
COLOR FLUX TUBES IN TWO-COLOR QCD AT LOW TEMPERATURE AND HIGH DENSITY

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with

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Probing the physics of high-density and low-temperature matter with ab initio calculations in 2-color QCD, YITP workshop, 3rd - 6th November, 2020, Online





Outline

- Introduction
- Two-color QCD with $N_f=2$
- Dual superconductor picture of the QCD vacuum
- Color flux tubes
- Summary

Introduction

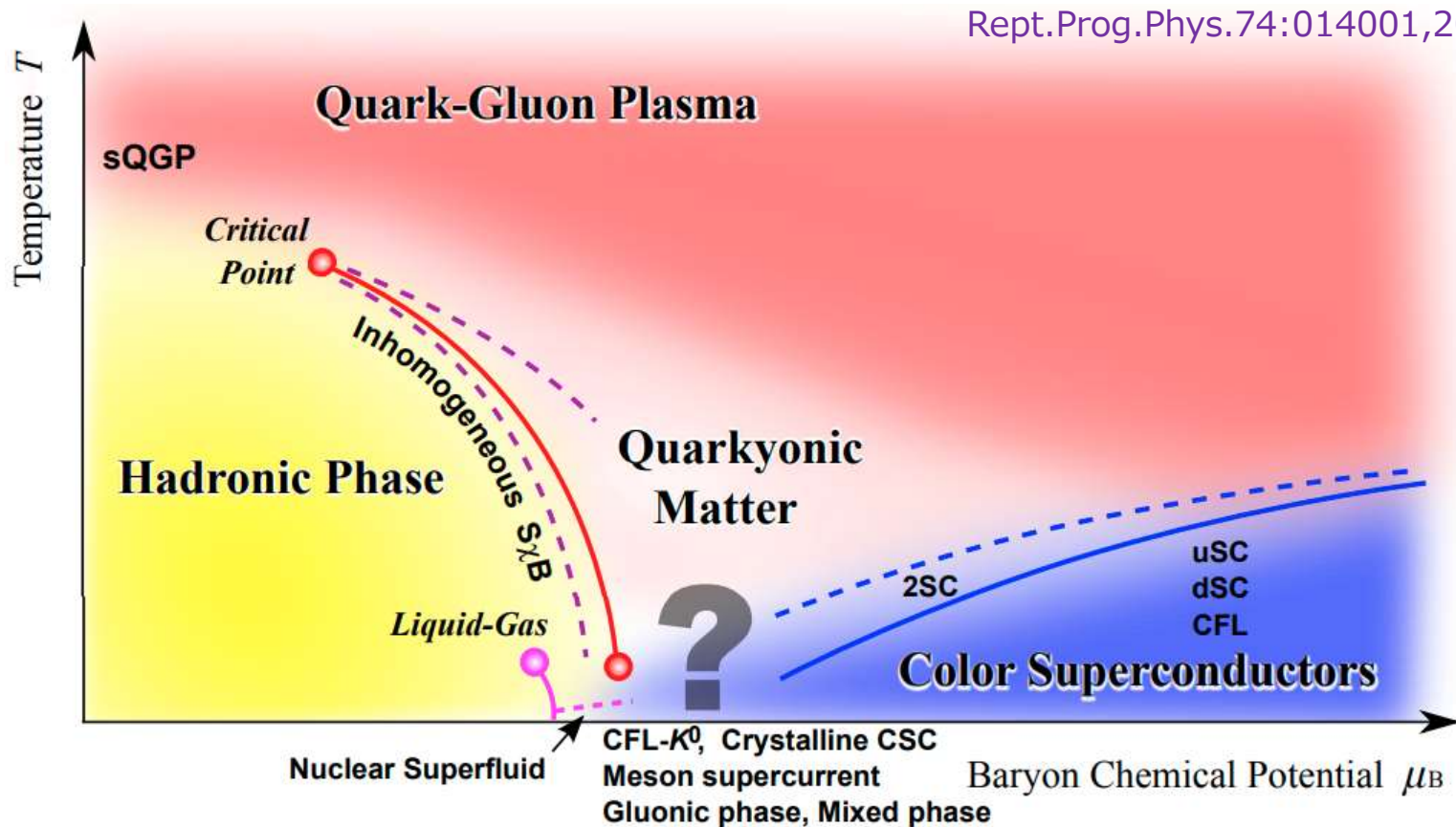
- Purpose
 - Understanding the physics of QCD in extreme conditions at finite temperature and density.
 - In the real world, we encountered the above situations.
 - For example, in neutron stars and relativistic heavy ion collisions.
 - What we want to investigate
 - Phase diagram on T- μ plane
 - Thermodynamic quantities
 - Nonperturbative properties (color confinement, χ_{SB} , ...)
- and others

- Many analytical calculations at finite temperature and density in QCD have been done in areas where perturbative calculations are useful, and various physics are known. However, perturbative calculations are limited to a part of the phase diagram in the T - μ plane.
- At zero baryon chemical potential, various non-perturbative quantities like the Polyakov loop and the chiral condensate can be calculated using lattice QCD simulations.
- Many lattice calculations have been carried out in the finite density regime, which is difficult to compute analytically. However, the applicability of this method is still limited due to the sign problem.

QCD phase diagram

- Schematic picture

K. Fukushima and T. Hatsuda
Rept.Prog.Phys.74:014001,2011



Two-color QCD with $N_f=2$

- In this work, we examine the temperature and density dependence of the phase and color flux tube structure of dense two-color QCD with two-flavor Wilson fermions by using a lattice simulation to avoid the sign problem.
- 2-color QCD has the same properties as 3-color QCD, e.g. color confinement and spontaneous chiral symmetry breaking.
- It is expected to provide insights for 3-color QCD.

Lattice setup

K. Iida, E. Itou and T.-G. Lee, JHEP01 (2020) 181

- Lattice action
 - Gauge part : Iwasaki gauge action

$$S_g = \beta \sum_x \left(c_0 \sum_{\substack{\mu < \nu \\ \mu, \nu=1}}^4 W_{\mu\nu}^{1 \times 1}(x) + c_1 \sum_{\substack{\mu \neq \nu \\ \mu, \nu=1}}^4 W_{\mu\nu}^{1 \times 2}(x) \right)$$

- Fermion part : Two-flavor Wilson fermion action including the quark number operator and the diquark source term

$$S_F = \bar{\psi}_1 \Delta(\mu) \psi_1 + \bar{\psi}_2 \Delta(\mu) \psi_2 - J \bar{\psi}_1 (C \gamma_5) \tau_2 \bar{\psi}_2^T + \bar{J} \psi_2^T (C \gamma_5) \tau_2 \psi_1$$

$$\Delta(\mu)_{x,y} = \delta_{x,y} - \kappa \sum_{i=1}^3 \left[(1 - \gamma_i) U_{x,i} \delta_{x+\hat{i},y} + (1 + \gamma_i) U_{y,i}^\dagger \delta_{x-\hat{i},y} \right] \\ - \kappa \left[e^{+\mu} (1 - \gamma_4) U_{x,4} \delta_{x+\hat{4},y} + e^{-\mu} (1 + \gamma_4) U_{y,4}^\dagger \delta_{x-\hat{4},y} \right]$$

Definition of the phases

K. Iida, E. Itou and T.-G. Lee
JHEP01 (2020) 181

	Hadronic		QGP	Superfluid	
		Hadronic matter		BEC	BCS
$\langle L \rangle$	zero	zero	non-zero		
$\langle qq \rangle$	zero	zero	zero	non-zero	$\propto \mu^2$
$\langle n_q \rangle$	$\langle n_q \rangle = 0$	$\langle n_q \rangle > 0$	$\langle n_q \rangle \geq 0$	non-zero	$\langle n_q \rangle / n_q^{\text{tree}} \approx 1$

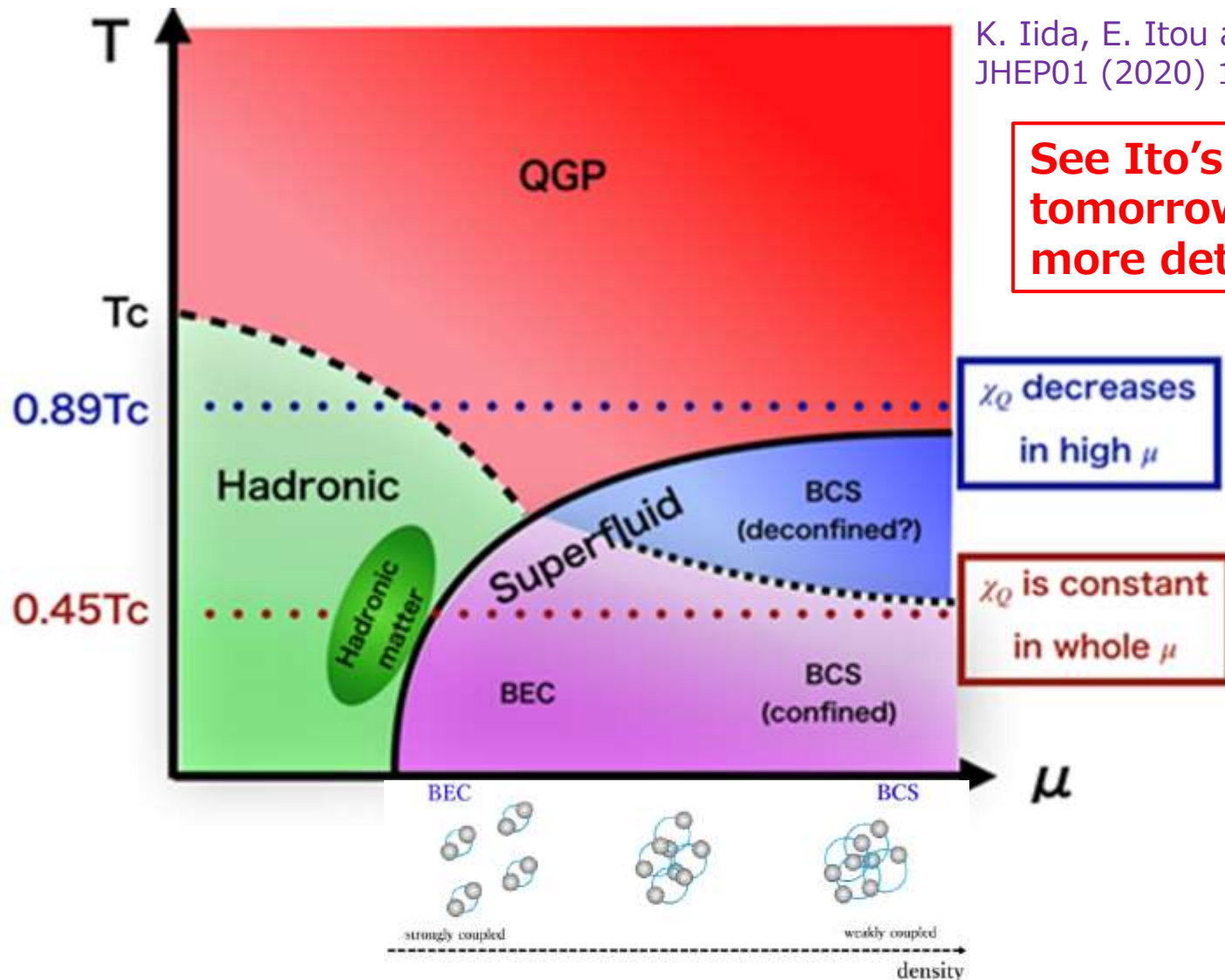
- $\langle |L| \rangle$: Polyakov loop
- $\langle qq \rangle$: diquark condensate
- $\langle n_q \rangle$: quark number density

$$L = \frac{1}{N_s^3} \sum_{\vec{x}} \prod_{\tau} U_4(\vec{x}, \tau)$$

$$\langle qq \rangle \equiv \frac{\kappa}{2} \langle \bar{\psi}_1 K \bar{\psi}_2^T - \psi_1 K \psi_2^T \rangle$$

$$a^3 n_q = \sum \kappa \langle \bar{\psi}_i(x) (\gamma_0 - \mathbb{I}_4) e^{\mu} U_4(x) \psi_i(x + \hat{4}) + \bar{\psi}_i(x) (\gamma_0 + \mathbb{I}_4) e^{-\mu} U_4^\dagger(x - \hat{4}) \psi_i(x - \hat{4}) \rangle.$$

