Metastable electroweak vacuum and inflationary Universe

based on: KK, T. Kobayashi, T.Takahashi, M.Yamaguchi, J.Yokoyama, PRD86(2012)023504 M.Asano, KK, O.Lebedev, & A.Westphal, in preparation

Kohei Kamada (DESY theory group)



基研研究会 素粒子物理学の進展 2013 @京都, 8/8/2013

PPP2012 Higaki-san's talk

ArXiv:1207. 2771 [hep-ph]

Higgs, Moduli Problem, Baryogenesis and Large Volume Compactifications

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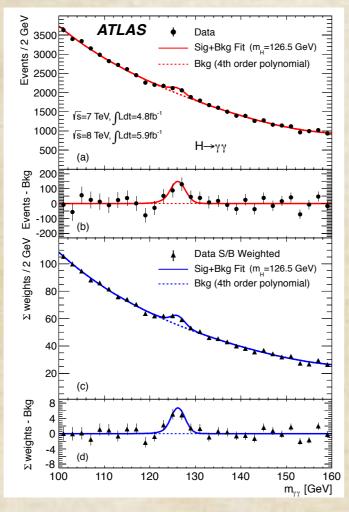
 $Deutsches\ Elektronen\text{-}Synchrotron\ DESY,$

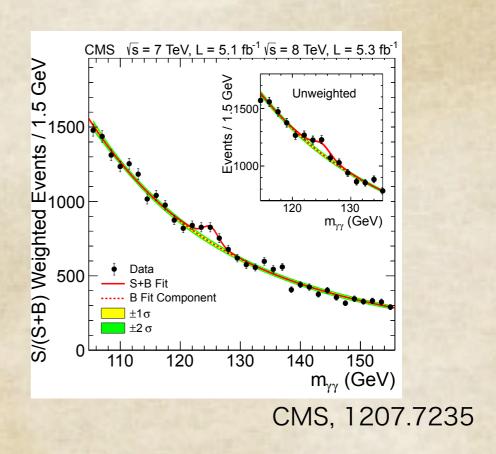
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Department of Physics, Tohoku University, Sendai 980-8578, Japan

July 2012, a (SM) Higgs boson with mass around 125 GeV is found at LHC!

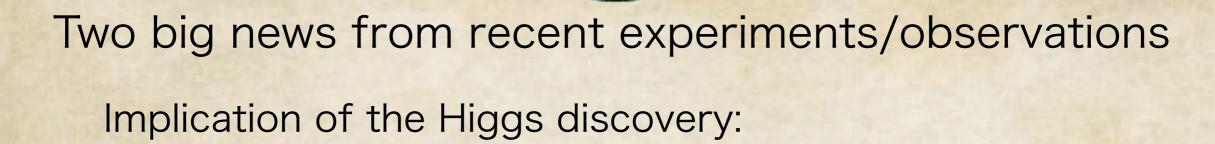




ATLAS, 1207.7214

...and the data supports the SM more and more strongly.

Any deviation from the SM has not been reported thus far.



- 1. Standard Model works well.
- Supersymmetry or other models beyond the standard model do not give any clues so far. (~10 TeV SUSY is favored? Where is naturalness?)

Strong coupling

e Higgs discovery:

1. Standard Model works well.

Scale u

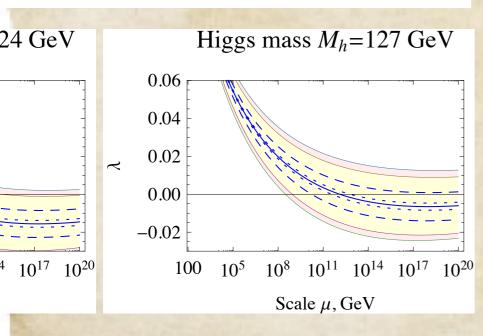
- 2 Supersymmetry or other models beyond the standard model do not give any clues so far.
 - (~10 TeV SUSY is favored? Where is naturalness?)

ne electroweak vacuum is at the boundary.

cf. Kawai-san & Iso-san's talk

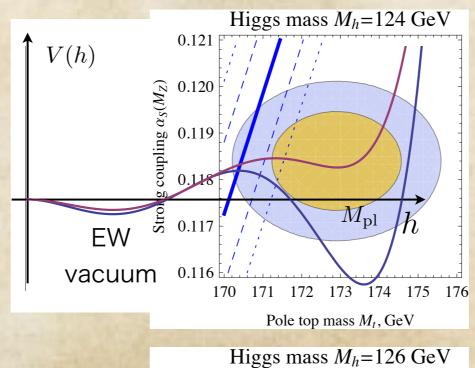
0.121

 \bigcirc 0.120

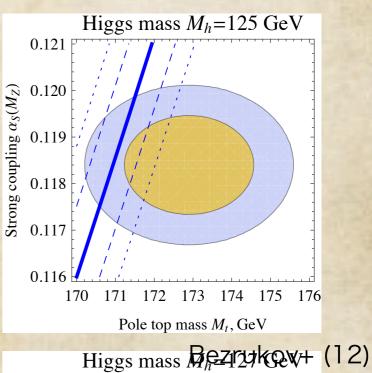


Zero

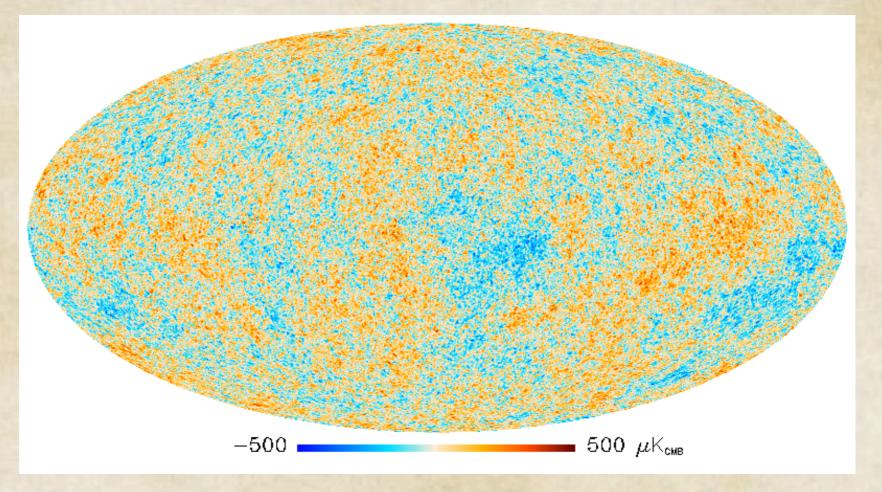
 $M_{\rm Planck}$

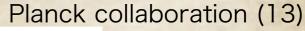


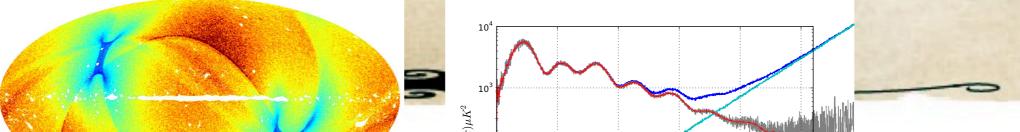
 \bigcirc 0.120



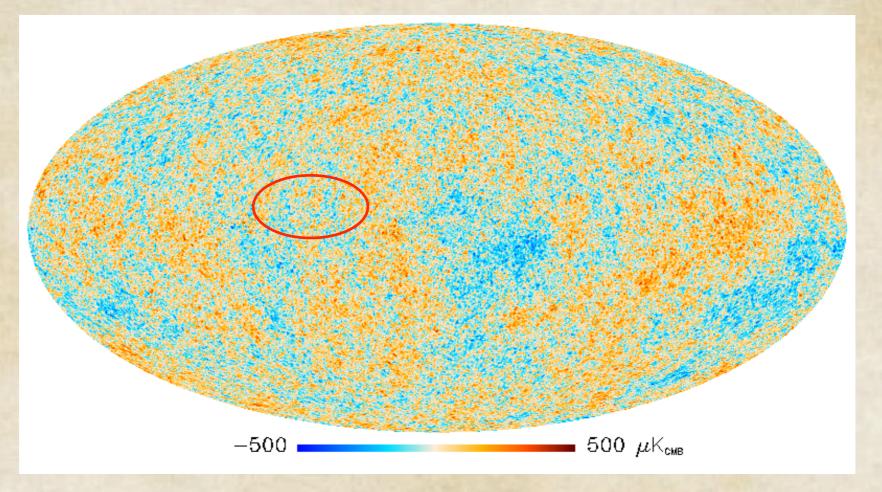
March, 2013, Planck Collaboration showed their first result and gave a detailed picture of CMB! cf. lchiki-san's talk

