### 一般的なNMSSMにおける Higgs混合と軽いHiggsino

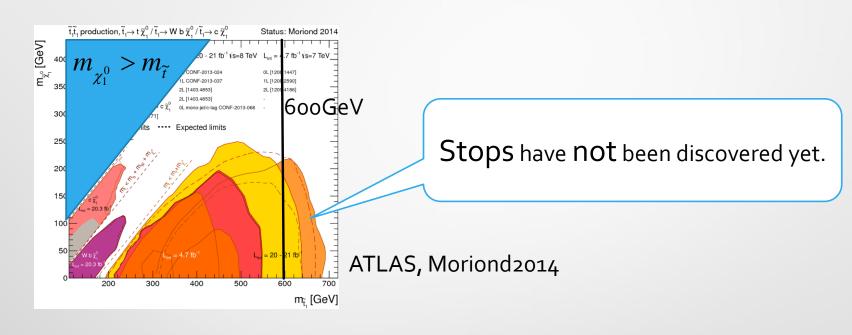
Yutaro SHOJI

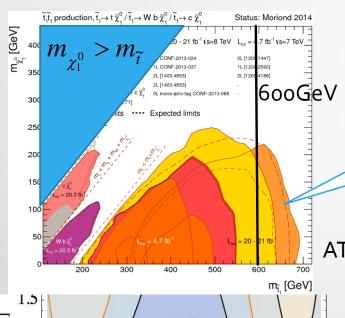
Tohoku University (Japan)

@Kyoto University, PPP2014

Collaborators: Kwang Sik Jeong (DESY), Masahiro Yamaguchi => IBS from Oct.

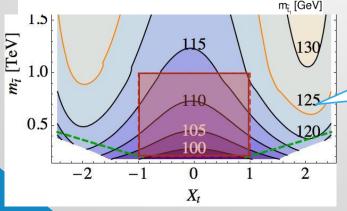
arXiv: 1407.0955





Stops have not been discovered yet.

ATLAS, Moriond2014

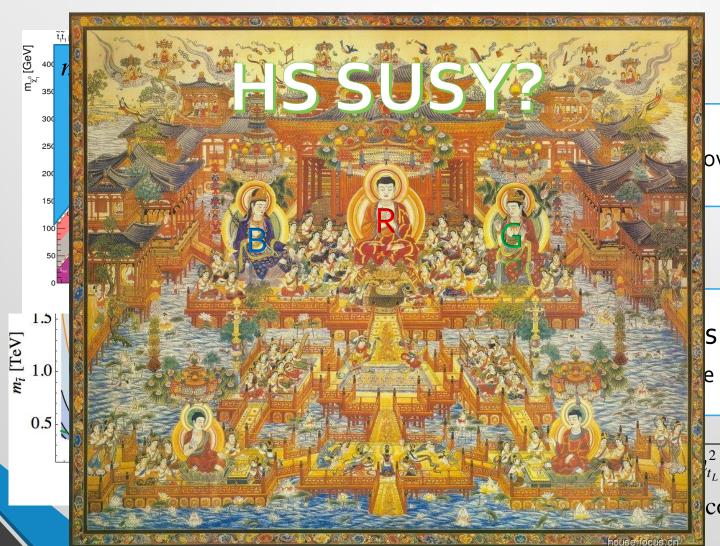


The Higgs boson mass requires heavy stop masses in the MSSM.

FeynHiggs 2.10.0

$$m_{\tilde{t}}^2 = \sqrt{m_{t_R}^2 m_{t_L}^2}$$

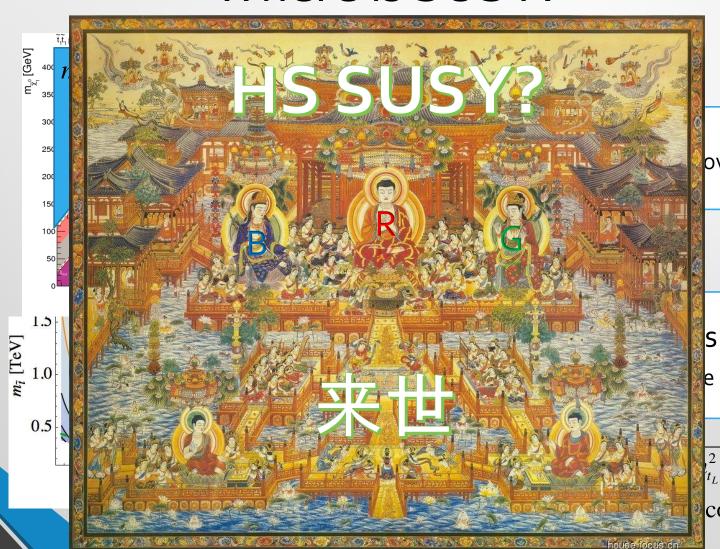
$$X_t = (A_t - \mu \cot \beta) / m_{\tilde{t}}$$



overed yet.

