散乱振幅で理論的に探る電弱対称性の破れ

The electroweak effective field theory from on-shell amplitudes

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基研研究会 素粒子物理学の進展2020

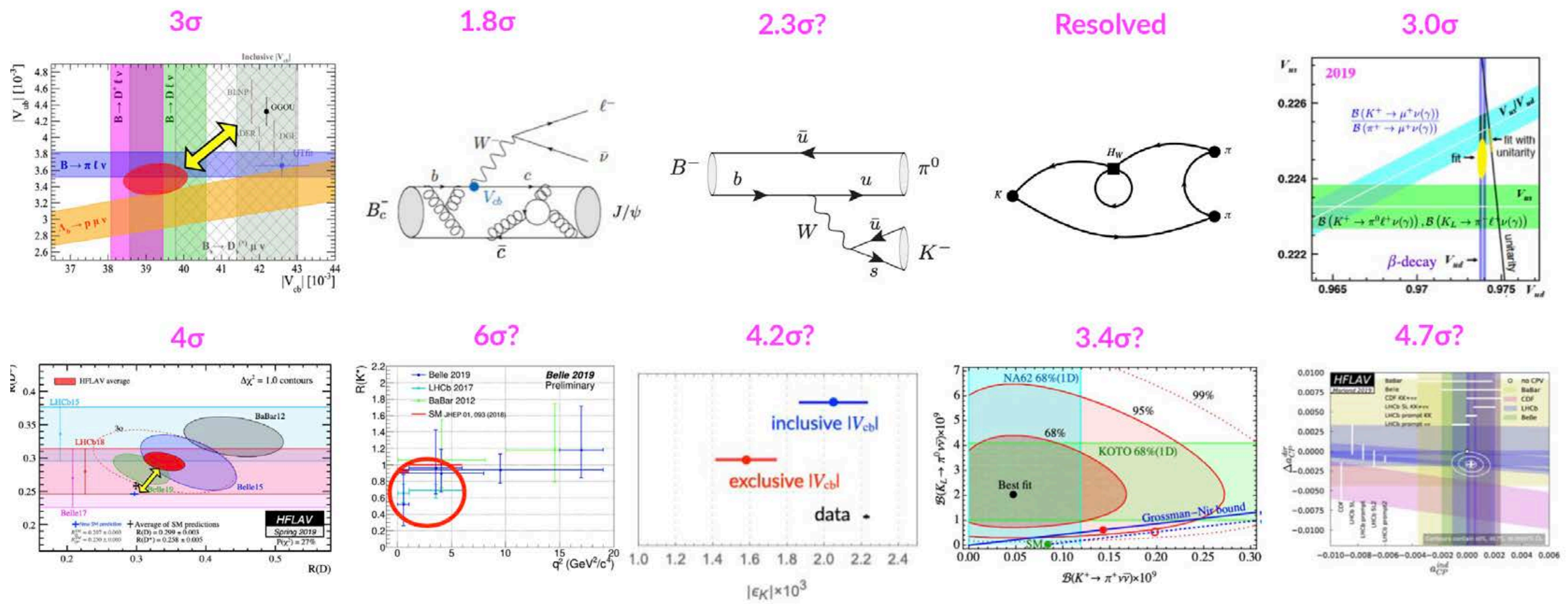
2020年9月4日, オンライン

本題に入る前に...

- フレーバーは出てきません
- グラフや実験結果は出てきません
- hep-phとhep-thの境界領域の研究です
- $D = 4$, Minkowski metric (+ , − , − , −)
- 興味のある方は一緒に共同研究しましょう

最新のフレーバーのレビュートークは
こちらをクリック

「中間子の精密測定におけるアノマリーの現状と新物理の識別」
於　物理学会第75回年次大会 (招待講演)，京都大学セミナー



Based on

Novel formalism

[1709.04891]
Nima Arkani-Hamed, Tzu-Chen Huang, Yu-tin Huang

Technion, scattering amplitudes group

[1809.09644]
Yael Shadmi, Yaniv Weiss

[1909.10551]
Gauthier Durieux, **TK**, Yael Shadmi, Yaniv Weiss

[2008.09652]
Gauthier Durieux, **TK**, Camila S. Machado, Yael Shadmi, Yaniv Weiss

Introduction (1/2)

- Effective field theory (EFT) can be generally constructed by assuming field contents and Lorentz, global and gauge symmetries, *e.g.*, SMEFT, HEFT, HQET, SCET, …
- EFT is bottom-up and natural approach (when one does not discover any new resonance)
- General problems of (effective) Lagrangian treatment:
  - Find nice operator basis: operator redundancy via field redefinitions and EOMs

    e.g., Warsaw basis (dimension-six SMEFT) [Grzadkowski, Iskrzynski, Misiak, Rosiek '10]
  - Gauge redundancy (=gauge-fixing dependence), which is canceled out at amplitude level (after the complicated calculations)

Introduction (2/2)