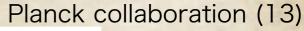


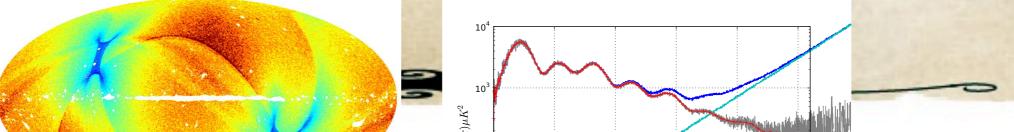




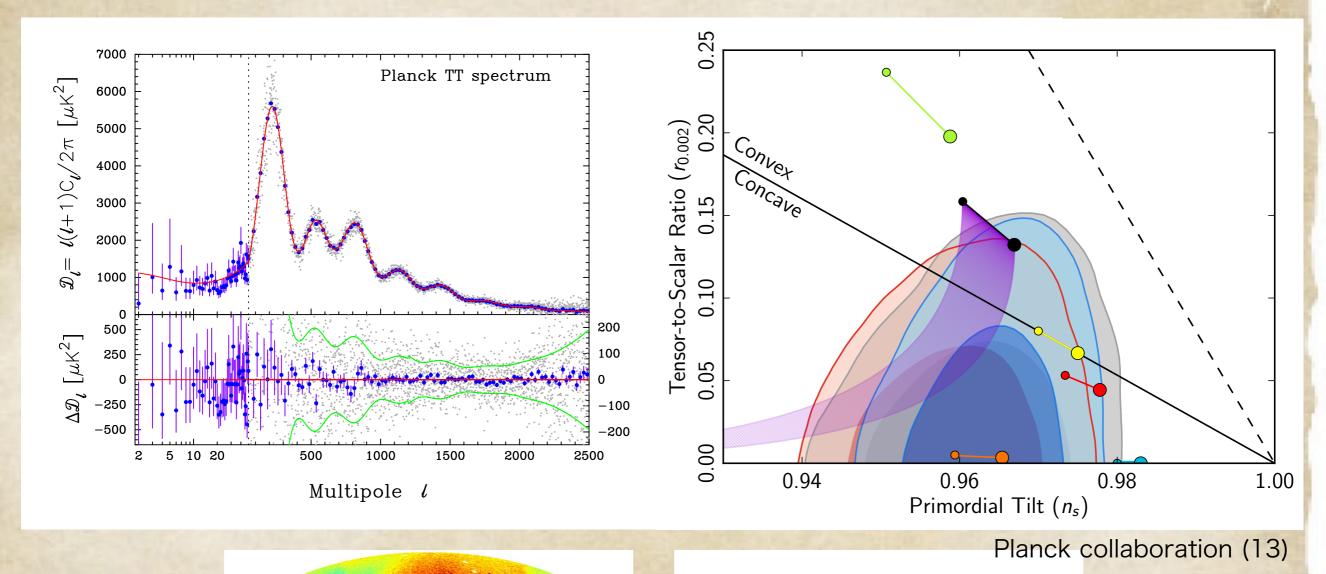
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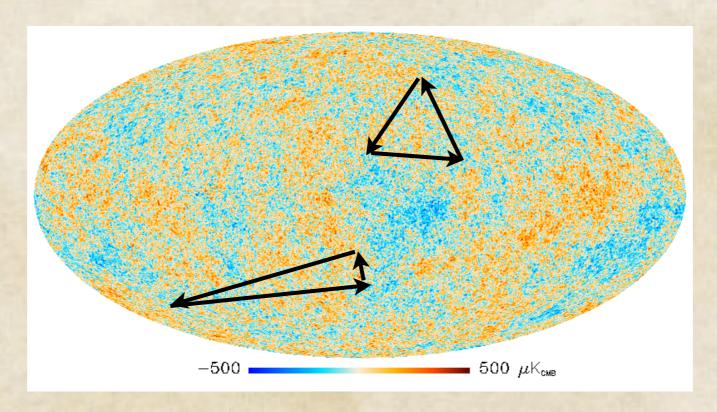
Two big news from recent experiments/observations

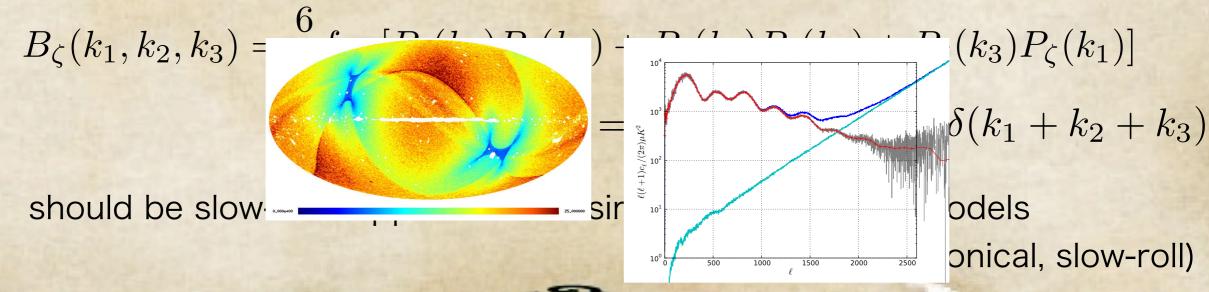
Planck results tells...

- 1. Simple type of inflation explains the observation very well.
- 2. No "anomalies" are confirmed.
 - Non-Gaussianity: $f_{
 m NL}^{
 m local}=2.7\pm5.8$ $f_{
 m NL}^{
 m equil}=-42\pm75$ $f_{
 m NL}^{
 m ortho}=-25\pm39$ (68% CL)
 - Dark radiation: $N_{
 m eff}=3.36^{+0.68}_{-0.64}$ (95% CL) # without BAO&H0
 - Gravitational waves: $r \equiv \frac{A_t}{A_s} < 0.11$ (95% CL)

Nongaussinaity

...Probe of nonminimal interaction of inflation and existence of light field







Dark radiation

...Probe of the existence of light particle at the present

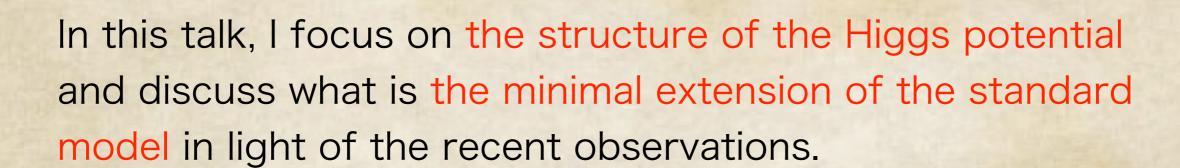
Gravitational waves cf. Jinno-san's talk

...Probe of the energy scale of inflation

$$A_T = \frac{8}{M_{\rm pl}^2} \left. \left(\frac{H_{\rm inf}}{2\pi} \right)^2 \right|_{H_{\rm inf} = k/a}$$

$$r \simeq 0.1 \leftrightarrow H_{\rm inf} \simeq 10^{13} {\rm GeV}$$

$$r \simeq 0.001 \leftrightarrow H_{\rm inf} \simeq 10^{12} {\rm GeV}$$



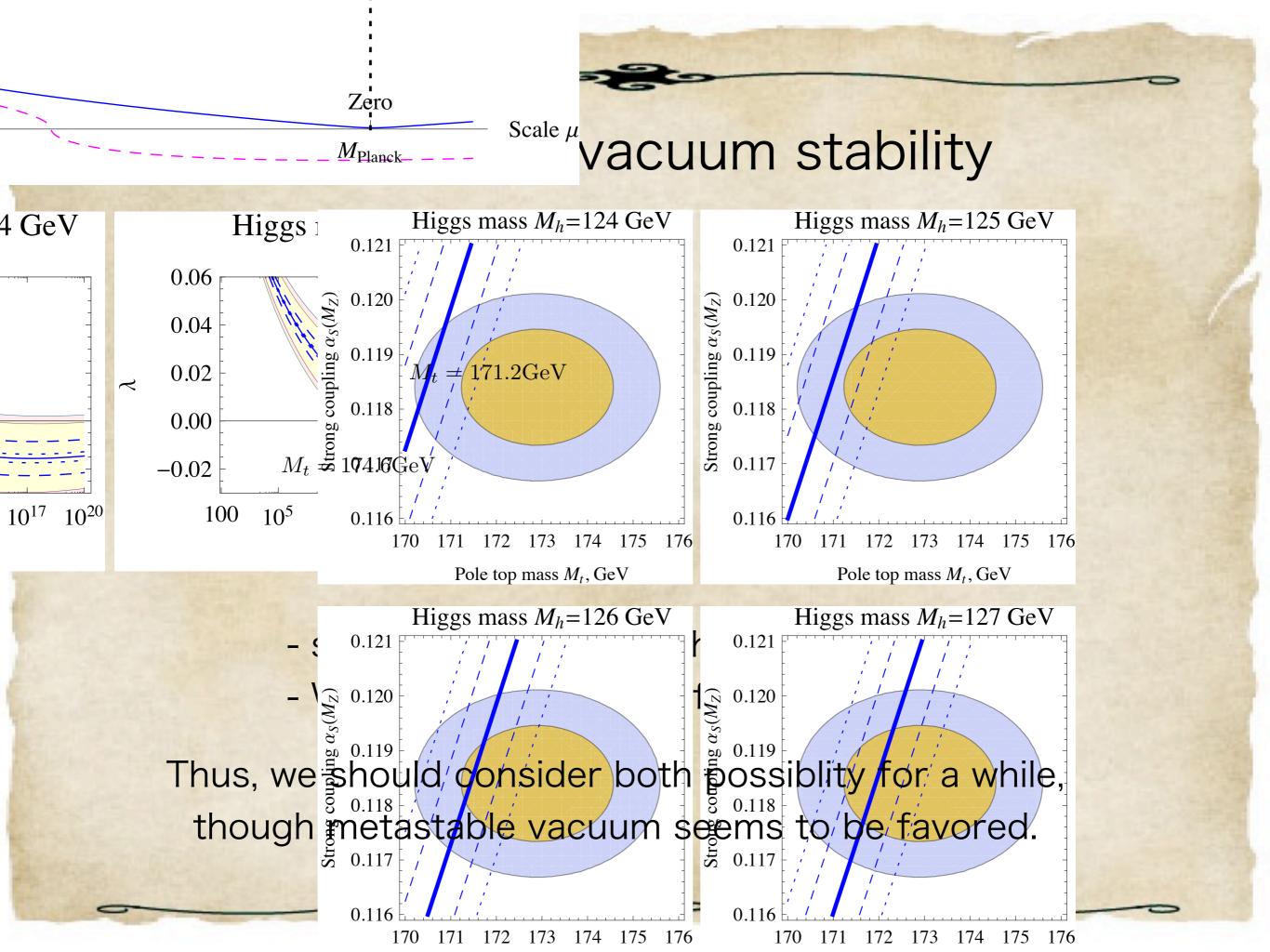
In particular,

Stable case: Higgs inflation

metastable case: stability during inflation

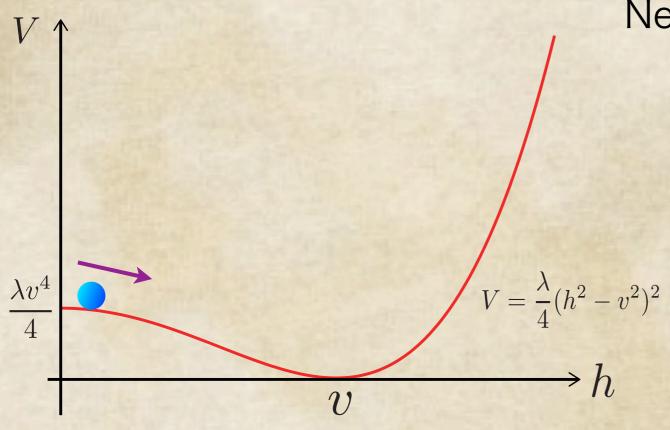
Here I mean "I require minimal extension of the physical degrees of freedom but allow nonrenormalizable interaction", by using the term "minimal".

Moreover, I focus on inflation and does not study dark matter, dark energy and baryogenesis.





In order to have an inflationary era, we need a scalar field that drives inflation, inflaton.



New inflation? ('82, Linde)

: impossible because the potential is too steep to realize accelerating expansion of the Universe.