Schematic picture of two-color QCD phase diagram



K. Iida, E. Itou and T.-G. Lee
JHEP01 (2020) 181

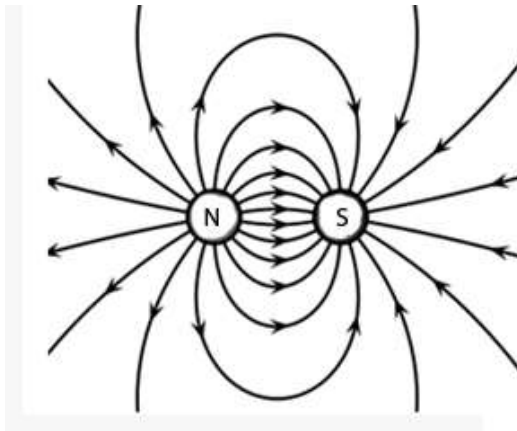
See Ito's talk tomorrow for more details!

Dual superconductor picture of the QCD vacuum

- In QCD vacuum, color magnetic monopoles condense instead of the formation of Cooper pairs in the BCS theory of normal superconductivity. (dual superconductor)
G. 'tHooft (1976), S.Mandelstam (1976)
- Color confinement is due to the dual Meissner effect by condensation of color magnetic monopoles.
- There are many numerical evidence for the dual superconductor picture of the QCD vacuum. Shiba-Suzuki (1994), Y.Matsubara et al.(1994), Cea-Cosmai(1995), G.S.Bali et al.(1998), A.Di Giacomo et al.(1999), ...
- Color electric field between quark and antiquark is squeezed into tube-like structure. The formation of color flux provides a linear potential between quark and antiquark.
- In this talk, we focus on the nature of color electric flux tube at low temperature and high density of two-color QCD.

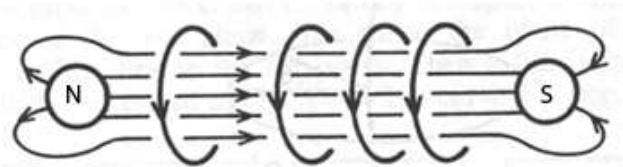
Flux tubes

- Magnetic field of a magnetic dipole



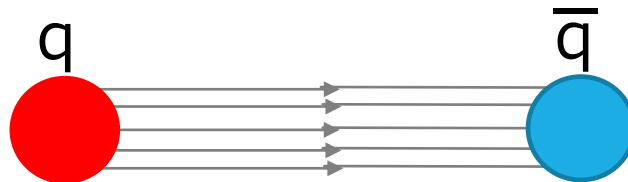
In empty space

K. Huang, *Int. J. Mod. Phys. A*30,1530056 (2015)

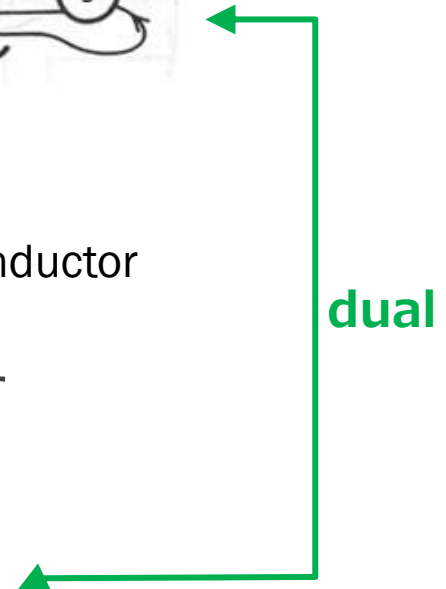


In a normal superconductor

- Color electric field of quark and antiquark pair



In the QCD vacuum



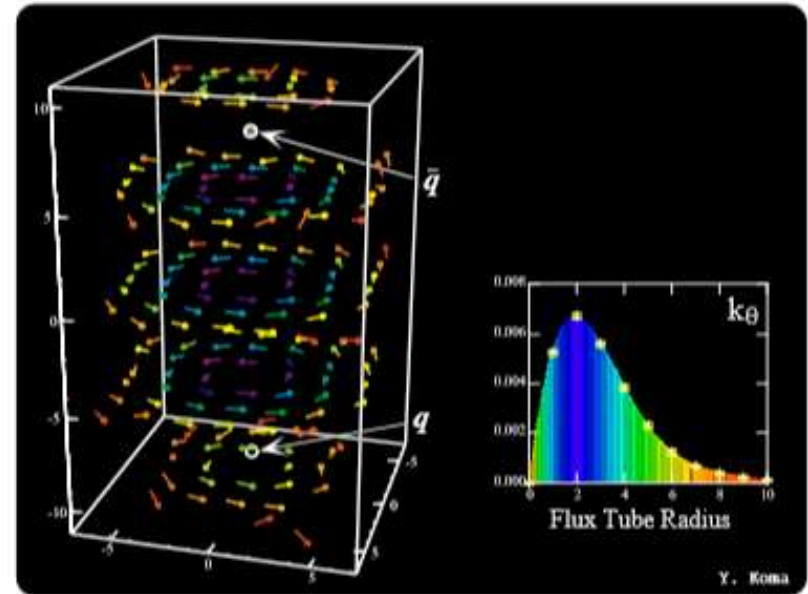
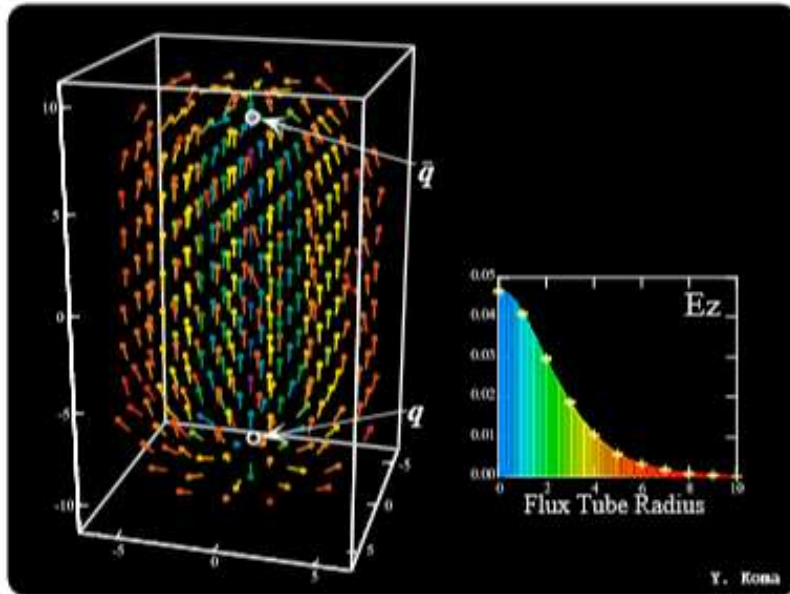
dual

Previous works of flux tubes in QCD

- quenched SU(2)
 - border between type 1 and type 2 (Y.Matsubara et al. '94, T.Suzuki et al. '07, P.Cea et al. '12)
 - Type 1 (G.S.Bali et al. '98)
 - Type 2 (P.Cea et al. '95)
 - weak Type 1 (S.Kato et al., '14)
- quenched SU(3)
 - border between type 1 and type 2 (Y.Matsubara '94)
 - Type 1 (P.Cea et al. '12)
- Finite temperature quenched SU(3) (P.Cea et al. '16)
- 2+1 flavor QCD
 - Type 1 (P.Cea et al. '17)
- 2+1 flavor QCD in external magnetic field (C.Bonati et al. '18)

The results for the type of vacuum are different depending on the physical quantity used and the analysis method.

Flux tube profiles in SU(2) gauge theory



- Quenched SU(2) QCD
- Maximally Abelian gauge
- Extract U(1) gauge fields and U(1) monopole
- E_z : U(1) electric field (longitudinal), k_θ : U(1) monopole

Y. Koma, M. Koma, E.-M. Ilgenfritz,
T. Suzuki and M. I. Polikarpov
Phys. Rev. D68, 094018 (2003).

Measuring the color fields on the lattice

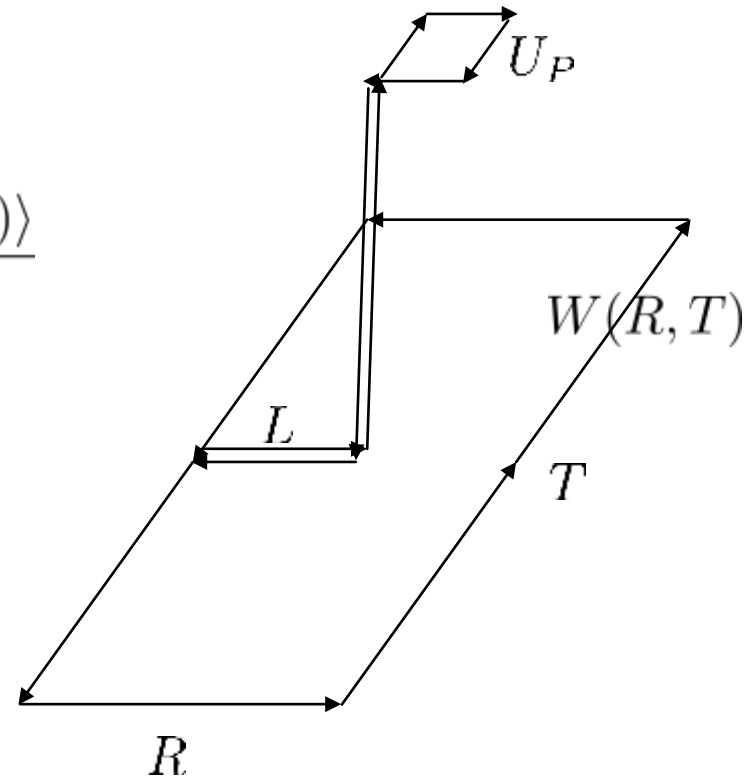
- Connected correlator

$$\rho_W = \frac{\langle \text{Tr}(W L U_P L^\dagger) \rangle}{\langle \text{Tr}(W) \rangle} - \frac{1}{N} \frac{\langle \text{Tr}(U_P) \text{Tr}(W) \rangle}{\langle \text{Tr}(W) \rangle}$$