- Scattering amplitude (on-shell amplitude, modern amplitude method, or spinor-helicity formalism) is an alternative way to EFTs (will explain at on after next slide)
- Scattering amplitudes can be bootstrapped from Lorentz symmetry, locality and unitarity
- Advantages:
  - No operator and gauge redundancies. Gauge invariance is manifest
  - Bypassing Lagrangian, operators, and Feynman rules/diagrams
  - Drastically simple results compared to Feynman methods

e.g., $gg \to ggg$ [Mangano, Parke '91]

$$\mathcal{M}_5(1_g^-, 2_g^-, 3_g^+, 4_g^+, 5_g^+) = i g_s^3 \frac{\langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 51 \rangle}$$

corresponds to sum of 25 diagrams.
n g is impossible by the Feynman methods

On-shell approach to the SMEFT

- Derive anomalous dimension matrix (one- and two-loop levels)

  [Cheung, Shen ’15; Bern, Parra-Martinez, Sawyer ’19, ’20; Elias Miro, Ingoldby, Riembau ’20; Jiang, Ma, Shu ’20]

- Derive non-interference theorem for the new physics operators

  [Azatov, Contino, Machado, Riva ’16; Craig, Jiang, Li, Sutherland ’20, Jiang, Shu, Xiao, Zheng ’20; Gu, Wang ‘20]

- Enumeration of independent massless operators (consistent with Hilbert series approach)

  [Shadmi, Weiss ’18; Ma, Shu, Xiao ’19; Falkowski ’19; Durieux, Machado ’19; Durieux, TK, Machado, Shadmi, Weiss ’20]

  Hilbert series [Henning, Lu, Melia, Murayama ’15, ‘17]

- Investigate the electroweak symmetry (relations from $SU(2)_L \times U(1)_Y$ SSB) using massive scattering amplitudes

  [Christensen, Field ’18; Aoude, Machado ‘19; Christensen, Field, Moore, Pinto ’19; Durieux, TK, Shadmi, Weiss ’19; Bachu, Yelleshpur ‘19]

This talk

Spinor-helicity formalism (massless scattering amplitudes) (1/2)

reviews *e.g.*, [Elvang, Huang '13, Dixon '13; Schwartz '14]

- Massless particle is an irreducible representations of the Poincaré group; particle $i = |p_i, h_i\rangle$

  $h = \pm 1/2, \pm 1$ is particle's helicity

- Massless *n*-pt amplitudes are given by $M_n(p_1^{h_1}, p_2^{h_2}, \ldots, p_n^{h_n})$ (all particles are incoming)

- Little-group (LG) is subgroup of the Lorentz group, which leaves p_i invariant; $p_i \to p_i$

- In $D = 4$, SO(2) $\simeq$ U(1) LG for massless particle

- Massless amplitudes are scaled by their helicities $\{h_1, h_2, \ldots\}$ under U(1) LG transformation

  Little group scaling; $M_n(p_1^{h_1}, \ldots, p_n^{h_n}) \to e^{2i\xi \sum h_i} M_n(p_1^{h_1}, \ldots, p_n^{h_n})$

Spinor-helicity formalism (massless scattering amplitudes) (2/2)

Lorentz group
irreducible representation

	symbol	(A, B) $\hat{A},\hat{B} = \frac{1}{2}(\hat{J} \pm i\hat{K})$	spinor-helicity formalism	
undotted spinor	$\lambda_{i,\alpha} = u_-(p_i), \bar{v}_-(p_i)$	**2**: (1/2, 0)	$\lvert i\rangle_\alpha \to e^{-i\xi} \lvert i\rangle_\alpha$ (under LG)	$\langle ij\rangle = -\langle ji\rangle$
dotted spinor	$\tilde{\lambda}_i^{\dot{\alpha}} = u_+(p_i), \bar{v}_+(p_i)$	**2***: (0, 1/2)	$\lvert i]^{\dot{\alpha}} \to e^{+i\xi} \lvert i]^{\dot{\alpha}}$ (under LG)	$\langle ii\rangle = [ii] = 0$
4-vector	p_i^μ	**2**×**2***: (1/2, 1/2)	$p_{i,\alpha\dot{\alpha}} = p_i^\mu \sigma_{\mu,\alpha\dot{\alpha}} = \lvert i\rangle_\alpha [i\rvert_{\dot{\alpha}}$	$\det p_{i,\alpha\dot{\alpha}} = p_i^2 = 0$
polarization vector	$\varepsilon_i^{\mu,\pm}$	constrained 4-vector $p_i \cdot \varepsilon_i^\pm = 0,\quad \varepsilon_i^\pm \cdot (\varepsilon_i^\pm)^* = -1$ $\sum_{\lambda=\pm} \varepsilon_i^{\mu,\lambda} (\varepsilon_i^{\nu,\lambda})^* = -\eta^{\mu\nu}$	$\varepsilon^+_{i,\alpha\dot{\alpha}} = \varepsilon_i^{\mu,+}\sigma_{\mu,\alpha\dot{\alpha}} = \sqrt{2}\frac{\lvert\zeta\rangle_\alpha [i\rvert_{\dot{\alpha}}}{\langle i\zeta\rangle}$ $\varepsilon^-_{i,\alpha\dot{\alpha}} = \varepsilon_i^{\mu,-}\sigma_{\mu,\alpha\dot{\alpha}} = \sqrt{2}\frac{\lvert i\rangle_\alpha [\zeta\rvert_{\dot{\alpha}}}{[i\zeta]}$	auxiliary spinor ζ
⋮	⋮		⋮	

massless → **massive**

[Kleiss, Stirling ’85; Dittmaier ’98; Cohen, Elvang, Kiermaier ‘10]