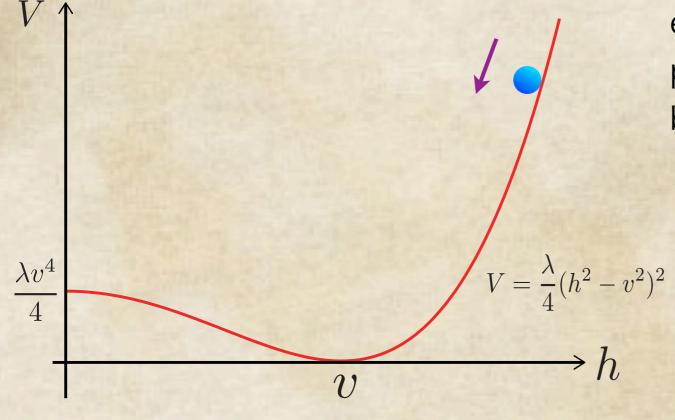
In order to have an inflationary era, we need a scalar field that drives inflation, inflaton.



expansion of the Universe, but the primordial density perturbation becomes too large.

$$\mathcal{P}_{\zeta} \sim 10^3 \lambda \;\; ext{for} \;\; V = rac{\lambda}{4} \phi^4$$

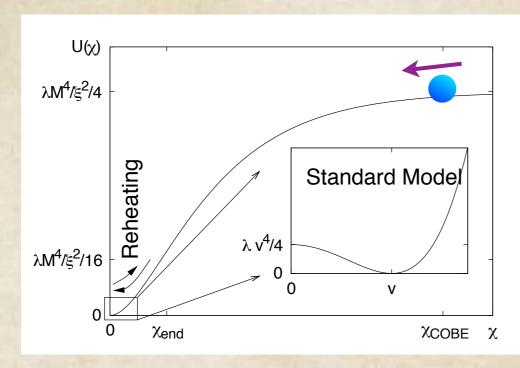
 $\lambda_{
m Higgs} \sim {\cal O}(1)$ is inconsistent with the observation ${\cal P}^{
m obs}_{\cal R} \simeq 2.4 imes 10^{-9}$ (WMAP('12))



In order to have an inflationary era, we need a scalar field that drives inflation, inflaton.

It is known that the SM Higgs can drive inflation that can explain our present Universe.

→ Higgs inflation



$$\Delta S = \int d^4x \sqrt{-g} \left[\frac{-\xi h^2}{2} R \right]$$

For $|\xi|h^2\gg 1$, the effective Planck mass becomes large

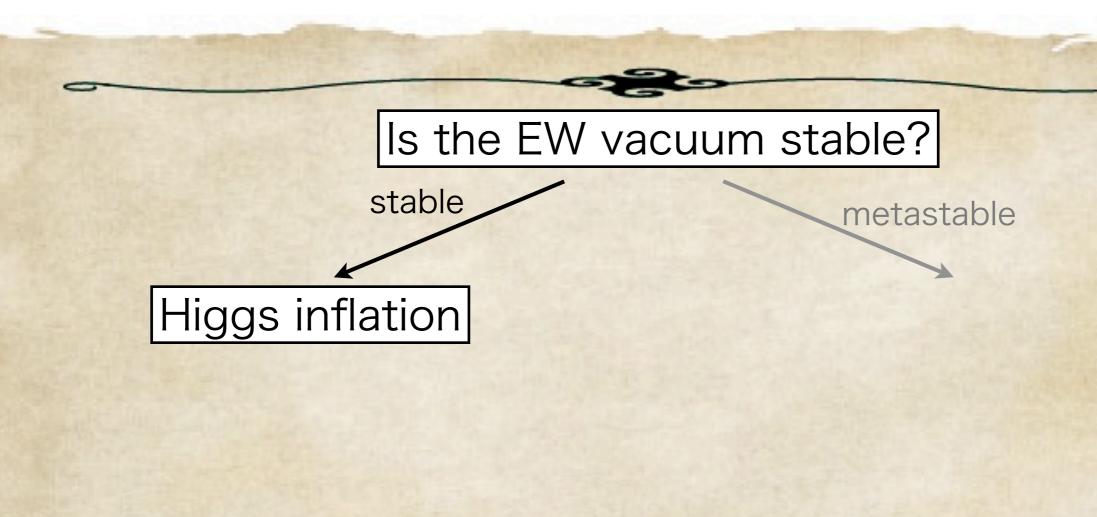
$$M_{\rm pl}^{\rm eff2} = M_{\rm pl}^2 + |\xi| h^2$$

and hence the primordial density fluctuation is suppressed.

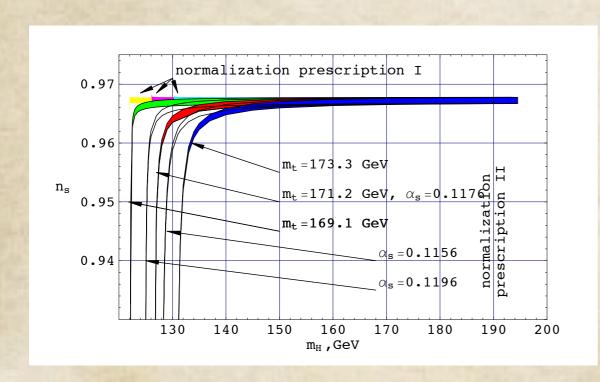
('08, Bezrukov & Shaposhnikov) (cf. '95, Cervantes-Cota+)

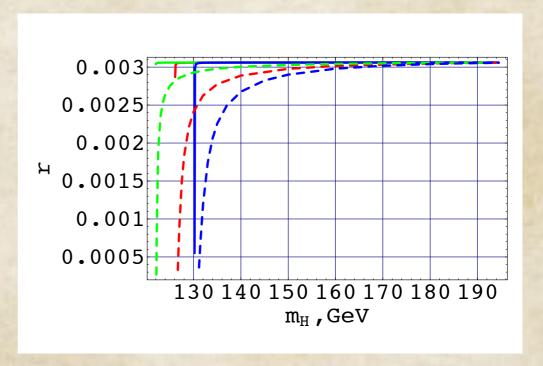
(In this notation, ξ is negative.)

This is the unique possiblity of the extension of the SM at renormalizable level.



Predictions of (nonminimal) Higgs inflation





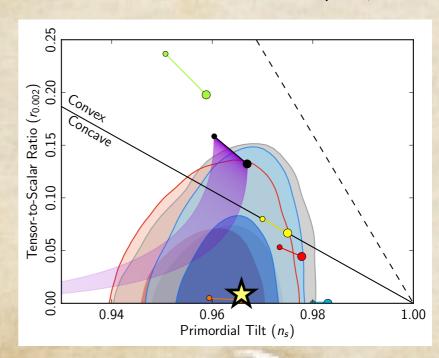
('09, Bezrukov+)

spectral tilt: $n_s \simeq 0.966$

gravitational waves: $r \simeq 0.002 \sim 3$

Non-Gaussianity: slow-roll suppressed.

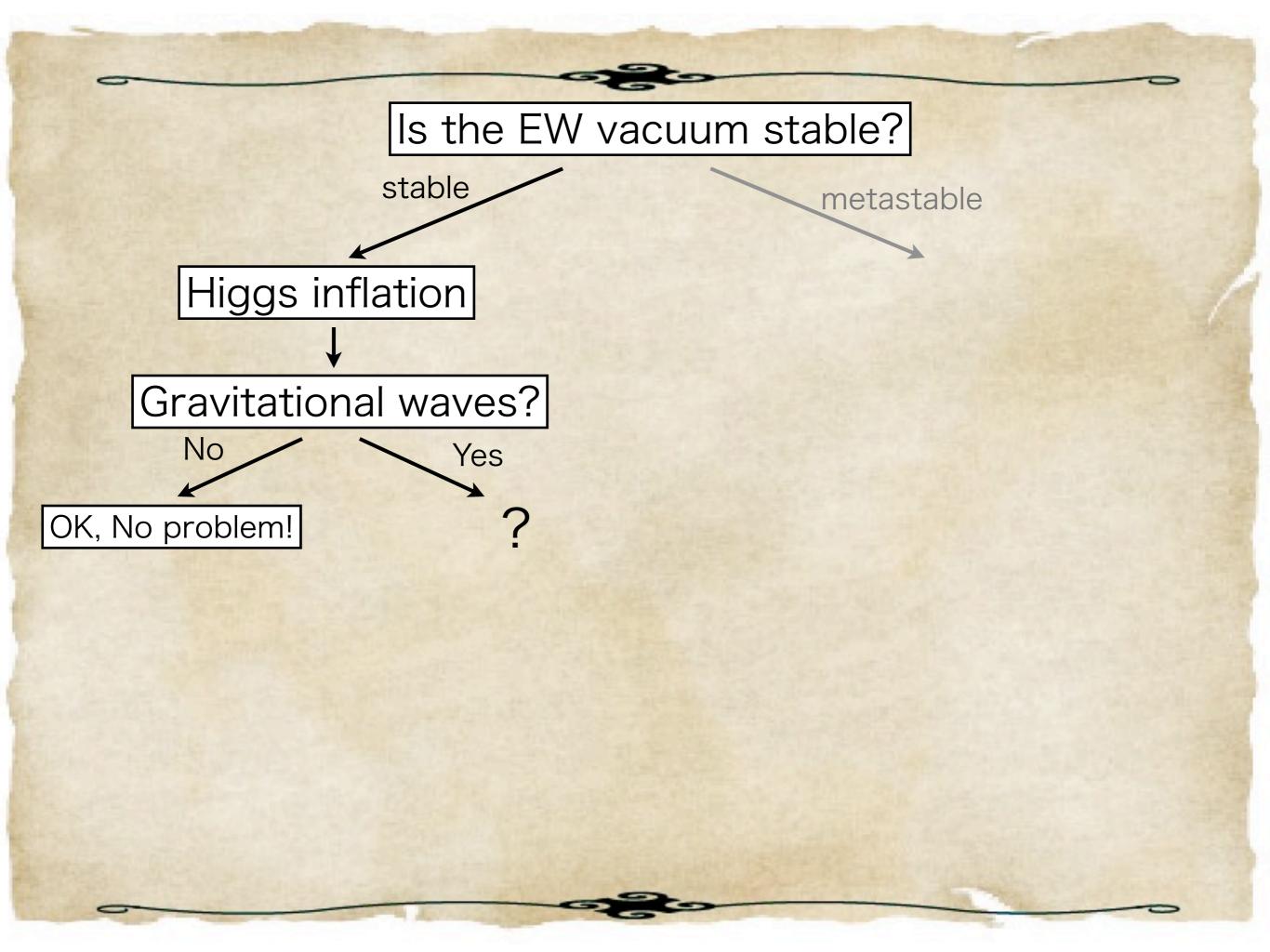
#depending on the top mass...