W : Wilson loop

L : Schwinger line

U_P : Plaquette

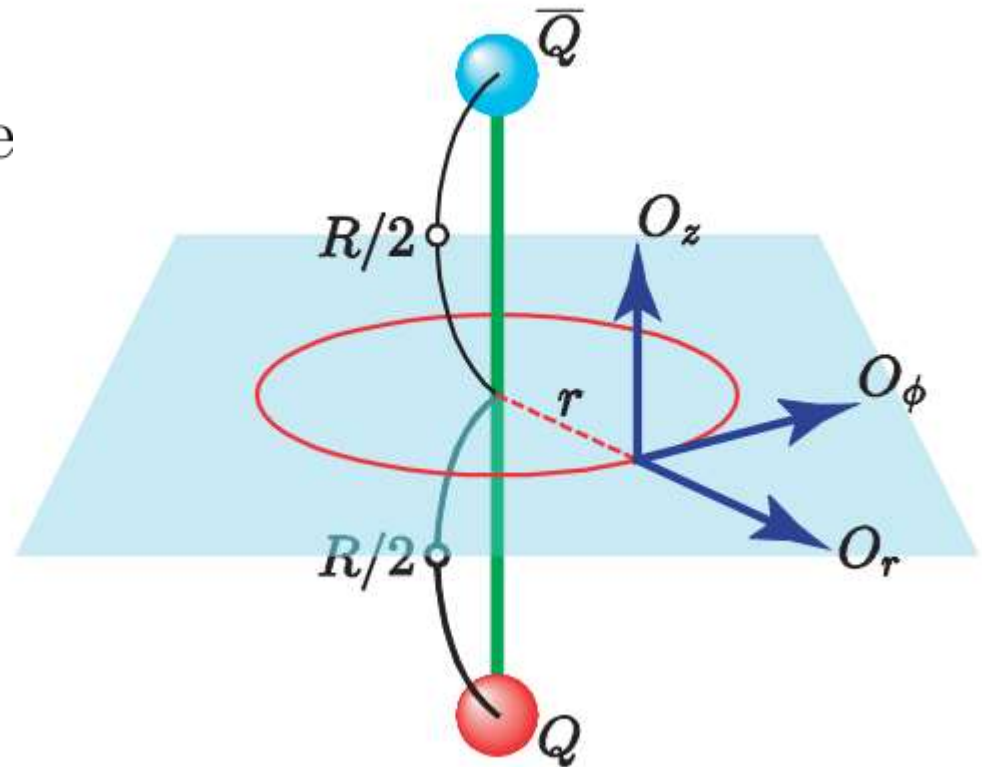


A. Di Giacomo, M. Maggiore, S. Olejnik, Nucl.Phys. B347 (1990) 441,
P. Cea, L. Cosmai, Phys.Rev. D52 (1995) 5152

Color flux tube

- Cylindrical coordinate system (r, ϕ, z)

- r : radial distance
- ϕ : azimuthal angle
- z : axial coordinate



Profile of E_z

- Exponential form (London theory)

$$E_z(r) = E_z(0) \exp\left(-\frac{r}{\lambda}\right)$$

- λ : London penetration length characteristic length of the exponential decrease of (electric) field in (dual) superconductor
- K_0 form (London theory for vortex as a line singularity : Abrikosov vortex)

$$E_z(r) = \frac{\phi}{2\pi\lambda^2} K_0\left(\frac{r}{\lambda}\right), \quad \lambda \gg \xi \quad r \rightarrow \infty \quad E(r) \propto \frac{1}{\sqrt{r}} \exp\left(-\frac{r}{\lambda}\right)$$

$K_n(r)$: n-th order modified
Bessel function of the second kind

$$r \rightarrow 0 \quad E(r) \propto \log \frac{\lambda}{r}$$

P. Cea and L. Cosmai, Phys. Rev. D52 (1995) 5152

G. S. Bali, K. Schilling, C. Schlichter, Phys. Rev. D51 (1995) 5165

- Clem form (cylindrical vortex with core)

$$E_z(r) = \frac{\phi}{2\pi\lambda^2\alpha} \frac{K_0[(r^2/\lambda^2 + \alpha^2)^{\frac{1}{2}}]}{K_1[\alpha]}$$

$$\frac{1}{\alpha} = \frac{\lambda}{\xi_v} \quad \kappa = \frac{\lambda}{\xi} = \frac{\sqrt{2}}{\alpha} \sqrt{1 - \frac{K_0^2(\alpha)}{K_1^2(\alpha)}}$$

ϕ : external flux

λ : London penetration length

ξ_v : variational core-radius

$\kappa = \lambda/\xi$: Ginzburg-Landau parameter

J. R. Clem, J. Low Temp. Phys 18, 427 (1975)
P. Cea, L. Cosmai and A. Papa, Phys. Rev. D86(2012)

Lattice setup

- Lattice volume

16^4 lattice

- Parameters

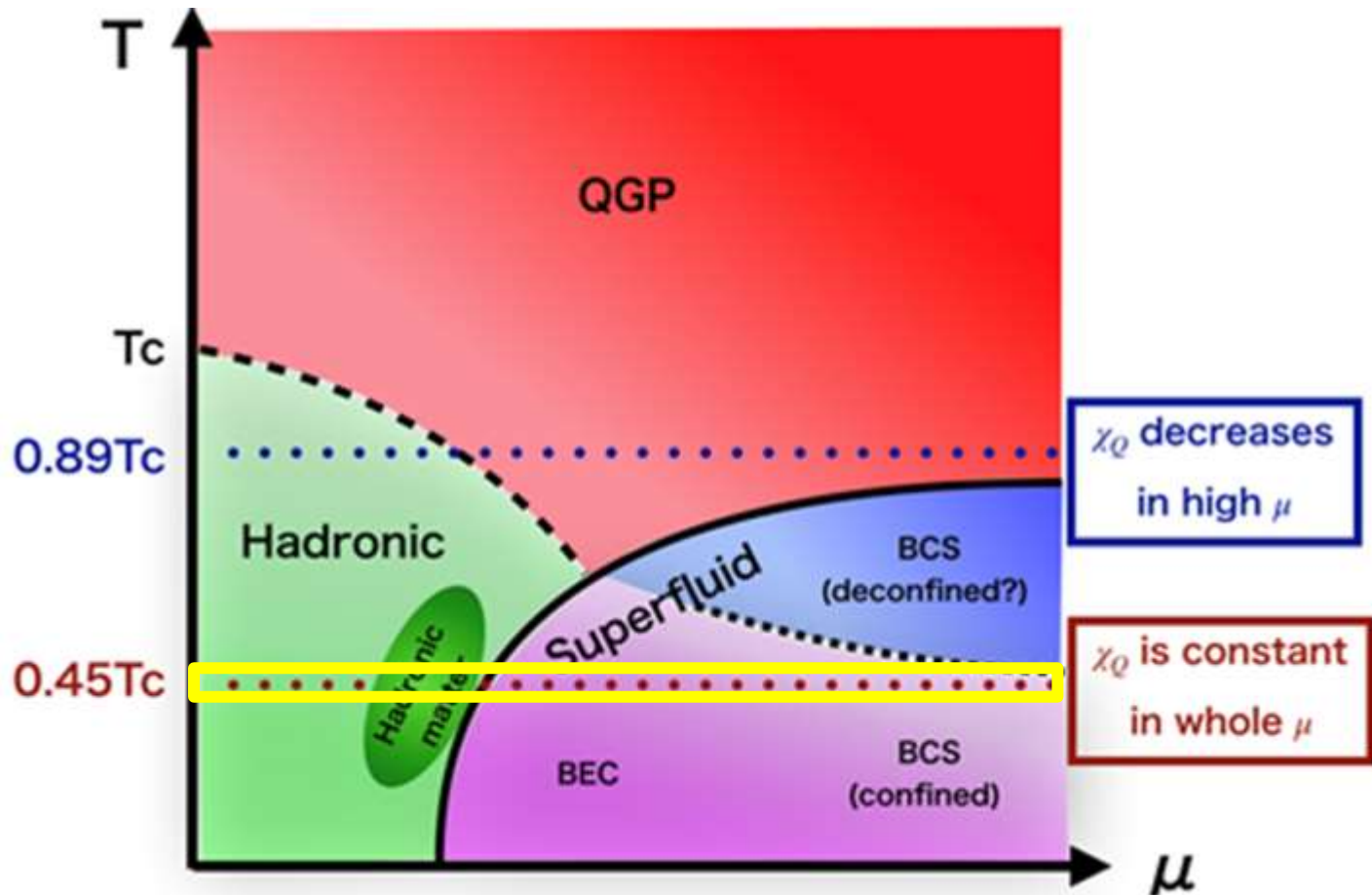
$$\beta = 0.800, \quad \kappa = 0.159, \quad T = 0.45T_c$$

$$m_{PS}/m_V = 0.823(9), \quad am_{PS} = 0.623(3),$$

(μ, j)	μ / m_{PS}	phase
(0.00, 0.00)	0.00	Hadronic
(0.25, 0.00)	0.40	near border
(0.35, 0.01)	0.56	BEC
(0.50, 0.01)	0.80	near border
(0.70, 0.01)	1.12	BCS

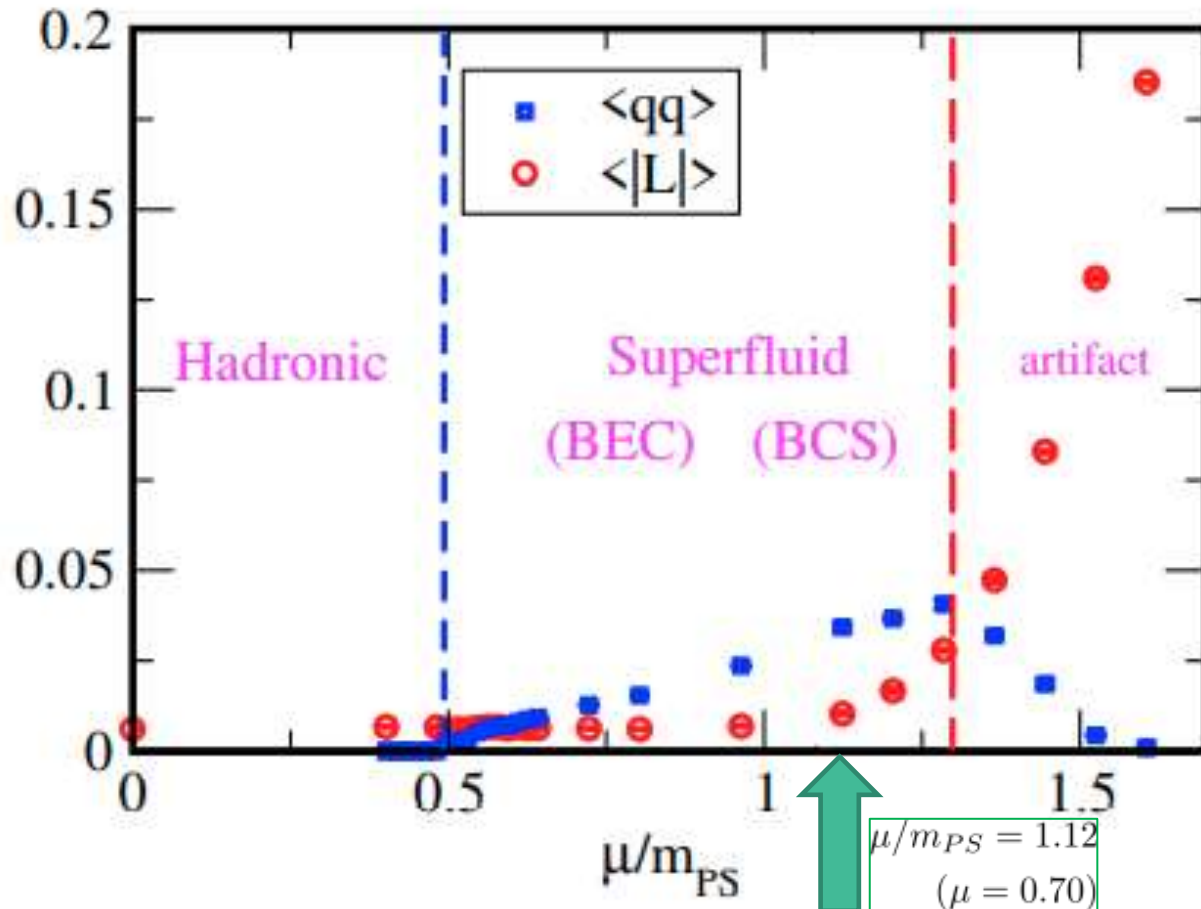
- Smoothing of the gauge configurations : APE smearings for spatial links , one hypercubic smearing for temporal links

