

Precise top-quark mass at hadron colliders

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(SeoulTech)

Based on

- SK and H. Yokoya, arXiv:1607.00990
- SK, Y. Shimizu, Y. Sumino and H. Yokoya, PLB741 (2015) 232-238

Outline

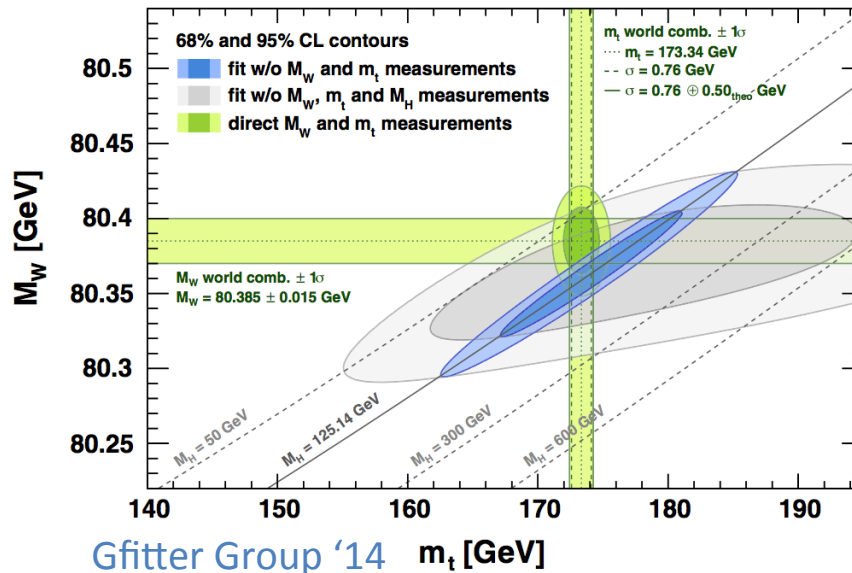
1. Introduction to m_t determination
2. Diphoton mass spectrum を用いる方法
3. Lepton energy distribution を用いる方法
4. Summary

1. Introduction

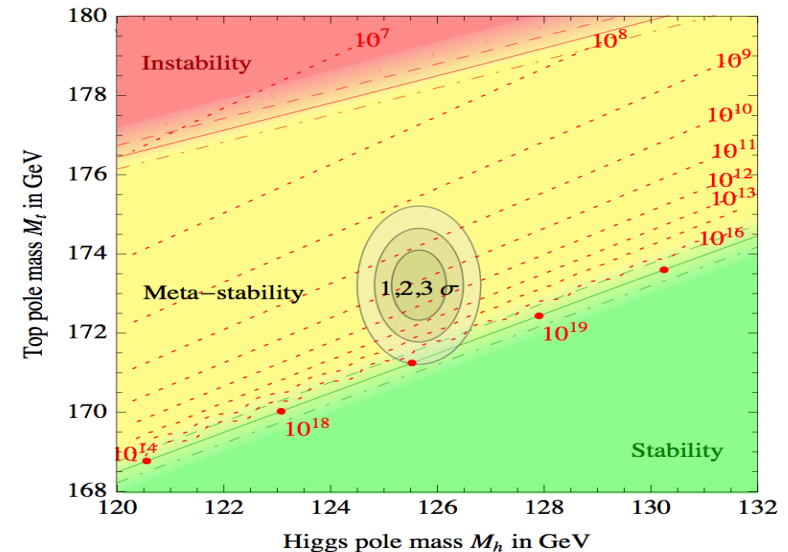
Motivation to measure m_t

- One of the fundamental parameters of the SM
- Important input to various physics

★ EW precision tests for SM



★ SM vacuum stability



Buttazzo et al. '13

★ Beyond SM

Current status at hadron colliders

Tevatron+LHC m_t combination [arXiv:1403.4427](https://arxiv.org/abs/1403.4427)

$m_t = 173.34 \pm 0.76 \text{ GeV}$ 0.4 % precision !

Current status at hadron colliders

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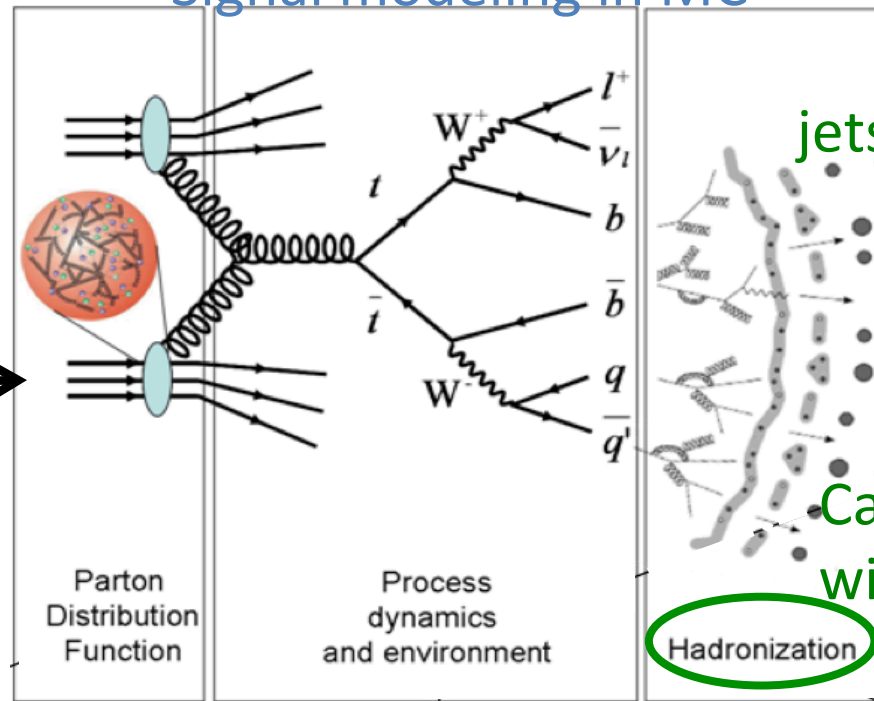
What kind of mass? $\neq m_t^{\text{pole}}, m_t^{\overline{\text{MS}}}$ “Monte-Carlo mass”