But we can expect for the detection of gravitational waves by Planck up to

 $r \gtrsim 0.05$

Should we give up Higgs inflation when Planck will detect gravitational waves?



We find that there are several possibilities of Higgs inflation if we allow nonrenormalizable (derivative) interaction for the Higgs field based on the generalized Galileon theory.

'11,'12, KK, T. Kobayashi, T.Takahashi, M.Yamaguchi, J.Yokoyama

$$\mathcal{L}_{2} = K(\phi, X),$$

$$\mathcal{L}_{3} = -G_{3}(\phi, X)\Box\phi,$$

$$S = \int d^{4}x\sqrt{-g}\sum_{i=2}^{5}\mathcal{L}_{i}$$

$$\mathcal{L}_{4} = G_{4}(\phi, X)R + G_{4X}\left[\left(\Box\phi\right)^{2} - \left(\nabla_{\mu}\nabla_{\nu}\phi\right)^{2}\right],$$

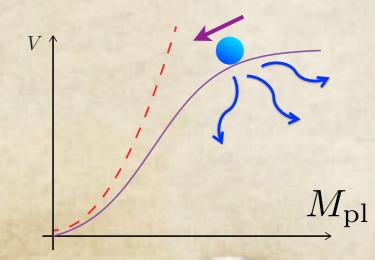
$$\mathcal{L}_{5} = G_{5}(\phi, X)G_{\mu\nu}\nabla^{\mu}\nabla^{\nu}\phi - \frac{1}{6}G_{5X}\left[\left(\Box\phi\right)^{3}\right]$$

$$-3\left(\Box\phi\right)\left(\nabla_{\mu}\nabla_{\nu}\phi\right)^{2} + 2\left(\nabla_{\mu}\nabla_{\nu}\phi\right)^{3}.$$

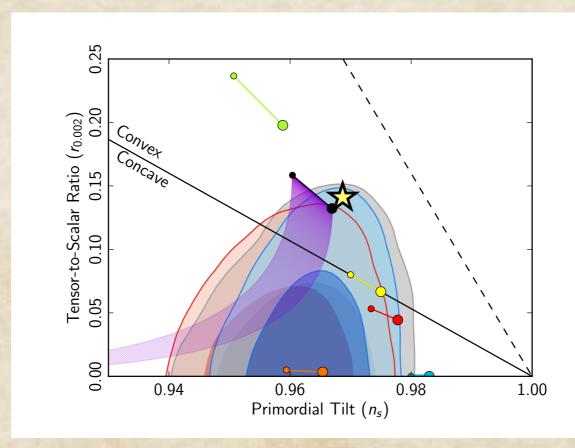
$$X = -\frac{1}{2}(\nabla\phi)^{2}, \quad G_{iX} = \frac{\partial G_{i}}{\partial X}$$

('74, Horndeski; '09, '10, '11, Deffayet+)

Derivative interaction enhance the friction term for the Higgs field and realize sub-Planck inflation

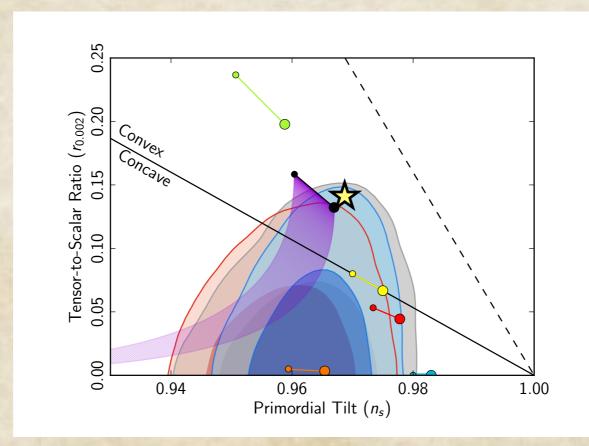


In the simplest case, introducing a term, $h\Box h\partial_{\mu}h\partial^{\mu}h$ we have $n_s\simeq 0.967, \quad r\simeq 0.14$ $M\sim 10^{13}{\rm GeV}, h_{\rm inf}\sim 10^{16}{\rm GeV}$



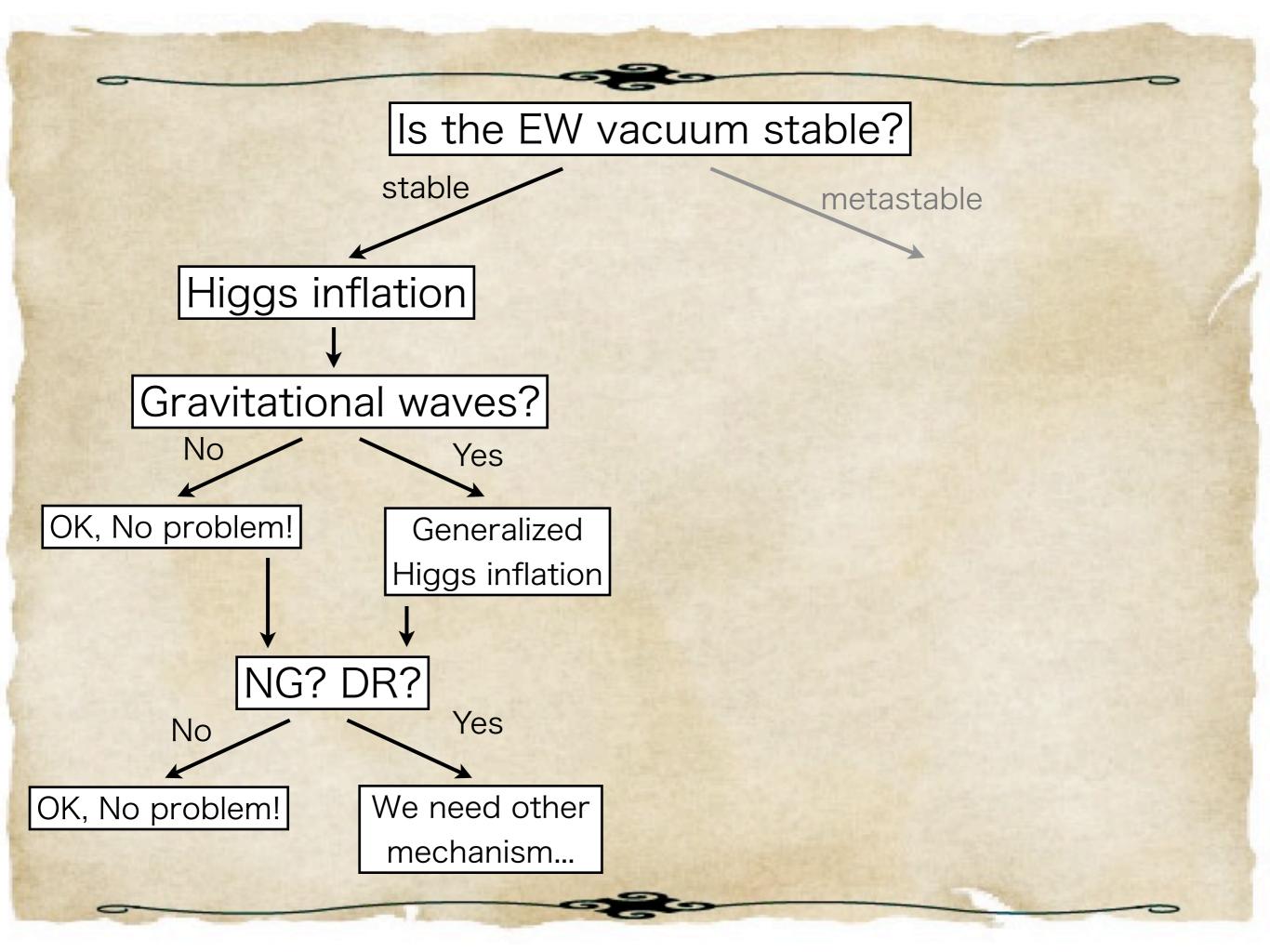
This model is found to have a problem in reheating ('12 Ohashi&Tsujikawa) but it can be solved in the context of Generalized Galileon (or generalized Higgs inflation). (See, Kunimitsu+ in preparation.)

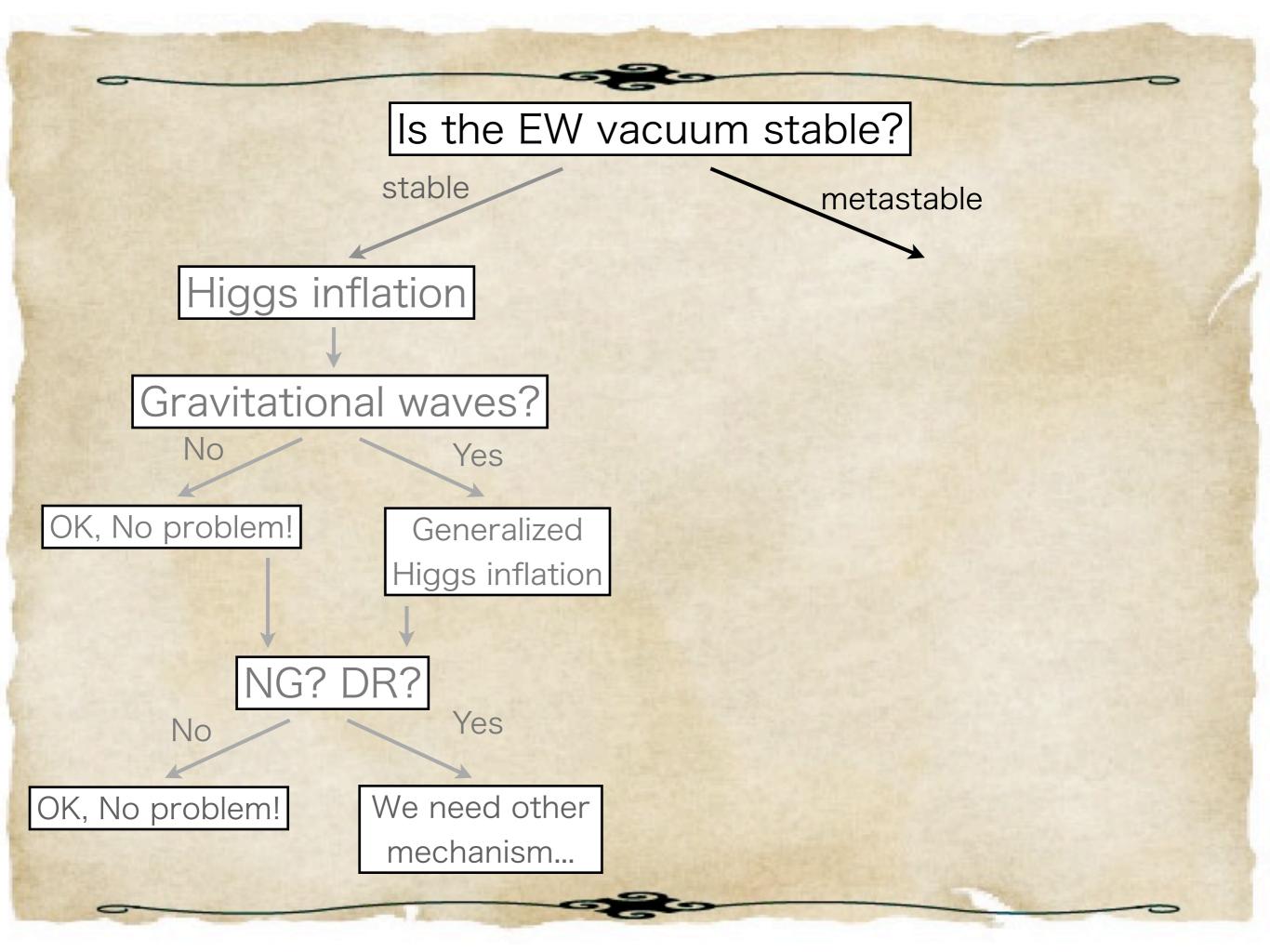
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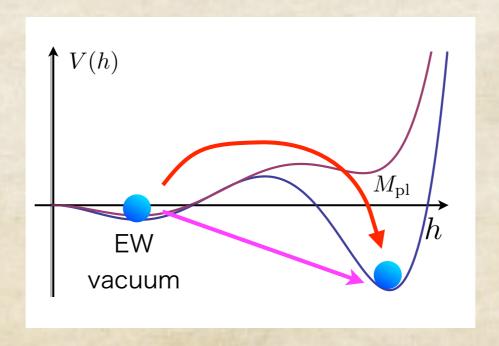
We can still hope for nongaussianity or dark radiation, but generalized Higgs inflation cannot explain it and it needs some other degrees of freedom.