The area used in this work



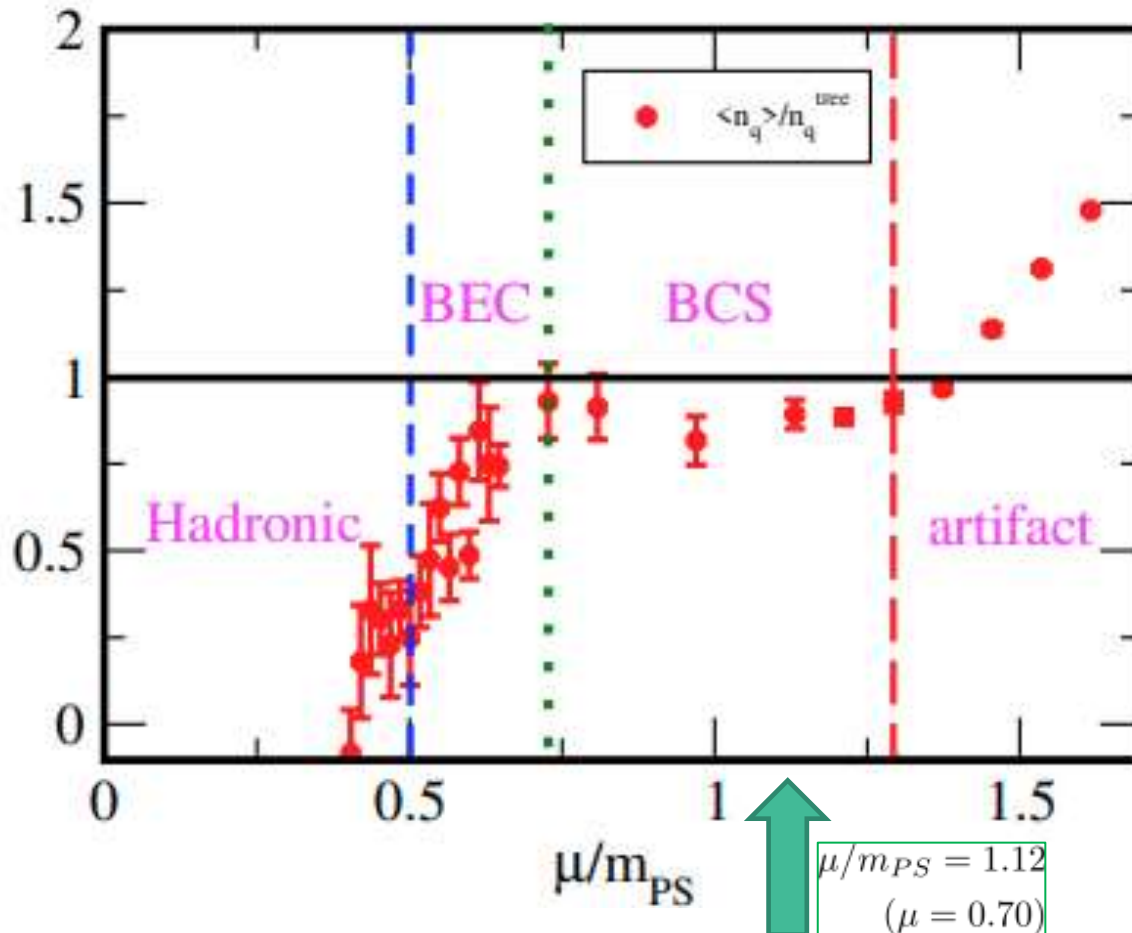
Polyakov loop and diquark condensate

- Confinement phase in all μ except for artifact region



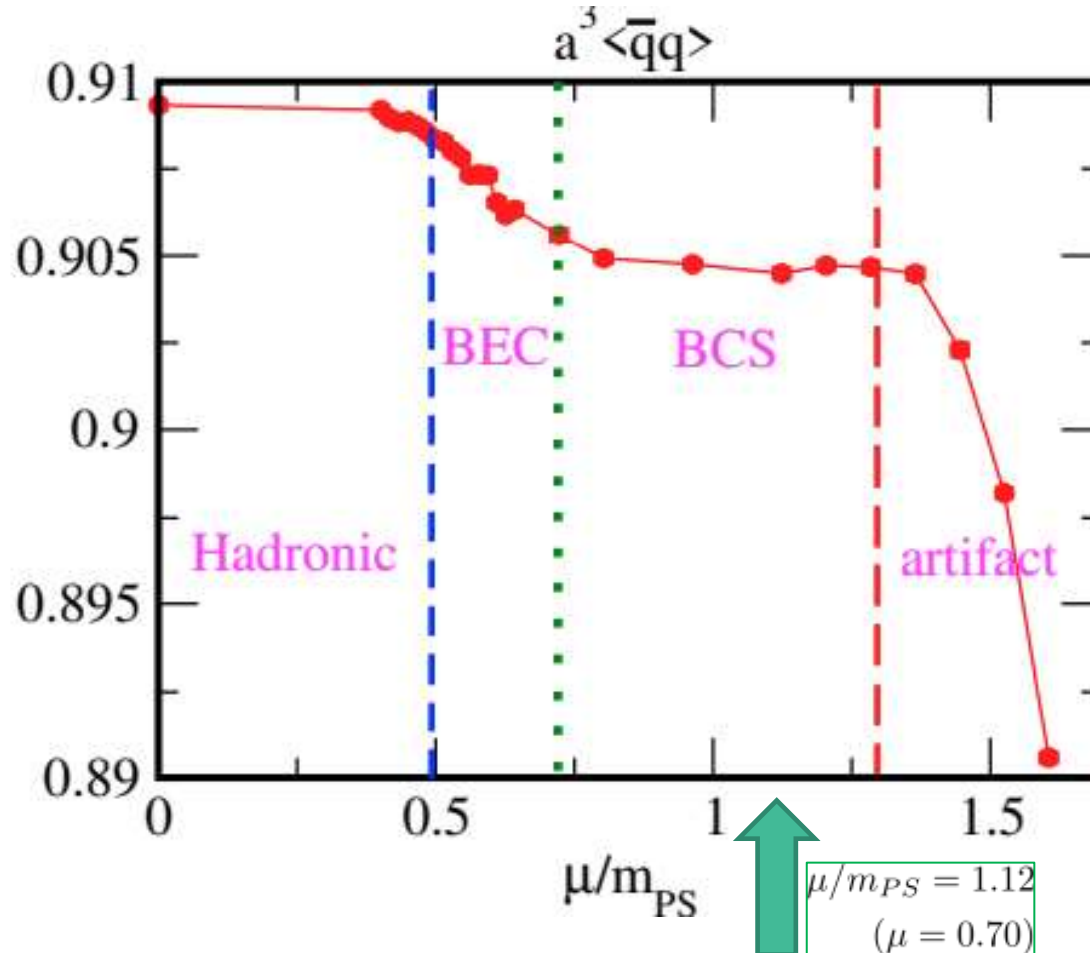
T=0.45T_c

Quark number density



$T=0.45T_c$

Chiral condensate

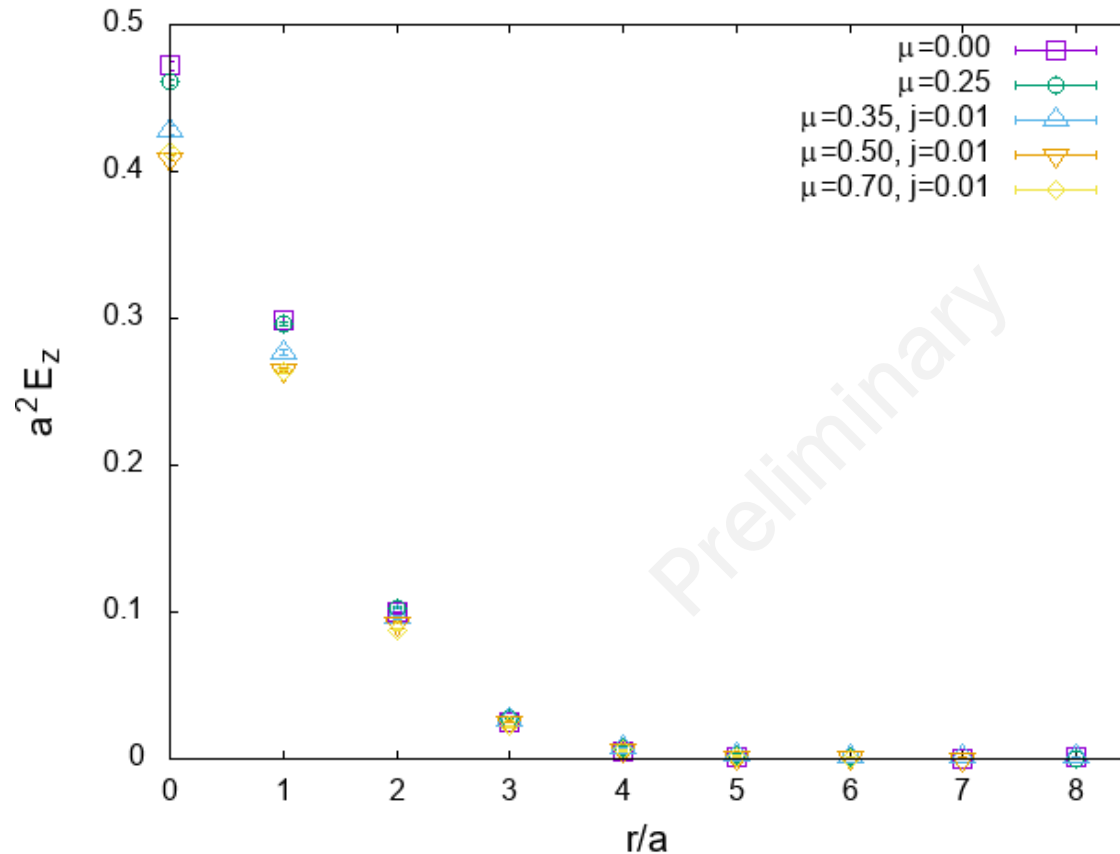


$T=0.45T_c$

Results (preliminary)

