

Search for light scalar boson in minimal $U(1)_{L\mu-L\tau}$ model at ILC experiment

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in collaboration with

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***“Light Z’ boson from scalar boson decay
at collider experiment in an $U(1)_{L\mu-L\tau}$ model”***

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August 7th, 2018 @ PPP2018

Introduction

- The discrepancy of muon (g-2) is a long standing problem in particle physics.

$$\Delta a_{\mu} \equiv \Delta a_{\mu}^{\text{exp}} - \Delta a_{\mu}^{\text{th}} = (28.8 \pm 8.0) \times 10^{-10},$$

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$$\Delta a_\mu^{NP} \sim \frac{f^2}{4\pi^2} \frac{m_\mu^2}{M^2} \sim 10^{-9}$$

M : mass of new particle
 f : coupling const with muon

$$\longrightarrow \frac{f}{M} \sim 2 \times 10^{-3} \text{ (/GeV)}$$

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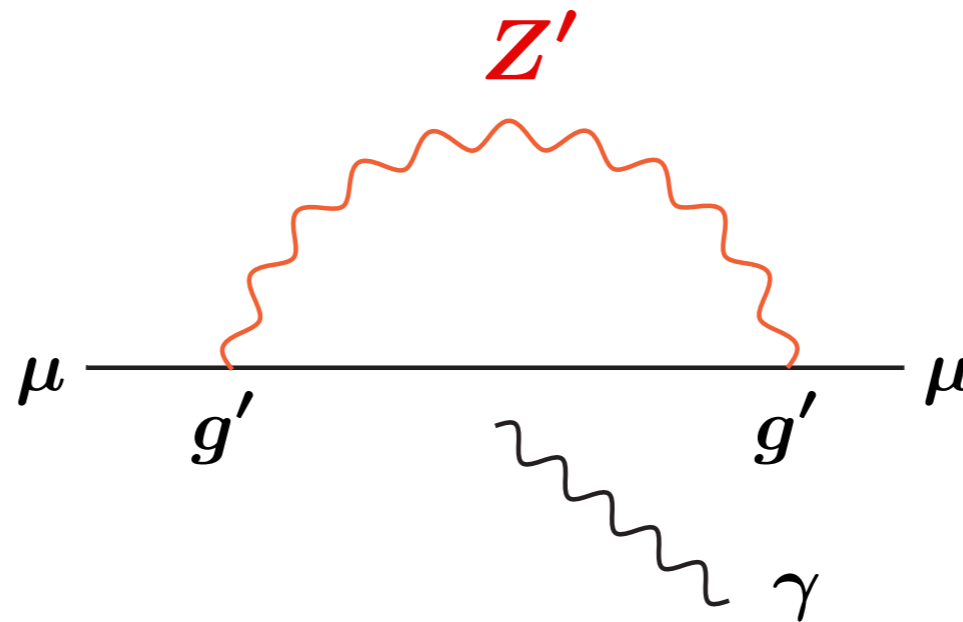
The scale of New Physics (NP)

- High scale NP, e.g. $M=1$ TeV and $f = O(1)$
- **Low scale NP, e.g. $M=100$ MeV and $f = O(10^{-4})$**

Introduction

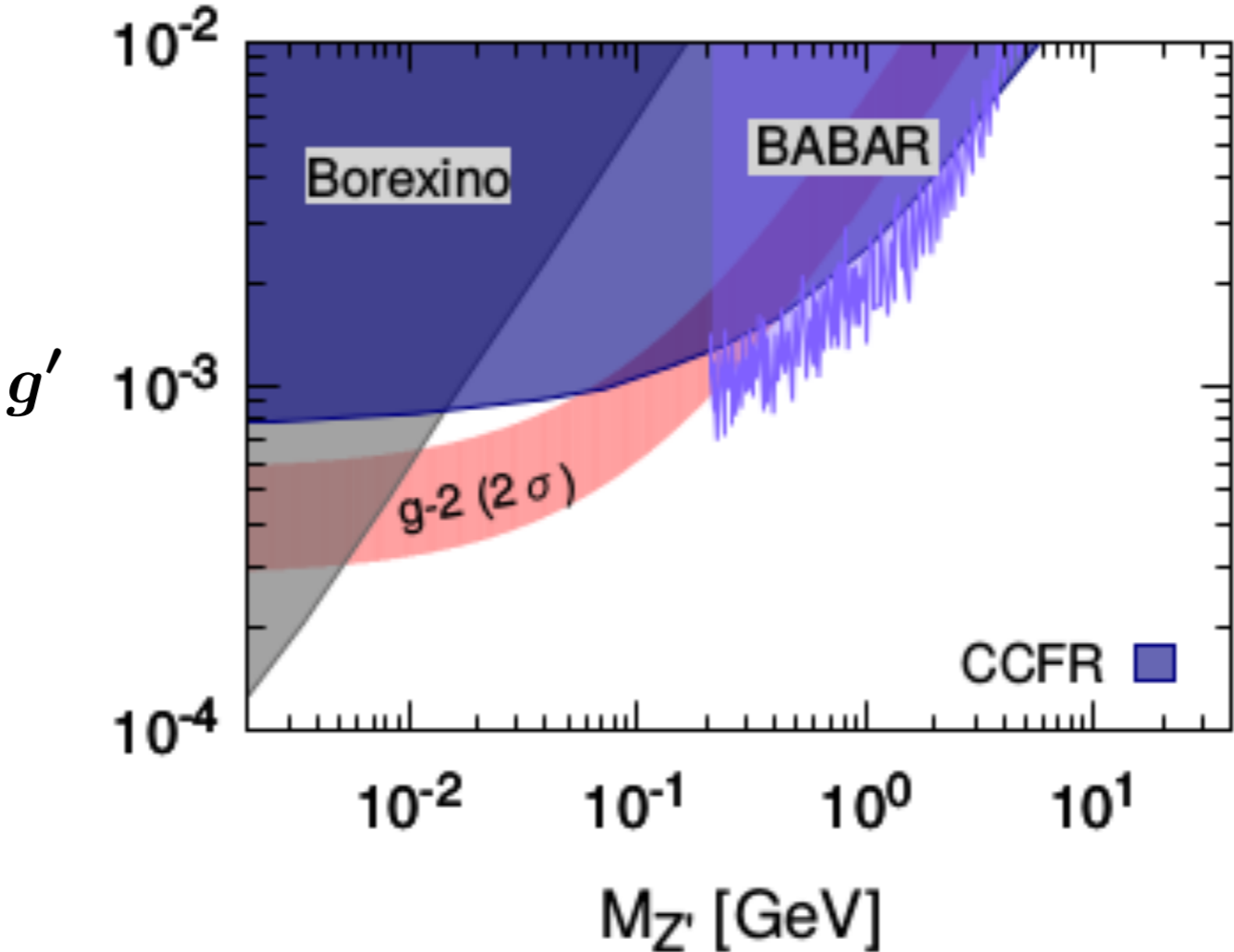
Gauged $U(1)_{L_\mu-L_\tau}$ model

Muon interacts with a gauge boson, Z' , of the symmetry



Introduction

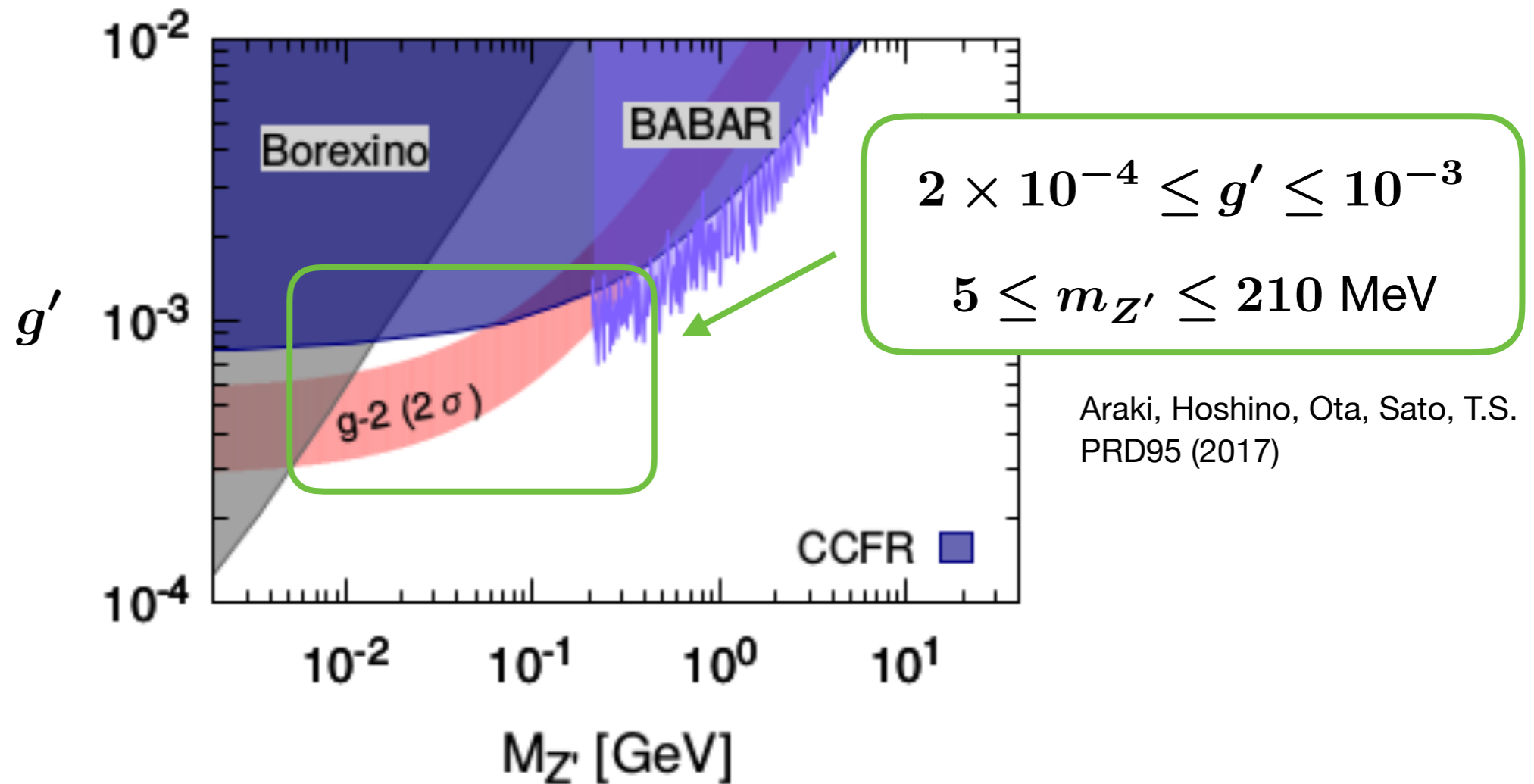
Gauged $U(1)_{L\mu-L\tau}$ model



Araki, Hoshino, Ota, Sato, T.S.
PRD95 (2017)

Introduction

Gauged $U(1)_{L\mu-L\tau}$ model



Introduction

The $U(1)_{L\mu-L\tau}$ symmetry should be broken so that Z' acquires mass.

A new scalar boson must exist to break the symmetry spontaneously.

- The vev of the scalar can be estimated as

$$v_\phi = \frac{m_{Z'}}{g'} \sim \mathcal{O}(10 - 100) \text{ GeV}$$

- The mass of the scalar will be the same order or less

$$m_\phi = \sqrt{\lambda} v_\phi \leq \mathcal{O}(100) \text{ GeV}$$

Purpose

There are two new particles at least

- Feebly int. gauge boson with mass $< O(100)$ MeV
Low energy with high luminosity, i.e. Belle-II
 - Araki, Hoshino, Ota, Sato, T.S. PRD95 (2017)
 - Kaneta, T.S. PTEP 2017
- Scalar boson with mass $< O(100)$ GeV
High energy with high luminosity, ILC!

Search for the symmetry breaking scalar at ILC

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2. Minimal gauged $U(1)_{L\mu-L\tau}$ model

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4. Search at ILC experiment

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Minimal gauged $U(1)_{L_\mu-L_\tau}$ model

Minimal Gauged $U(1)_{L_\mu-L_\tau}$ model

We introduce one complex scalar to $U(1)_{L_\mu-L_\tau}$ model

- Anomaly free
- neutrino mass and mixing
- minimal setup

Choubey, Rodejohann, Eur.Phys.J, (2005)

Ota, Rodejohann, Phys.Lett. (2006)

Asai, Hamaguchi, Nagata, 1705.00419

	Scalar		Lepton					
	H	φ	L_e	L_μ	L_τ	e_R	μ_R	τ_R
$SU(2)_L$	2	1	2	2	2	1	1	1
$U(1)_Y$	$\frac{1}{2}$	0	$-\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	-1	-1	-1
$U(1)_{L_\mu-L_\tau}$	0	1	0	1	-1	0	1	-1

The Lagrangian

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + |D_\mu \varphi|^2 - \underline{V} - \frac{1}{4} Z'_{\mu\nu} Z'^{\mu\nu} - \frac{\epsilon}{2} \underline{B_{\mu\nu} Z'^{\mu\nu}} + \underline{g' Z'_\mu J_{Z'}^\mu}$$

where

gauge kinetic mixing

$$J_{Z'}^\mu = \bar{L}_\mu \gamma^\mu L_\mu + \bar{\mu}_R \gamma^\mu \mu_R - \bar{L}_\tau \gamma^\mu L_\tau - \bar{\tau}_R \gamma^\mu \tau_R$$

$$V = -\mu_H^2 H^\dagger H - \mu_\varphi^2 \varphi^* \varphi + \frac{\lambda_H}{2} (H^\dagger H)^2 + \frac{\lambda_\varphi}{2} (\varphi^* \varphi)^2 + \lambda_{H\varphi} (H^\dagger H) (\varphi^* \varphi)$$

scalar mixing

* We omit RH neutrino sector. Neutrino mass generation is discussed in Asai, Hamaguchi, Nagata

Mass eigenstates

After the EW and $U(1)_{L\mu-L\tau}$ symmetry breaking,

scalar bosons:

$$\begin{aligned} h &= \cos \alpha H + \sin \alpha \phi, \\ \phi &= -\sin \alpha H + \cos \alpha \phi, \end{aligned} \quad \tan 2\alpha = \frac{2\lambda_{H\phi} v v_\phi}{\lambda_H v^2 - \lambda_\phi v_\phi^2}$$

where v and v_ϕ are the vev of H and ϕ , respectively.

gauge boson:

$$Z_3 = Z' - \epsilon \sin \theta_W Z$$

Mass eigenstates

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The Yukawa and gauge int.

$$\begin{aligned} \mathcal{L} \supset & \sin \alpha \phi \left[\sum_f \frac{m_f}{v} \bar{f} f + \frac{m_Z^2}{v} Z_\mu Z^\mu + \frac{2m_W^2}{v} W_\mu^+ W^{-\mu} \right] \\ & + \frac{m_{Z'}^2}{v_\phi} \cos \alpha \phi Z'_\mu Z'^\mu + Z'_\mu (-e\epsilon \cos \theta_W J_{EM}^\mu + g' J_{Z'}^\mu) + \mathcal{O}(\epsilon^2) \end{aligned}$$

J_{EM} : Electromagnetic current

Decay of ϕ

Assuming $2m_t > m_\phi \gg m_{Z',f}$, the decay widths are

$$\Gamma_{\phi \rightarrow Z' Z'} \simeq \frac{g'^2 \cos^2 \alpha}{32\pi} \frac{m_\phi^2}{m_{Z'}^2} m_\phi$$

$$\Gamma_{\phi \rightarrow f \bar{f}} \simeq \frac{\sin^2 \alpha}{8\pi} \left(\frac{m_f}{v} \right)^2 m_\phi$$

Decay of ϕ

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$$\Gamma_{\phi \rightarrow Z'Z'} \simeq \frac{g'^2 \cos^2 \alpha}{32\pi} \frac{m_\phi^2}{m_{Z'}^2} m_\phi \quad \leftarrow \text{enhancement by longitudinal mode}$$
$$\Gamma_{\phi \rightarrow f\bar{f}} \simeq \frac{\sin^2 \alpha}{8\pi} \left(\frac{m_f}{v} \right)^2 m_\phi$$