↓ formalize/generalize for any mass and spin particles

[1709.04891]
Arkani-Hamed, Huang, Huang

Massive-spinor formalism (1/4) [Arkani-Hamed, Huang, Huang '17]

$$\det p_{i,\alpha\dot\alpha} = \det p_i\cdot\sigma = \overbrace{\begin{vmatrix} p_i^0+p_i^3 & p_i^1-ip_i^2 \\ p_i^1+ip_i^2 & p_i^0-p_i^3\end{vmatrix}}^{p_{i,\alpha\dot\alpha}} = (p_i^0)^2-(p_i^1)^2-(p_i^2)^2-(p_i^3)^2$$

$$= p_i^2 = 0 \longrightarrow \; = m^2 > 0$$

$p_{i,\alpha\dot\alpha}$: rank 1 → product of two vectors

$$p_{i,\alpha\dot\alpha} = |i\rangle_\alpha [i|_{\dot\alpha}$$

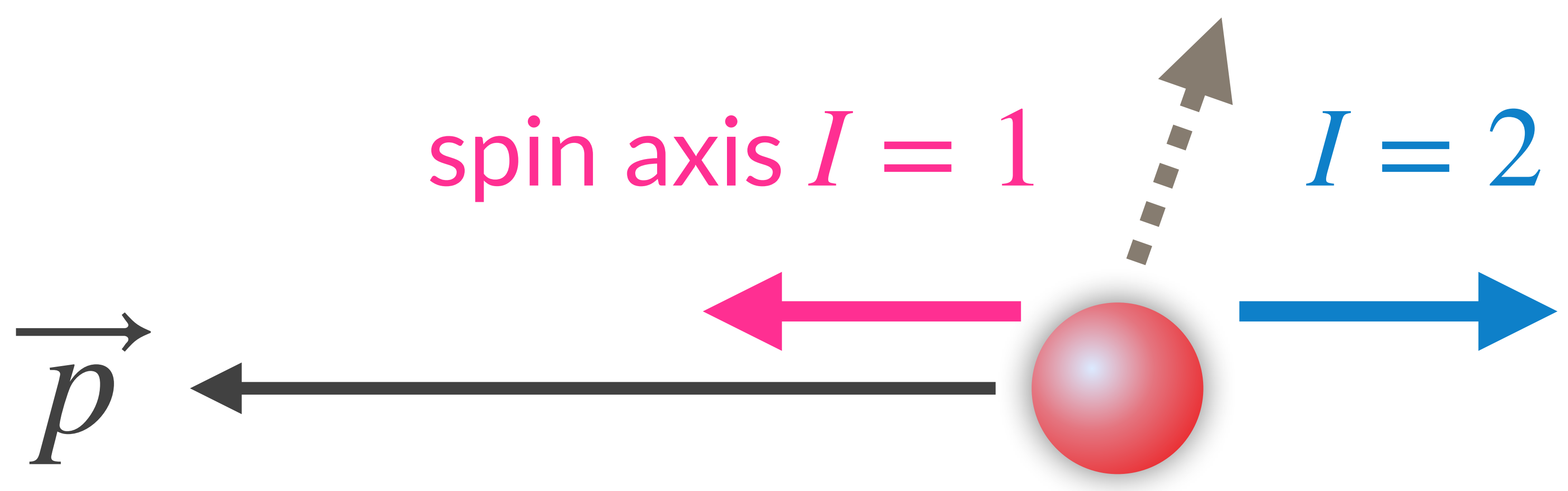
rank 2 → sum of two products of two vectors

$$p_{i,\alpha\dot\alpha} = |i^1\rangle_\alpha [i_1|_{\dot\alpha} + |i^2\rangle_\alpha [i_2|_{\dot\alpha} \equiv \sum_{I=1,2} |\mathbf{i}^I\rangle_\alpha [\mathbf{i}_I|_{\dot\alpha}$$

- In $D=4$, SO(3) ≃ SU(2) LG for massive particles; leaves $p_{i,\alpha\dot\alpha}$ invariant; $p_{i,\alpha\dot\alpha} \to p_{i,\alpha\dot\alpha}$
- Amplitudes are transformed by SU(2) LGs (for massive external particles)
- **Bold spinors** $|\mathbf{i}^I\rangle, |\mathbf{i}^I]$ carry the SU(2) LG index $I = 1,2$