S requires e MSSM.

 $\cot \beta$ )/ $m_{\tilde{t}}$ 

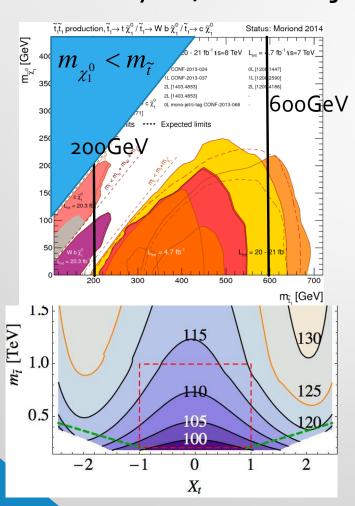


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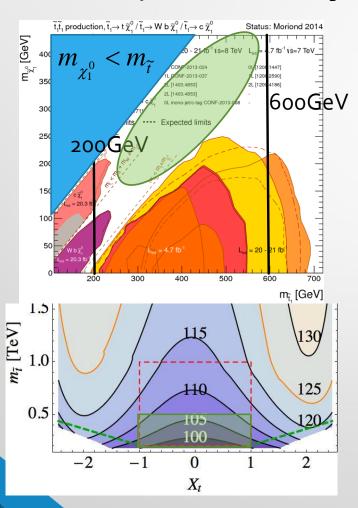
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### But, don't worry! Maybe, "she" is just around the corner!



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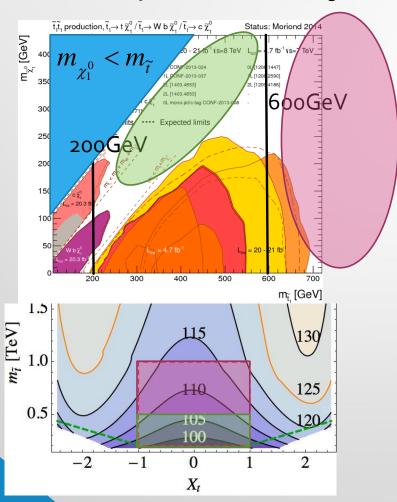
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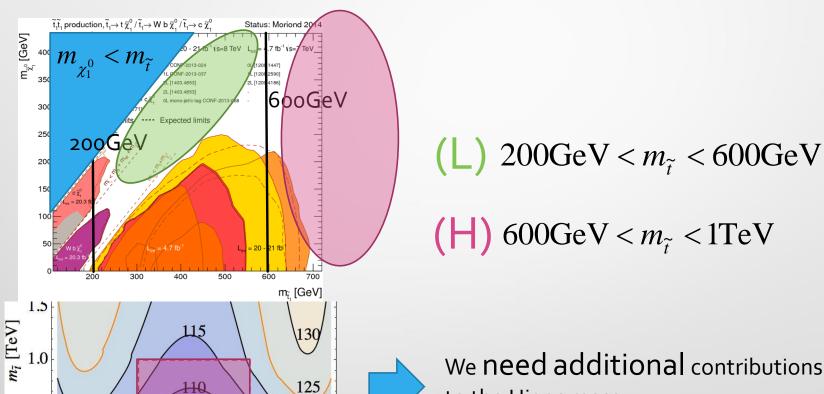


(L) 
$$200 \text{GeV} < m_{\tilde{t}} < 600 \text{GeV}$$

(H) 
$$600 \text{GeV} < m_{\tilde{t}} < 1 \text{TeV}$$

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0.5

0  $X_t$  to the Higgs mass.

Add a SM gauge singlet supermultiplet!

Z3-invariant NMSSM, nMSSM, PQ-NMSSM, U(1)-extended MSSM, ...