Signal modeling in MC

Experiment

m_t measurement

Theory (MC)



Cannot be treated within pert. QCD

Current status at hadron colliders

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What kind of mass? $\neq m_t^{\text{pole}}, m_t^{\overline{\text{MS}}}$ “Monte-Carlo mass”

$m_t^{\text{pole}}, m_t^{\overline{\text{MS}}}$ measurements

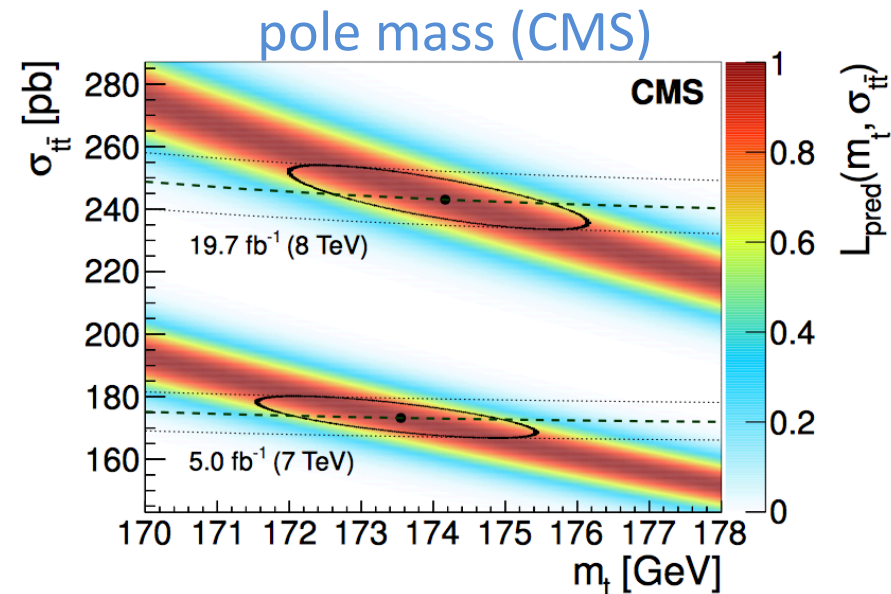
◆ $m_t^{\text{pole}} = 173.8 \pm 2 \text{ GeV}$

CMS, JHEP 08 (2016) 029

◆ $m_t^{\overline{\text{MS}}} = 160.0^{+5.1}_{-4.5} \text{ GeV}$

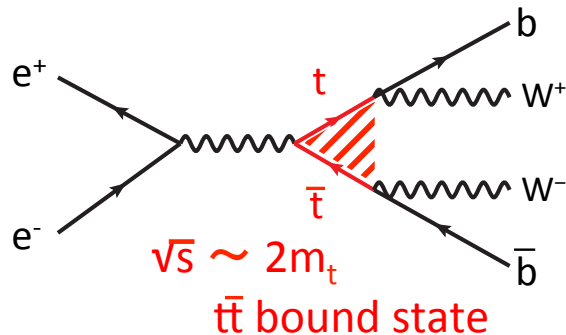
D0, PLB 703 (2011) 422

The errors are still large.

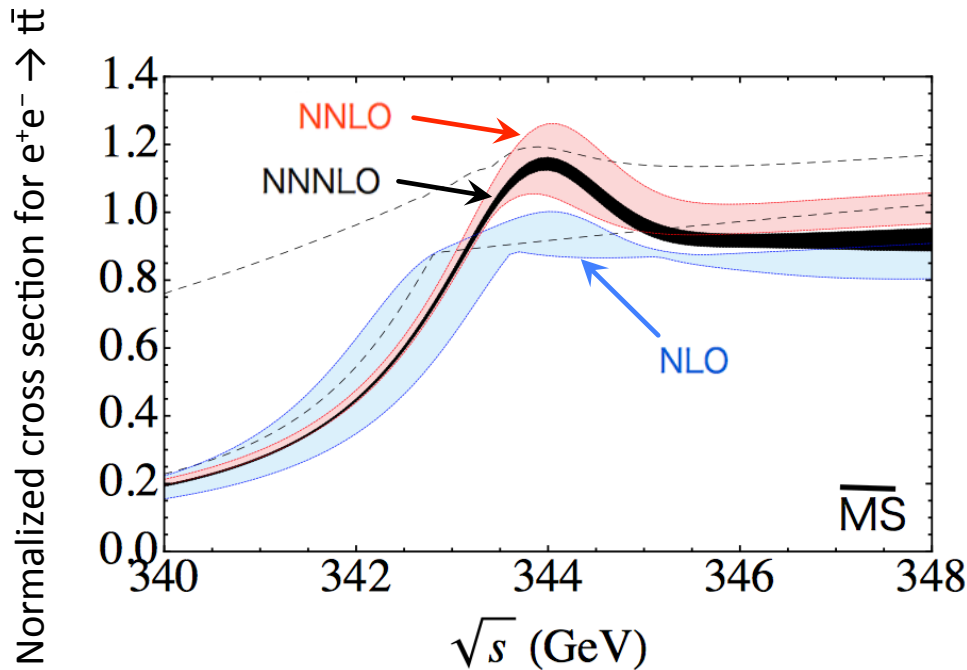


m_t measurement at future e^+e^- colliders

Threshold scan of $t\bar{t}$ production



- N³LO: Beneke, Kiyo, Marquard, Penin, Piclum, Steinhauser, PRL 115 (2015)192001
- ILC study: Horiguchi, Ishikawa, Suehara, Fujii, Sumino, Kiyo, Yamamoto, arXiv:1303.3758
- CLIC study: Seidel, Simon, Tesar, Poss, Eur.Phys.J.C73 (2013)2530