To resolve $(g-2)_\mu$, $m_\phi/m_{Z'} \sim 10^3$ for $m_\phi=100$ GeV and $m_{Z'}=100$ MeV

$$\Gamma_{\phi \rightarrow Z'Z'} \gg \Gamma_{\phi \rightarrow f\bar{f}}$$

Thus,

ϕ dominantly decays into Z' pair

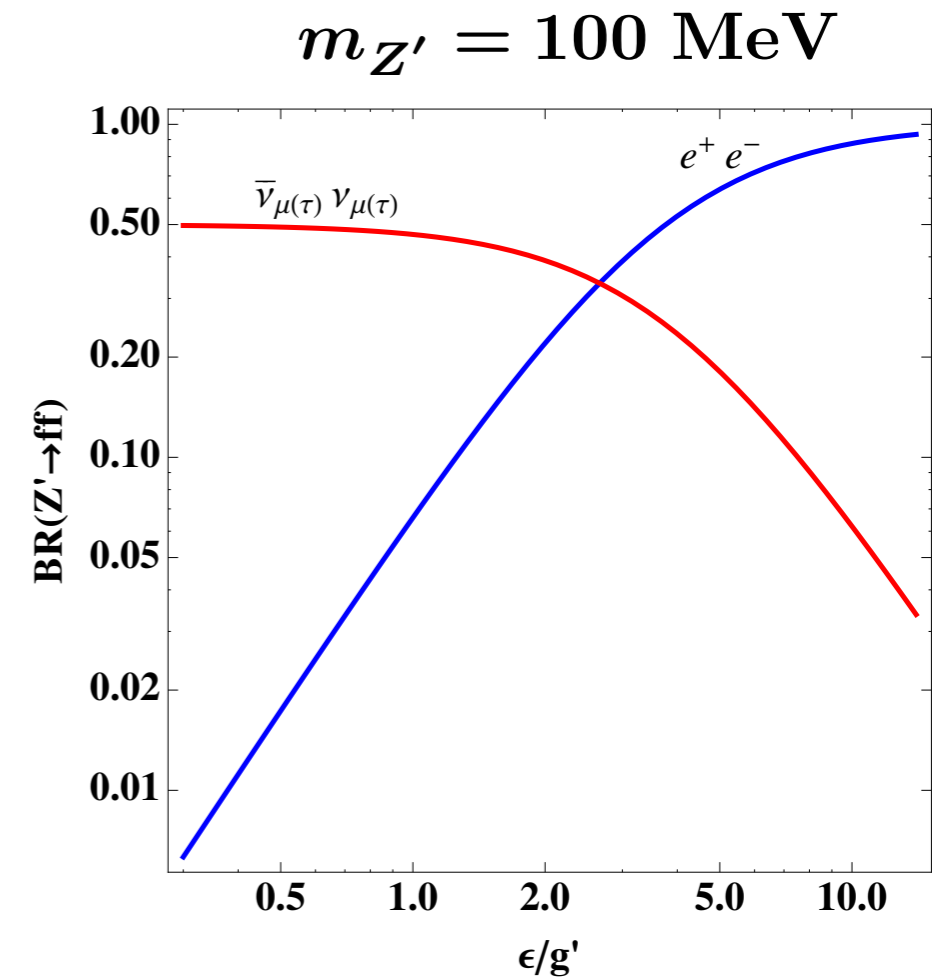
Decay of Z'

$$\Gamma_{Z' \rightarrow \nu \bar{\nu}} = \frac{g'^2}{24\pi} m_{Z'}$$

$$\Gamma_{Z' \rightarrow e^+ e^-} \simeq \frac{e^2 \epsilon^2 \cos^2 \theta_W}{12\pi} m_{Z'}$$

For $\epsilon/g' < 1$,

Z' mainly decays into ν



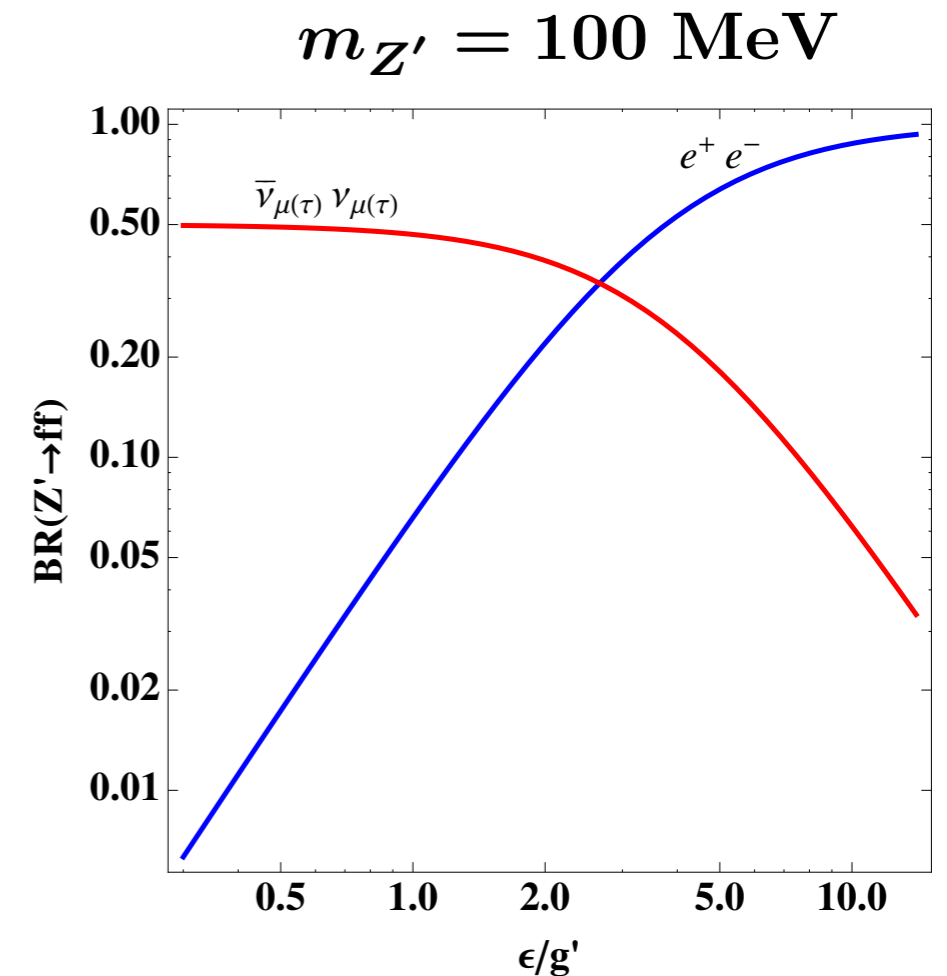
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Therefore the main decay mode of ϕ is

$$\phi \rightarrow Z' Z' \rightarrow 4\nu$$

- **Invisible** at collider experiments
- **different** from the SM Higgs

Allowed Parameter Region

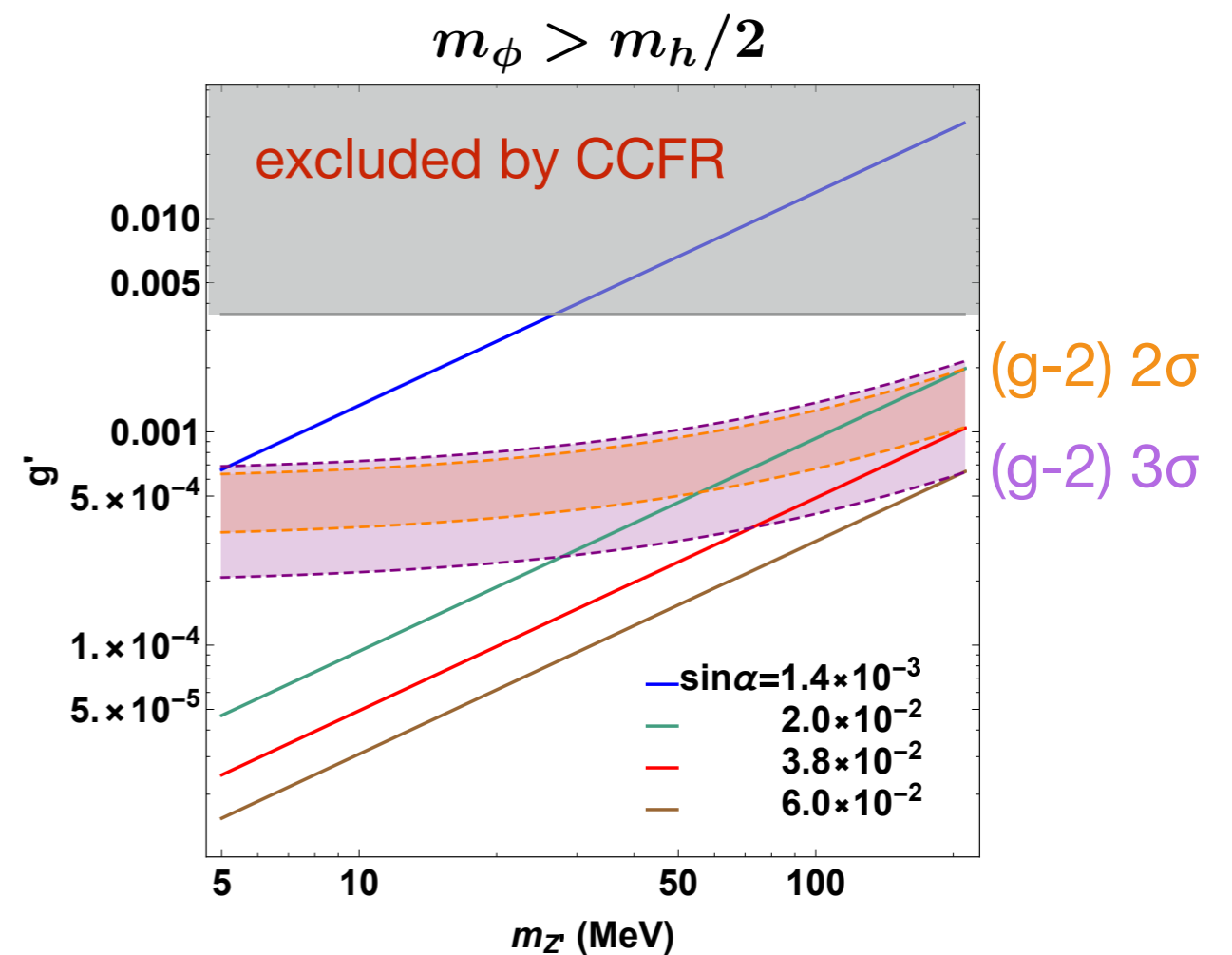
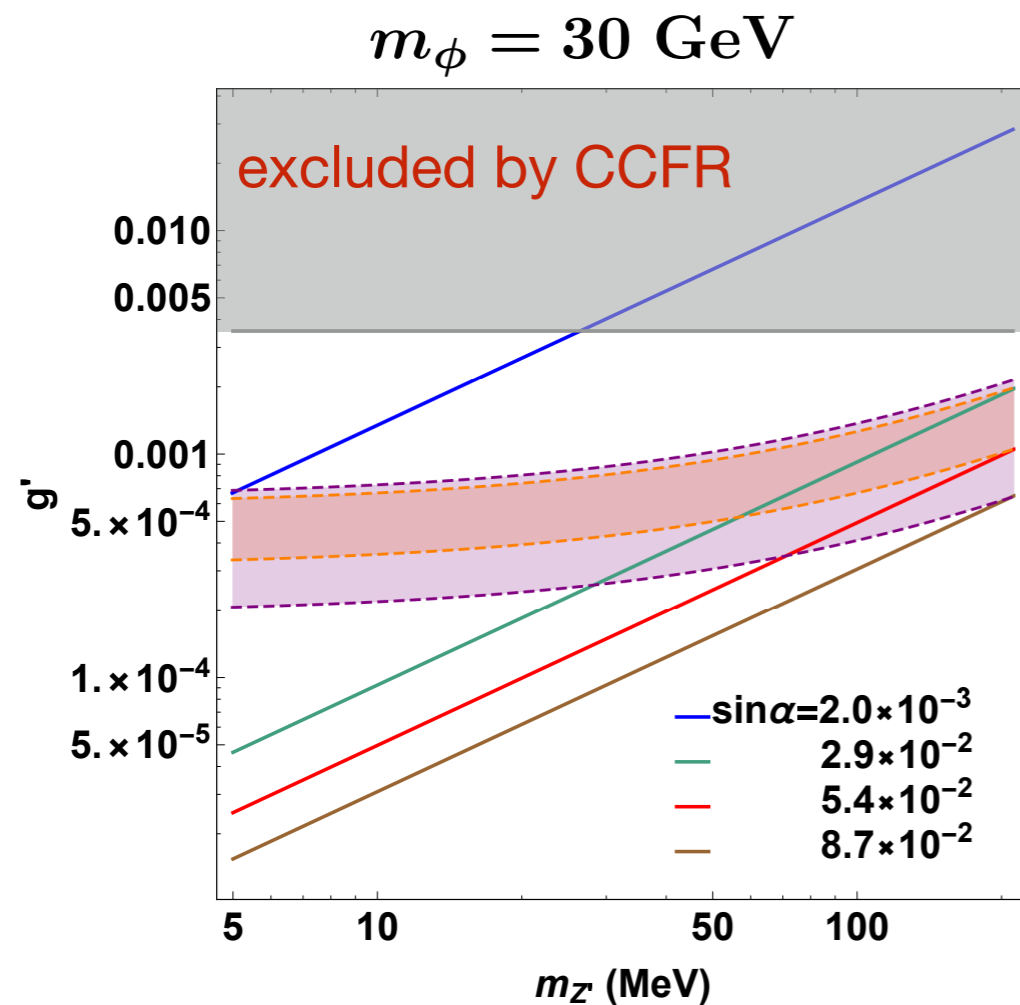
Allowed parameter region