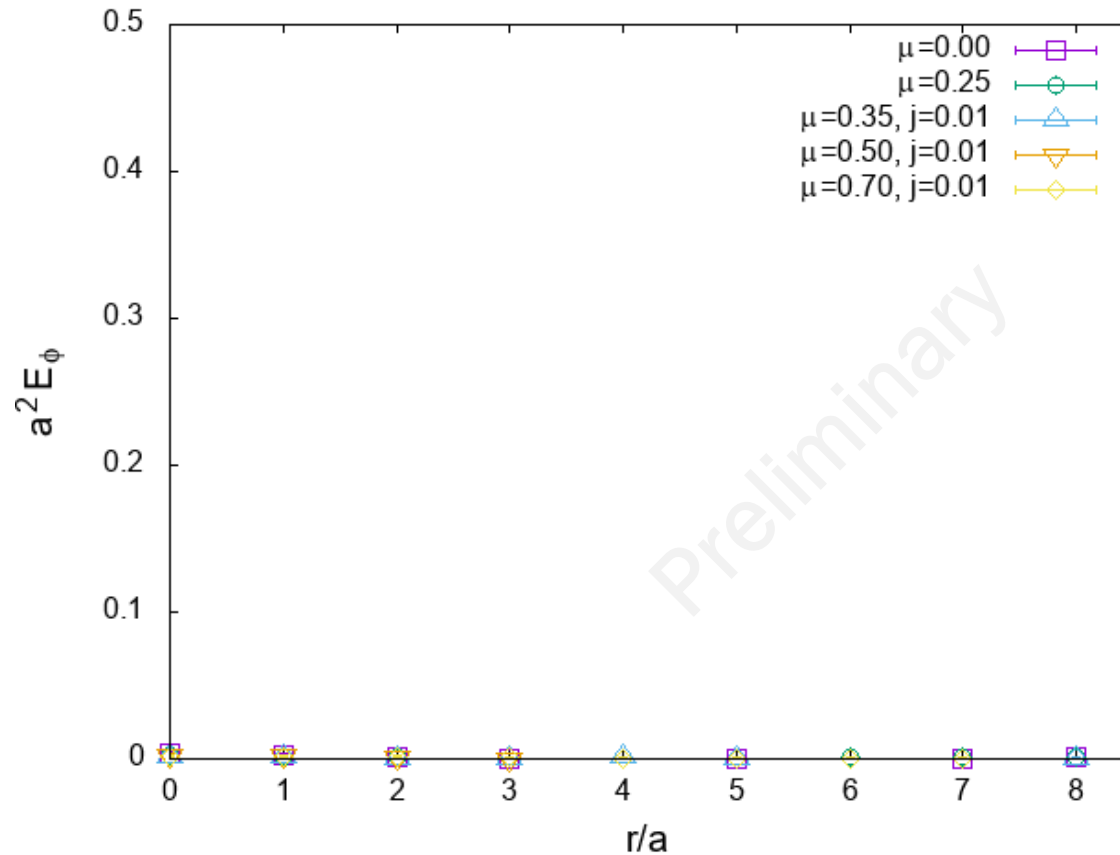
- μ dependence of color electric fields
- Penetration length
 - Exponential form
 - K_0 form
 - Clem form
- GL parameter
- Static potential between quark and antiquark

Ez (color electric field)



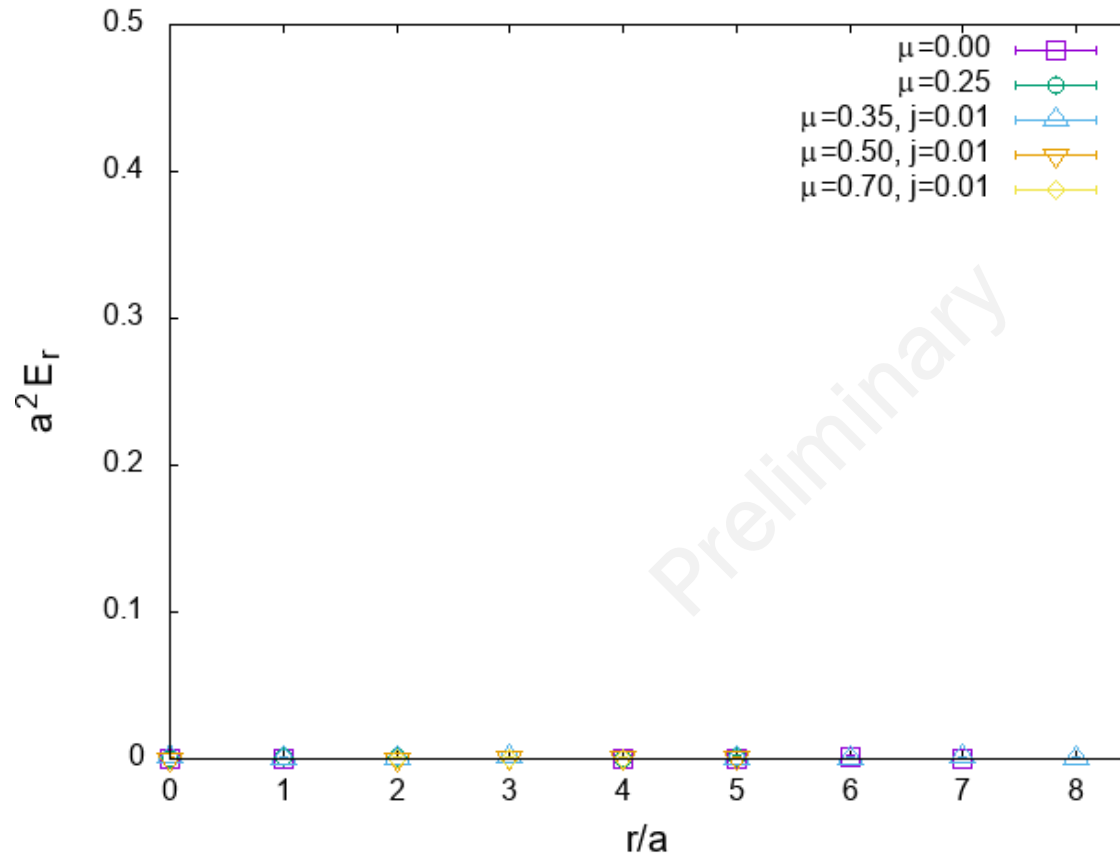
- Flux tubes are formed and there is a slight density dependence.

E_ϕ



- The azimuthal components of the color electric field are almost zero.

Er



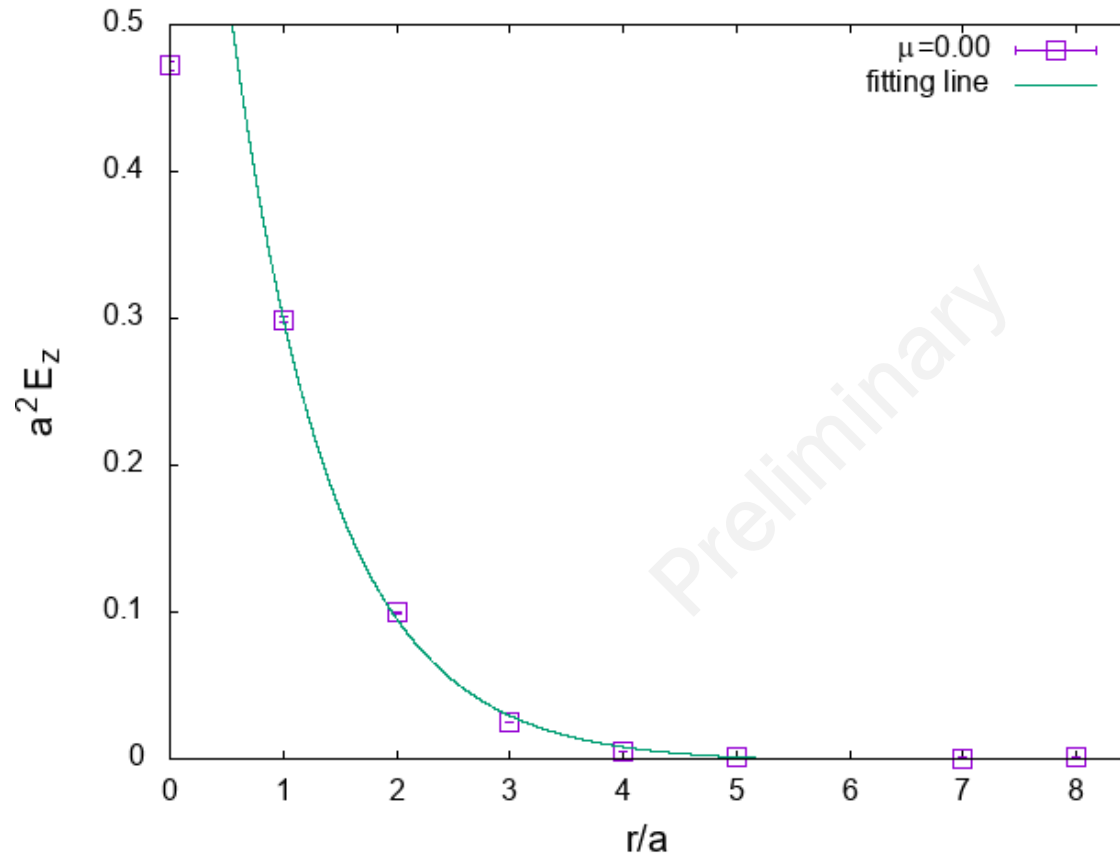
- The radial components of the color electric field are almost zero.

-
- E_ϕ and E_r are zero and only E_z has a non-zero value in all μ region.
 - The color electric field is squeezed into a linear shape to form a flux tube.



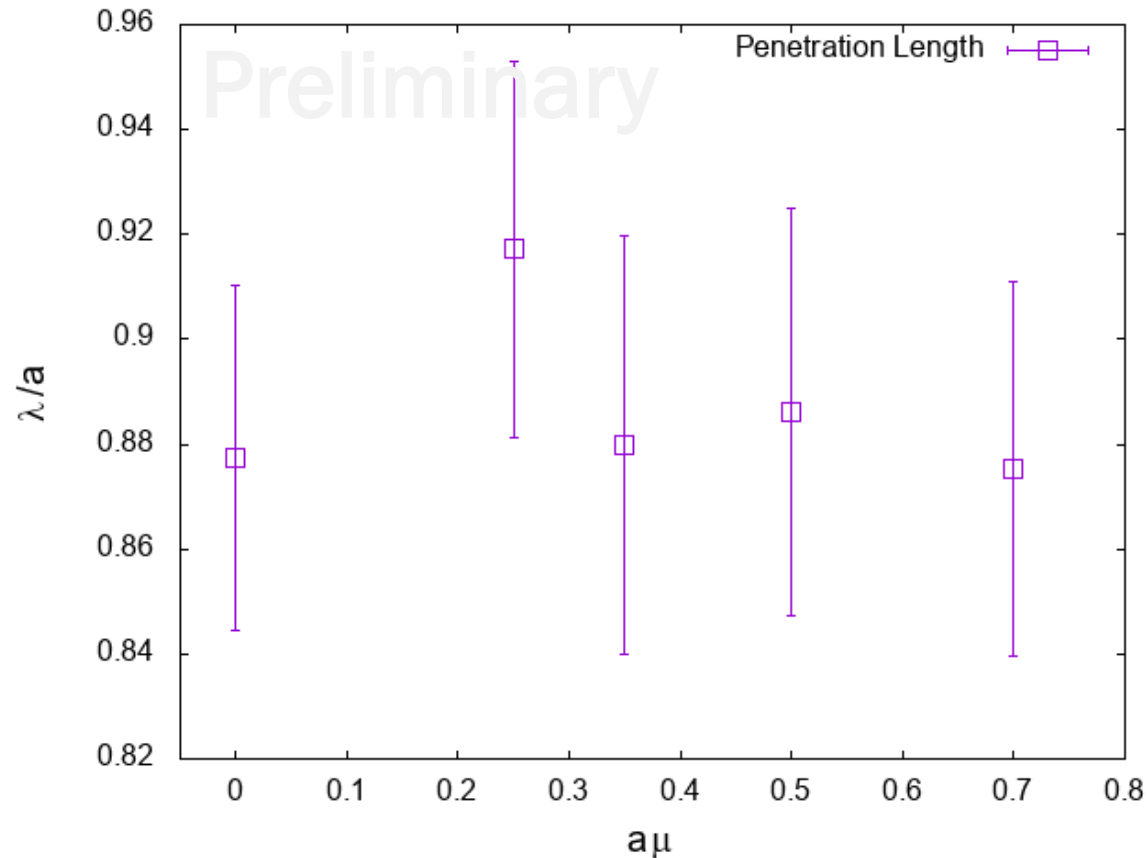
- Investigate the electric field E_z by doing a fit using the three forms described above.