Massive-spinor formalism (2/4)

- One can use the SU(2) LG rotation for the spin-quantization axis
- Convenient choice (for any spin particles):



Arbitrary spin polarization can be given by two opposite spin states

$$\begin{pmatrix} a \\ b \end{pmatrix} = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix} = a|+_z\rangle + b|-_z\rangle$$

- In this choice, in high energy limit, $I = 1$ ($I = 2$) spinor corresponds to positive (negative) helicities
- Any choice of spin-quantization axis is possible in general (“SU(2) LG covariant”)

Massive-spinor formalism (3/4)

	symbol	massive-spinor formalism	
undotted spinor	$\lambda^s_{i,\alpha} = P_L u^I(p_i), \bar{v}^I(p_i) P_L$	$\vert \mathbf{i}^I\rangle_\alpha \to W^I_J \vert \mathbf{i}^J\rangle_\alpha$ (under LG)	$\langle \mathbf{i}^I \mathbf{j}^J \rangle = -\langle \mathbf{j}^J \mathbf{i}^I \rangle$
dotted spinor	$\tilde{\lambda}^{s,\dot{\alpha}}_i = P_R u^I(p_i), \bar{v}^I(p_i) P_R$	$\vert \mathbf{i}^I]^{\dot{\alpha}} \to \left(W^{-1}\right)^I_J \vert \mathbf{i}^J]^{\dot{\alpha}}$ (under LG)	$\langle \overbracket{\mathbf{i}^I \mathbf{i}^J} \rangle = [\overbracket{\mathbf{i}^I \mathbf{i}^J}] = 0$
4-vector	p^μ_i	$p_{i,\alpha\dot{\alpha}} = p^\mu_i \sigma_{\mu,\alpha\dot{\alpha}} = \sum_{I=1,2} \vert \mathbf{i}^I\rangle_\alpha [\mathbf{i}_I \vert_{\dot{\alpha}}$	$\det p_{i,\alpha\dot{\alpha}} = p_i^2 = m^2$
polarization vector	$\varepsilon^{\mu,\pm,L}_i$	$\varepsilon^{IJ}_{i,\alpha\dot{\alpha}} = \varepsilon^{\mu,\pm,L}_i \sigma_{\mu,\alpha\dot{\alpha}} = \sqrt{2}\dfrac{\vert \mathbf{i}^I\rangle_\alpha [\mathbf{i}^J \vert_{\dot{\alpha}}}{m}$	no auxiliary spinor
⋮	⋮	⋮	

Massive-spinor formalism (4/4)

- Equations of motion (EOM) $\sim$ "chirality flip"

$$\bar{p}_i|\mathbf{i}^I\rangle = m|\mathbf{i}^I], \;\; p_i|\mathbf{i}^I] = m|\mathbf{i}^I\rangle, \;\; \langle\mathbf{i}^I|p_i = -m[\mathbf{i}^I|, \;\; [\mathbf{i}^I|\bar{p}_i = -m\langle\mathbf{i}^I|$$

- Massive polarization vectors [Gauthier Durieux, **TK**, Yael Shadmi, Yaniv Weiss '19]

$$\varepsilon^{IJ}_{i,\alpha\dot\alpha} = \sqrt{2}\frac{|\mathbf{i}^I\rangle_\alpha[\mathbf{i}^J|_{\dot\alpha}}{m} \quad \begin{cases} \varepsilon^{+}_{i,\alpha\dot\alpha} = \varepsilon^{11}_{i,\alpha\dot\alpha} = \sqrt{2}\frac{|\mathbf{i}^1\rangle_\alpha[\mathbf{i}^1|_{\dot\alpha}}{m} \xrightarrow{m\to 0} \sqrt{2}\frac{|\zeta\rangle_\alpha[i|_{\dot\alpha}}{\langle i\zeta\rangle} = \varepsilon^{+}_{i,\alpha\dot\alpha} \\ \varepsilon^{-}_{i,\alpha\dot\alpha} = \varepsilon^{22}_{i,\alpha\dot\alpha} = \sqrt{2}\frac{|\mathbf{i}^2\rangle_\alpha[\mathbf{i}^2|_{\dot\alpha}}{m} \longrightarrow \sqrt{2}\frac{|i\rangle_\alpha[\zeta|_{\dot\alpha}}{[i\zeta]} = \varepsilon^{-}_{i,\alpha\dot\alpha} \\ \varepsilon^{L}_{i,\alpha\dot\alpha} = \varepsilon^{\sqcap 12}_{i,\alpha\dot\alpha} = \frac{|\mathbf{i}^1\rangle_\alpha[\mathbf{i}^2|_{\dot\alpha} + |\mathbf{i}^2\rangle_\alpha[\mathbf{i}^1|_{\dot\alpha}}{m} \longrightarrow \sim\frac{p_{i,\alpha\dot\alpha}}{m} = \mathcal{O}\left(\frac{E}{m}\right) \end{cases}$$