Kiyo, Top WS@Valencia 2015

Peak position
 $\sim 1\sigma$ resonance mass

\overline{MS} mass

$\Delta m_t^{\overline{MS}} \sim 30$ MeV in principle

Kiyo, Mishima, Sumino, JHEP11 (2015) 084

Aim of this study

e^+e^- collider と違って、hadron collider で
摂動QCD における定義の明確な top mass を
高精度で決定することは困難



この困難を克服する2つの方法を提案する

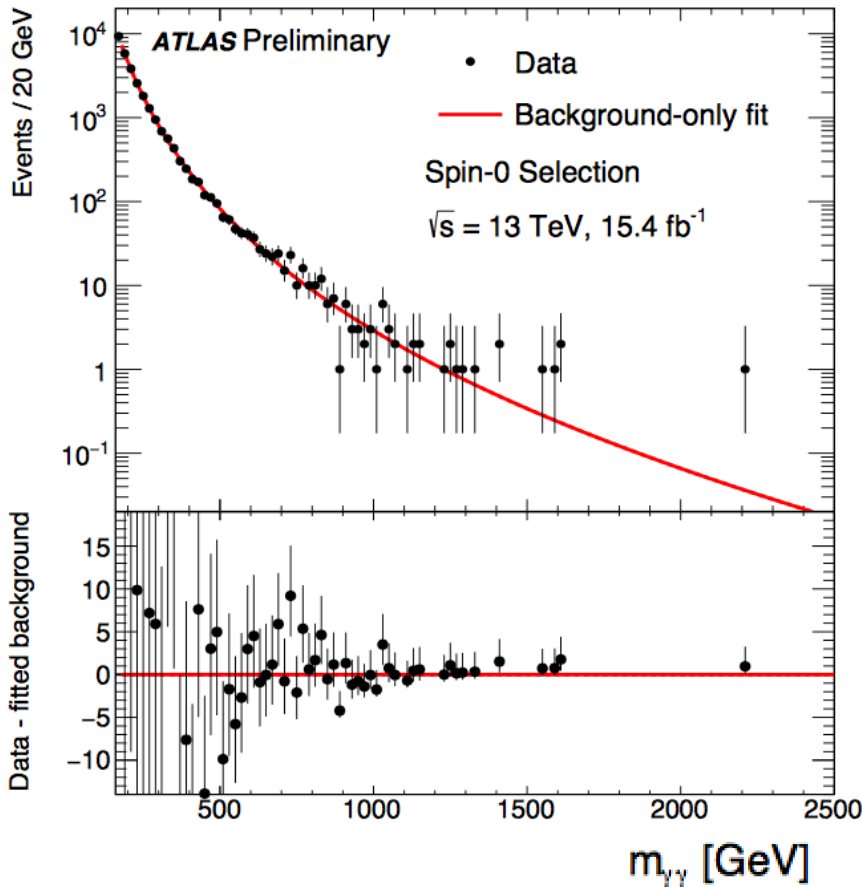
- **Diphoton mass spectrum を用いる方法**
SK and H. Yokoya, arXiv:1607.00990
- **Lepton energy distribution を用いる方法**
SK, Y. Shimizu, Y. Sumino and H. Yokoya,
PLB741 (2015) 232-238

2. Diphoton mass spectrum を 用いる方法

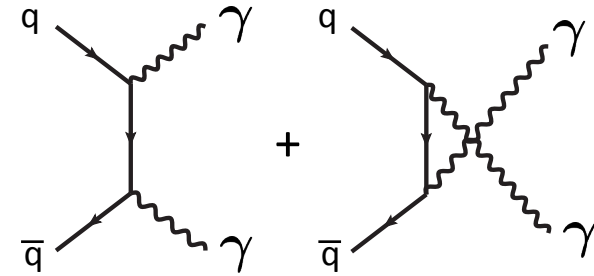
SK and H. Yokoya, arXiv:1607.00990

Diphoton mass spectrum

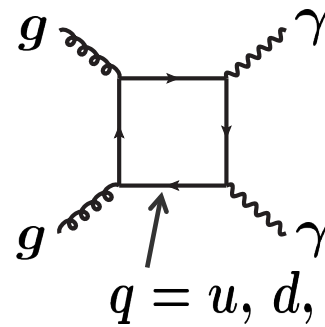
- $H \rightarrow \gamma\gamma$
- New resonance search