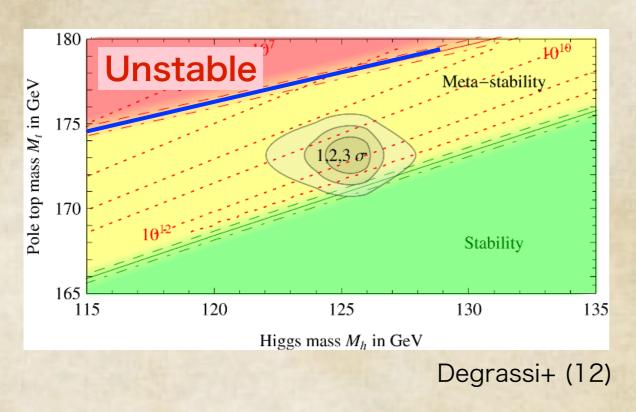
Metastable EW vacuum Can we live in a metastable electroweak vacuum?

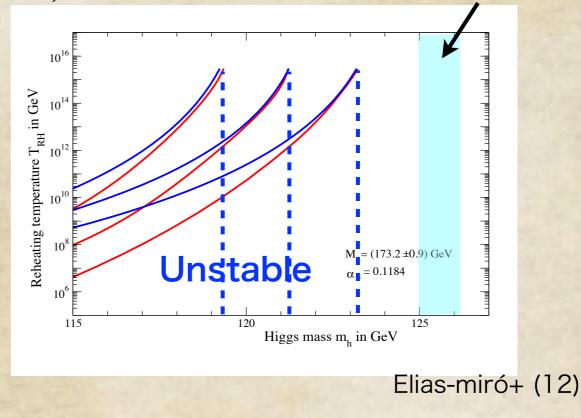


Metastable EW vacuum Can we live in a metastable electroweak vacuum?

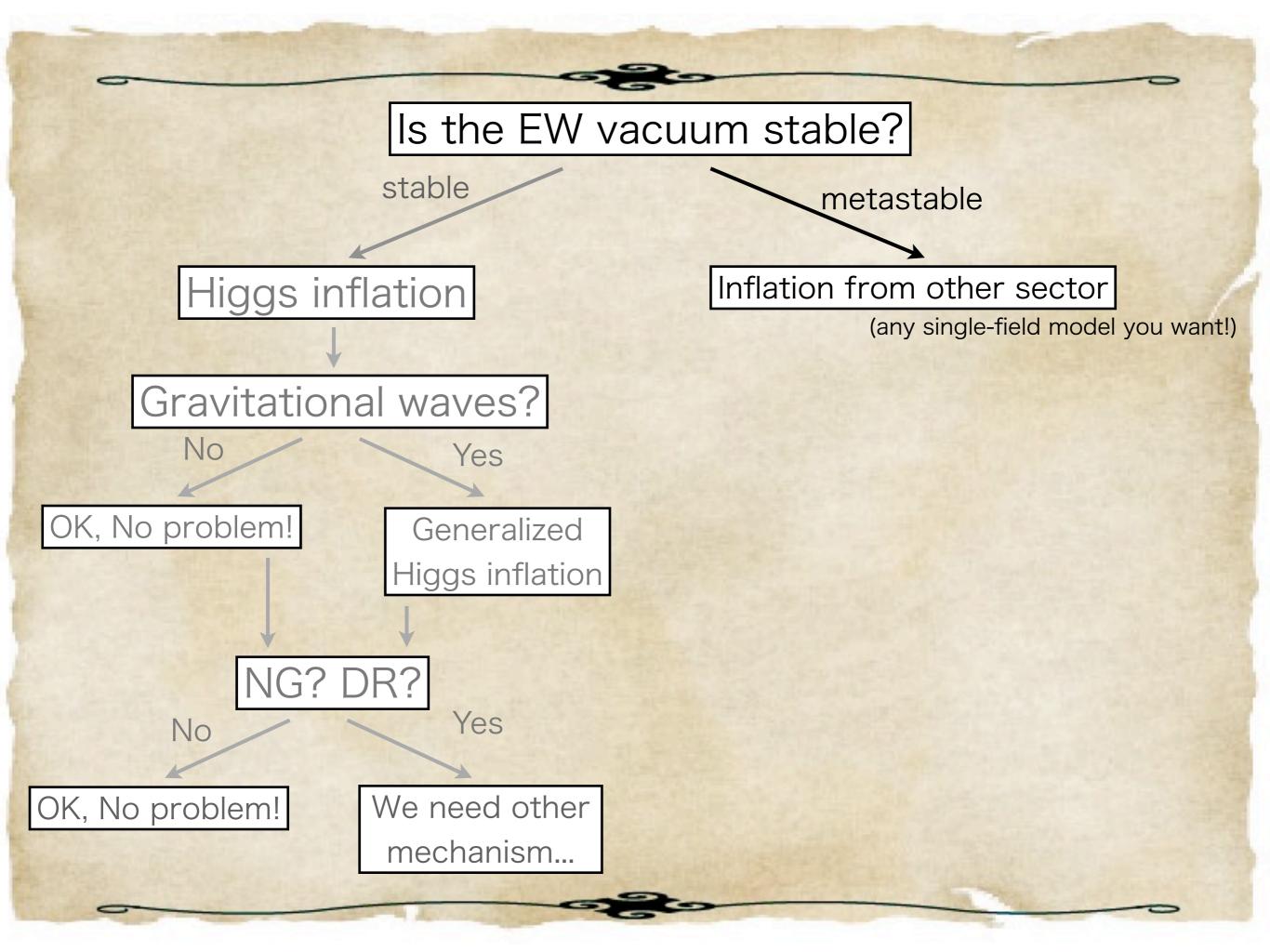
-zero temperature decay : $p \sim \max \{\tau^4 h^4 \exp[-8/3|\lambda(h)|]\}$

-thermal decay : $\Gamma(T) \sim T^4 \left(\frac{S_3(T)}{2\pi T}\right)^{3/2} \exp[-S_3(T)/T]$ LHC data





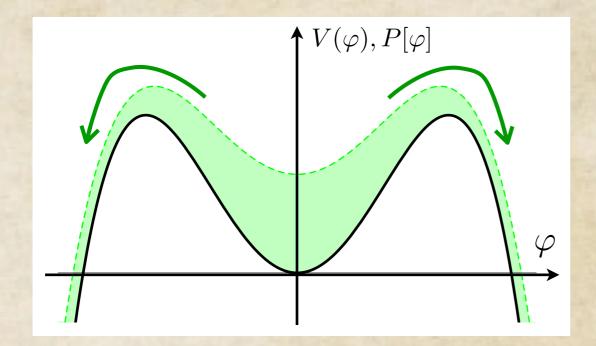
Current data suggests that we live in a "safe" meta-stable vacuum.

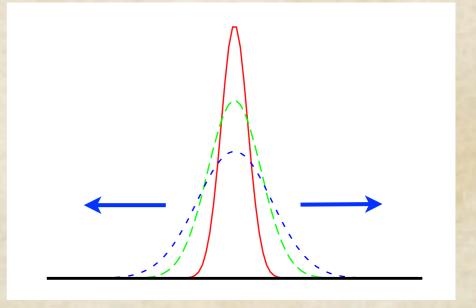


Are we all right?

No! During inflation, the expectation value of the light (massless) scalar field evolves as

$$\langle \varphi^2 \rangle = \frac{H_{\rm inf}^2}{4\pi^2} \mathcal{N}_e$$



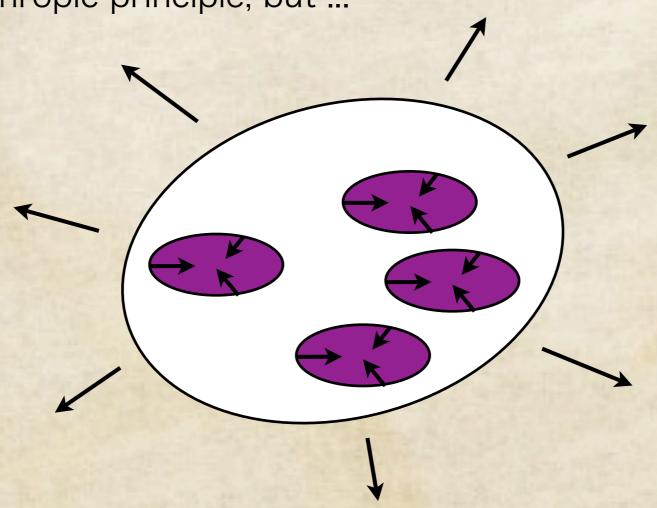




See Starobinsky & Yokoyama (93)



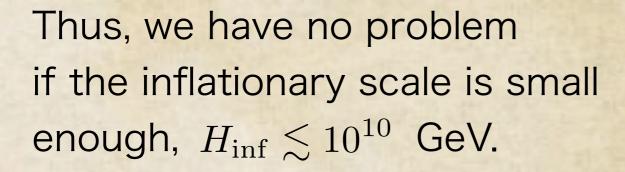
The region where the Higgs falls down to the true vacuum will collapse and inflation does not continue any more in such a region. It may be a matter of anthropic principle, but ...