massless polarizations

well-known energy growth

Factor $1/\sqrt{2}$ (in *L* mode) corresponds to Clebsch-Gordan; we modify the original formalism

$$\Rightarrow \quad p_i\cdot\varepsilon^{\pm}_i = 0, \qquad \varepsilon^{\pm,L}_i\cdot(\varepsilon^{\pm,L}_i)^* = -1, \qquad \sum_{\lambda=\pm,L}\varepsilon^{\mu,\lambda}_i(\varepsilon^{\nu,\lambda}_i)^* = -\left(\eta^{\mu\nu} - \frac{p_{i,\mu}p_{i,\nu}}{m^2}\right)$$

corresponds to "unitary gauge"

Our several results

Our strategy

- Spectrum: different masses + massless photon

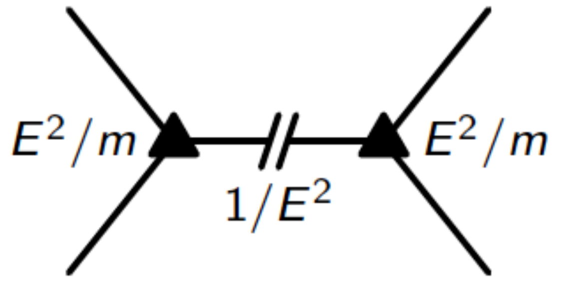
$$\psi\,(\psi^c), Z, W^\pm, h \quad +\gamma$$

- We do not impose $SU(2)_L \times U(1)_Y$ symmetry, but impose only $U(1)_{EM}$

- [LGs ⊂ Lorentz ⊂ Poincaré] + [locality] [Arkani-Hamed, Huang, Huang '17]

  + [perturbative unitarity ⊂ unitarity] [Gauthier Durieux, **TK**, Yael Shadmi, Yaniv Weiss '19]

  - For three-pt amplitudes,

    E^2/m has to be forbidden; 

$\sim E^2/m^2$ unacceptable energy growth

  - For full four-pt amplitudes, E/m has to be forbidden

Note that: there is no longitudinal mode in massless scattering amplitudes

Three-point: hhZ

- Result (LGs + locality): [Durieux, **TK**, Shadmi, Weiss '19]

SU(2) LG indices I, J are implicit

$$\mathcal{M}_3(\mathbf{1}_h, \mathbf{2}_h, \mathbf{3}_Z) \propto \langle \mathbf{3}(\mathbf{1}-\mathbf{2})\mathbf{3}] \quad = \langle \mathbf{3}|(p_1 - p_2)|\mathbf{3}] \text{ (notation)}$$

The scalars $\mathbf{1}$ and $\mathbf{2}$ have to be asymmetric:
when the scalars $\mathbf{1}$ and $\mathbf{2}$ are identical, this amplitude must vanish at the all order

One-line proof to "why $\rho^0 \to 2\pi^0$ is forbidden in our world"

A good application of massive scattering amplitude!

Three-point: W^+W^-Z (1/4)

- Result (LGs + locality): 11 spinor structures [Arkani-Hamed, Huang, Huang '17]

  Schouten identity, and momentum conservation → 8 spinor structures

  $|\mathbf{i}\rangle\langle\mathbf{jk}\rangle + |\mathbf{j}\rangle\langle\mathbf{ki}\rangle + |\mathbf{k}\rangle\langle\mathbf{ij}\rangle = 0$ $\quad p_1 + p_2 + p_3 = 0$

- Furthermore, we observe a non-trivial massive spinor identity [Durieux, **TK**, Shadmi, Weiss '19, + Machado '20]

$$m_1\langle\mathbf{12}\rangle\langle\mathbf{13}\rangle[\mathbf{23}] + m_2\langle\mathbf{12}\rangle[\mathbf{13}]\langle\mathbf{23}\rangle + m_3[\mathbf{12}]\langle\mathbf{13}\rangle\langle\mathbf{23}\rangle = m_1[\mathbf{12}][\mathbf{13}]\langle\mathbf{23}\rangle + m_2[\mathbf{12}]\langle\mathbf{13}\rangle[\mathbf{23}] + m_3\langle\mathbf{12}\rangle[\mathbf{13}][\mathbf{23}]$$

  → 7 spinor structures (final)

- Angular momentum conservation (in three-pt amplitudes): [Costa, Penedones, Poland, Rychkov '11]

  # of irreps of sum of three spins = # of independent spinors in three-pt amplitudes

  → 7 combinations is expected $\quad 3 \otimes 3 \otimes 3 = 1 \oplus 3 \oplus 3 \oplus 3 \oplus 5 \oplus 5 \oplus 7$

- 7 form factors for general WWZ coupling [Hagiwara, Peccei, Zepenfeld, Hikasa '86]

Three-point: W^+W^-Z (2/4)