$$pp \rightarrow \gamma\gamma$$



$$\mathcal{O}(\alpha^2)$$

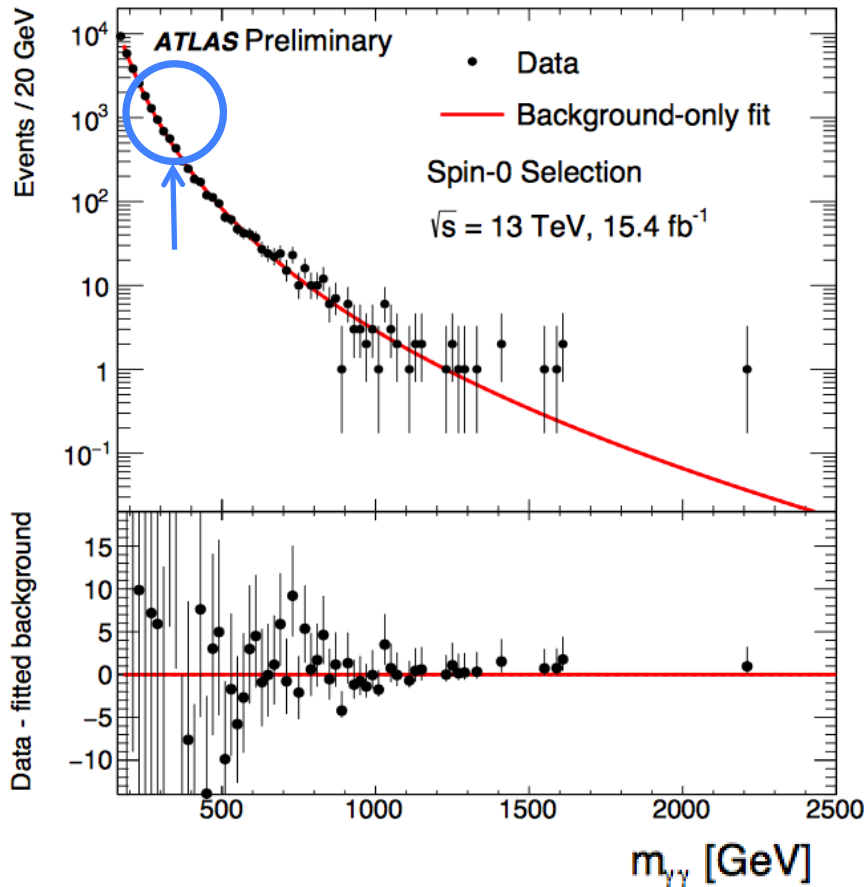


$$\mathcal{O}(\alpha^2 \alpha_s^2)$$

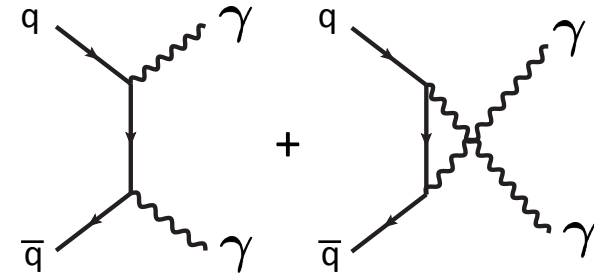
But large gluon PDF

Diphoton mass spectrum

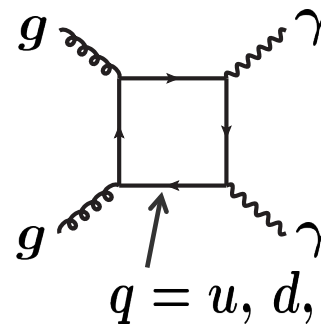
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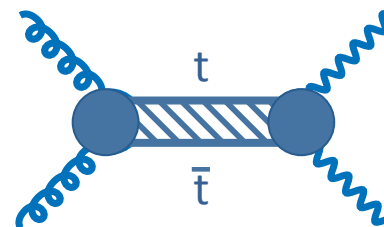
$$\mathcal{O}(\alpha^2)$$



$$\mathcal{O}(\alpha^2 \alpha_s^2)$$

But large gluon PDF

$$m_{\gamma\gamma} \sim 2 m_t$$



Bound-state effects

gg \rightarrow $\gamma\gamma$ amplitude

• Two-loop amplitude for massless-quark loop Bern, De Freitas, Dixon, JHEP0109 (2001) 037

• One-loop amplitude for 5 massless-quark + top-quark loop

Dicus, Willenbrock, PRD37 (1988) 1801

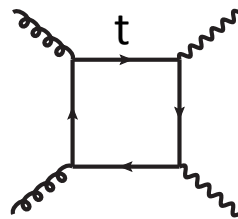
Campbell, Ellis, Li, Williams, arXiv:1603.02663



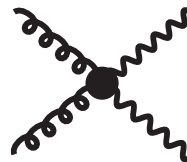
Bound-state effectsを含める (this work)

Near $t\bar{t}$ threshold

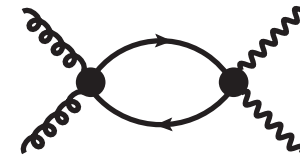
$$v = \sqrt{1 - 4m_t^2/m_{\gamma\gamma}^2} \text{ で展開}$$



=



+



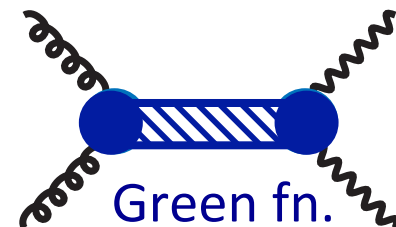
+ $\mathcal{O}(v^2)$



Schrodinger eq.

$$\left[\left\{ -\frac{\nabla^2}{m_t} + V_{\text{QCD}}(r) \right\} - \mathcal{E} \right] G(\vec{r}; \mathcal{E}) = \delta^3(\vec{r})$$

$$\text{Green fn.} = \text{loop} + \alpha_s/v \text{ loop} + (\alpha_s/v)^2 \text{ loop} + (\alpha_s/v)^3 \text{ loop} + \dots$$



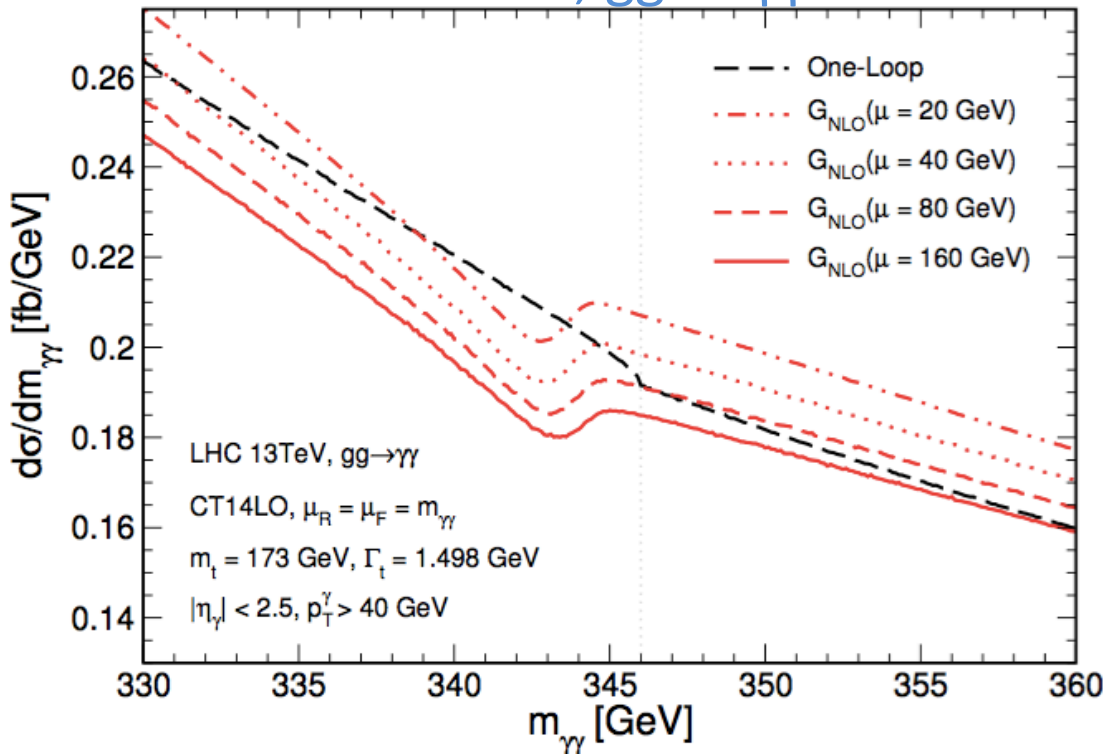
Green fn.
in NRQCD

Diphoton mass spectrum near threshold

$$\left| \begin{aligned}
 & \text{Light quarks} + \text{Crossed} + \text{Re} \left(\text{t-loop} \right) + i \text{Im} \left(\text{t-loop} \right) + \mathcal{O}(v^2) \end{aligned} \right|^2$$