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F-term potential of S

But it can be sizable **only** when

$$2 < \tan \beta < 4$$
 (perturbative bounds)

$$\tan \beta = \langle H_{\rm u}^{\,0} \rangle / \langle H_{\rm d}^{\,0} \rangle$$

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With a light singlet boson



Mixing with a singlet

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- But, it is still alive.

(perturbative bounds)

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Mixing with a singlet

$$m_{S} < m_{h} = 126 \text{GeV}$$

$$h = \hat{h}\cos\theta_1\cos\theta_2 - \hat{H}\sin\theta_1 - \hat{s}\cos\theta_1\sin\theta_2$$

observed

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observed

SM-like

MSSM Higgs

singlet

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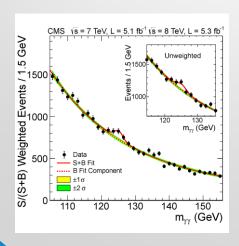
observed

SM-like

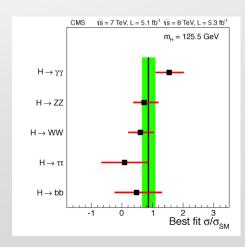
MSSM Higgs

singlet

$$m_h = 126 \text{GeV}$$



#### $h \approx \hat{h}$

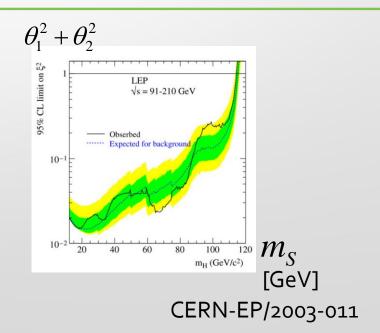


```
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H = \cdots
s = \cdots
```

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The LEP Higgs search



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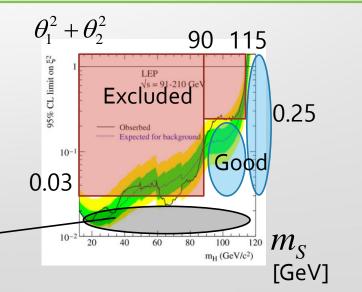
$$H = \cdots$$

$$s = \cdots$$

The LEP Higgs search

$$90 - 115 \text{GeV} < m_S$$

Not enough to raise the Higgs mass

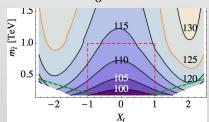


CERN-EP/2003-011

$$\hat{h} \qquad \hat{S} \\
\begin{pmatrix}
m_0^2 + (\lambda^2 v^2 - m_Z^2)\sin^2 2\beta & -(\lambda^2 v^2 - m_Z^2)\sin 2\beta \cos 2\beta & \lambda v(2\mu - \Lambda \sin 2\beta) \\
- & -(\lambda^2 v^2 - m_Z^2)\sin^2 2\beta + \frac{2b}{\sin 2\beta} & \lambda v\Lambda \cos 2\beta \\
- & - & m_{\hat{S}}^2
\end{pmatrix}$$

 $b,\Lambda,m_{\hat{S}}^2$  : parameters depending on the model

 $m_0^2$ : Z boson mass + Q.C. (top, stop, ...)



$$W = \lambda SH_u H_d + \cdots$$

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- \qquad \qquad -(\lambda^2 v^2) - m_Z^2 \end{pmatrix} \sin^2 2\beta + \frac{2b}{\sin 2\beta} \qquad \lambda v \wedge \cos 2\beta \\
- \qquad \qquad - \qquad \qquad m_{\hat{S}}^2$$

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The higgsino mass  $\mu$  is constrained by the Higgs results!

$$W = \lambda SH_u H_d + \cdots$$

$$W = (MSSM - Yukawas) + \lambda SH_uH_d + f(S)$$

Anything is OK.

No CPV is assumed.