- + perturbative unitarity [Durieux, **TK**, Shadmi, Weiss '19]

  $\bar{\Lambda}$ dependence of 7 spin structures is fully determined

  c_{WWZ} : dimensionless

$$
\begin{aligned}
\mathcal{M}(\mathbf{1}_{W^+}, \mathbf{2}_{W^-}, \mathbf{3}_Z) =& 2\frac{c_{WWZ}}{m_Z m_W}\left(\frac{m_Z}{m_W}\langle\mathbf{12}\rangle[\mathbf{13}][\mathbf{23}] + [\mathbf{12}]\langle\mathbf{13}\rangle[\mathbf{23}] + [\mathbf{12}][\mathbf{13}]\langle\mathbf{23}\rangle\right) \\
&+ \frac{c_{WWZ}^{[L0]0}}{m_Z\bar{\Lambda}}\langle\mathbf{12}\rangle\left(\langle\mathbf{13}\rangle[\mathbf{23}] - [\mathbf{13}]\langle\mathbf{23}\rangle\right) + \frac{c_{WWZ}^{\{L0\}0}}{m_Z\bar{\Lambda}}\langle\mathbf{12}\rangle\left(\langle\mathbf{13}\rangle[\mathbf{23}] + [\mathbf{13}]\langle\mathbf{23}\rangle\right) \\
&+ \frac{c_{WWZ}^{[R0]0}}{m_Z\bar{\Lambda}}[\mathbf{12}]\left(\langle\mathbf{13}\rangle[\mathbf{23}] - [\mathbf{13}]\langle\mathbf{23}\rangle\right) + \frac{c_{WWZ}^{\{R0\}0}}{m_Z\bar{\Lambda}}[\mathbf{12}]\left(\langle\mathbf{13}\rangle[\mathbf{23}] + [\mathbf{13}]\langle\mathbf{23}\rangle\right) \\
&+ \frac{c_{WWZ}^{RRR}}{\bar{\Lambda}^2}[\mathbf{12}][\mathbf{13}][\mathbf{23}] + \frac{c_{WWZ}^{LLL}}{\bar{\Lambda}^2}\langle\mathbf{12}\rangle\langle\mathbf{13}\rangle\langle\mathbf{23}\rangle\,.
\end{aligned}
$$

non-trivial single renormalizable structure

$\mathbf{1} \leftrightarrow \mathbf{2}$ symmetric: $\not{C}$

- $m_Z \to 0$ limit provides $M_3(\mathbf{1}_{W^+}, \mathbf{2}_{W^-}, 3_\gamma^\pm)$ with 5 spin structures

  consistent with angular momentum analysis: $3 \otimes 3 \otimes 2 = 2 \oplus 2 \oplus 4 \oplus 4 \oplus 6$

Three-point: W^+W^-Z (3/4)

- Moreover, we match the massive scattering amplitudes onto the SMEFT in the broken phase.

compare our massive amplitudes to the SMEFT

$$\mathcal{L}_{\rm SMEFT} = \mathcal{L}_{\rm SM} + \sum_i \frac{C_i}{\Lambda^2}\mathcal{O}_i$$

Warsaw basis (dimension-six SMEFT)
[Grzadkowski, Iskrzynski, Misiak, Rosiek '10]

Warsaw basis in the broken phase
[Dedes, Materkowska, Paraskevas, Rosiek, Suxho '17]

result of 7 coefficients

$$2c_{WWZ} = -\sqrt{2}\frac{\bar{g}^2}{\sqrt{\bar{g}^2+\bar{g}'^2}} + \sqrt{2}\frac{\bar{g}^3\bar{g}'}{(\bar{g}^2+\bar{g}'^2)^{3/2}}\frac{v^2}{\Lambda^2}C_{\varphi WB}\,,$$

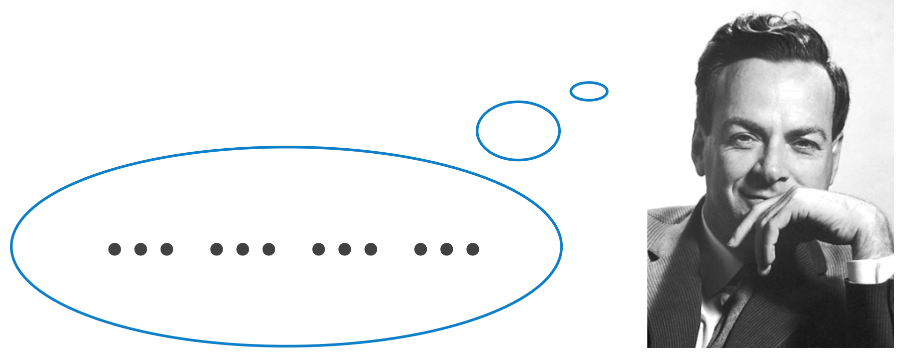
$$c^{[L0]0}_{WWZ} = c^{[R0]0}_{WWZ} = 0\,, \quad \not{C}$$