$e^+e^- \rightarrow t\bar{t}$ の場合

LHC13TeV, $gg \rightarrow \gamma\gamma$



Theoretical uncertainties

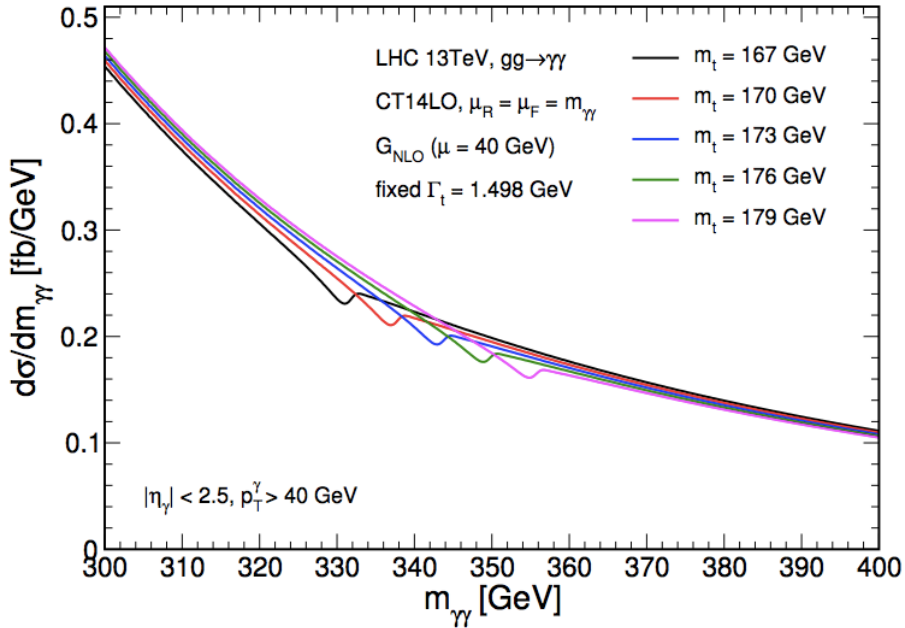
◆ Green fn. scale

- Overall normalization $\sim 10\%$
- Positions of dip and bump \sim sub-GeV level

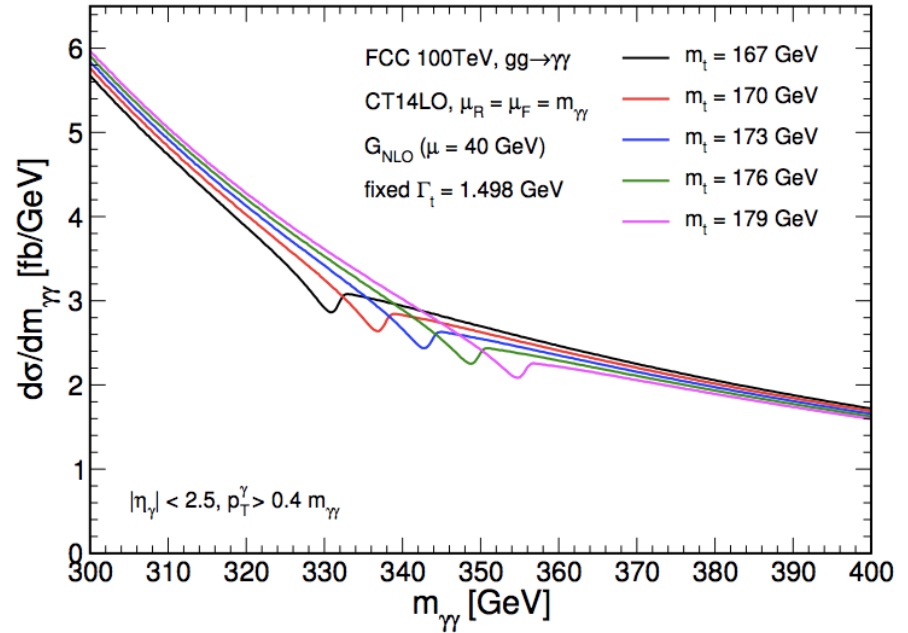
◆ PDF and α_s scale $\sim 20\%$

m_t measurement at LHC and FCC

LHC 13TeV



FCC 100TeV



Simulation

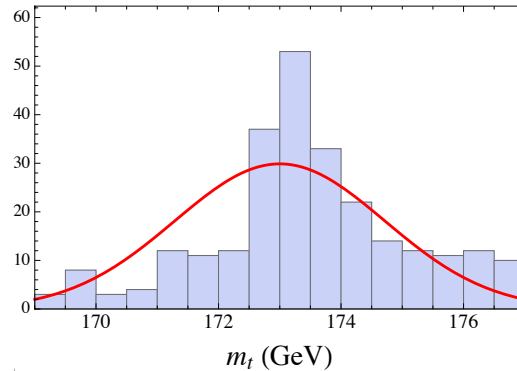
- $gg \rightarrow \gamma\gamma$ generated according to our calculation
 - $q\bar{q} \rightarrow \gamma\gamma$
 - Fragmentation contributions
- } Diphox (LO)

Fitting fn : (gg prediction) + analytic smooth fn.