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 $0.01 < \lambda < 1$ 

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$$b \rightarrow s \gamma$$

$$m_S < m_h = 126 \text{GeV}$$
  
350 GeV  $< m_H < 1 \text{TeV}$ 

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Let us take 
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m GeV}$$

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(L) 
$$200 \text{GeV} < m_{\tilde{t}} < 600 \text{GeV}$$

 $100 \text{GeV} < m_0 < 115 \text{GeV}$ 

(H) 
$$600 \text{GeV} < m_{\tilde{t}} < 1 \text{TeV}$$

 $105 \text{GeV} < m_0 < 120 \text{GeV}$ 

### Why does µ get small?

When m\_H is rather heavy,

$$|\lambda\mu| \cong \frac{1}{2\nu} \left[ \frac{2|\theta_1\theta_2|}{\theta_1^2 + \theta_2^2} \frac{1}{\tan\beta} (m_H^2 - m_S^2) + |\theta_2| (m_h^2 - m_S^2) \right]$$

$$\lambda^{2} \cong \frac{m_{Z}^{2}}{v^{2}} + \frac{\tan^{2} \beta}{4v^{2}} \left[ \left( m_{h}^{2} - m_{0}^{2} \right) - \left( \theta_{1}^{2} + \theta_{2}^{2} \right) \left( m_{h}^{2} - m_{S}^{2} \right) \right]$$

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Lower bound

Bounded above

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**Bounded below** 

Upper bound



Thus  $\mu$  is bounded above!

# How small is μ?Simple example -

(L) If 
$$\sqrt{\theta_1^2 + \theta_2^2} < 0.58$$
 (H) If  $\sqrt{\theta_1^2 + \theta_2^2} < 0.44$ 

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$$> 0$$

$$|\lambda \mu| \approx \frac{1}{2v} \left[ \frac{2 |\theta_1 \theta_2|}{\theta_1^2 + \theta_2^2} \frac{1}{\tan \beta} (m_H^2 - m_s^2) + |\theta_2| (m_h^2 - m_s^2) \right]$$

$$|\mu| < 350 \text{GeV} \quad (\tan \beta = 10)$$

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- Additional contributions to the effective couplings
  - are dominated by those of the higgsinos and the stops
  - are taken freely, but consistently with our parameter space

Then we obtain the  $1\sigma$  and  $2\sigma$  preferred regions

- using the chi-squared distribution with 5 or 6 DOF
- taking the parameters of higgsinos and stops that minimize the chi-squared.

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	WW/ggF	ZZ/ggF	bb/VH-VBF	ττ/VH-VBF	γγ/Χ	γγ/Υ
ATLAS	0.99±0.30	1.43±0.38	$1.09 \pm 0.34$		1.66±0.37	1.72±0.79
CMS	$0.68 \pm 0.20$	$0.92 \pm 0.28$	$1.15 \pm 0.62$	$1.10 \pm 0.41$	$0.65 \pm 0.34$	$1.80 \pm 0.89$

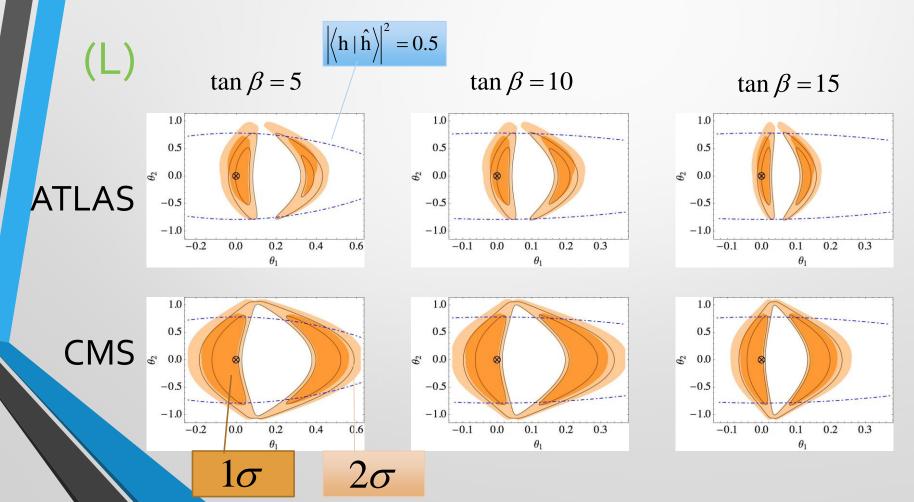
$$\mu^{X} - 1 = (\mu^{ggF} - 1)\cos\varphi + (\mu^{VH/VBF} - 1)\sin\varphi$$