- Extra invisible decays of the SM Higgs into Z' and ϕ pairs,

$$h \rightarrow Z'Z', \phi\phi$$

- Constraints from invisible Higgs decay measured at LHC,

$$\sin \alpha < 0.3, \quad BR_{\text{invisible}} < 0.25$$



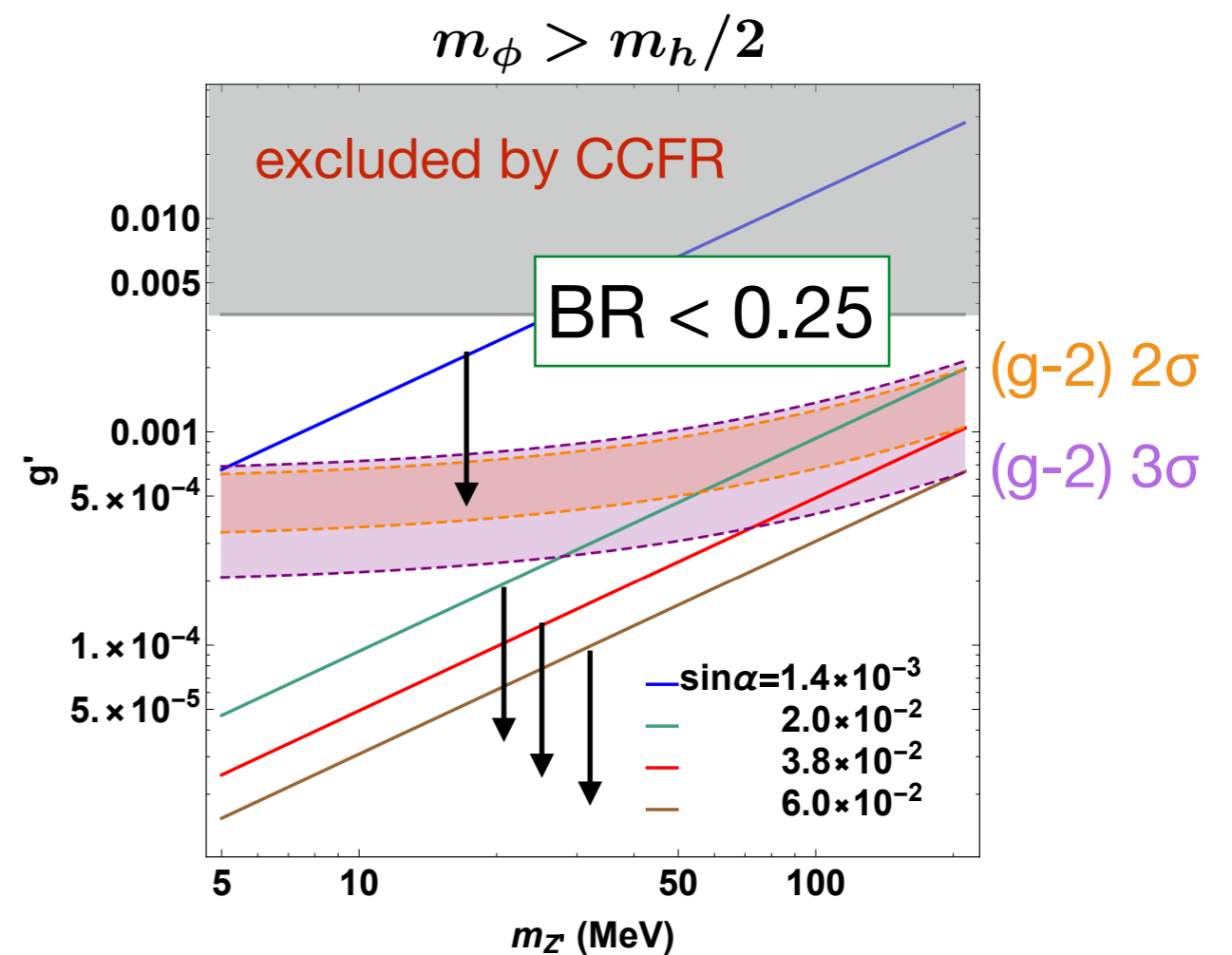
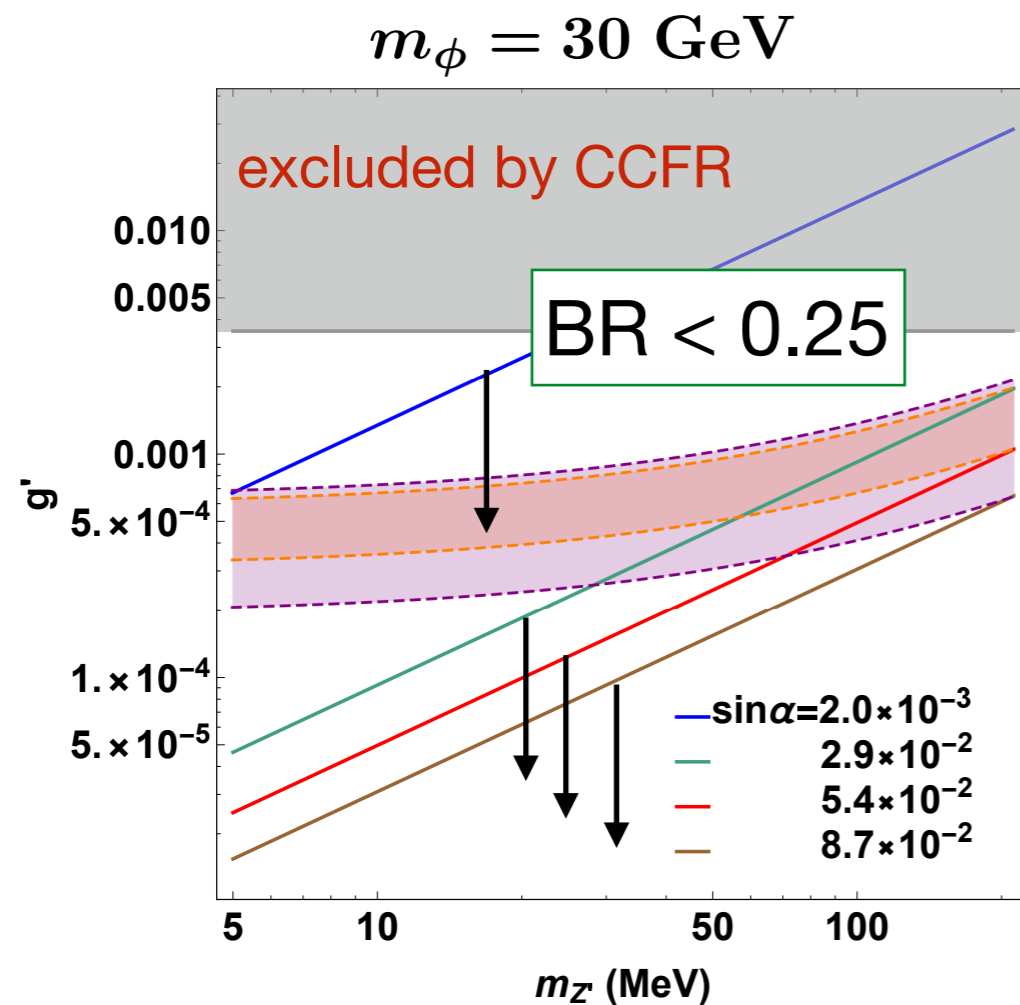
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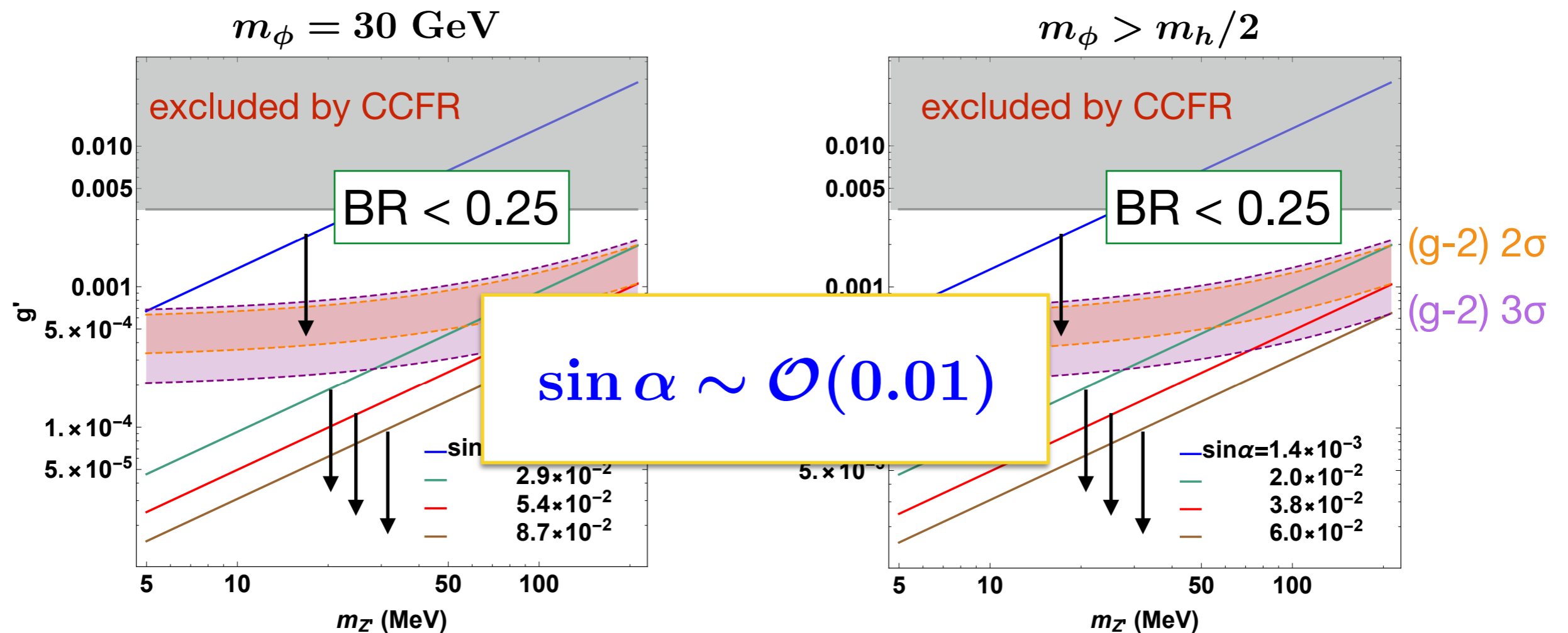
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Allowed parameter region

- Z' search in meson decays at NA64: $Z' \rightarrow e^+e^-$

$$\epsilon/g' \leq 2 \text{ (0.6)} \quad \text{for} \quad m_{Z'} = 100 \text{ (5)} \text{ MeV}$$

- EW precision: ρ parameter

$$\epsilon \leq 7 \times 10^{-4}$$

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In the following discussion, we fix

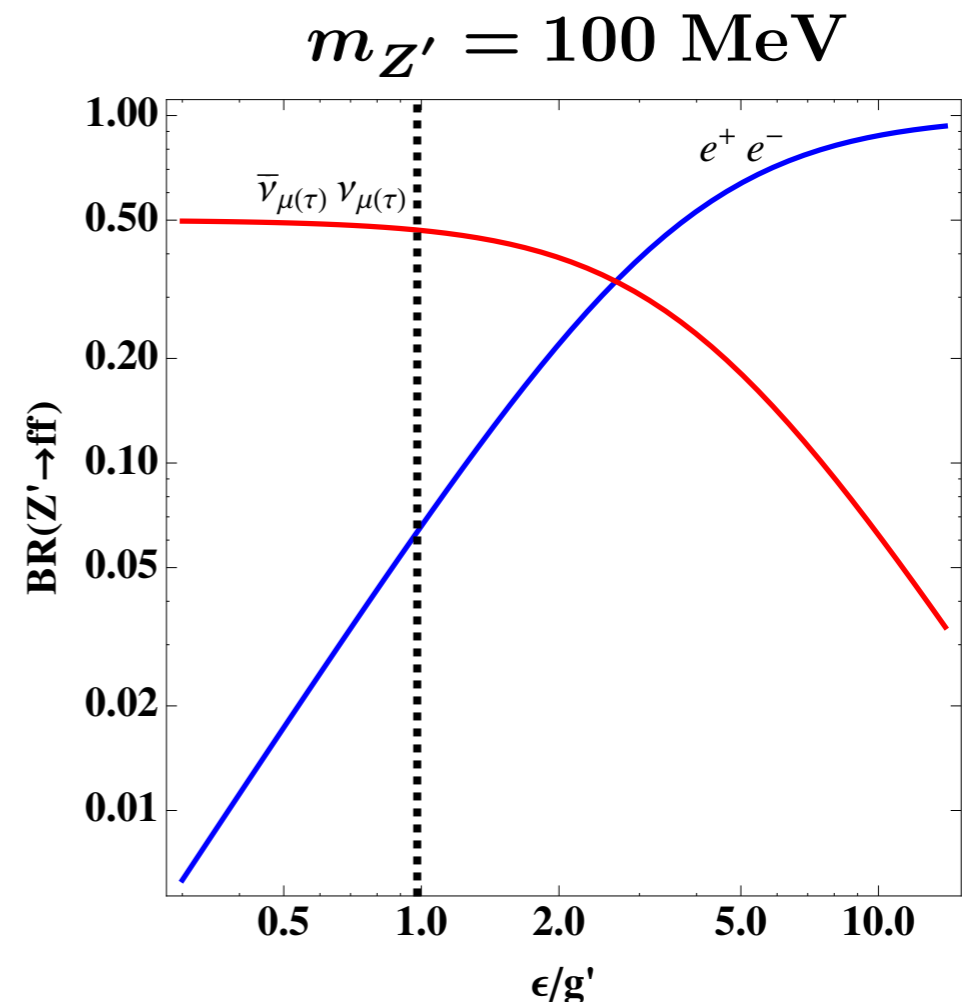
$$m_{Z'} = 100 \text{ MeV}$$

$$\sin \alpha = 0.05$$

$$\epsilon/g' = 1$$

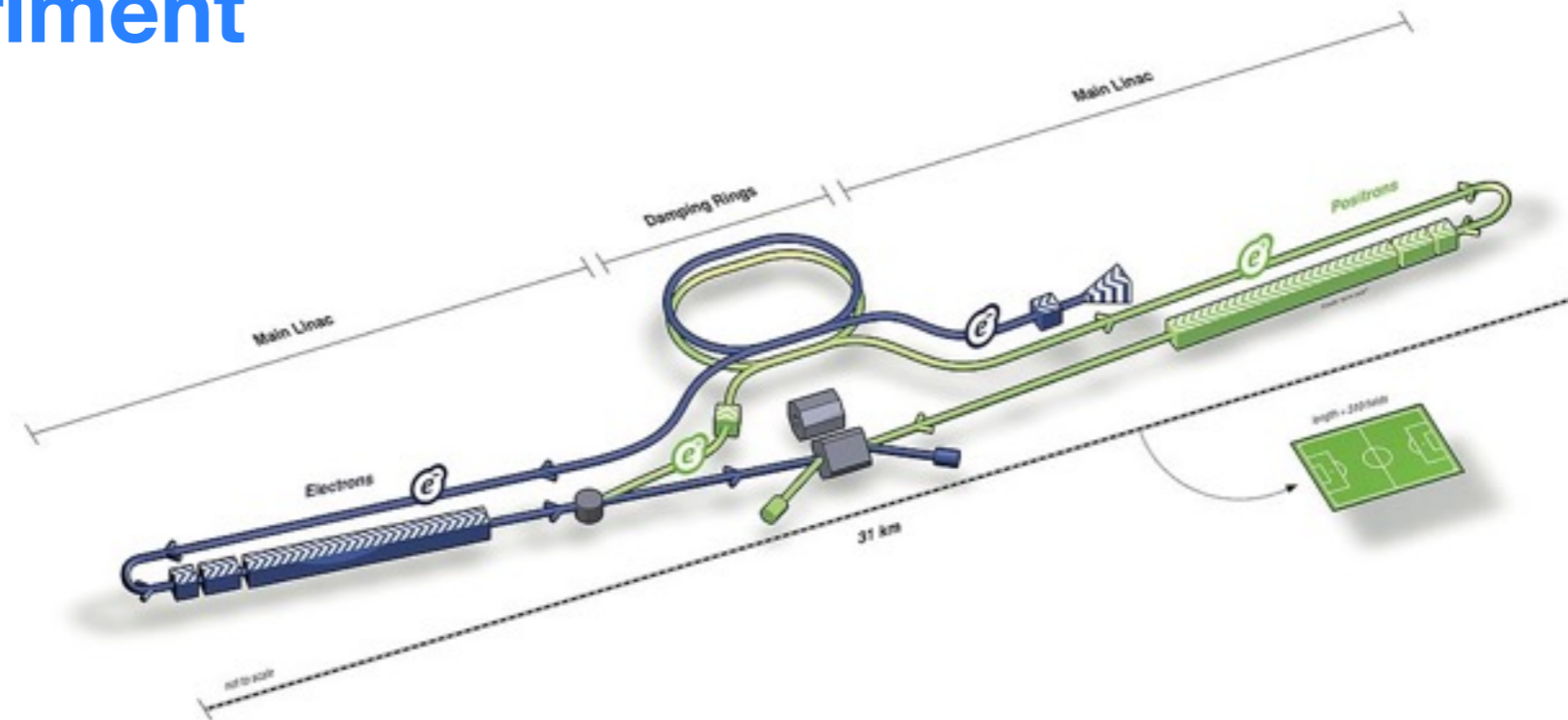


$$BR(Z' \rightarrow \nu\bar{\nu}) = 0.97$$



Search at ILC experiment

ILC experiment



- Electron-Positron collider with $\sqrt{s}=250$ GeV
- Beam polarization (+ −, − +, ++, −−) for (e⁺, e⁻)
- Total integrated luminosity, 2000 fb⁻¹, (45%, 45%, 5%, 5%)
- Clean BG and invariant mass reconstruction available