Sensitivity

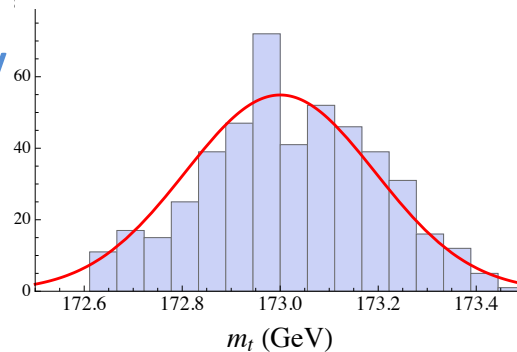
Estimate of statistical error

LHC $\sqrt{s} = 13\text{TeV}$
 3 ab^{-1}



$\Delta m_t \sim 2\text{ GeV}$ for 3 ab^{-1}

FCC $\sqrt{s} = 100\text{TeV}$
 $1\text{ ab}^{-1}, 10\text{ ab}^{-1}$



$\Delta m_t = 0.2\text{ GeV}$ for 1 ab^{-1}

$\Delta m_t = 0.06\text{ GeV}$ for 10 ab^{-1}

- At FCC, systematic uncertainties should dominate.
Photon energy scale, isolation cuts, ...
- Higher-order QCD corrections for $pp \rightarrow \gamma\gamma X$ would be important.
NLO $q\bar{q}$, LO and NLO qg , NLO gg

3. Lepton energy distribution を 用いる方法

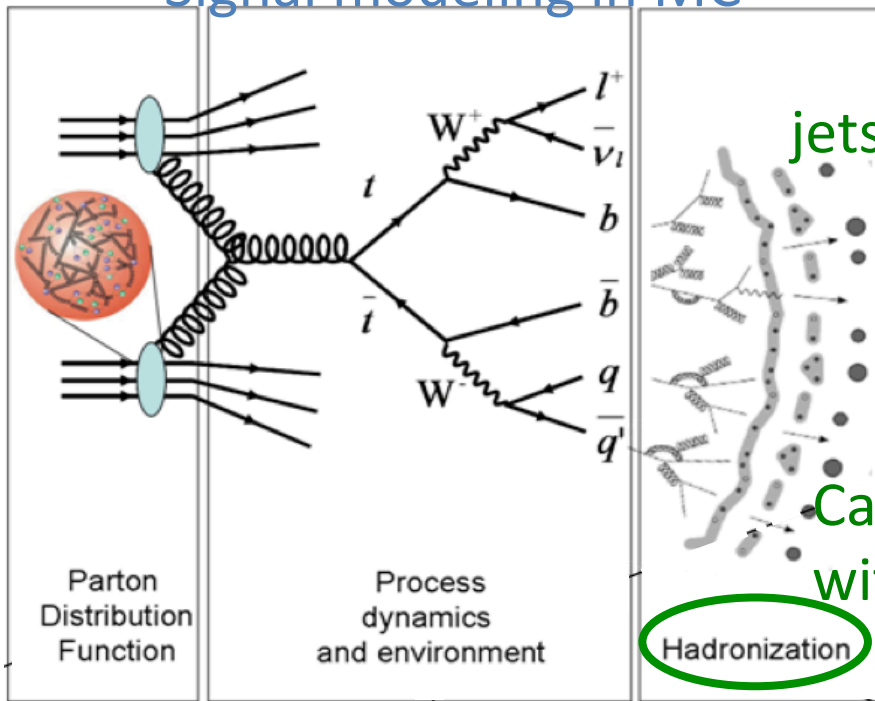
SK, Y. Shimizu, Y. Sumino and H. Yokoya,
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Weight function method

SK, Y.Shimizu, Y.Sumino, H.Yokoya, PLB 710, 658 (2012)
SK, Y.Shimizu, Y.Sumino, H.Yokoya, JHEP 08, 129 (2013)

- Only **lepton energy distribution** is needed
- **Independent** of top-quark velocity distribution

Signal modeling in MC



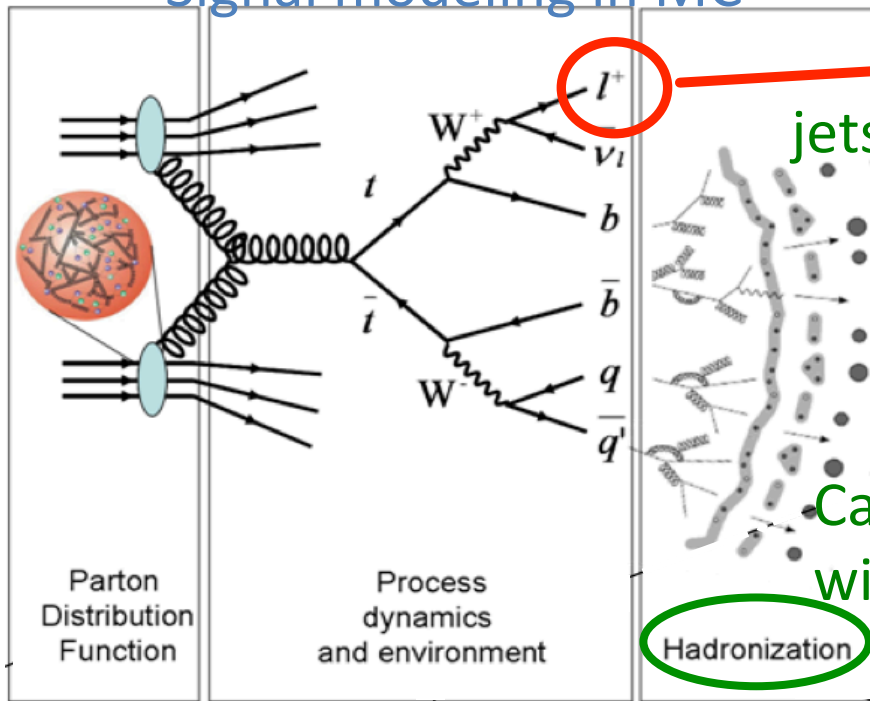
Cannot be treated
within pert. QCD

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- Only **lepton energy distribution** is needed
- **Independent** of top-quark velocity distribution