$$\frac{c^{\{R0\}0}_{WWZ}}{m_Z\Lambda} = -\frac{1}{\sqrt{2}m_W m_Z}\frac{\bar{g}\bar{g}'}{\sqrt{\bar{g}^2+\bar{g}'^2}}\frac{v^2}{\Lambda^2}\left(C_{\varphi WB}+iC_{\varphi\tilde{W}B}\right)\,,$$

$$c^{\{L0\}0}_{WWZ} = \left(c^{\{R0\}0}_{WWZ}\right)^*\,,$$

$$\frac{c^{RRR}_{WWZ}}{\Lambda^2} = -3\sqrt{2}\frac{\bar{g}}{\sqrt{\bar{g}^2+\bar{g}'^2}}\frac{1}{\Lambda^2}(C_W+iC_{\tilde{W}})\,,$$

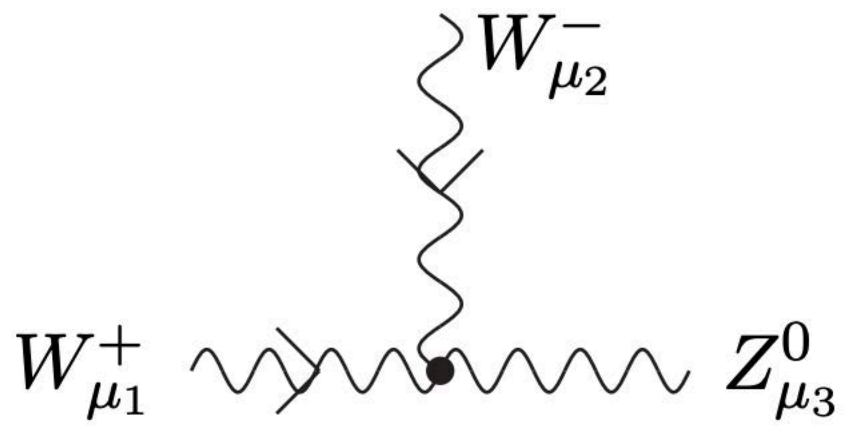
$$c^{LLL}_{WWZ} = \left(c^{RRR}_{WWZ}\right)^*\,.$$



Three-point: W^+W^-Z (4/4)

- One example

dimension-six operator: $\varepsilon^{IJK} W_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$



Feynman rule for this operator

$$-\frac{6i\bar{g}}{\sqrt{\bar{g}^2+\bar{g}'^2}}\frac{C_W}{\Lambda^2}\left(p_3^{\mu_1}p_1^{\mu_2}p_2^{\mu_3} - p_2^{\mu_1}p_3^{\mu_2}p_1^{\mu_3} + \eta_{\mu_1\mu_2}\left(p_1^{\mu_3}p_2\cdot p_3 - p_2^{\mu_3}p_1\cdot p_3\right)\right.$$
$$\left.+\eta_{\mu_2\mu_3}\left(p_2^{\mu_1}p_1\cdot p_3 - p_3^{\mu_1}p_1\cdot p_2\right) + \eta_{\mu_3\mu_1}\left(p_3^{\mu_2}p_1\cdot p_2 - p_1^{\mu_2}p_2\cdot p_3\right)\right)$$

massive-spinor formalism

$$-3\sqrt{2}\frac{\bar{g}}{\sqrt{\bar{g}^2+\bar{g}'^2}}\frac{C_W}{\Lambda^2}$$
$$\times\left([\mathbf{12}][\mathbf{13}][\mathbf{23}] + \langle\mathbf{12}\rangle\langle\mathbf{13}\rangle\langle\mathbf{23}\rangle\right)$$

@Wikipedia

All EW three-points are bootstrapped and mapped

[Durieux, **TK**, Shadmi, Weiss '19]

LGs
+
Locality
+
Unitarity
+
$U(1)_{EM}$

bootstrap

our work

$\psi^c \psi Z$

$\psi^c \psi \gamma$

$\psi^c \psi' W$

$\psi^c \psi h$

$ZZh,\ WWh$

$\gamma\gamma h,\ \gamma Zh$

$hhZ,\ hh\gamma$ (example 1)

hhh

WWZ (example 2)

$WW\gamma$

$ZZZ,\ ZZ\gamma,\ Z\gamma\gamma,\ \gamma\gamma\gamma$

map

our work

[Dedes, Materkowska, Paraskevas, Rosiek, Suxho '17]

broken phase

$v \neq 0$

canonically normalized

Warsaw basis

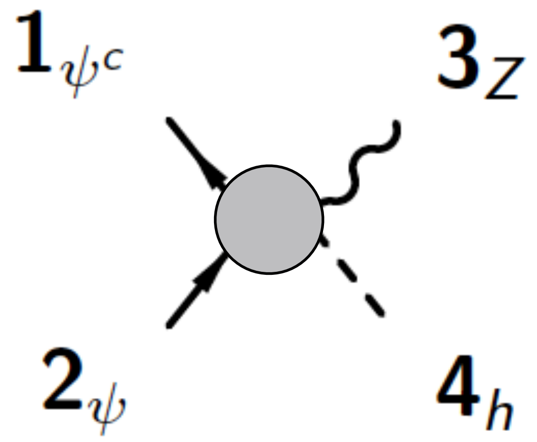
$$\mathcal{L}_{\rm SMEFT} = \mathcal{L}_{\rm SM} + \sum_i \frac{C_i}{\Lambda^2}\mathcal{O}_i$$