Exponential fitting of E_z $E_z(r) = E_z(0) \exp\left(-\frac{r}{\lambda}\right)$



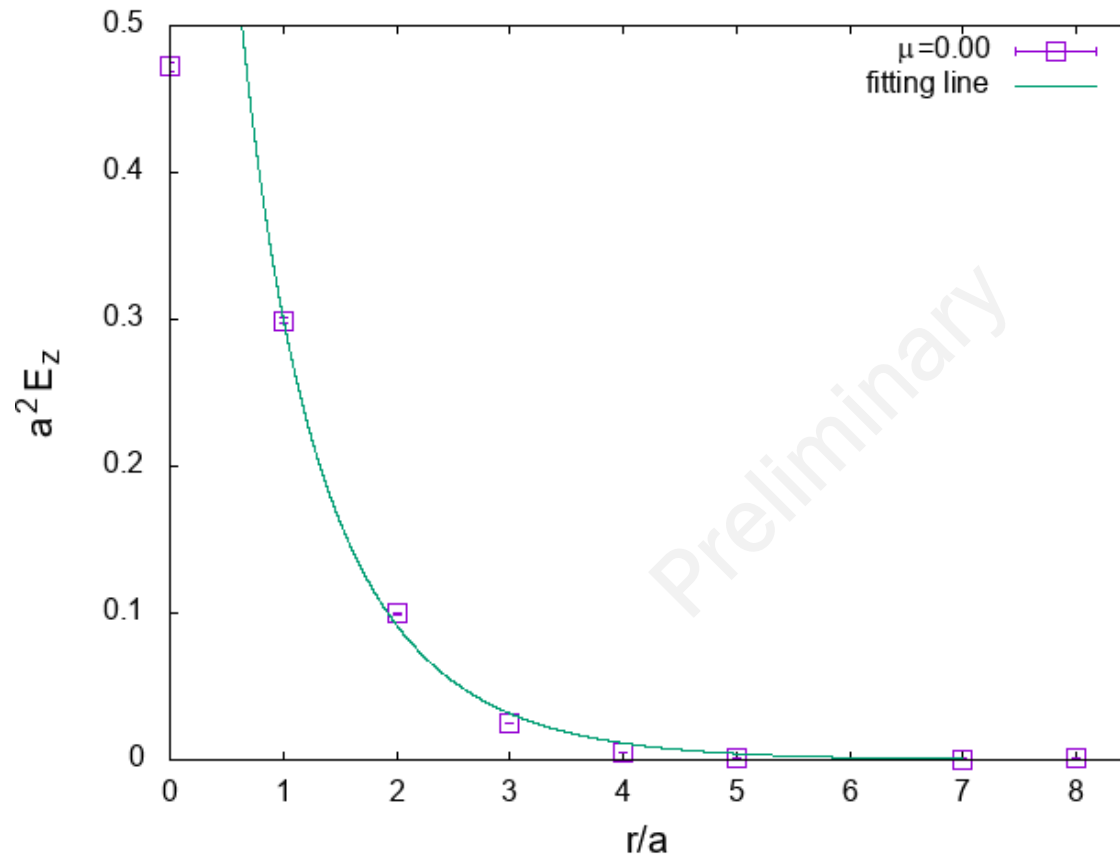
- The exponential fitting works, except for $r=0$.

Penetration length (Exponential fitting)



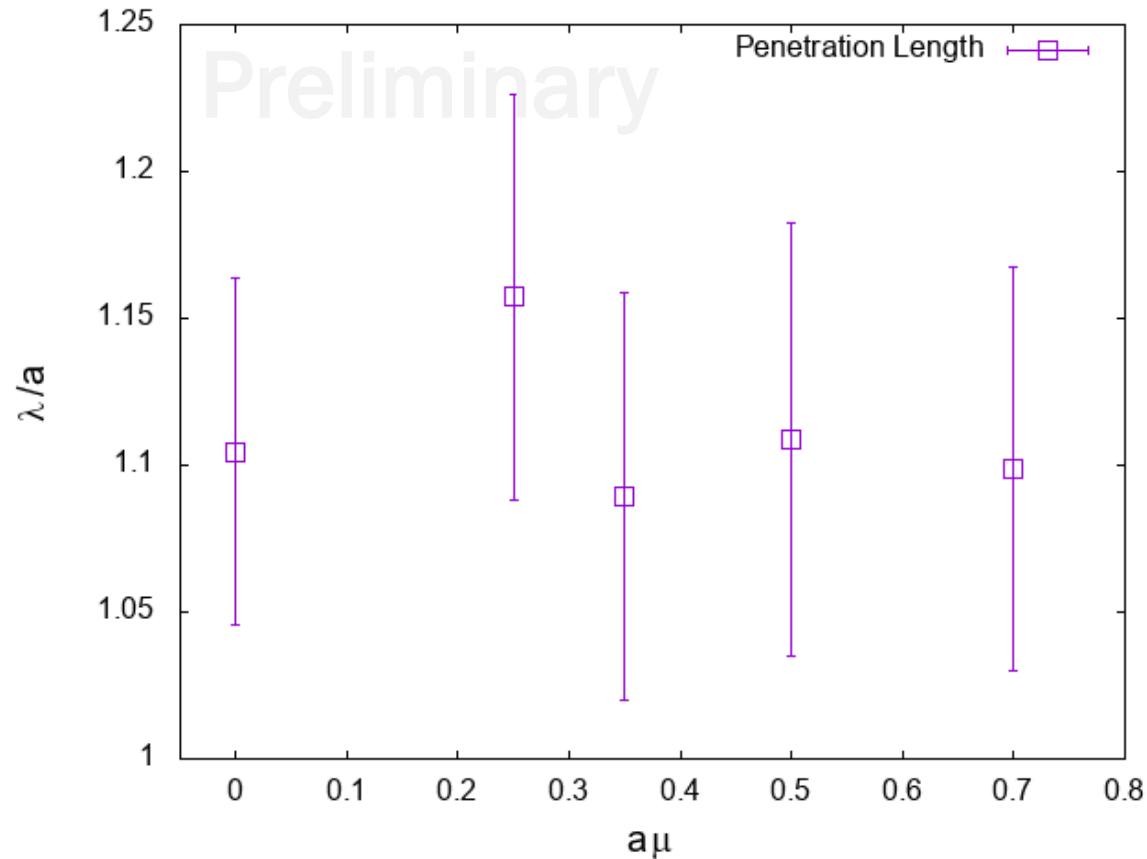
- No density dependence on penetration length

K_0 form fitting $E_z(r) = \frac{\phi}{2\pi\lambda^2} K_0\left(\frac{r}{\lambda}\right)$



- The K_0 form fitting works, except for $r=0$.

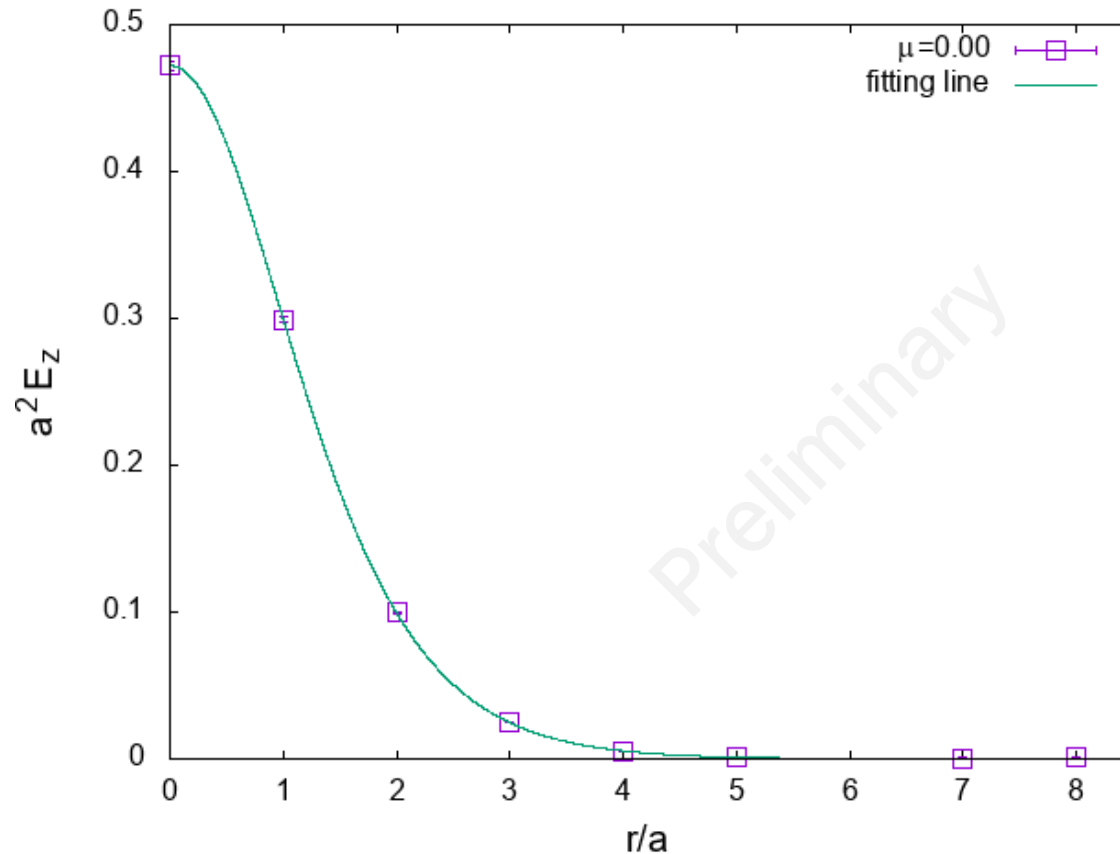
Penetration length (K_0 form fitting)