Signal modeling in MC



Free from ambiguity in jets



Theoretically well-defined m_t

Cannot be treated
within pert. QCD

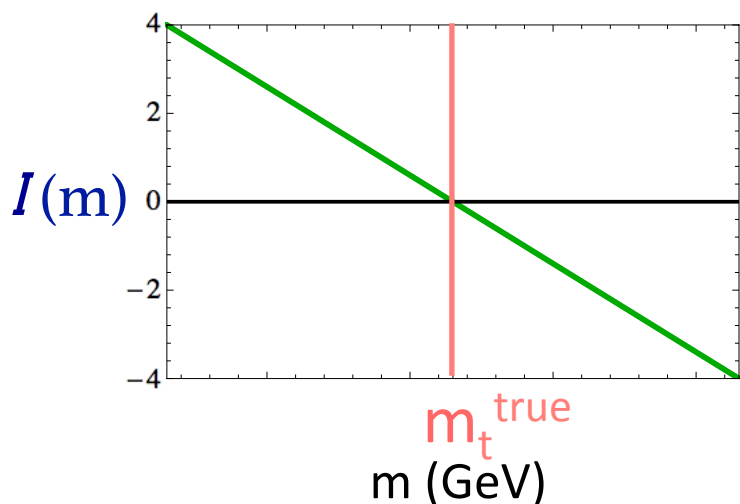
Method

1. Compute a weight function $W(E_\ell, m)$

$$W(E_\ell, m) = \int dE \mathcal{D}_0(E; m) \frac{1}{EE_\ell} (\text{odd func. of } \rho) \Big|_{e^\rho = E_\ell/E}$$

Lepton energy dist. in the **rest frame** of top quark (with mass m)

2. Using lepton energy dist. **from experiment** $D(E_\ell)$, perform its weighted integration



$$I(m) \equiv \int dE_\ell D(E_\ell) W(E_\ell, m)$$

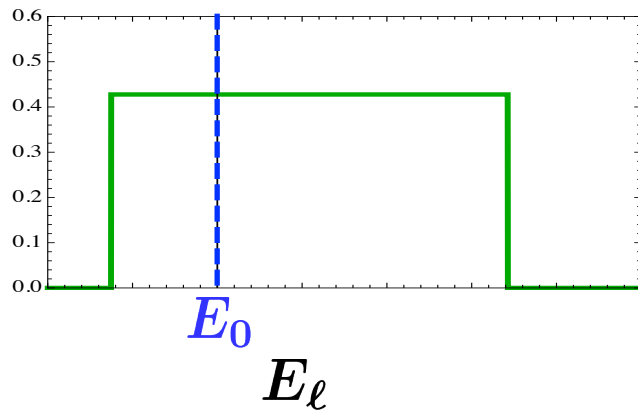
3. Obtain m_t^{true} from

$$I(m = m_t^{\text{true}}) = 0$$

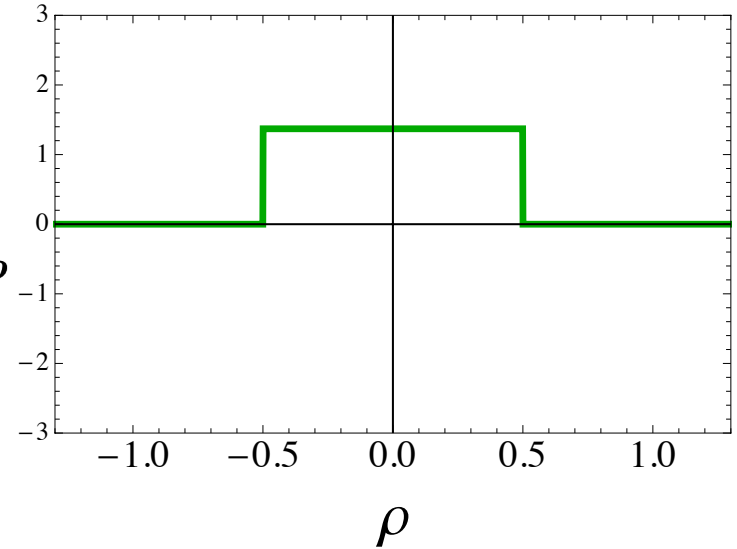
Construction of weight function

For a two-body decay : $X \rightarrow \ell + Y$ (X is unpolarized)

Lepton energy distribution



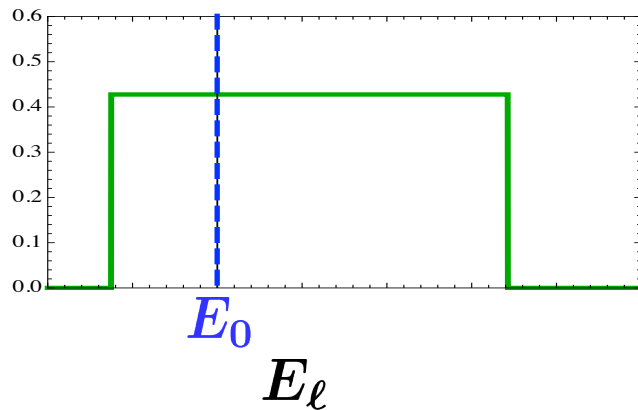
$$E_\ell / E_0 = e^\rho$$



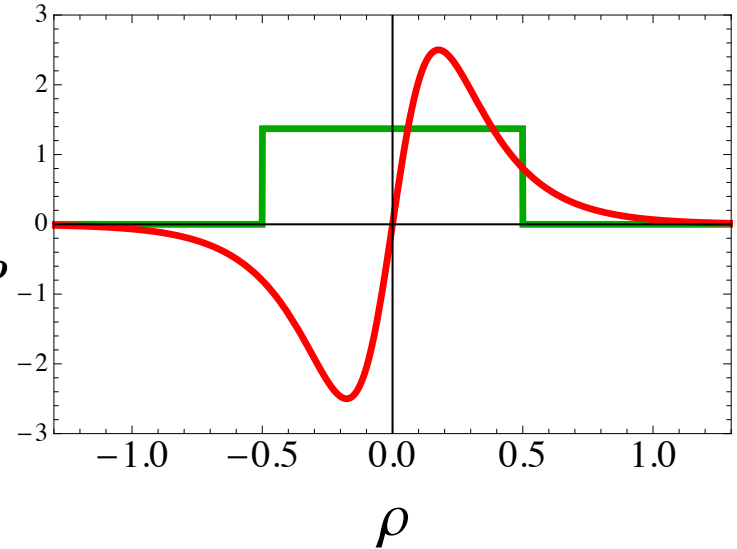
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$$E_\ell / E_0 = e^\rho$$



$$\int dE_\ell D(E_\ell) W(E_\ell, m_X^{\text{true}}) = 0 \iff \int d\rho (\text{even fn. of } \rho)(\text{odd fn. of } \rho) = 0$$