$$\mu^{Y} - 1 = -(\mu^{ggF} - 1)\sin\varphi + (\mu^{VH/VBF} - 1)\cos\varphi$$

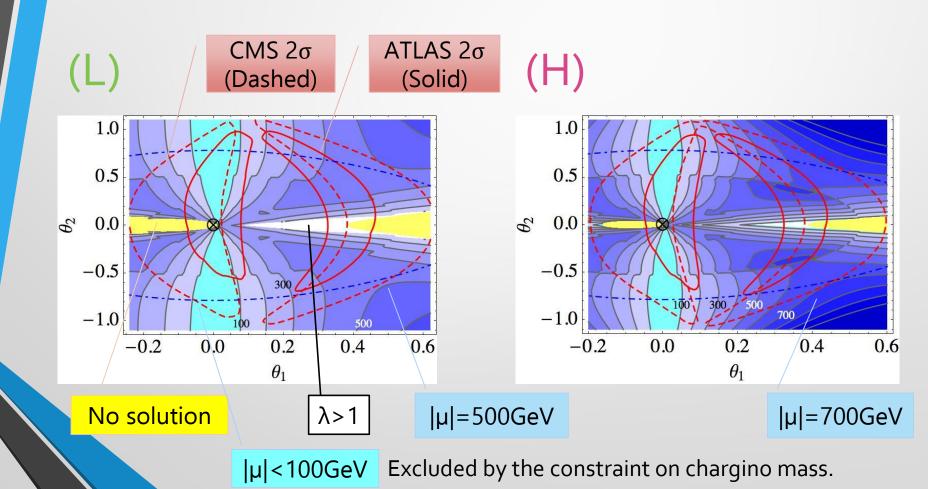
$$\cos\varphi = 0.98(ATLAS), 0.97(CMS)$$

hep-ex/1307.1427, ATLAS-CONF-2014-009, CMS-PAS-HIG-13-005

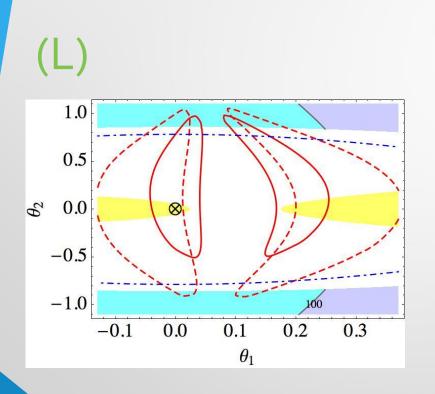
# How large $\theta_1$ and $\theta_2$ can be? -the Higgs signal strengths-

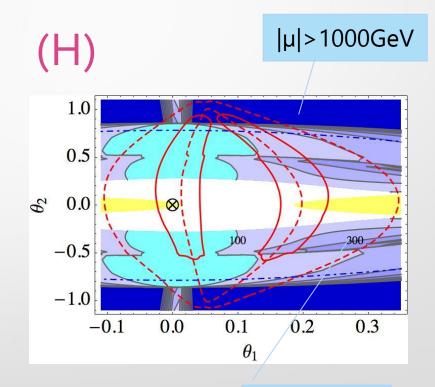


#### Numerical results ( $tan\beta=5$ )



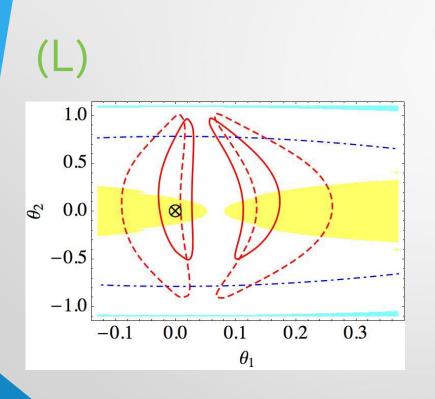
#### Numerical results (tanβ=10)

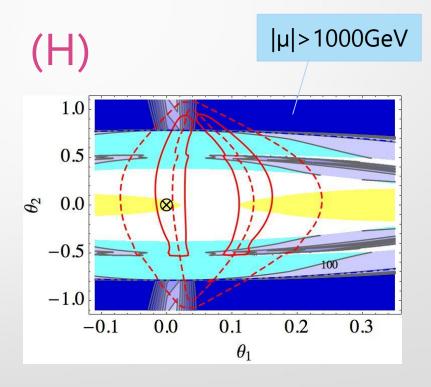




 $|\mu| = 300 \text{GeV}$ 

# Numerical results (tanβ=15)





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Then, in much of the parameter region, the higgsinos tend to be light!

Find the stop, find the deviation, then the higgsinos are just beside you!

Backups

#### gamma gamma channels