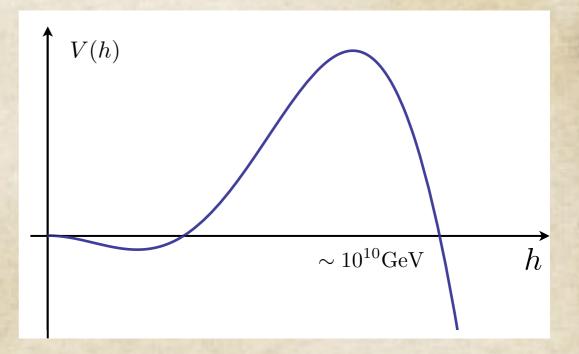


The probability that our Universe has experienced inflation with e-fold ~50 is very rare and it would be better to require that the square root of the Higgs expectation value squared does not overwhelm the barrier of the potential.

Higgs potential, typically, has a potential barrier at $h \sim 10^{10}$ GeV.

For the nearly ϕ^4 potential, the expectation value during inflation reads, $\sqrt{\langle h^2 \rangle} \sim H$.

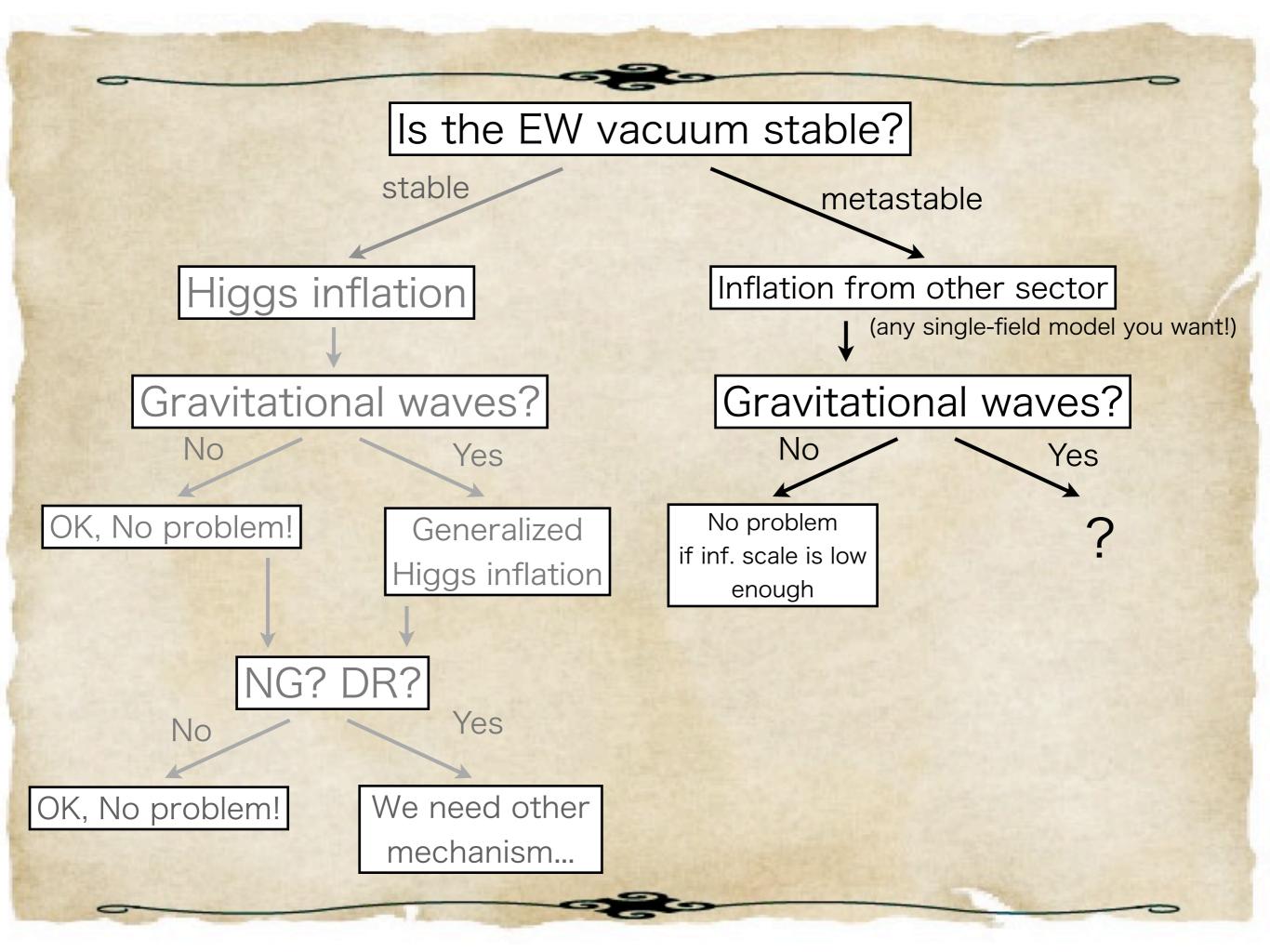


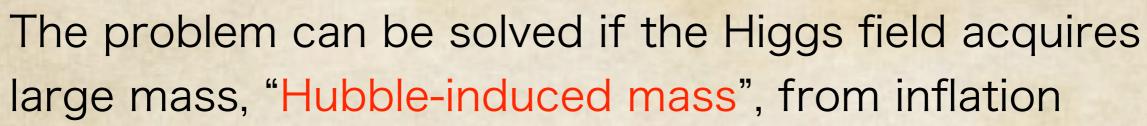


But, again, we can expect for the detection of gravitational waves by Planck up to

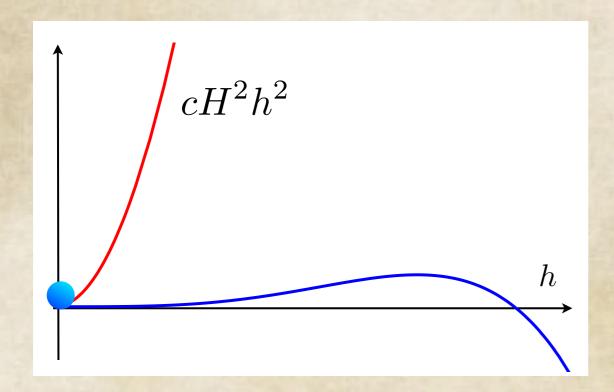
 $r \gtrsim 0.05$

that means the Hubble parameter during inflation is $H_{\rm inf} \gtrsim 0.7 \times 10^{13}$ GeV.





'13 Lebedev&Westphal



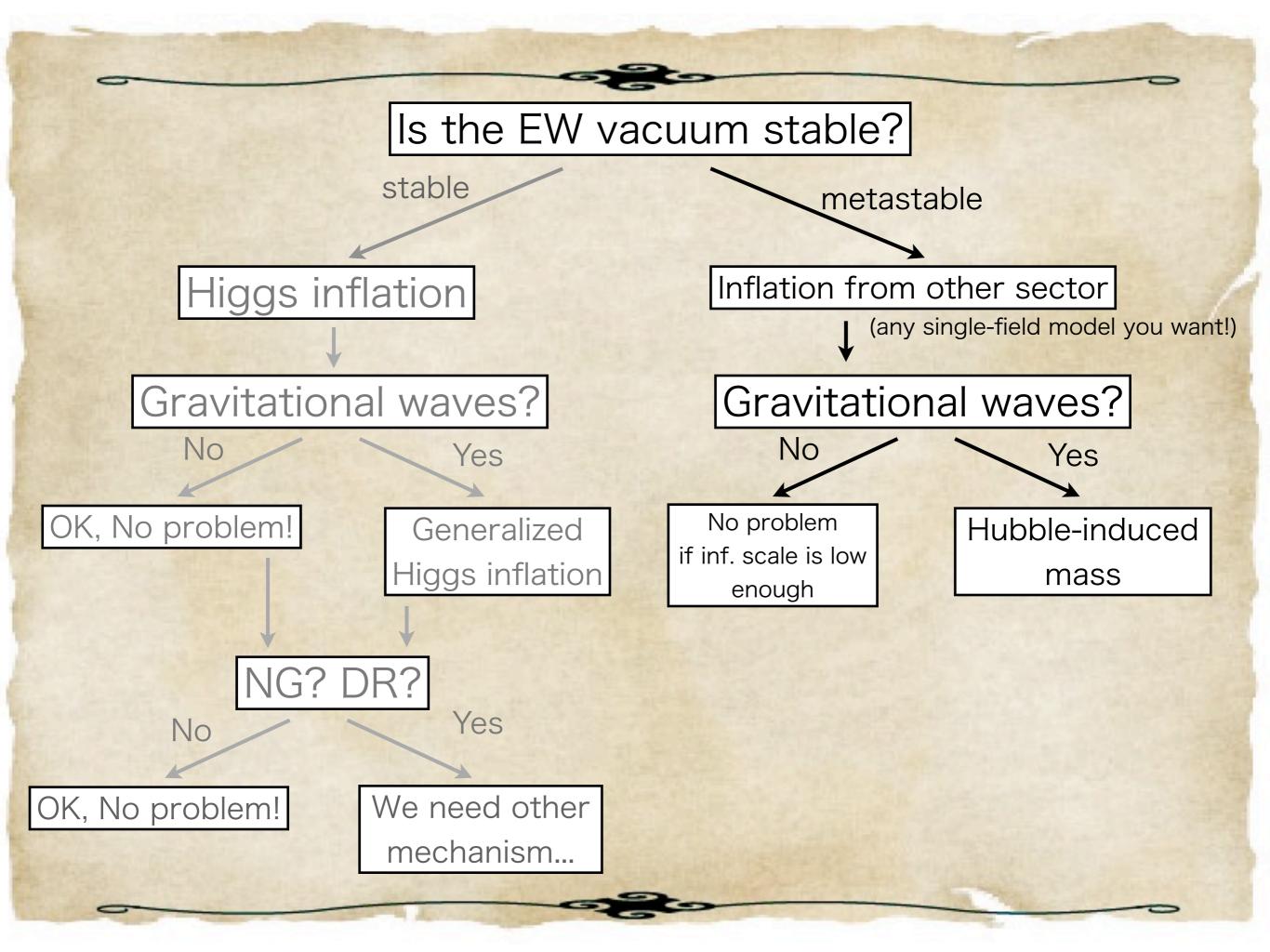
Possible origins;