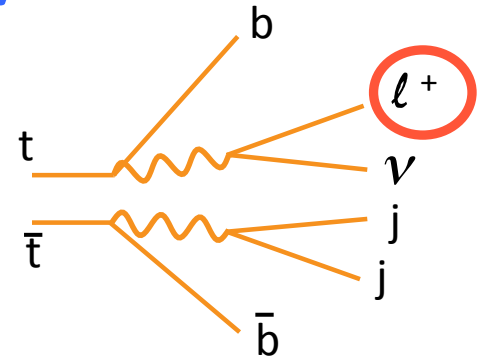
$$dE_\ell \propto d\rho e^\rho$$



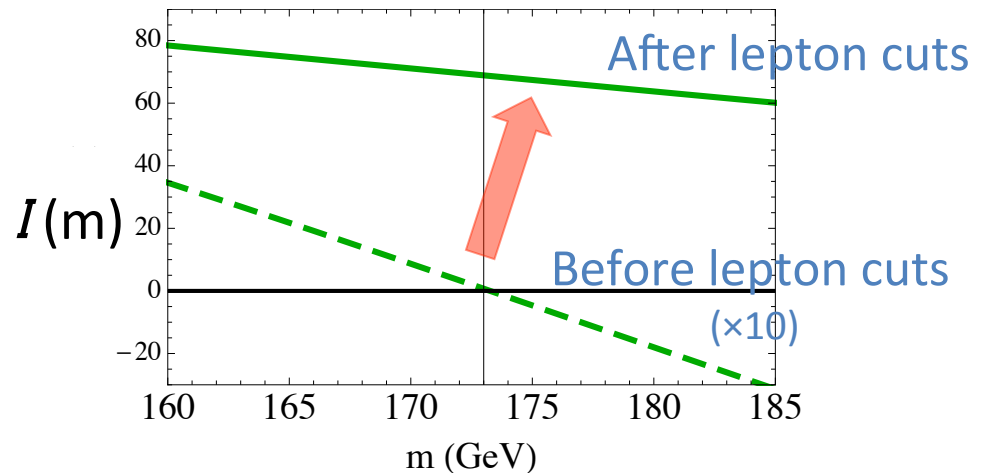
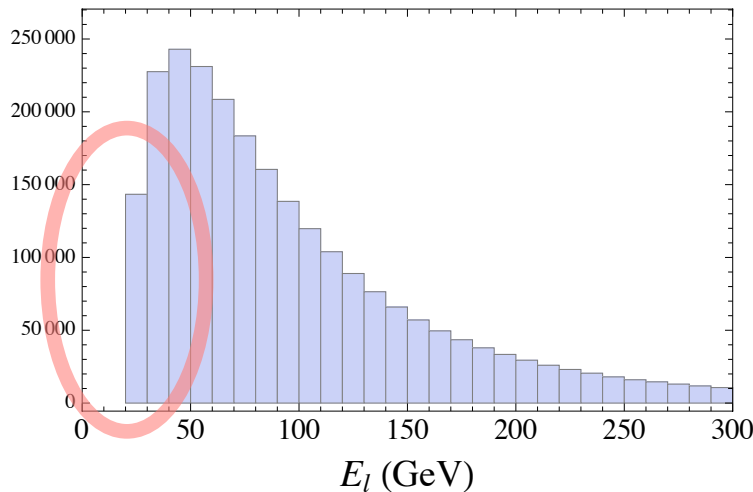
$$W(E_\ell, m_X^{\text{true}}) = e^{-\rho} (\text{odd func. of } \rho) \Big|_{e^\rho = E_\ell / E_0}$$

Simulation analysis

- LHC $\sqrt{s} = 14$ TeV
- $t\bar{t}$ events, Lepton+jets channel



Event selection cuts and background deform lepton distribution.



Parton level の lepton distribution に戻す方針

Sensitivity of m_t determination

Uncertainties [GeV] (LO analysis)

| | |
|--------------------|-----------|
| Signal stat. error | 0.4 |
| μ_F scale | +1.5/-1.4 |
| PDF | 0.6 |
| Jet energy scale | +0.2/-0.0 |
| BG stat. error | 0.4 |

← At 100 fb⁻¹, Lepton+jets channel

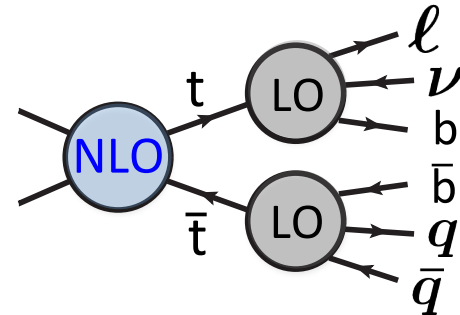
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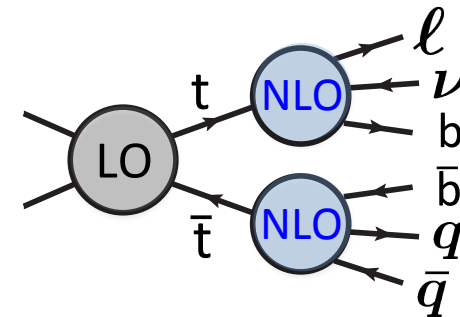
← At 100 fb⁻¹, Lepton+jets channel

← Can be improved by including NLO



Scale uncertainty in NLO top decay

+0.3/-0.2 GeV preliminary



We aim for $\Delta m_t^{\text{pole}} < 1$ GeV

4. Summary

- Precise determination of a theoretically well-defined m_t is demanded.
- We proposed two methods for precise m_t measurement at hadron colliders.

- Diphoton mass spectrum を用いる方法

$$\Delta^{\text{stat}} m_t = 0.06 \text{ GeV for } 10 \text{ ab}^{-1} \text{ at FCC}$$

- Lepton energy distribution を用いる方法

$$\Delta m_t^{\text{pole}} < 1 \text{ GeV is probable at LHC}$$