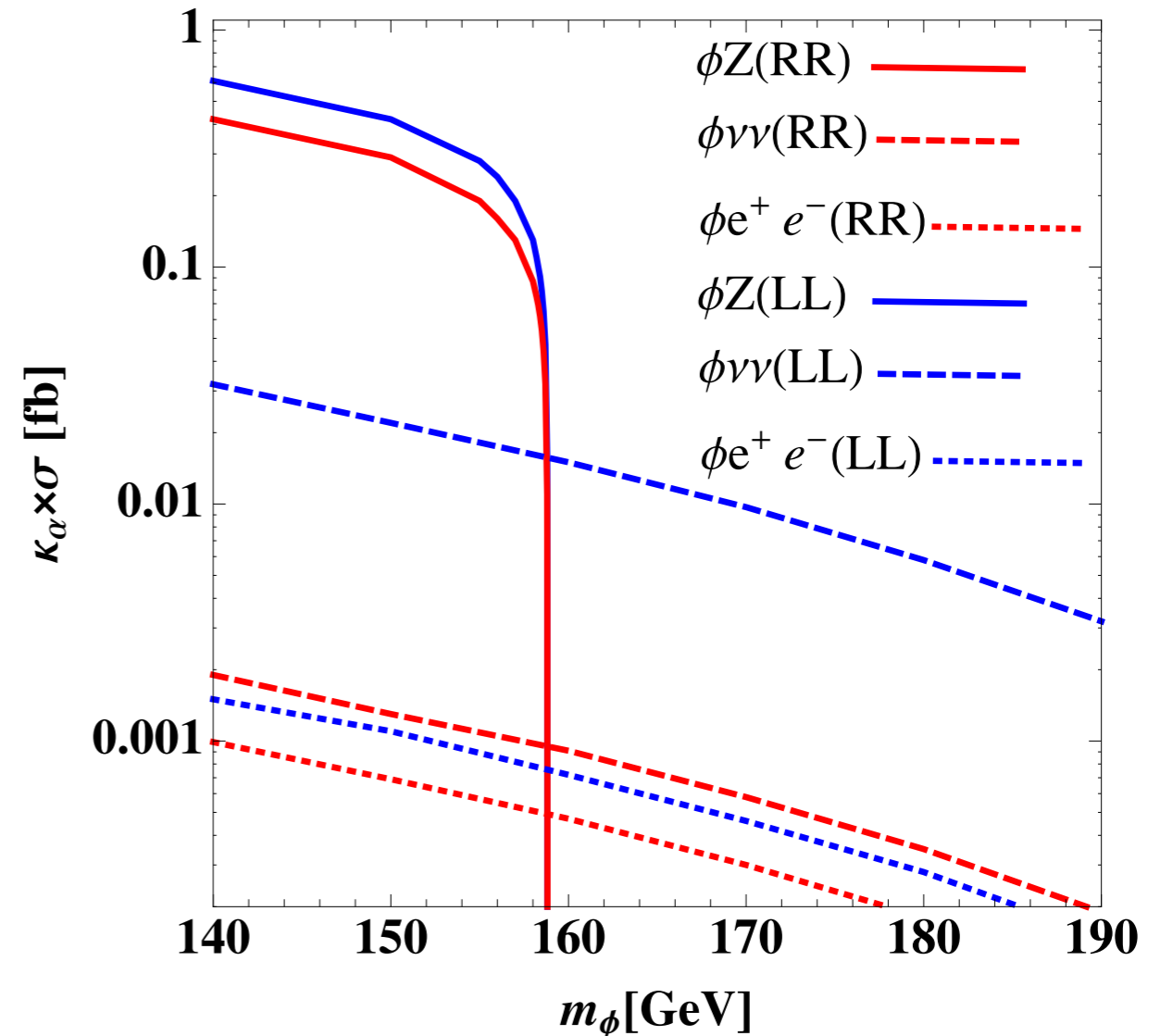
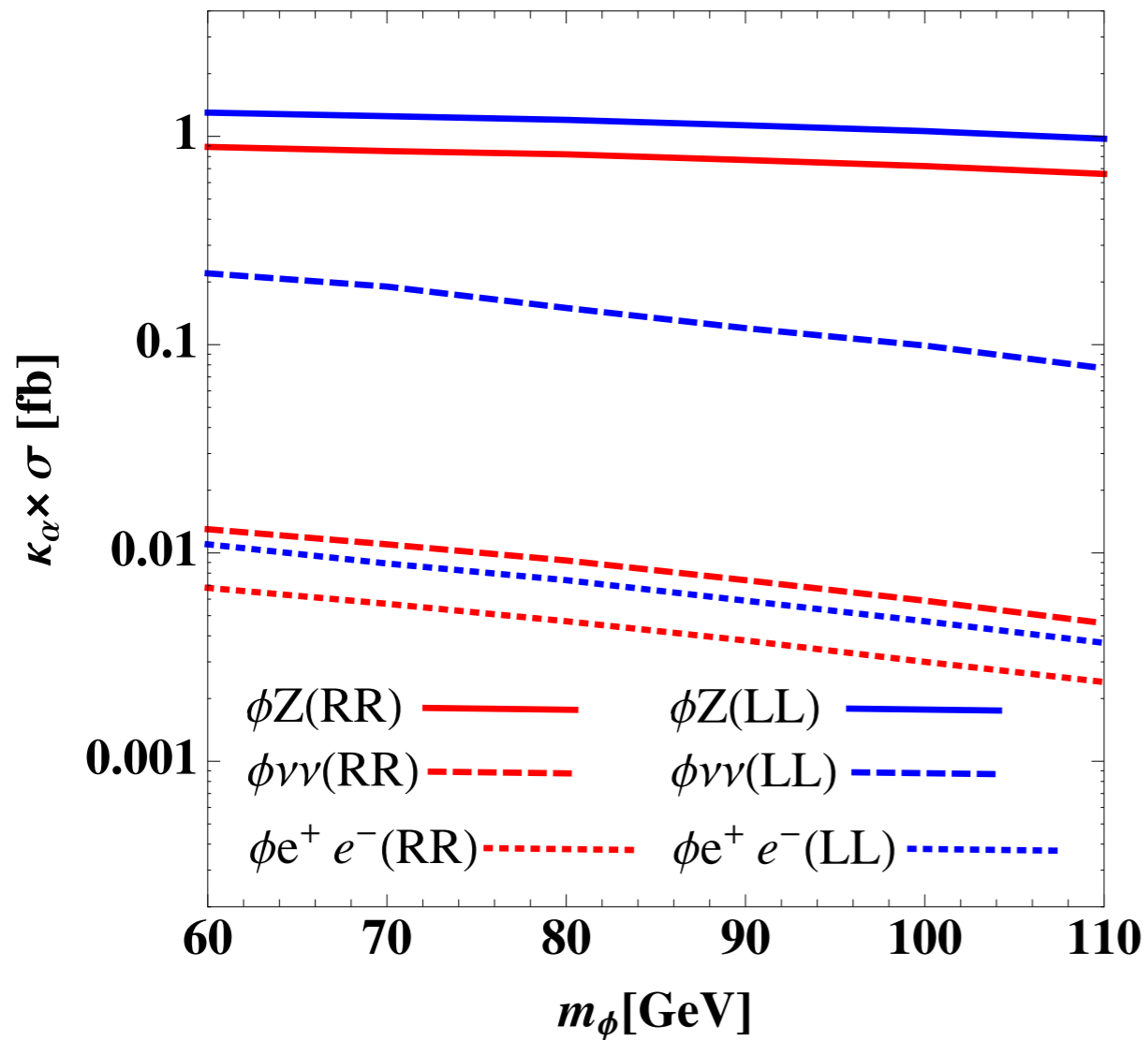
In the following analysis, we apply

$$LL : (e^-_L, e^+_R) = (-80\%, 30\%) \text{ with } L=900\text{fb}^{-1}$$

$$RR : (e^-_L, e^+_R) = (80\%, -30\%) \text{ with } L=900\text{fb}^{-1}$$

ILC experiment

Production cross section of ϕ at ILC experiment

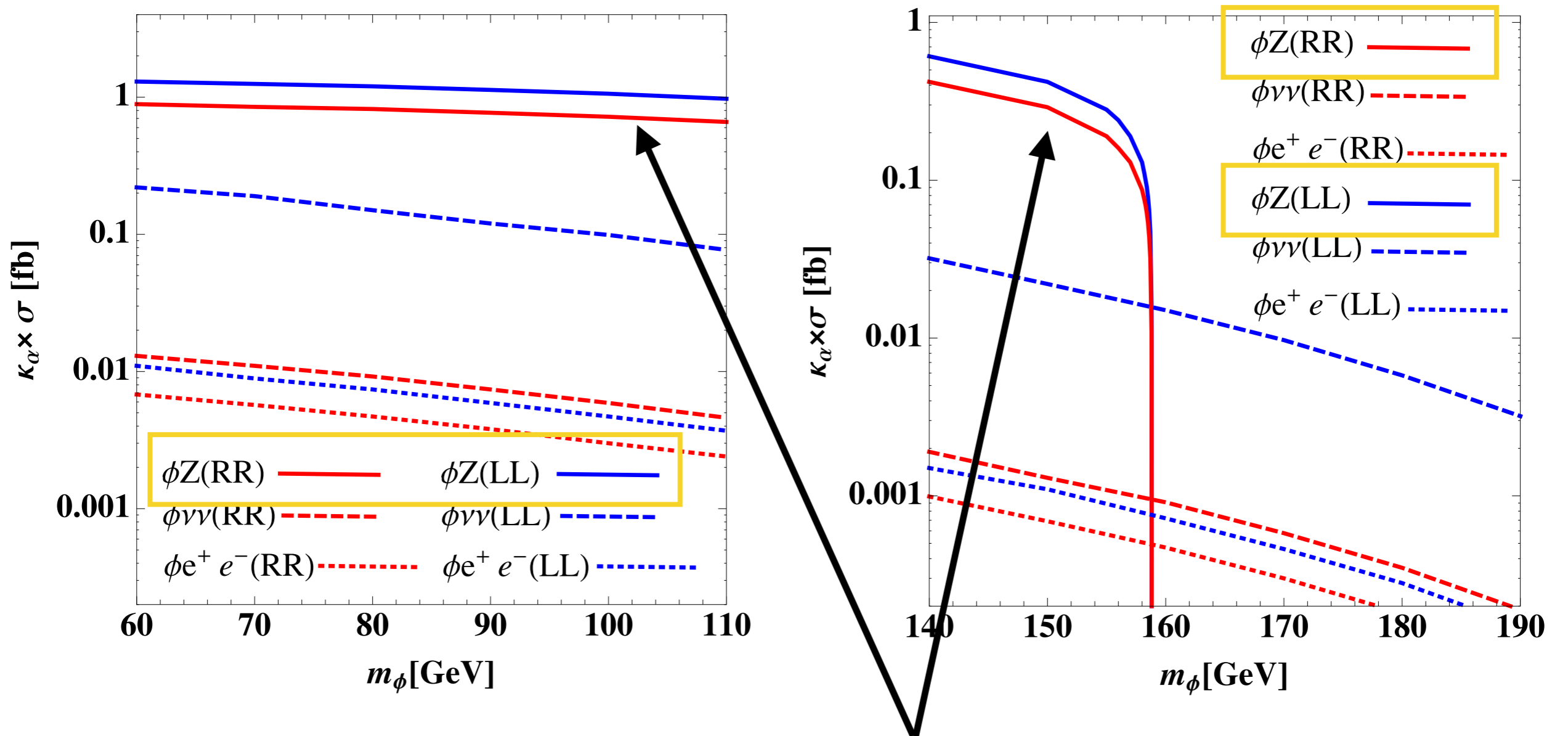


normalized by

$$\kappa_\alpha = \left(\frac{0.05}{\sin \alpha} \right)^2$$

ILC experiment

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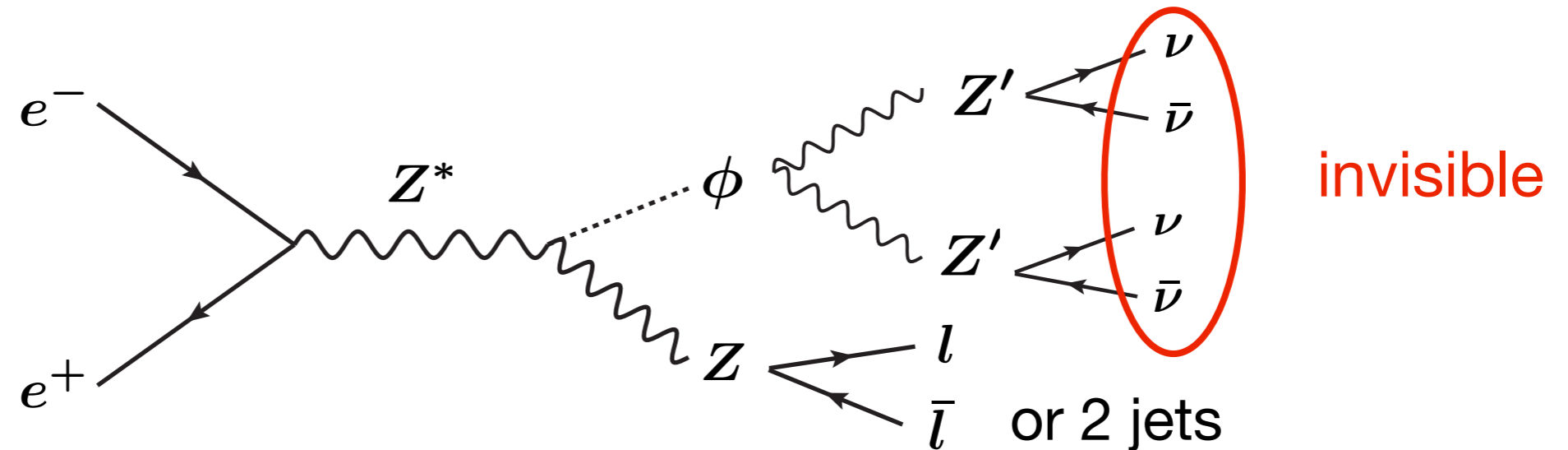
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For $m_\phi < 160$ GeV,
Higgs-Strahlung is dominant

ILC experiment

Main production process of ϕ at ILC with $\sqrt{s} = 250$ GeV



The signal

- 1) $e^+e^- \rightarrow l^+l^- + \cancel{E}$
- 2) $e^+e^- \rightarrow jj + \cancel{E}$

The backgrounds

- 1) $e^+e^- \rightarrow l^+l^-\nu\bar{\nu}, \tau^+\tau^-$ followed by leptonic decay
- 2) $e^+e^- \rightarrow jj\nu\bar{\nu}, \tau^+\tau^-$ followed by hadronic decay

ILC experiment

We perform a simulation study by using MADGRAPH/MADEVENT with PYTHIA6, assuming $L=900/\text{fb}$ and $m_\phi=30, 65 \text{ GeV}$.

Basic Cuts

Reduce BGs



- 1) $p_T(\ell^\pm) > 7 \text{ GeV}$, $\eta(\ell^\pm) < 2.5$ for leptons
- 2) $p_T(j) > 20 \text{ GeV}$, $\eta(j) < 5.0$ for jets

Invariant Mass Cuts

Reduce BGs
in M^{rec} dist.



- 1) $m_Z - 10 \text{ GeV} < M_{\ell^+\ell^-} < m_Z + 10 \text{ GeV}$ for leptons
- 2) $m_Z - 20 \text{ GeV} < M_{jj} < m_Z + 5 \text{ GeV}$ for jets

reconstructed mass

$$M_\phi^{\text{rec}} = \sqrt{s + m_Z^2 - 2(E_{l/j_1} + E_{l/j_2})\sqrt{s}}$$

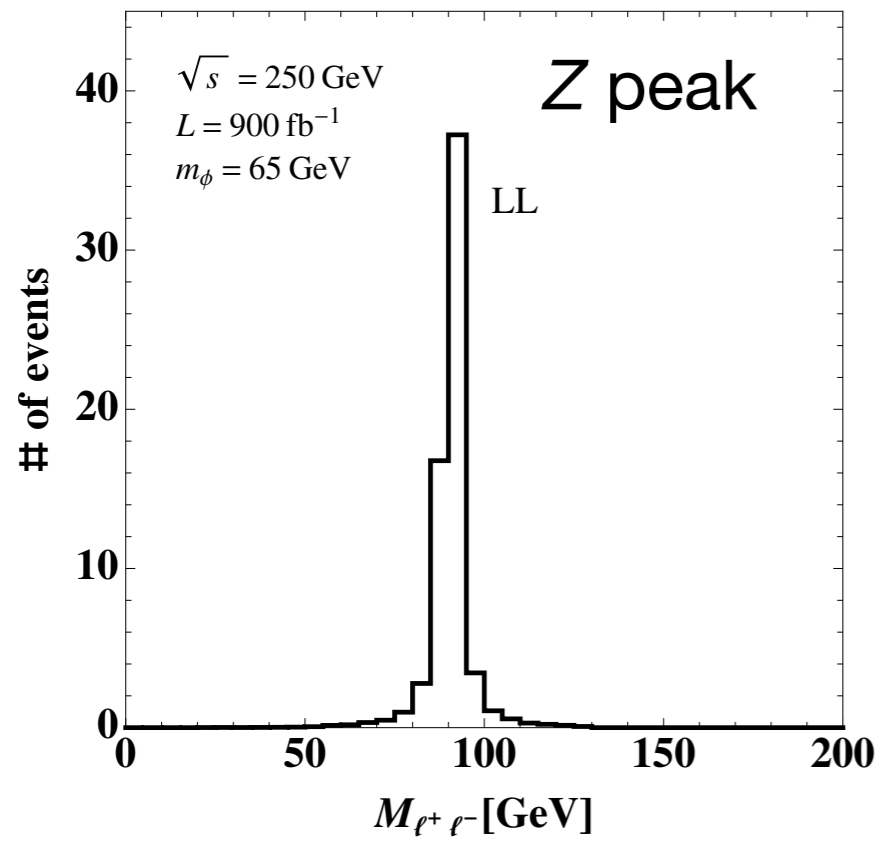
Kinematical Cuts

- 1) $m_\phi - 10 \text{ GeV} < M_{\phi_\ell}^{\text{rec}} < m_\phi + 10 \text{ GeV}$ for leptons
- 2) $m_\phi - 15(10) \text{ GeV} < M_{\phi_j}^{\text{rec}} < m_\phi + 25(50) \text{ GeV}$ for jets