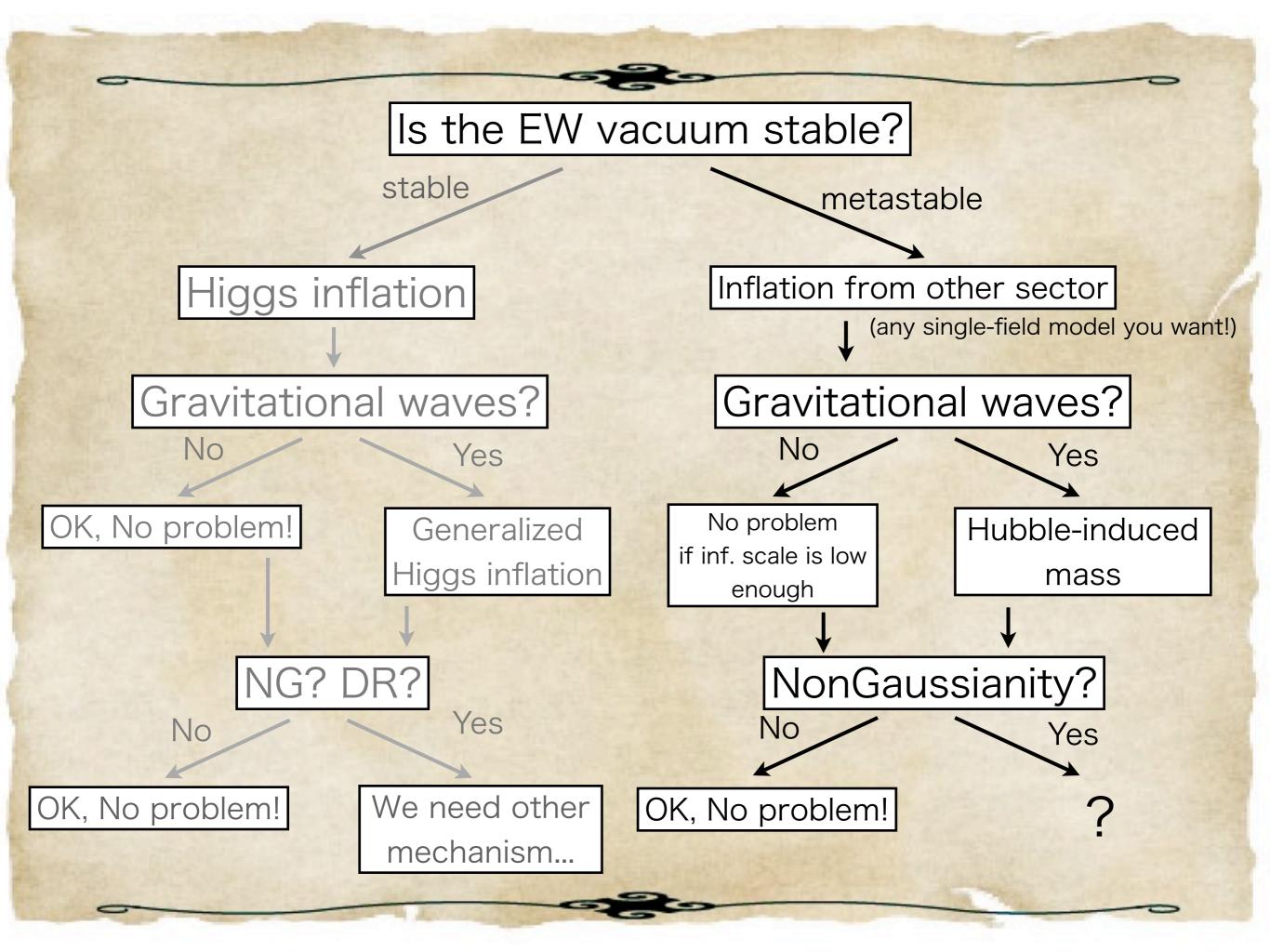
direct coupling to inflaton

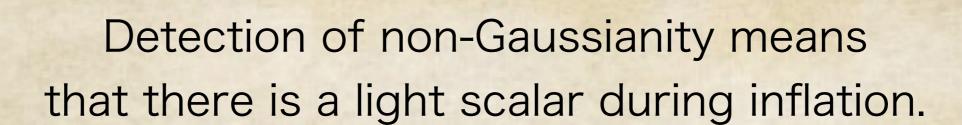
$$\lambda_i h^2 \phi_{\rm inf}^2 \to \frac{m_{h, \rm eff}^2}{2} h^2$$

nonminimal coupling to gravity

$$\xi Rh^2 \to 12\xi H^2h^2$$

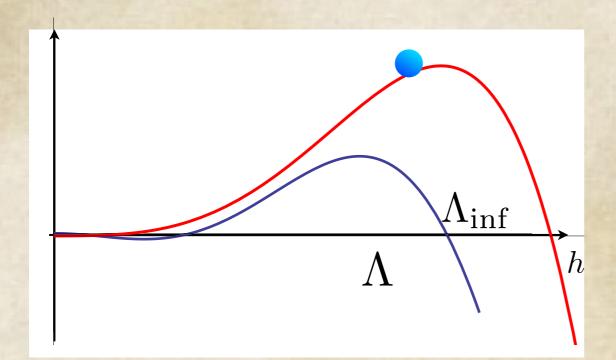




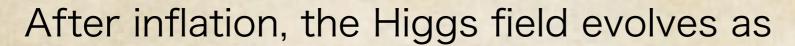


Detection of non-Gaussianity means that there is a light scalar during inflation.

We find that the Hubble induced mass can be relatively small during inflation, $m_H = c_{\rm inf} H_{\rm inf} \lesssim H_{\rm inf}$.



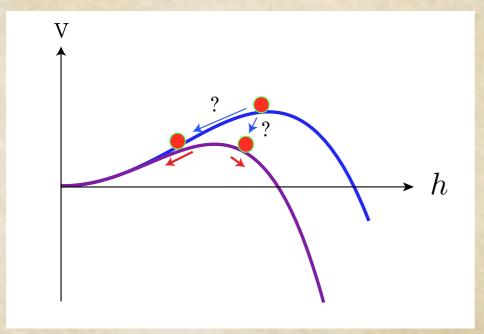
$$c_{\text{inf}} > \sqrt{\frac{-3\tilde{\lambda}}{8\pi^2}} \simeq 1.9 \times 10^{-2} \left(\frac{\tilde{\lambda}}{-0.01}\right)^{1/2}$$



$$\ddot{h} + 3H(t)\dot{h} + \frac{\partial V(H(t), h)}{\partial h} = 0$$

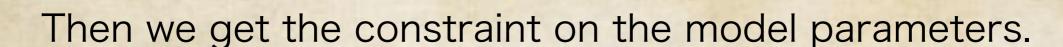
with an initial condition, typically, $h_{\rm ini} = \langle h^2 \rangle_{\rm inf}^{1/2} = \frac{\sqrt{3} H_{\rm inf}}{2 \sqrt{2} \pi \sqrt{c_{\rm inf}}}$

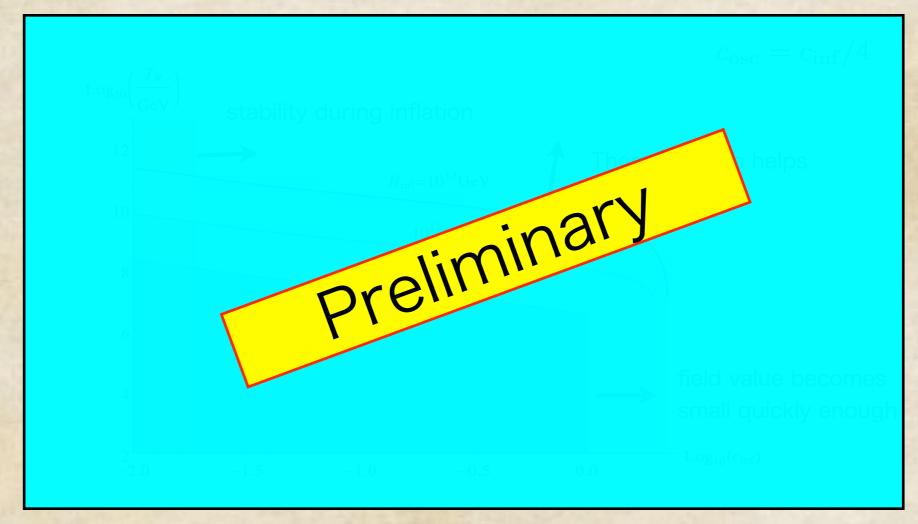
For small coefficient, $c_{\rm osc} < 9/16$, the Higgs field decreases much slower than the potential barrier and may be taken over by it.



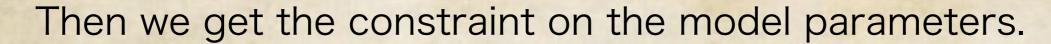
We are safe if...

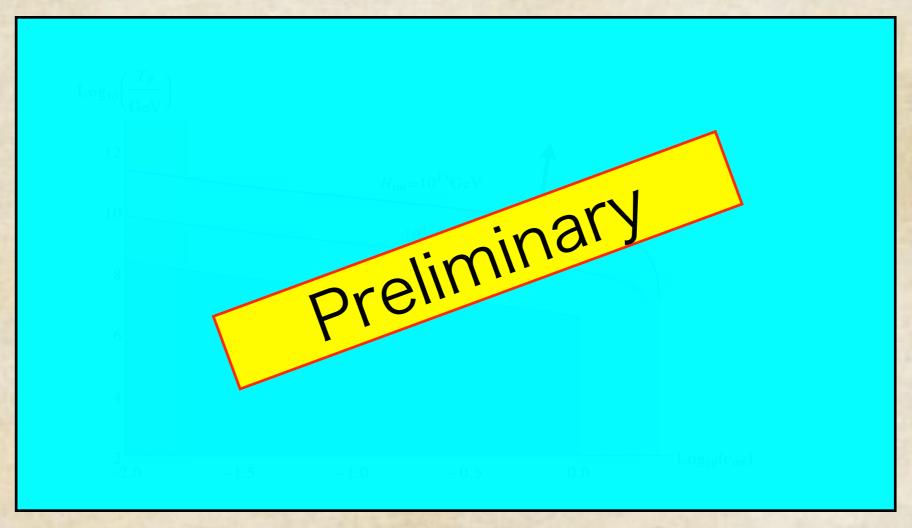
- Thermalization takes place earlier.
- The Higgs field value becomes small enough, $h(t) < \Lambda_0$ sufficiently quickly.