They are important building blocks for the SMEFT computations with on-shell techniques

Four-point: $\psi^c\psi Zh$

factorizable contribution

non-factorizable contribution (contact term)



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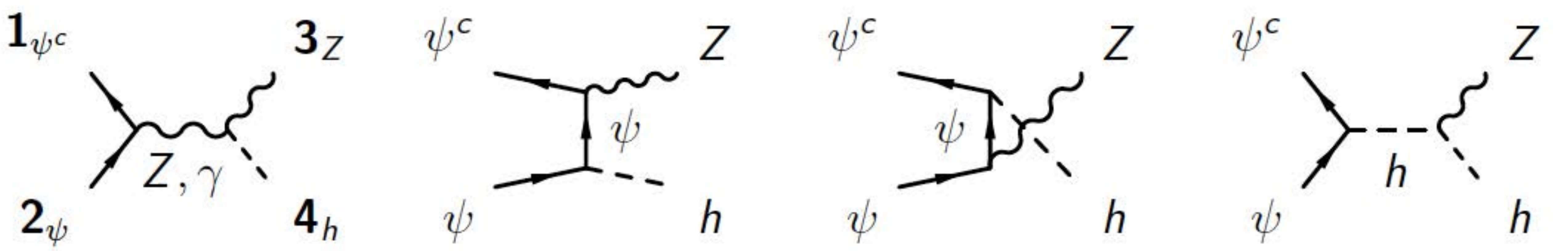

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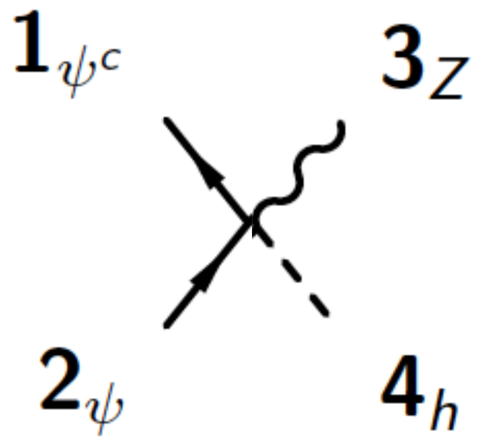

+ perturbative unitarity requires [Durieux, **TK**, Shadmi, Weiss '19]

$$(--00): \quad -\langle 12\rangle \left(c^{RL0}_{\psi^c\psi Z} - c^{LR0}_{\psi^c\psi Z}\right)\left(c^{00}_{ZZh}\, m_\psi/2m_Z - c^{LL}_{\psi^c\psi h}\right)/\sqrt{2}m_Z \quad = 0 + \mathcal{O}(m/\bar{\Lambda})$$

$$(++00): \quad +[12] \left(c^{RL0}_{\psi^c\psi Z} - c^{LR0}_{\psi^c\psi Z}\right)\left(c^{00}_{ZZh}\, m_\psi/2m_Z - c^{RR}_{\psi^c\psi h}\right)/\sqrt{2}m_Z \quad = 0 + \mathcal{O}(m/\bar{\Lambda})$$

either vector-like fermion: $c^{RL0}_{\psi^c\psi Z} = c^{LR0}_{\psi^c\psi Z}$

or Higgs mechanism: $c^{RR}_{\psi^c\psi h} = c^{00}_{ZZh} m_\psi/2m_Z = c^{LL}_{\psi^c\psi h}$

up to $\mathcal{O}(m/\bar{\Lambda})$ consistent with study for $t\bar{t}Zh$ amplitude [Maltoni, Mantani, Mimasu '19]

single non-trivial identity is observed; 12 independent spinors are found

Soft Higgs limit recovers $\psi^c\psi Z$ amplitudes

Outlook

- Map EW four-point amplitudes onto the SMEFT
- Renormalization group evolution, running coupling in massive scattering amplitudes?
- An application: infrared photon/gluon corrections?

[Soft Matters, or the Recursions with Massive Spinors, Falkowski, Machado '20]

Conclusions

- The powerful scattering amplitude approach avoids gauge redundancy and operator redundancy
- We clarified a few details in the massive-spinor formalism, and bootstrapped all the EW three-point amplitudes, as well as the four-point amplitudes
- We mapped all EW three-point amplitudes onto the SMEFT
- We observed the emergence of the EW relations from the perturbative unitarity

- We paved the way for the SMEFT computations in the on-shell formalism

Backup