Lepton

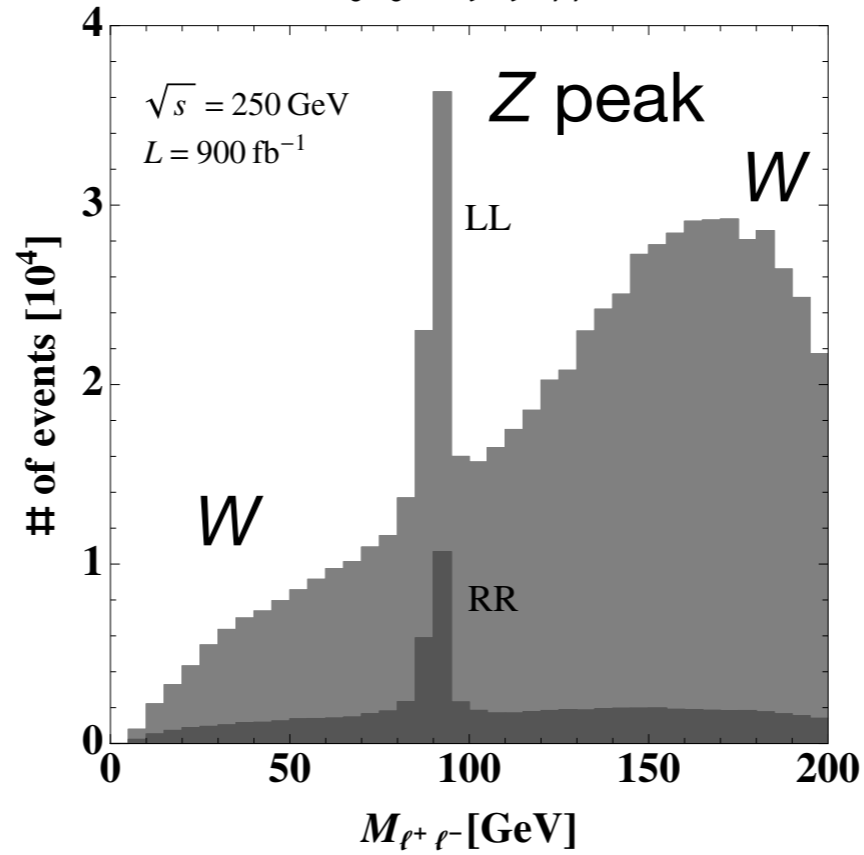
signal

$$e^+e^- \rightarrow Z\phi \rightarrow \ell^+\ell^- + 4\nu$$



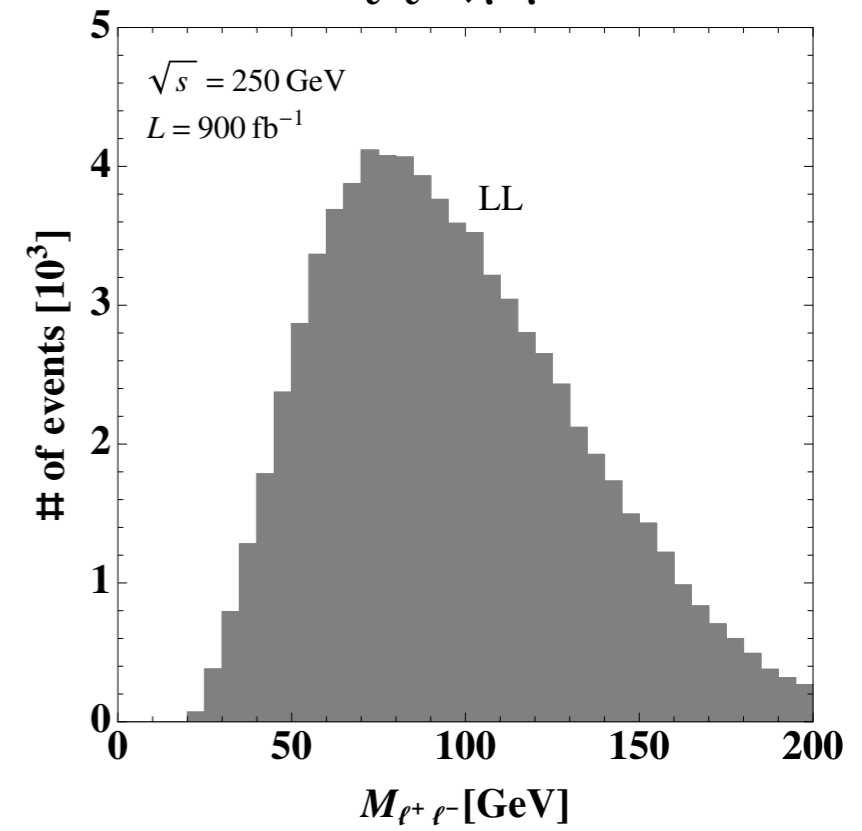
lepton BG

$$e^+e^- \rightarrow \ell^+\ell^- \nu\bar{\nu}$$



tau BG

$$e^+e^- \rightarrow \tau^+\tau^-$$

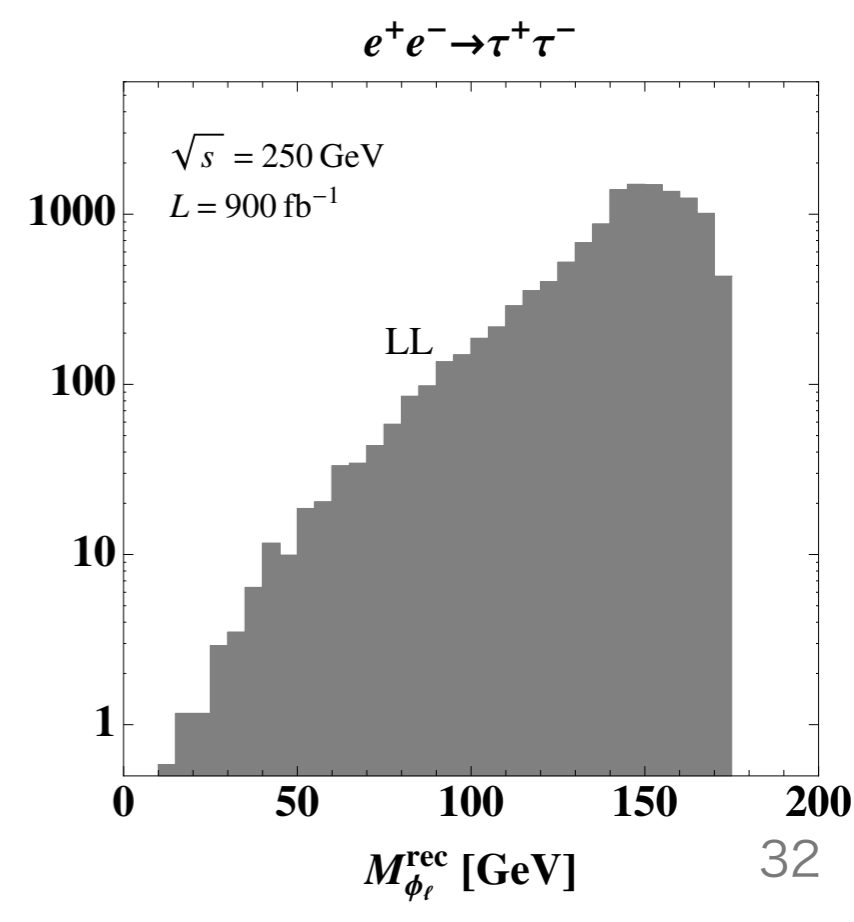
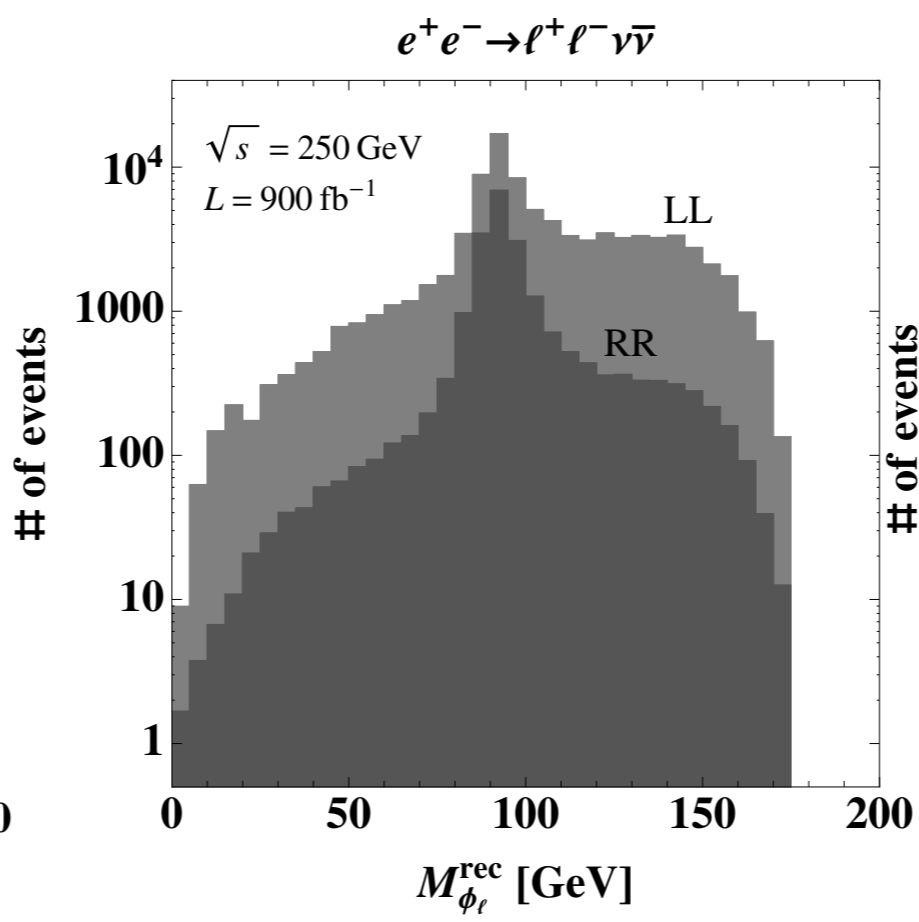
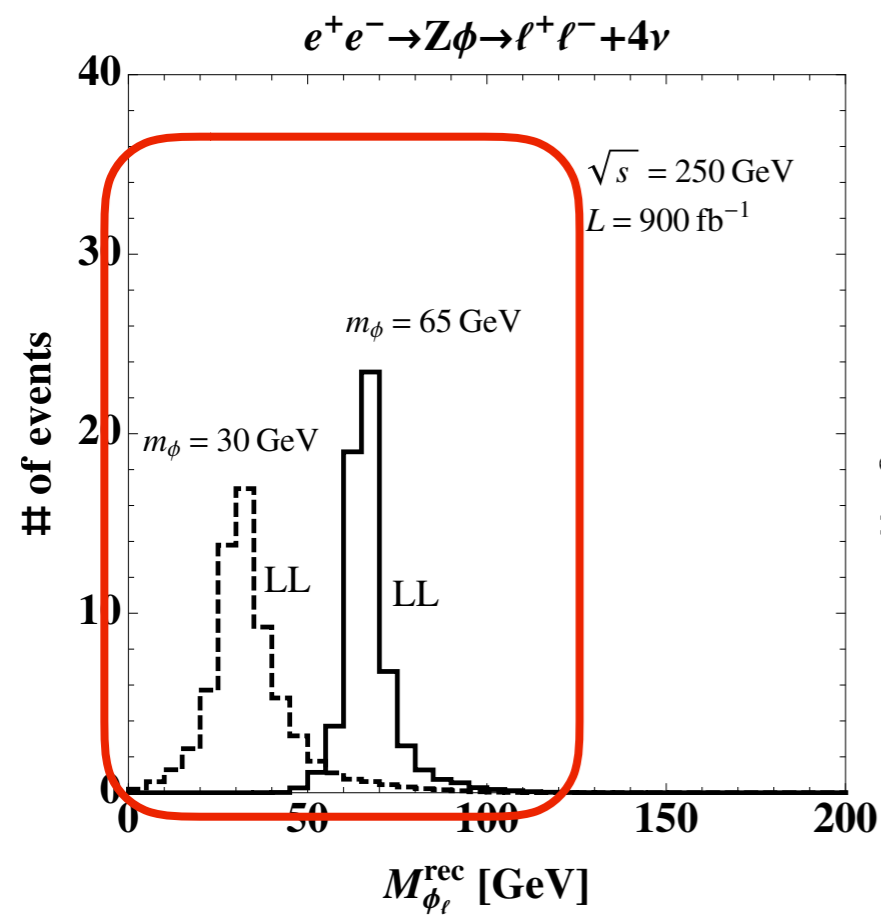
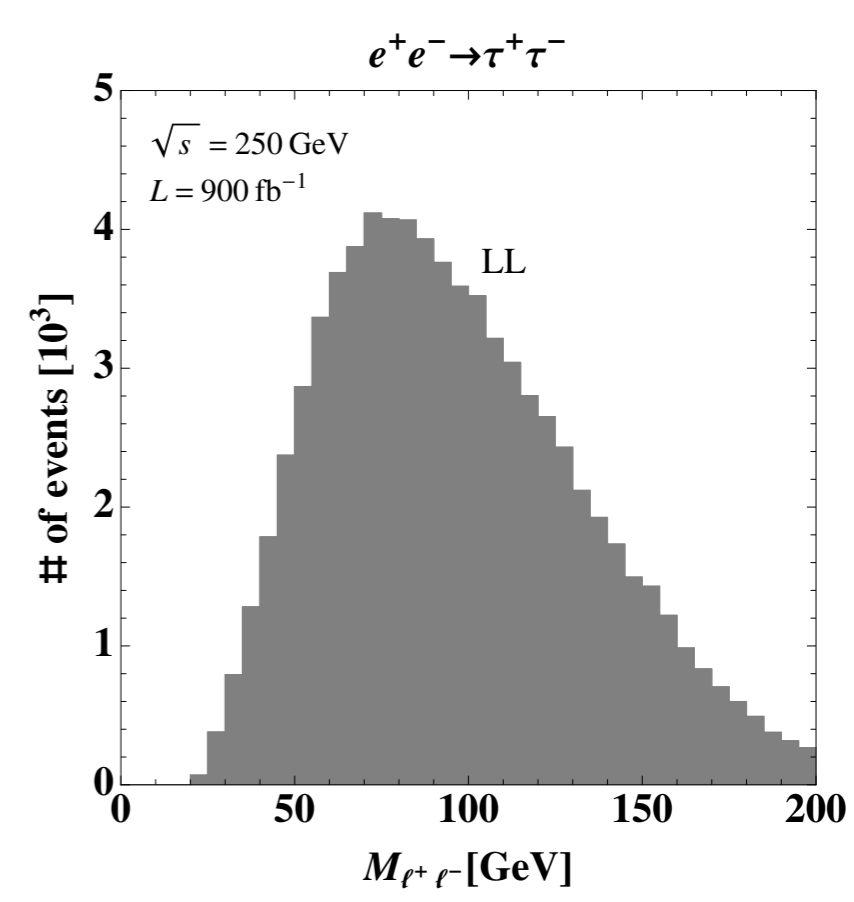
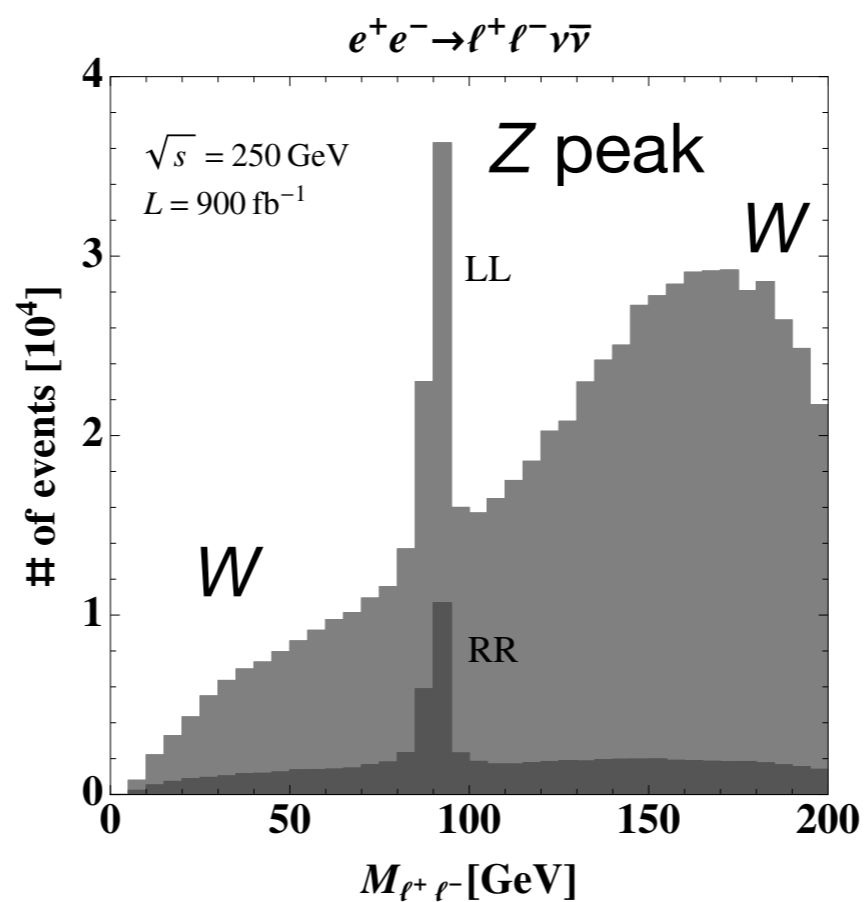
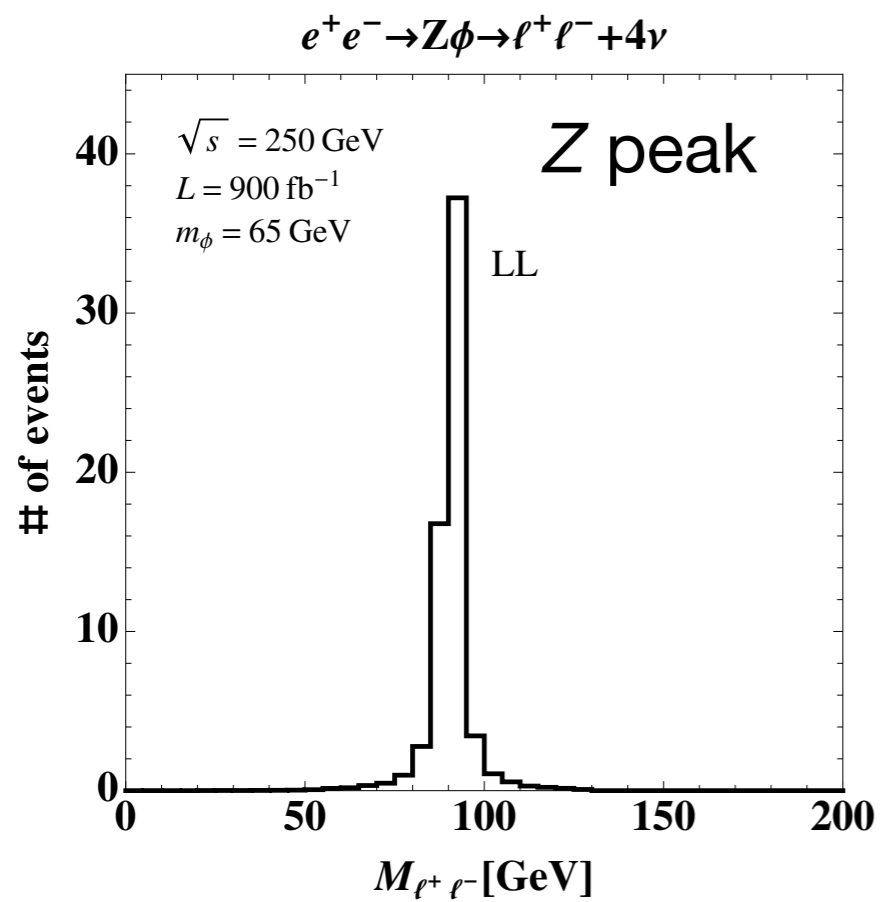