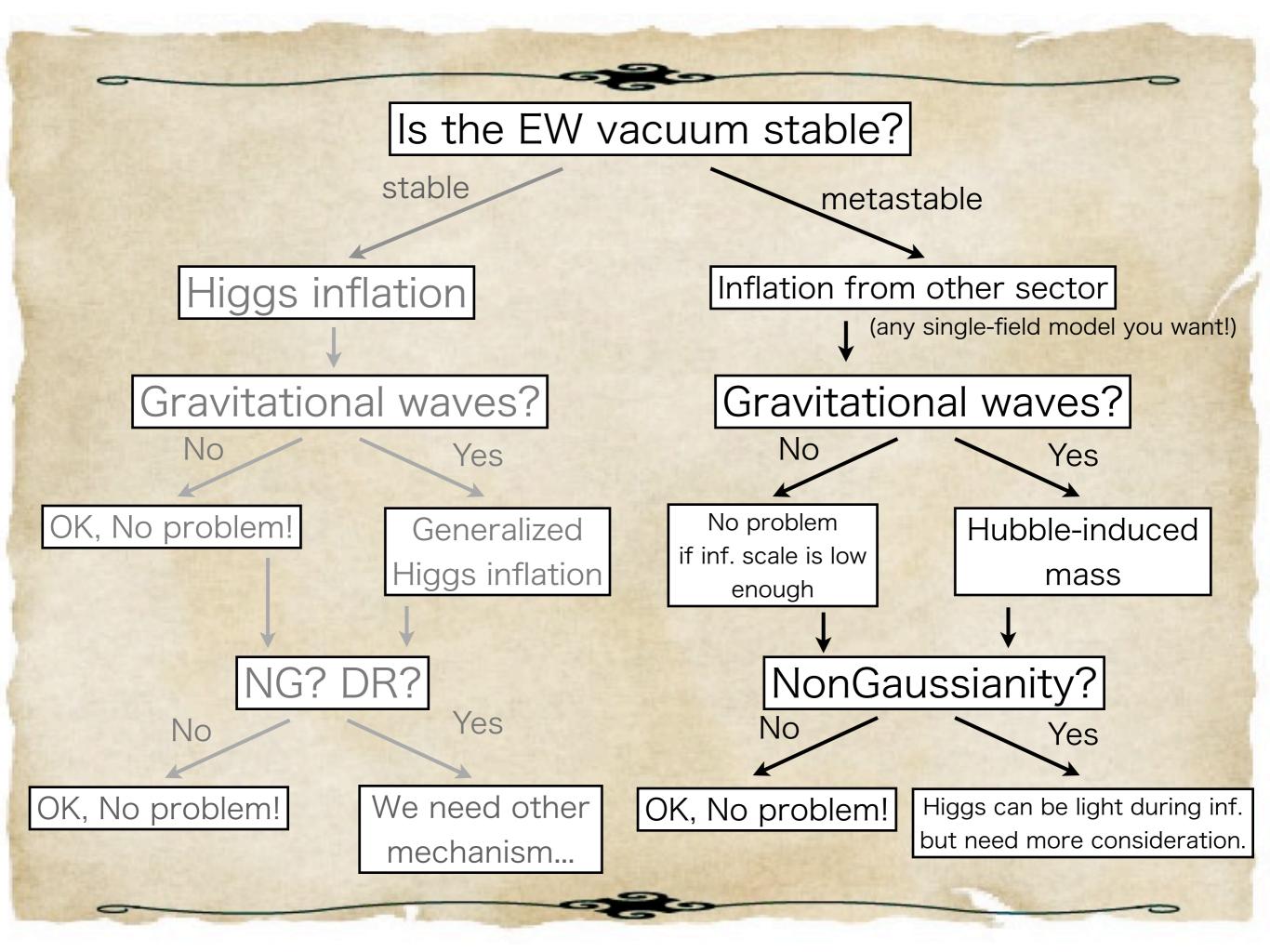
Relatively large reheating temperature is required, which can be tested by future gravitational wave experiments.



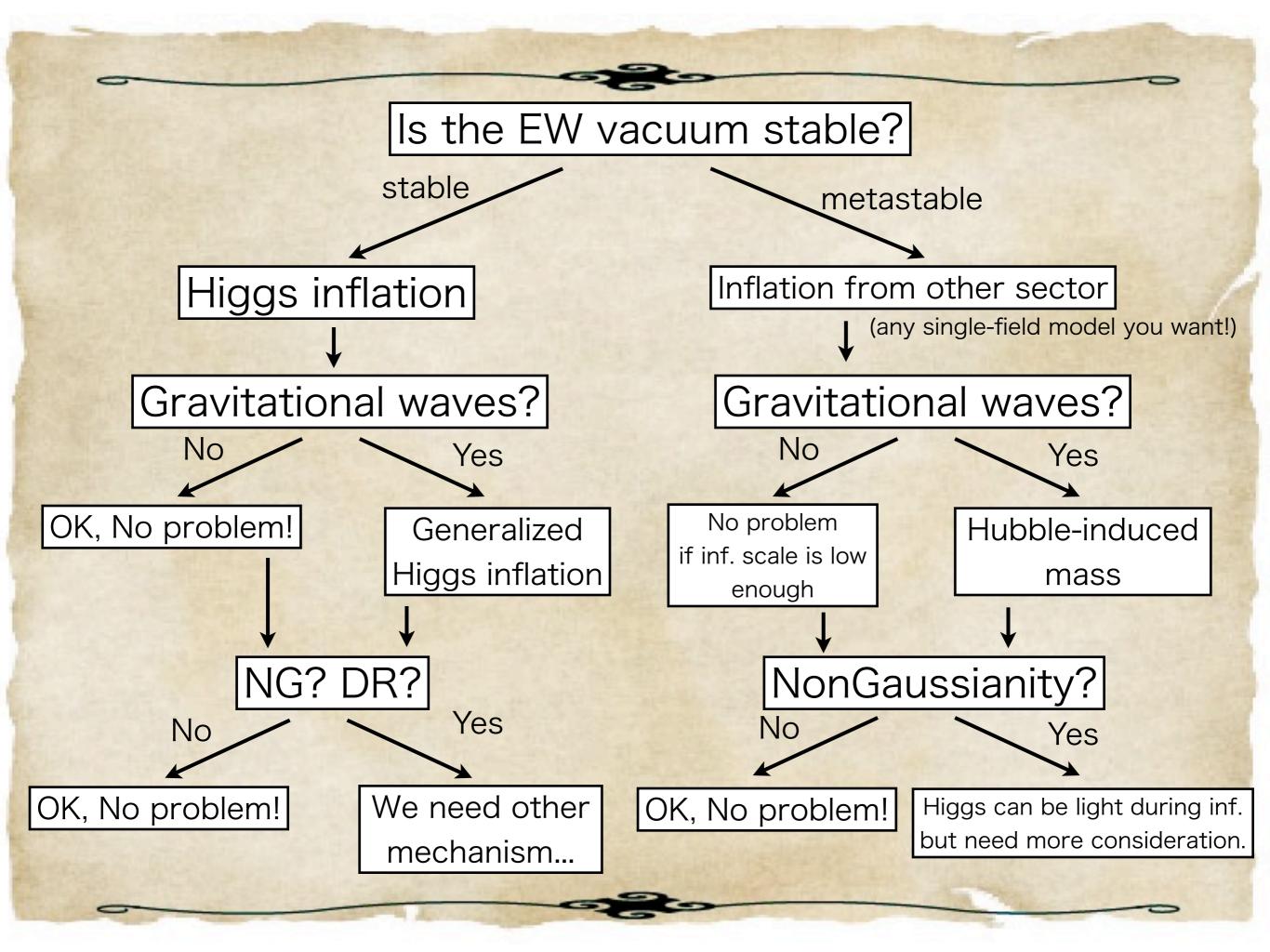


Relatively large reheating temperature is required, which can be tested by future gravitational wave experiments.

But we need some more technique to generate non-Gaussianity from light Higgs.

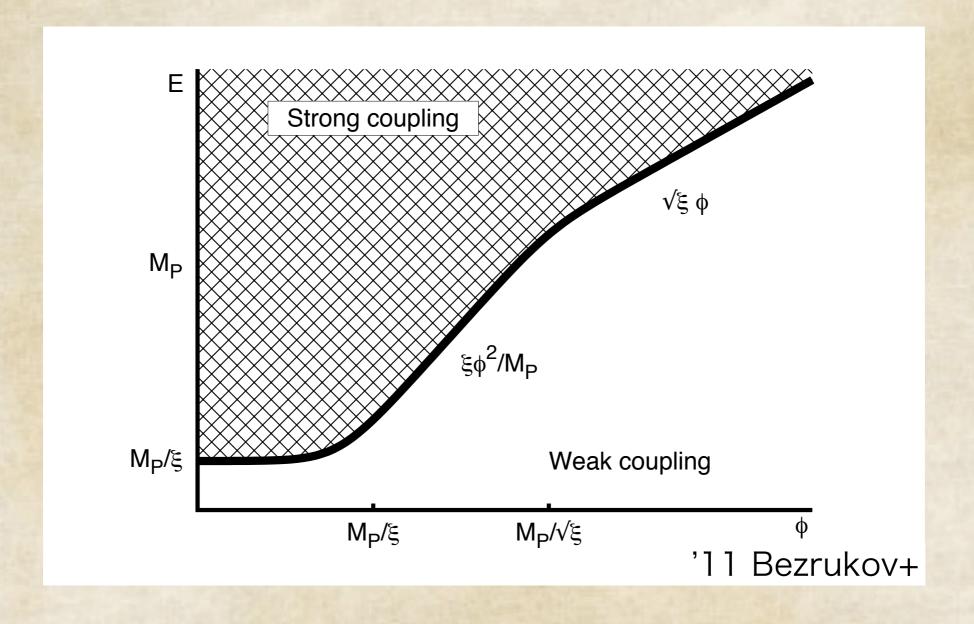








Unitarity problem in Higgs inflation



10⁻²² 10⁻³ 10⁻² 10⁻¹ 10⁰

Determination of reheating temperature by GWB

