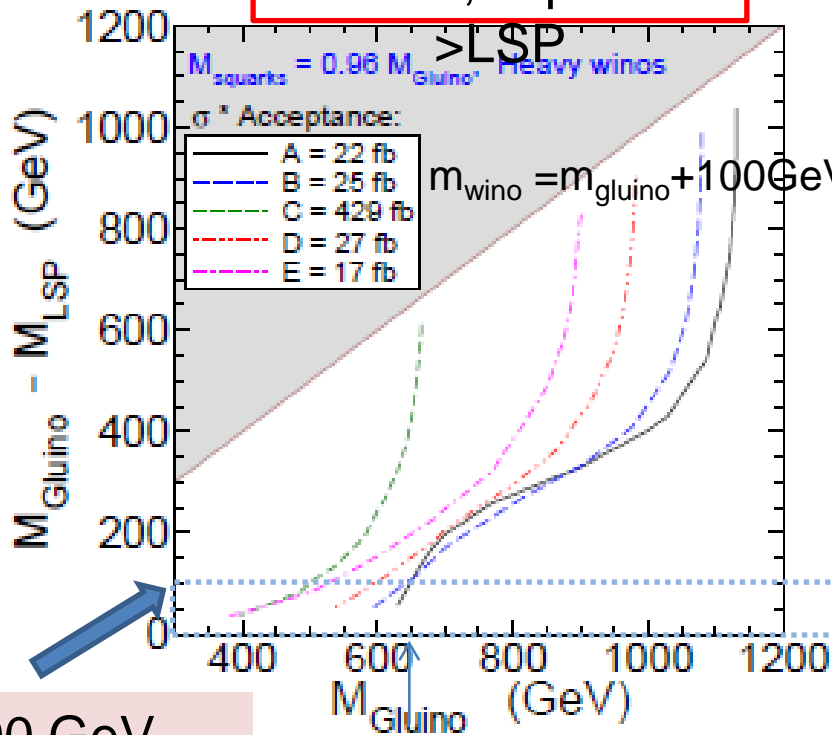
Squarks+Gluino+LSP  
model

$M_{\text{squark}} = M_{\text{gluino}} \sim 400 \text{ GeV}$

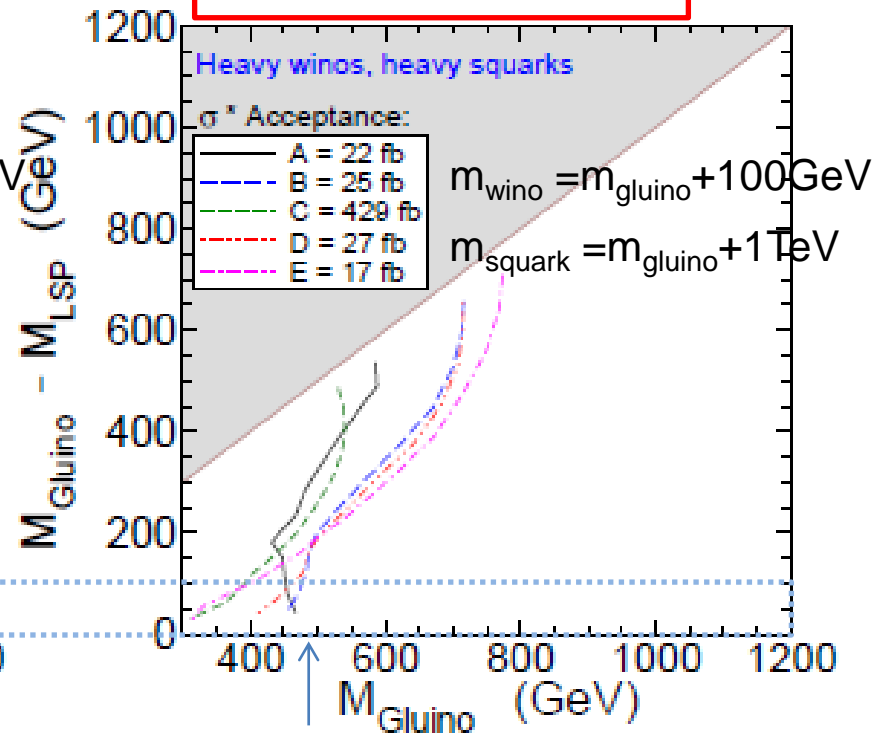
with >5GeV splitting!



Glauino, Squark -



Glauino ->LSP



100 GeV splitting

$$M_1 = \left( \frac{1+5c}{6} \right) M_{\tilde{g}}$$

$$M_2 = \left( \frac{1+2c}{3} \right) M_{\tilde{g}}$$