Lepton

signal

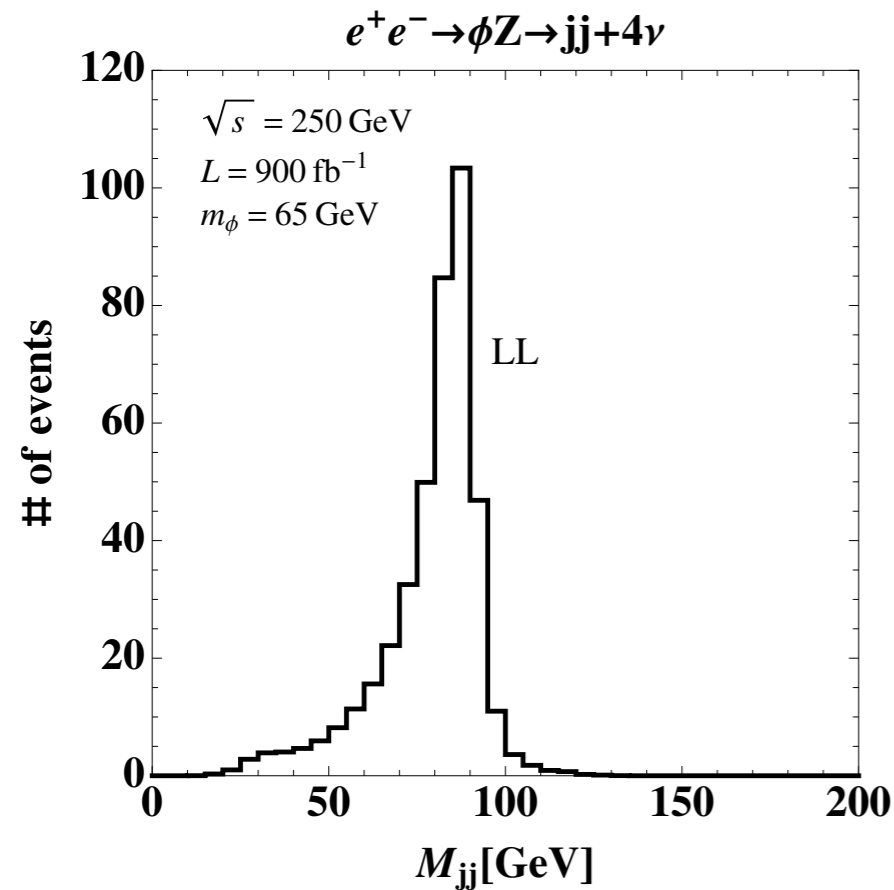
lepton BG

tau BG

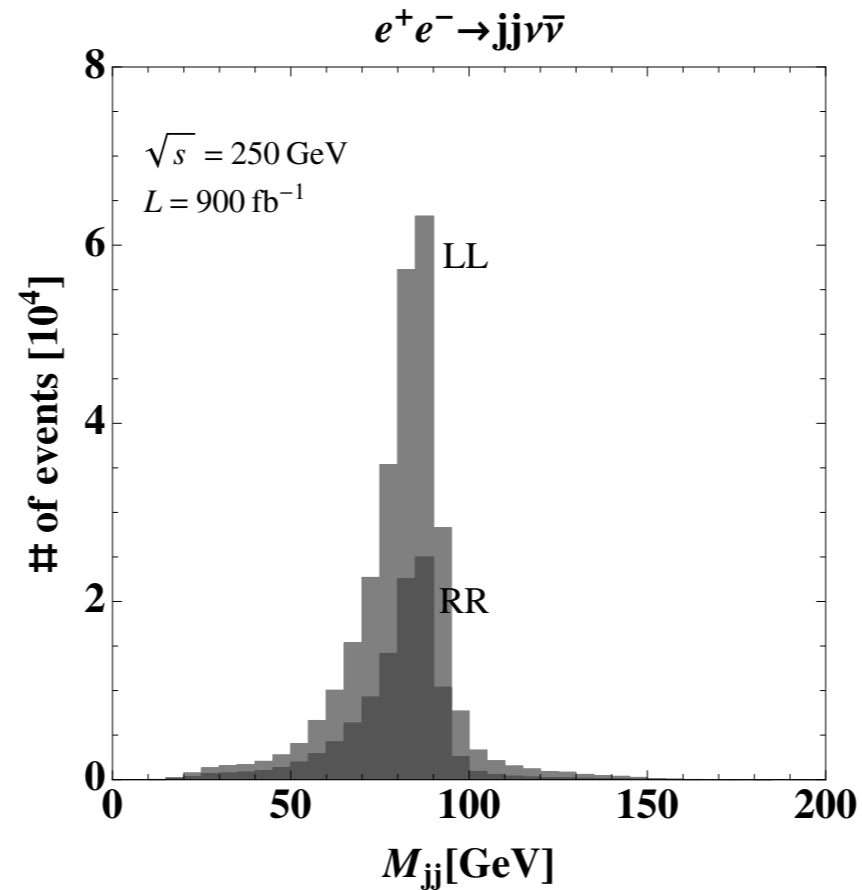


Jets

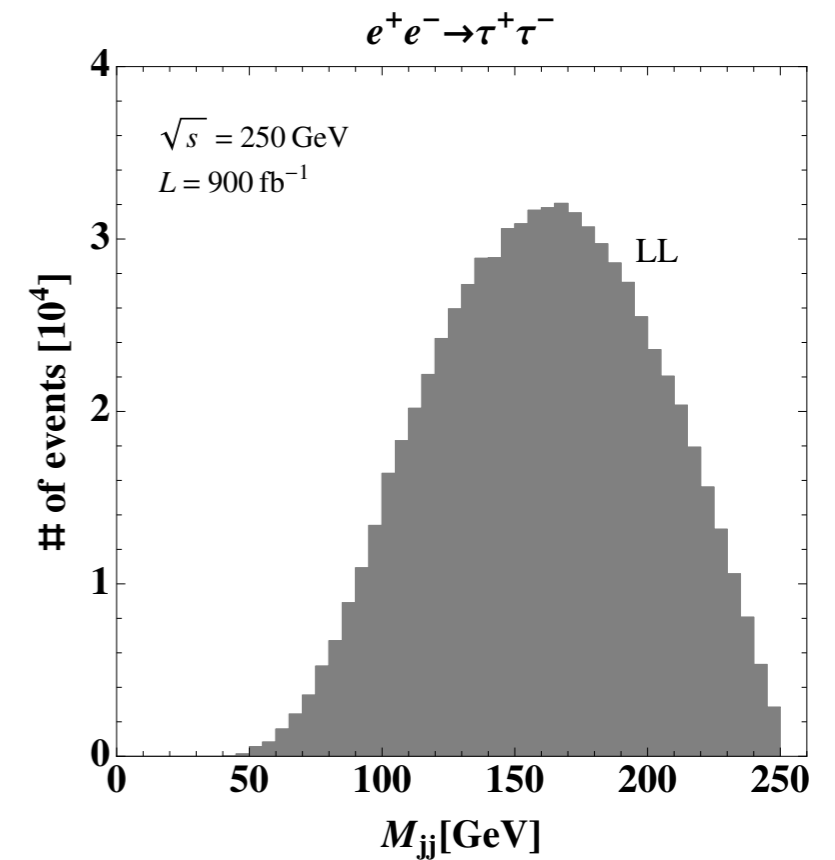
signal



lepton BG



tau BG



Slightly broader Z peak
due to jet energy resolution

Continuous BG suppressed at
low mass region

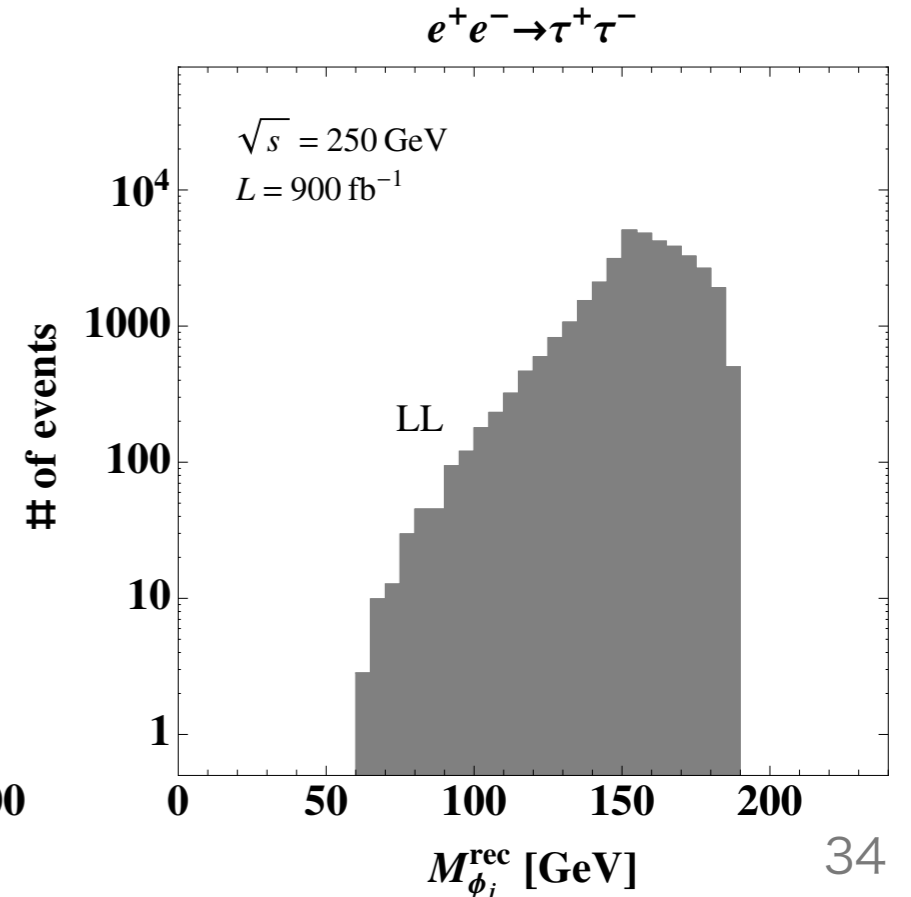
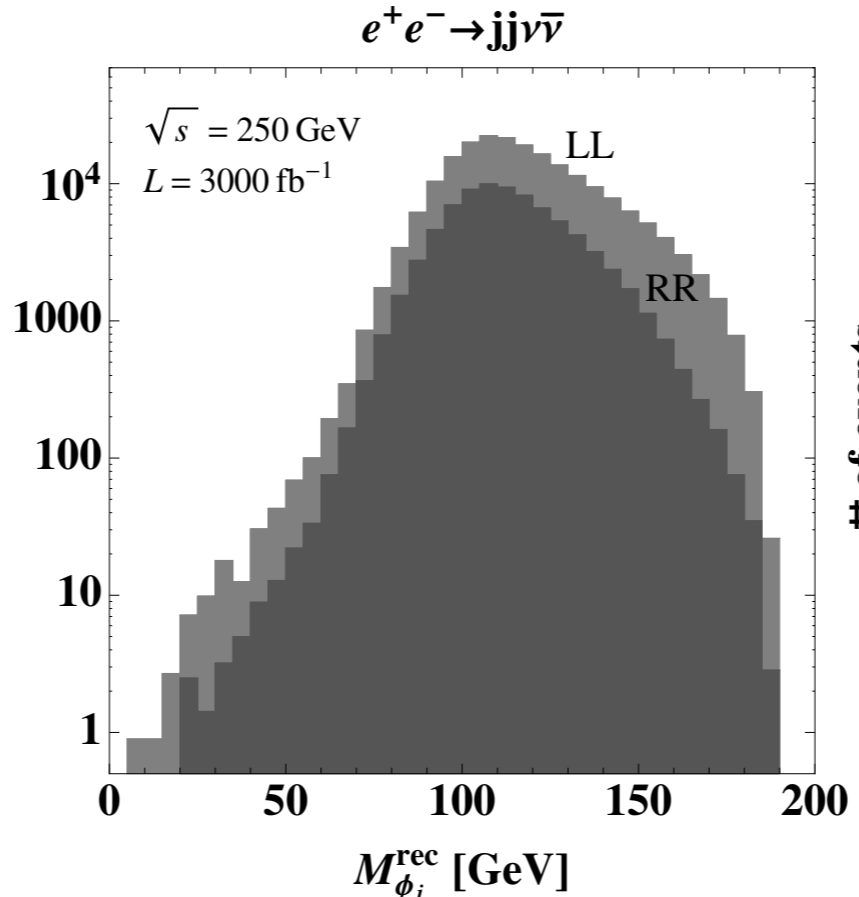
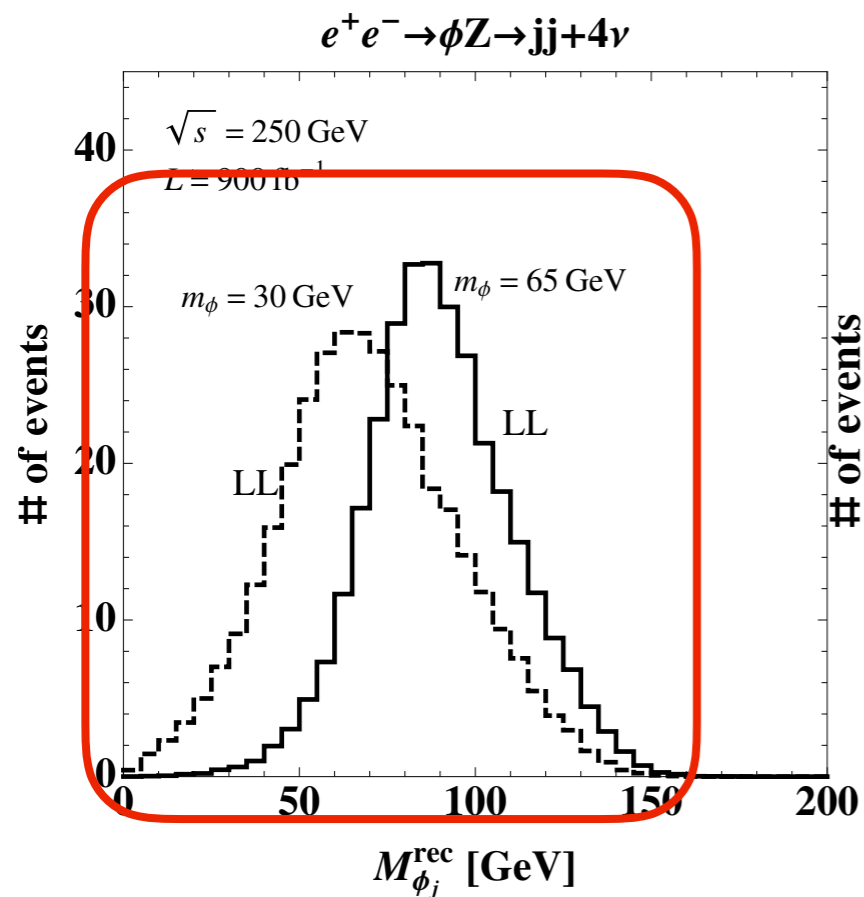
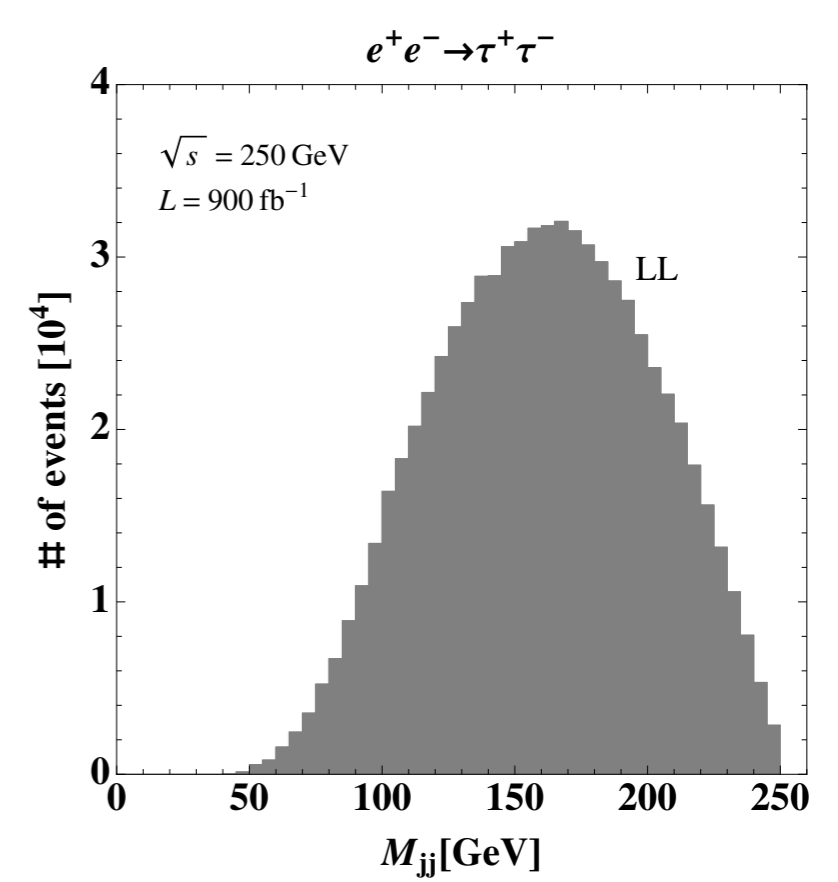
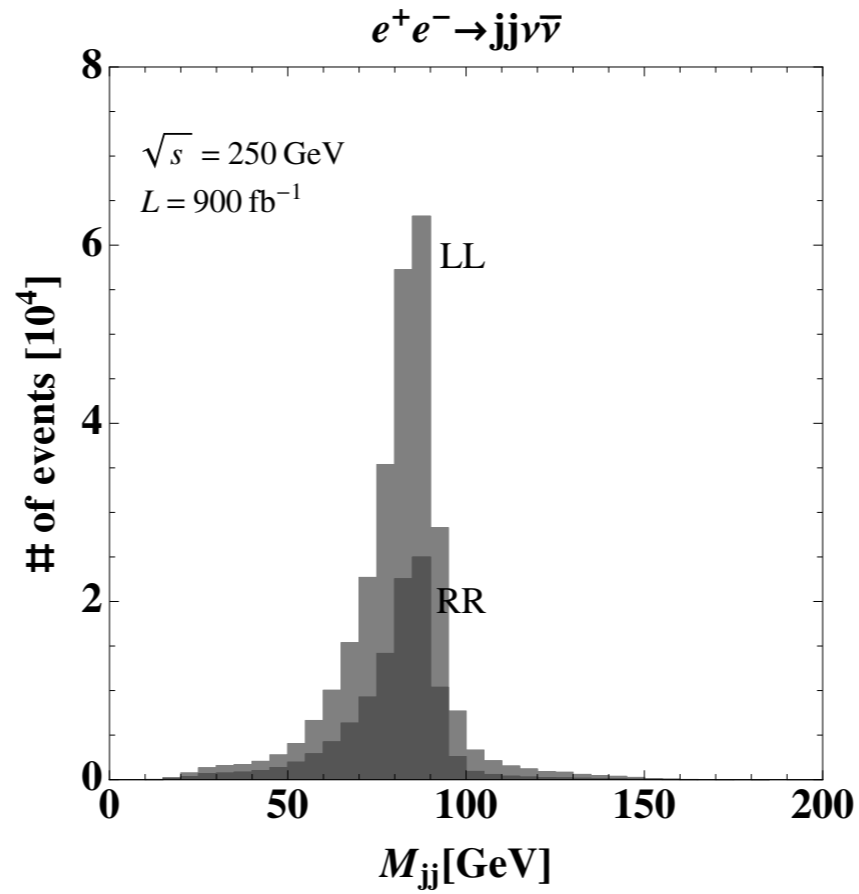
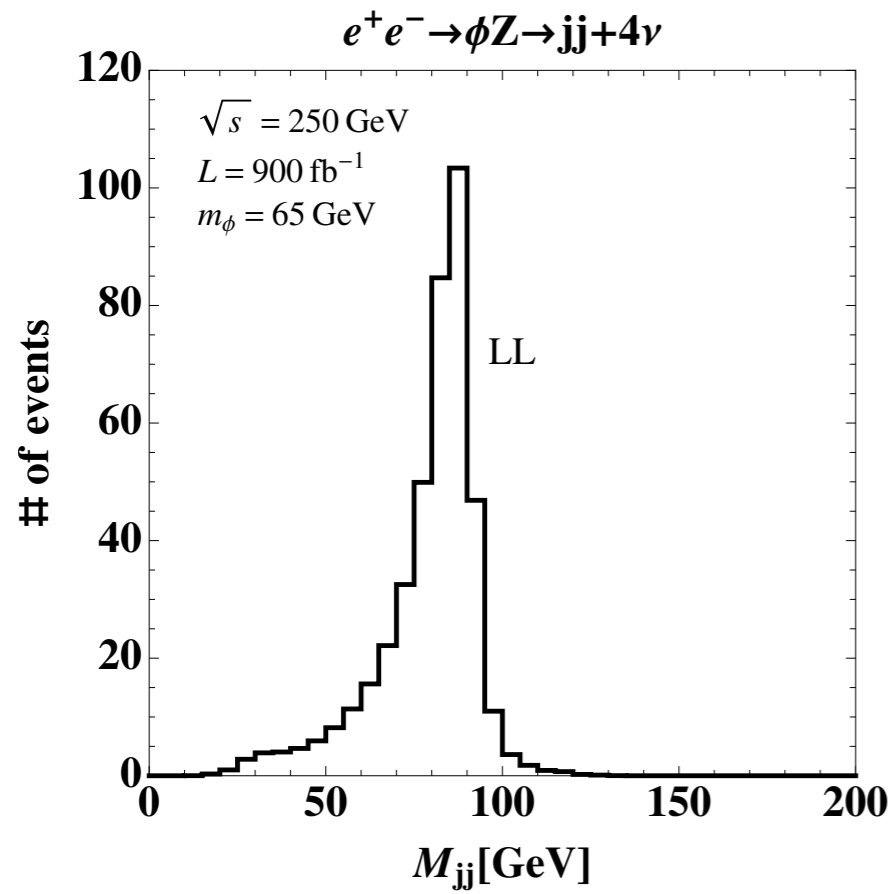
1. Lepton BG comes from ZZ followed by Z decays into jets/neutrinos.
2. Tau BG shifts to higher mass region due to miss-identification of τ -jet with missing energy.

Jets

signal

lepton BG

tau BG



lepton pair event

$$\kappa_\alpha = (0.05 / \sin \alpha)^2$$

	$\kappa_\alpha N_S^{\kappa_\alpha=1}; m_\phi = 65(30) \text{ GeV}$	$N_{BG}^{\ell^+ \ell^- \nu \bar{\nu}}$	$N_{BG}^{\tau\tau\tau}$	$\kappa_\alpha S_{cl}^{\kappa_\alpha=1}$
<i>RR</i>	42.(34.)	$2.2(1.7) \times 10^2$	$1.3(0.14) \times 10^2$	2.2(2.5)
<i>LL</i>	53.(47.)	$4.7(1.7) \times 10^3$	$1.6(0.15) \times 10^2$	0.75(1.1)
