

Partial Breaking of $\mathcal{N} = 2$ Supersymmetry and Decoupling Limit of Nambu-Goldstone Fermion in $U(N)$ Gauge Model

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The fermionic shift symmetry is a key point of the proof of the Dijkgraaf-Vafa conjecture which assert that non-perturbative quantities in a $\mathcal{N} = 1$ supersymmetric gauge theory can be computed by a matrix model . Thanks to this symmetry, effective superpotential is written as $W_{\text{eff}} = \int d^2\chi \mathcal{F}_p$, for some function \mathcal{F}_p which is related to the free energy of the matrix model. The fermionic shift symmetry is due to a free fermion and is expected to come from a second, spontaneously broken supersymmetry. The $U(N)$ gauge model which breaks $\mathcal{N} = 2$ supersymmetry to $\mathcal{N} = 1$ spontaneously was studied in [1]. The Nambu-Goldstone fermion appears in the overall $U(1)$ part of $U(N)$ gauge group and is coupled to the $SU(N)$ sector because of the fact that the 3rd derivatives of the prepotential \mathcal{F} are non-vanishing .

In [2], We examine a decoupling limit of the Nambu-Goldstone fermion. How can we take such limit with partial breaking of supersymmetry ($\mathcal{N} = 2 \rightarrow 1$) in $U(N)$ gauge model? In order to decouple the Nambu-Goldstone fermion, we should make the prepotential \mathcal{F} be a second order polynomial. However the order of the prepotential is greater than or equal to 3 because of the condition for partial breaking of $\mathcal{N} = 2$ supersymmetry. This problem can be solved by a large limit of the parameters (e, m, ξ) of the electric and magnetic FI terms. Let us reparametrize

$$(e, m, \xi) = (\Lambda e', \Lambda m', \Lambda \xi'), \quad \mathcal{F} = \sum_{k=0}^n \text{tr} \frac{g_k}{k!} \Phi^k = \text{tr} \left(g_0 \mathbf{1} + g_1 \Phi + \frac{g_2}{2} \Phi^2 \right) + \frac{1}{\Lambda} \sum_{k=3}^n \text{tr} \frac{g'_k}{k!} \Phi^k$$

After taking the limit $\Lambda \rightarrow \infty$, the Nambu-Goldstone fermion is decoupled from other fields, while partial breaking of $\mathcal{N} = 2$ supersymmetry is realized as before. We get a general $\mathcal{N} = 1$ action which include the free fermion. It shows that the fermionic shift symmetry is due to the free Nambu-Goldstone fermion.

References

- [1] K. Fujiwara, H. Itoyama and M. Sakaguchi, “Supersymmetric $U(N)$ gauge model and partial breaking of $\mathcal{N} = 2$ supersymmetry,” *Prog. Theor. Phys.* **113** (2005) 429 [arXiv:hep-th/0409060]; “Partial breaking of $\mathcal{N} = 2$ supersymmetry and of gauge symmetry in the $U(N)$ gauge model,” *Nucl. Phys. B* **723** (2005) 33 [arXiv:hep-th/0503113]; “Partial supersymmetry breaking and $\mathcal{N} = 2$ $U(N(c))$ gauge model with hypermultiplets in harmonic superspace,” *Nucl. Phys. B* **740** (2006) 58 [arXiv:hep-th/0510255].
- [2] K. Fujiwara, “Partial breaking of $\mathcal{N} = 2$ supersymmetry and decoupling limit of Nambu-Goldstone fermion in $U(N)$ gauge model,” arXiv:hep-th/0609039.