- No density dependence on penetration length

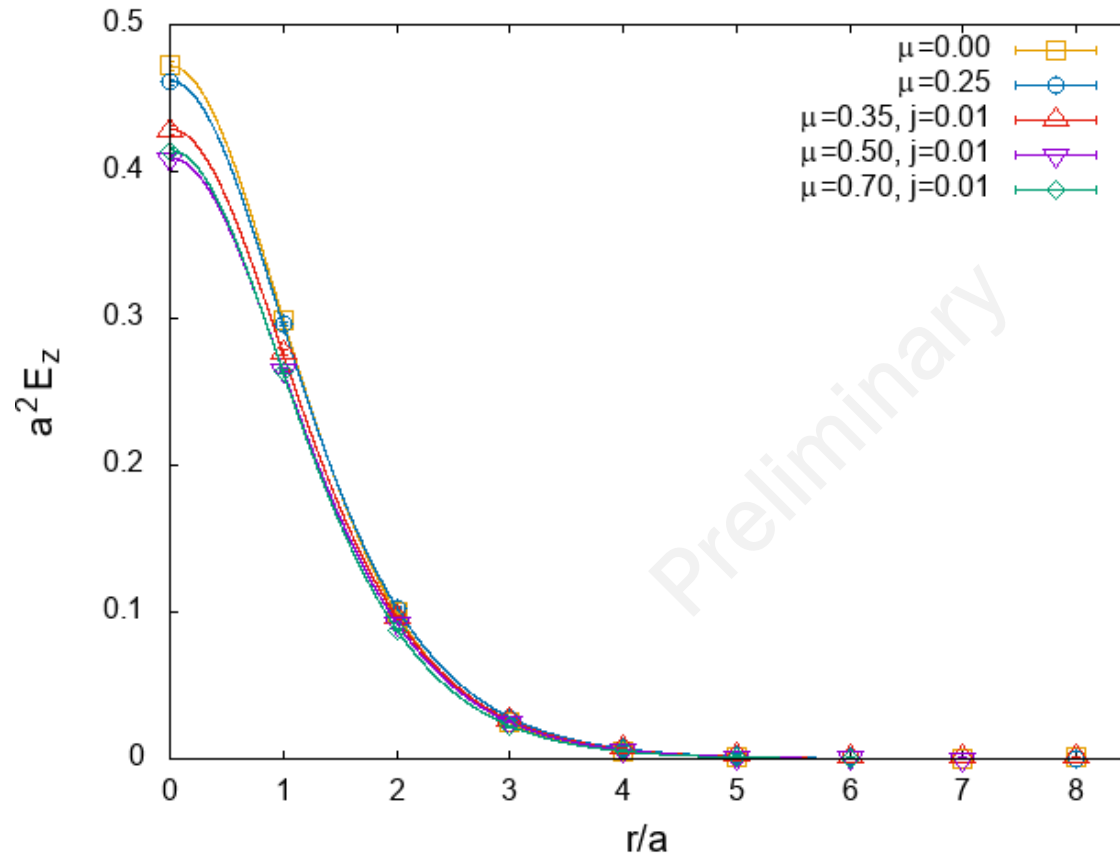
Clem form

$$E_z(r) = \frac{\phi}{2\pi\lambda^2\alpha} \frac{K_0[(r^2/\lambda^2 + \alpha^2)^{1/2}]}{K_1[\alpha]}$$



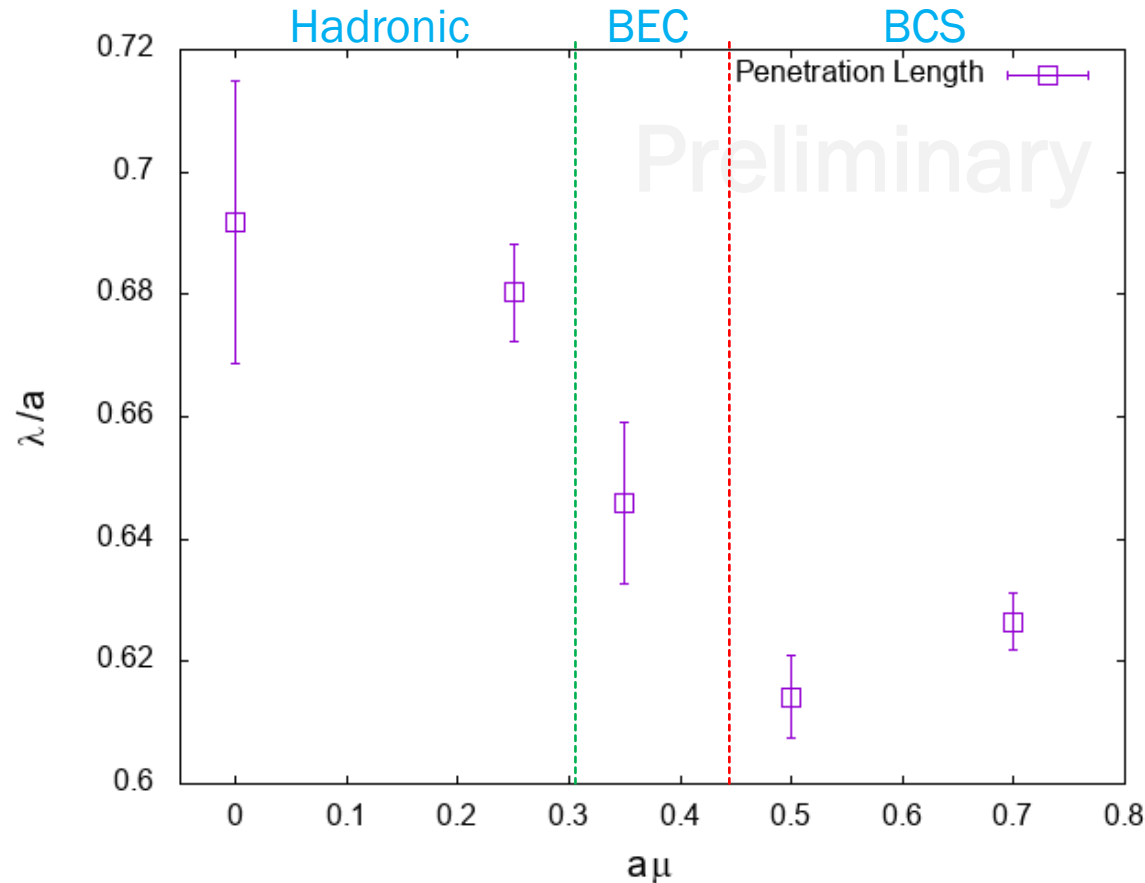
- Fitting by the Clem form is working, including $r=0$.

Density dependence of E_z



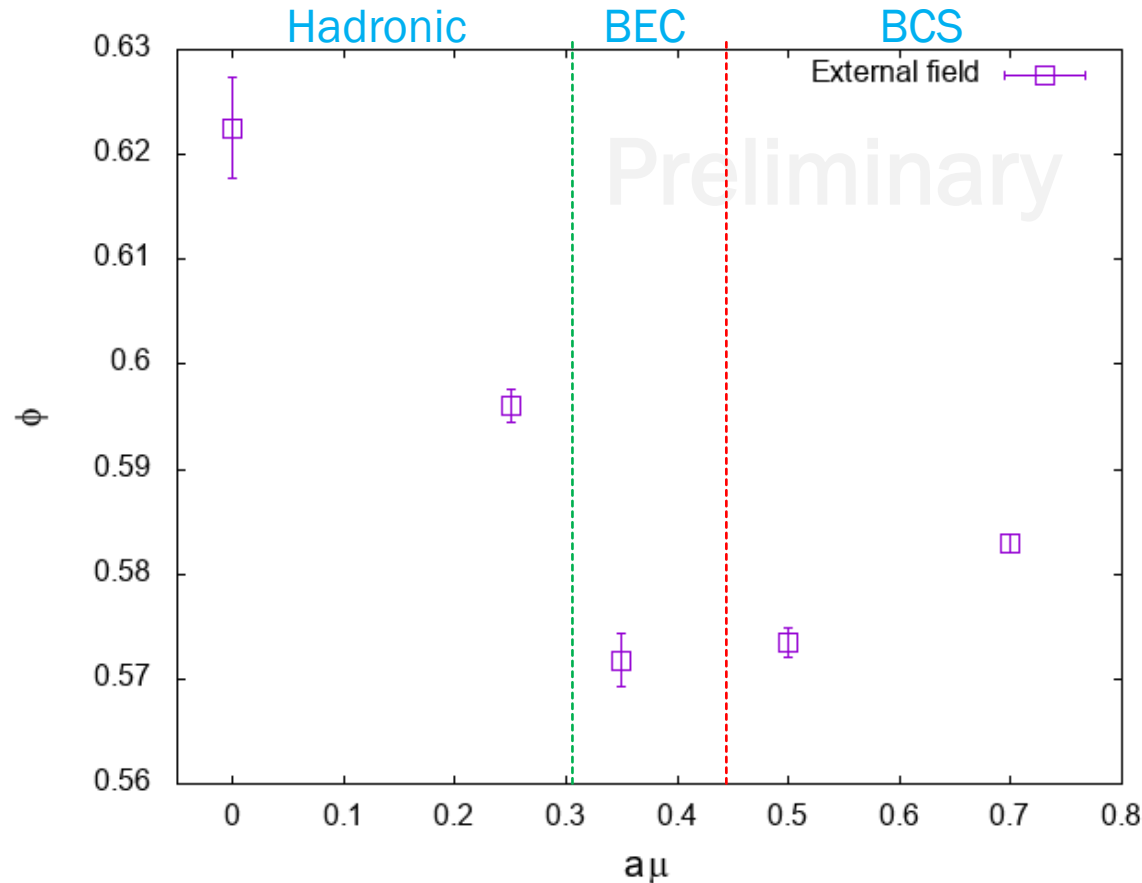
- Fittings by the Clem form is working in all μ region.

Penetration length



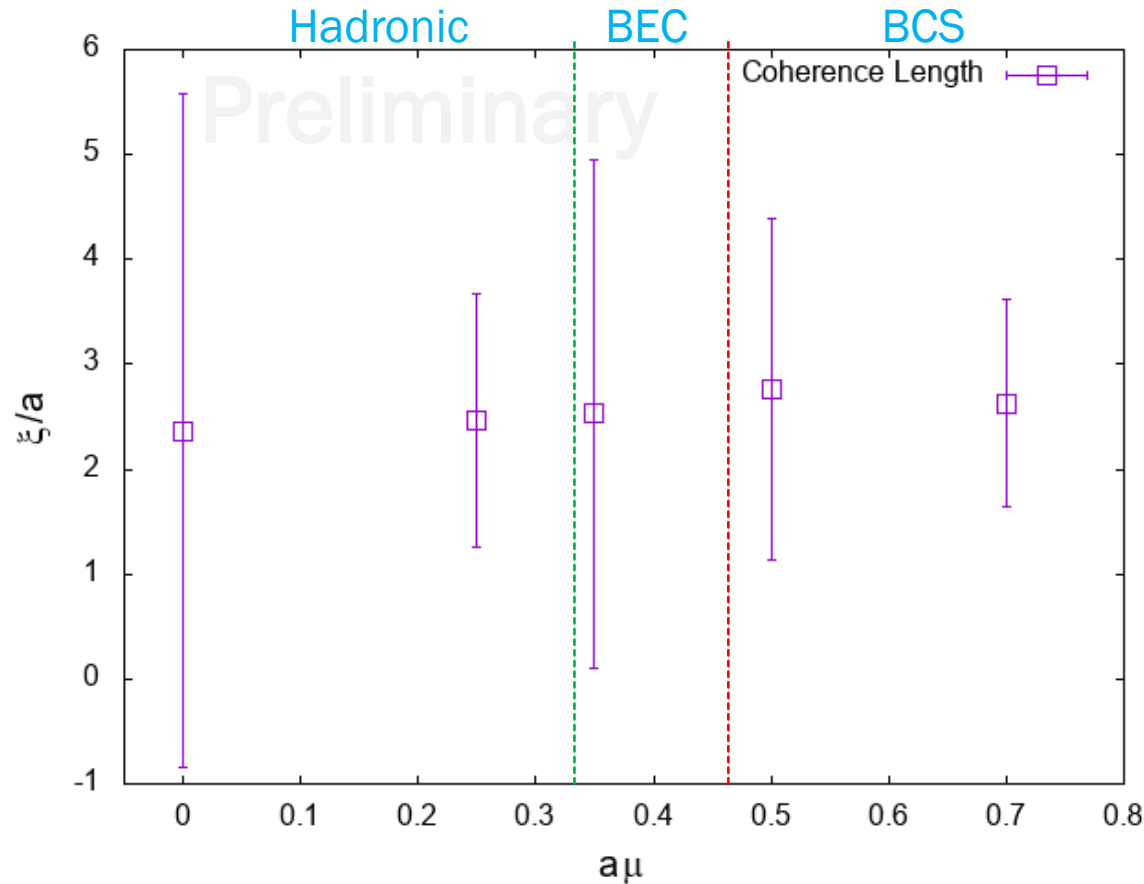
- The penetration length fitted by the Clem form shows a density dependence. The penetration length gradually decreases with μ .