<i>LL + RR</i>	95.(81.)	$4.9(1.9) \times 10^3$	$2.9(0.29) \times 10^2$	1.3(1.8)

jet event

	$\kappa_\alpha N_S^{\kappa_\alpha=1}; m_\phi = 65(30) \text{ GeV}$	$N_{BG}^{jj\nu\bar{\nu}}$	$N_{BG}^{\tau\tau\tau}$	$\kappa_\alpha S_{cl}^{\kappa_\alpha=1}$
<i>RR</i>	$1.3(1.5) \times 10^2$	$5.6(0.33) \times 10^3$	$1.3(0.064) \times 10^2$	1.6(8.3)
<i>LL</i>	$1.6(1.9) \times 10^2$	$1.3(0.085) \times 10^4$	$2.0(0.13) \times 10^2$	1.4(6.5)
<i>LL + RR</i>	$2.9(3.4) \times 10^2$	$1.9(0.12) \times 10^4$	$3.3(0.19) \times 10^2$	2.1(9.7)

Summary

Summary

We have studied the scalar boson search which breaks gauged $U(1)_{L\mu-L\tau}$ symmetry spontaneously at ILC250.

From (g-2) and the invisible Higgs decay,

- The gauge boson mass and gauge coupling are $O(100)$ MeV and 10^{-4} , respectively.
- The scalar mass should be $<O(100)$ GeV, and its mixing should be $O(0.01)$
- The scalar boson is invisible because its main decay mode is $\phi \rightarrow Z'Z' \rightarrow 4\nu$.

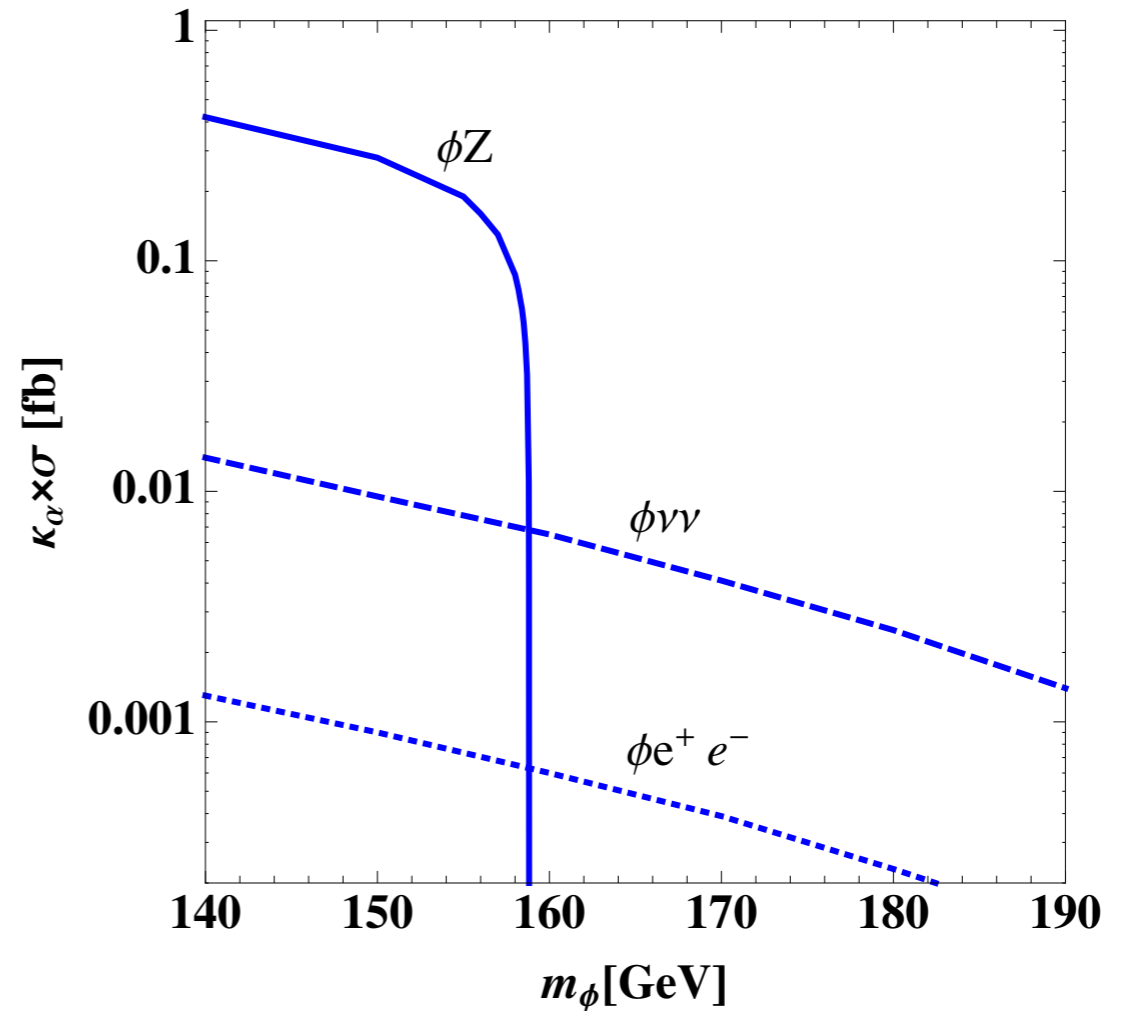
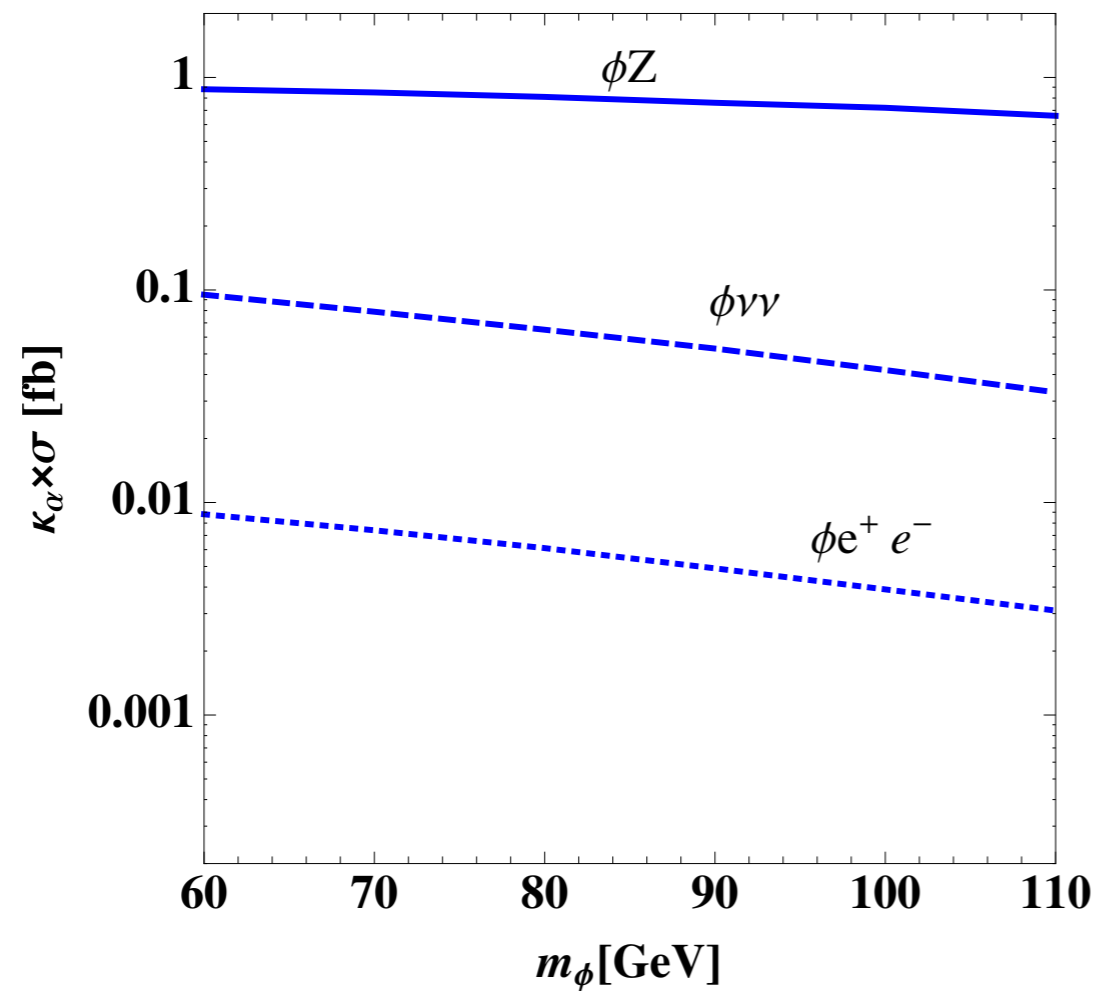
Then, at the ILC experiment,

- ϕ mass can be reconstructed from lepton pair and jet events.
- Signal significance is obtained as 2.2 (2.5) for lepton pair in RR , and 2.1 (9.7) for jets in $LL+RR$ with $m_\phi=65$ (30) GeV after cuts.

Back-up

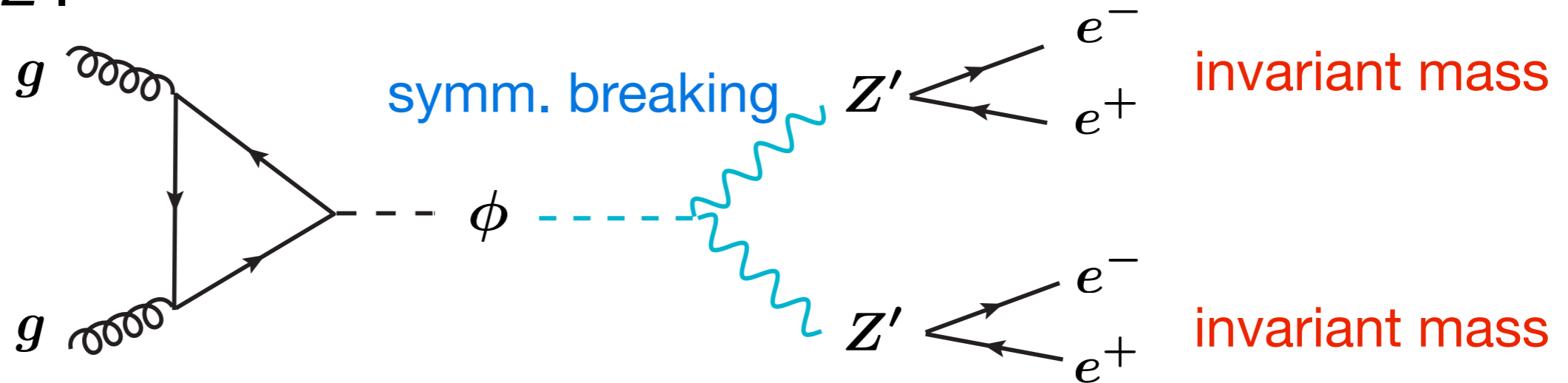
ILC experiment

Production cross section at ILC250



LHC experiment

The main production mode of ϕ is gluon fusion, followed by decay into $Z'Z'$.



In principle, Z' and ϕ mass can be reconstructed by measuring e^+e^- energy and momentum.

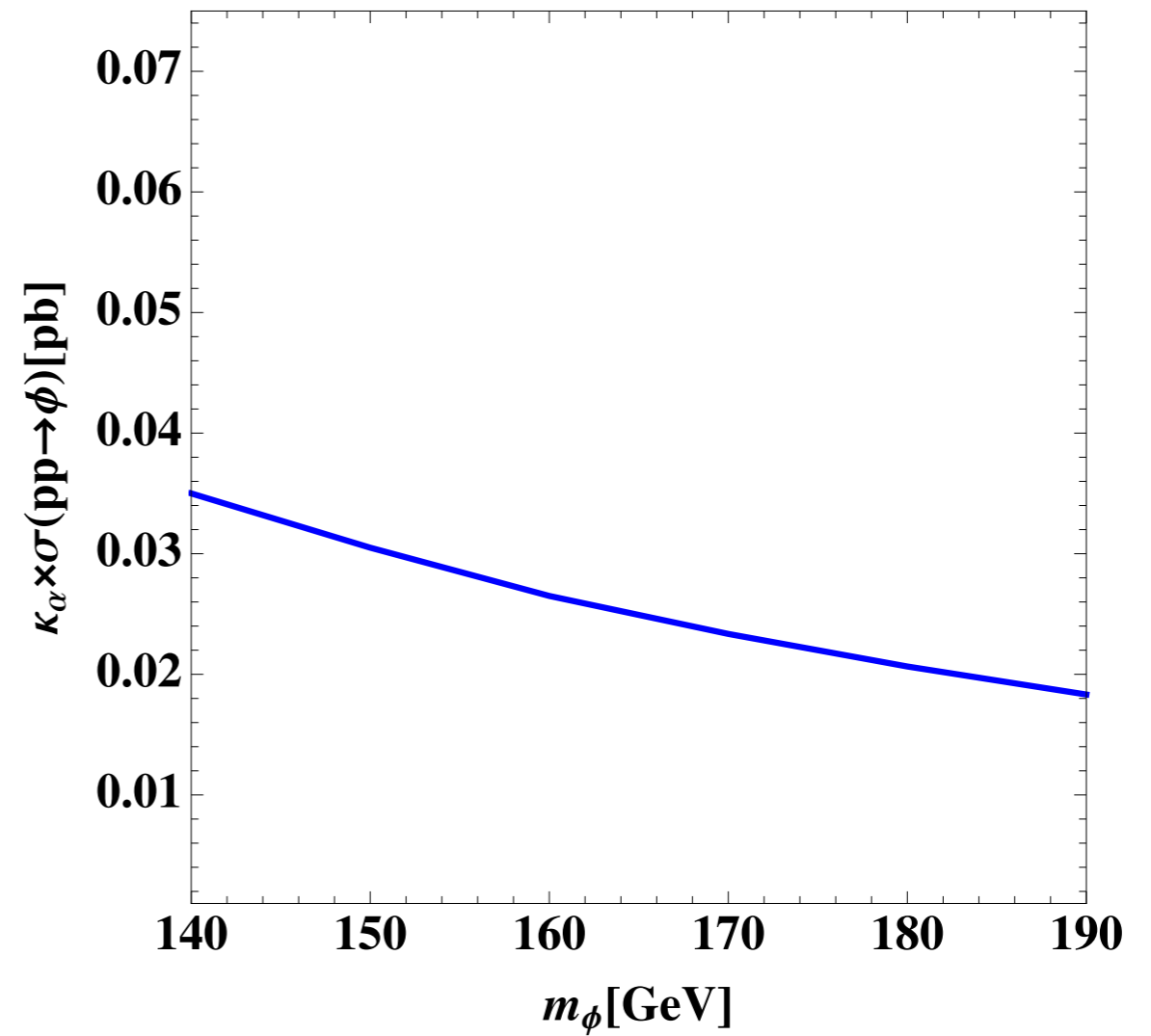
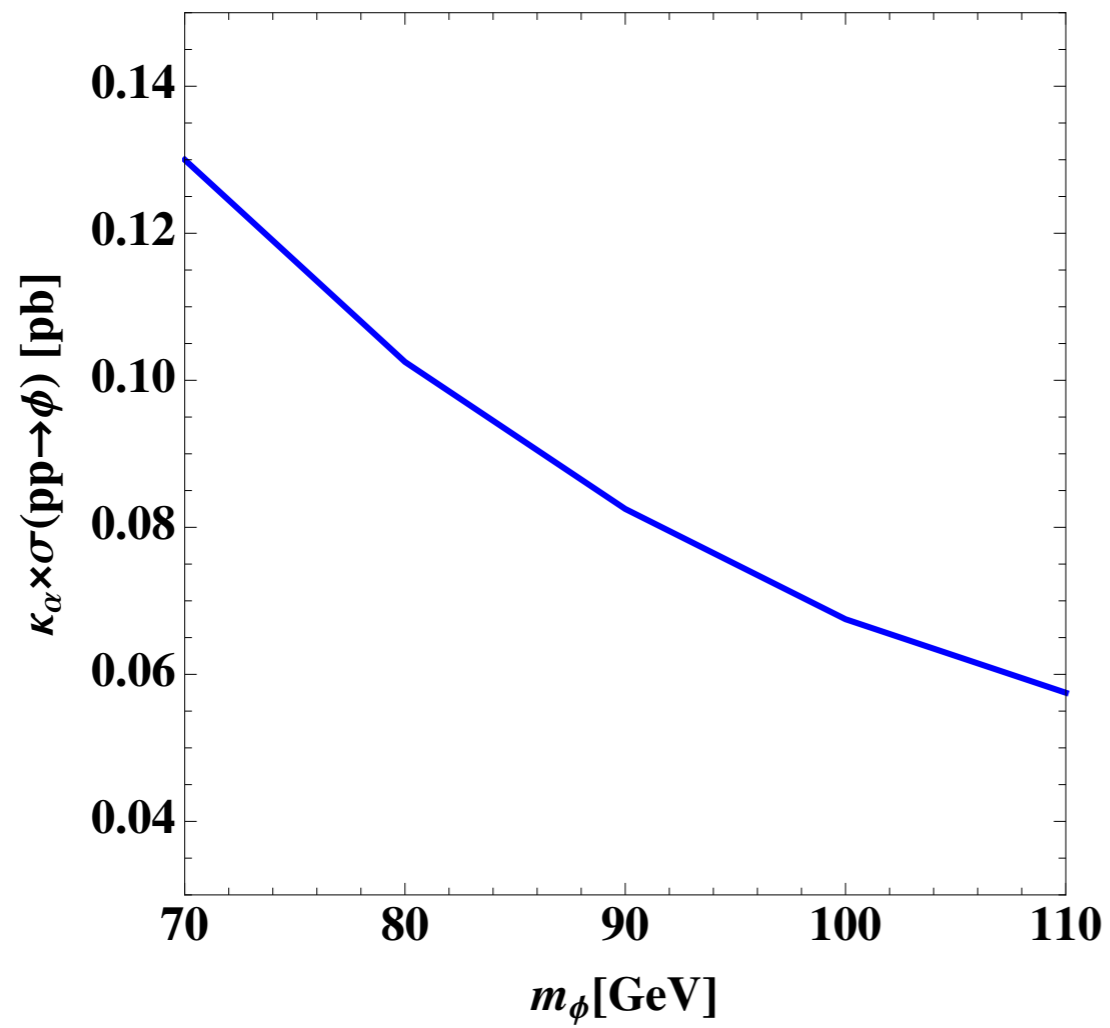
$$m_{Z'} = \sqrt{M_{e^+e^-}^2}$$
$$m_{\phi} = \sqrt{(p_{Z'} + p'_{Z'})^2}$$

The SM BGs are ZZ production and Drell-Yang with ISR/FSR e^+e^-

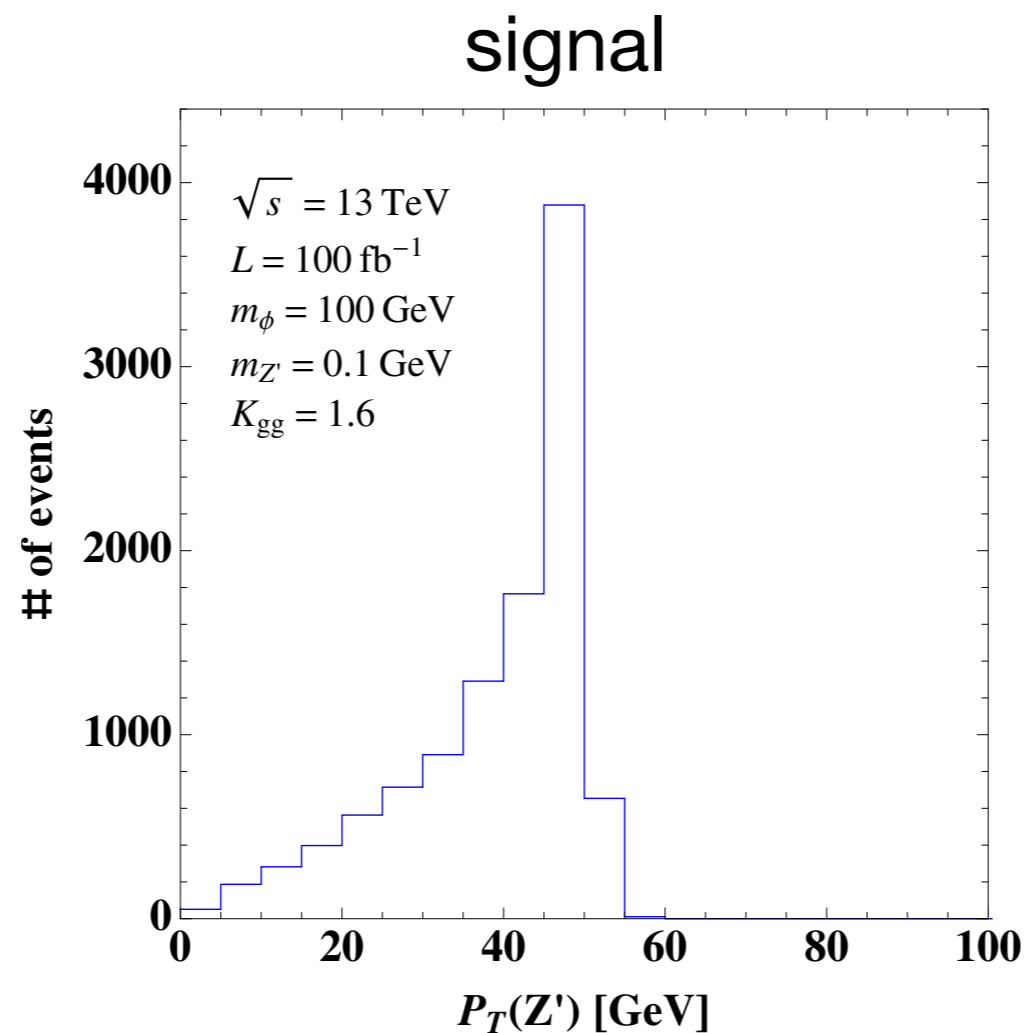
The momentum distrib. of e^+e^- in BG are different from the signal.
Kinematical cuts will reduce the BGs

LHC experiment

Production cross section at LHC13



LHC experiment



$O(10^3)$ events expected
even $BR(Z' \rightarrow e^+e^-) = 0.07$

However, e^+e^- pair is highly collimated, with the angle $\theta \sim 0.5^\circ$

Analysis of such a collimated e^+e^- is challenging at present