External field Φ



- External field fitted in the Clem form shows a density dependence. Φ decreases until $\mu=0.35$ and appears to remain constant thereafter.

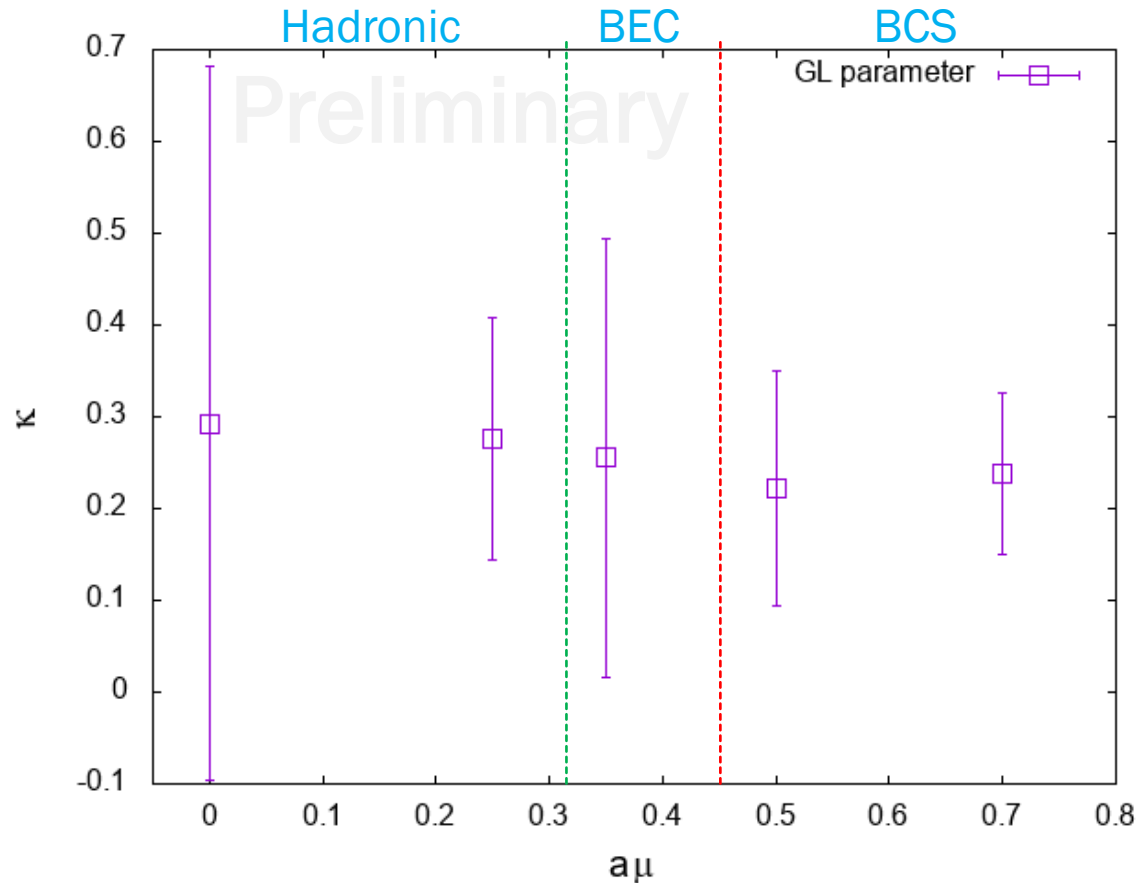
Coherence length



- The values of coherence length are almost constant in all μ region, although the error bars are large.

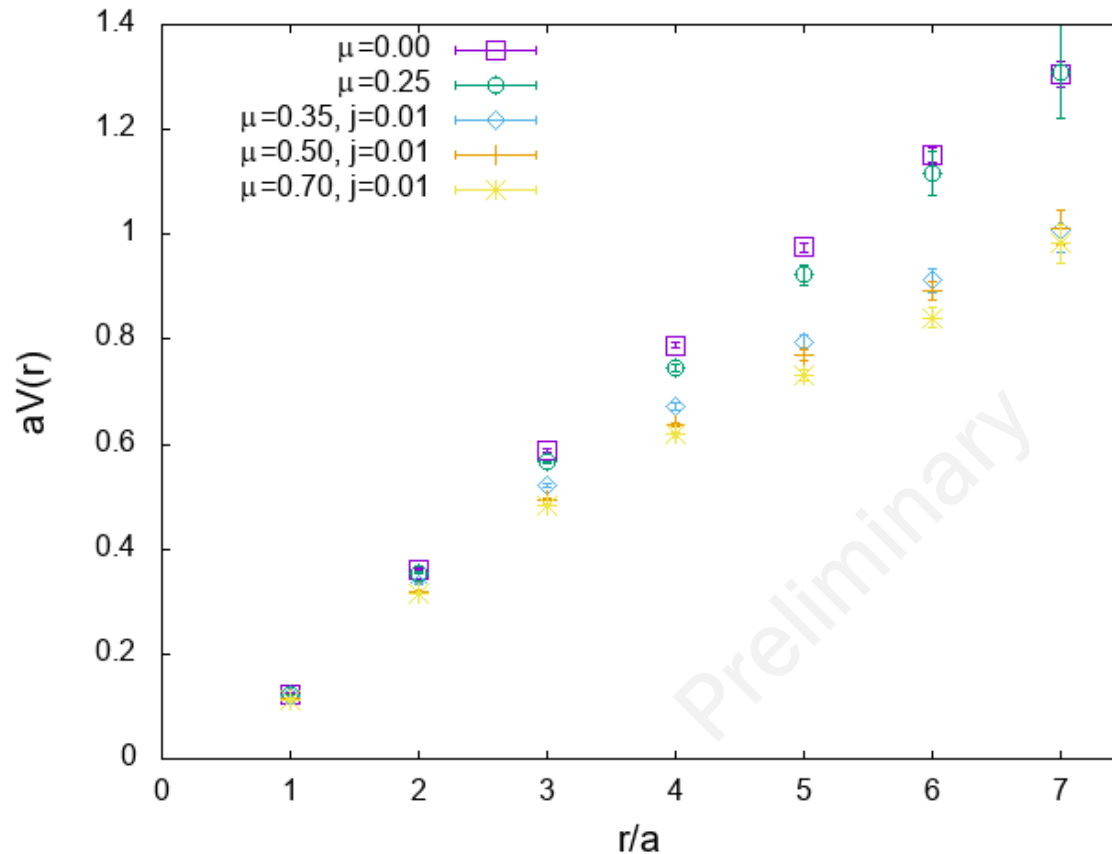
Ginzburg-Landau parameter

$$\kappa = \lambda/\xi$$



- Although the error bars are large, the values of Ginzburg-Landau parameter show almost constant or slightly decreasing with μ . This result suggests that the vacuum of two-color QCD is a type 1 dual superconductor.

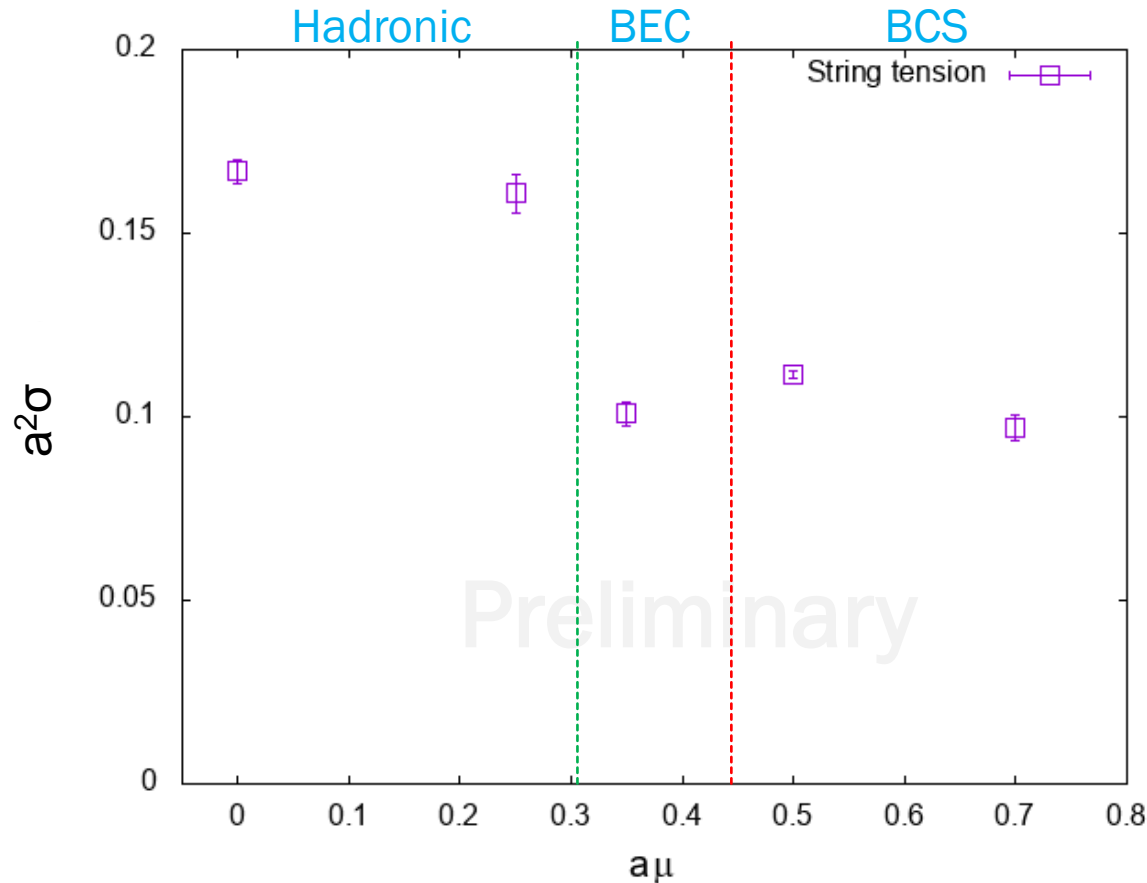
Static potential



- Static potentials from Wilson loop operator and their density dependence .

String tension

$$V(r) = \sigma r + \alpha/r + c$$



- The value of string tension σ varies around $\mu=0.35$ and has a non-zero value even in the high μ region. This indicates that the system at $T=0.45T_c$ is in a confined phase for all μ region within the scope of this study.

Summary

- Calculate the μ dependence of the color electric fields in $N_f=2$ two-color QCD at $T=0.45T_c$.
- Penetration length of E_z
 - No μ dependence in exponential and K_0 form
 - Slightly μ dependence in the Clem form
- External flux
 - μ dependence
- GL parameter
 - No μ dependence or slightly μ dependence
 - Suggest type 1 dual superconductor

- Remain flux tube squeezing in all μ region at $T=0.45T_c$
- Non-zero string tension in all μ region
- The system at $T=0.45T_c$ is in a confined phase for all μ region within the scope of this study.
- Consistent for the previous results of Polyakov loop and topological susceptibility
- In future works
 - Large lattice
 - Different lattice spacings
 - Different temperatures ($0.89T_c, \dots$)