

Instanton Effects in Orbifold ABJM Theory

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1. Introduction
2. ABJM
3. Generalizations

Message

The Exact Instanton Expansion
of The ABJM Partition Function
Has Been Found!

Caution:

Not Solved Completely Logically
(Partial Assist From Numerical Method)

Why ABJM Interesting?

ABJM Theory Describes Multiple M2-branes on

$$\mathbb{C}^4/\mathbb{Z}_k \rightarrow \mathbb{R}_{>0} \times S^1 \times \mathbb{C}\mathbb{P}^3 \quad (k \rightarrow \infty)$$

- $\log(\text{Partition Function}) = \text{M2 DOF} = N^{3/2} + \dots$
[Drukker-Marino-Putrov]
- Understanding **Membrane** Interactions
- A prototype of **AdS/CFT**

Why ABJM Exact PF Interesting?

- A completion of $N^{3/2}$ DOF
(Understanding Membrane Interaction)
- Perturbative Corrections Sum Up To Airy
 $Z(N) \sim \text{Ai}[N] \sim \exp(N^{3/2})$ [Fuji-Hirano-M]
- Poles from Coefficients of Two Instantons
Cancel Between Themselves [Hatsuda-M-Okuyama]
- Completely by Refined Topological Strings
[Hatsuda-Marino-M-Okuyama]

Question

So Far So Good for ABJM

- How General, Properties of ABJM?
 - More Instantons, More Cancellations?
 - Exact Expansion for General Theories?

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ABJM Matrix Model

- Partition Function of ABJM Theory
- Due to SUSY, Localized to Matrix Model

$$Z(N) = \frac{1}{N_1! N_2!} \int \prod_{i=1}^{N_1} \frac{d\mu_i}{2\pi} \prod_{k=1}^{N_2} \frac{d\nu_k}{2\pi} e^{-(\sum \mu_i^2 - \sum \nu_k^2)/2g_s}$$
$$\times \frac{\prod_{i < j} \left(2 \sinh \frac{\mu_i - \mu_j}{2} \right)^2 \prod_{k < l} \left(2 \sinh \frac{\nu_k - \nu_l}{2} \right)^2}{\prod_{i,k} \left(2 \cosh \frac{\mu_i - \nu_k}{2} \right)^2}$$

[Pestun, Kapustin-Willett-Yaakov]

't Hooft Expansion

- $N \rightarrow \infty$ Expansion with $\lambda = N/k$ Fixed
- Genus Expansion [Drukker-Marino-Putrov]

$$\log[Z(N)] = N^2 F_0(\lambda) + N^0 F_1(\lambda) + N^{-2} F_2(\lambda) + N^{-4} F_3(\lambda) + \dots$$

$$= N^2 [F_0^{\text{pert}}(\lambda) + \# e^{-2\pi\sqrt{2}\lambda} + \# e^{-4\pi\sqrt{2}\lambda} + \dots]$$

$$+ N^0 [F_1^{\text{pert}}(\lambda) + \# e^{-2\pi\sqrt{2}\lambda} + \# e^{-4\pi\sqrt{2}\lambda} + \dots]$$

+

Sum Up To Airy Function!

[Fuji-Hirano-M, Honda et al]

Worldsheet Instanton

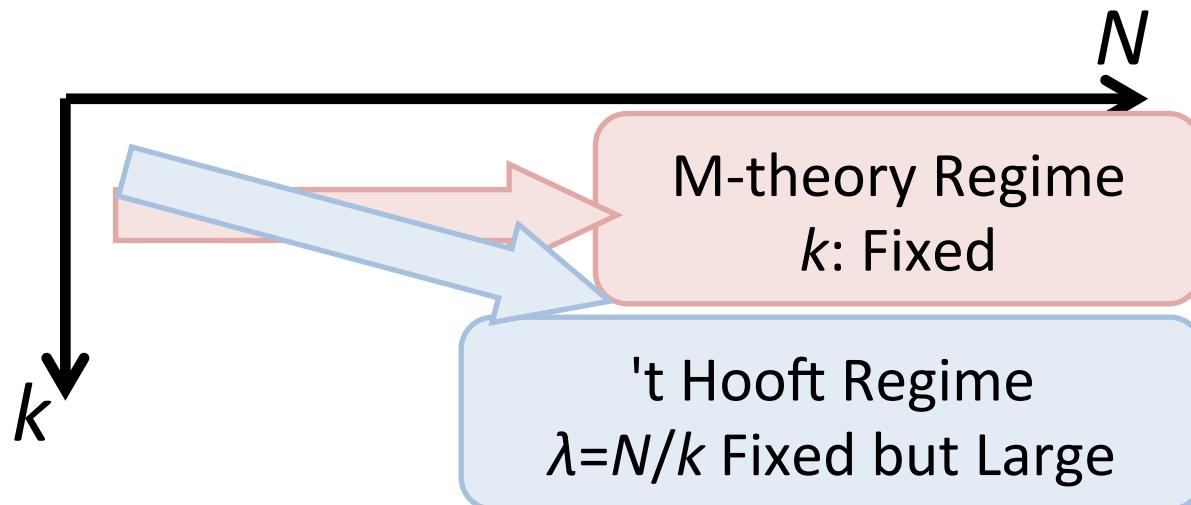
$$e^{-2\pi\sqrt{2}\lambda} = \exp[-T_{F1} \text{Area}(\mathbf{CP}^1)]$$

F-String Wrapping $\mathbf{CP}^1 \subset \mathbf{CP}^3$

[Cagnazzo-Sorokin-Wulff, DMP]

M-theory Expansion [Herzog-Klebanov-Pufu-Tesileanu]

- M-theory Background: C^4/Z_k
- $N \rightarrow \infty$ with k Fixed
- Close To 't Hooft Limit But Not Exactly



Fermi Gas Formalism [Marino-Putrov]

- Rewriting Partition Function
- Regarding as Fermi Gas with Hamiltonian e^{-H}

$$Z(N) = (N!)^{-1} \sum_{\sigma \in S(N)} (-1)^{\sigma} \int \prod_i dq_i \langle q_i | e^{-H} | q_{\sigma(i)} \rangle$$

$$e^{-H} \sim (2 \cosh q/2)^{-1} (2 \cosh p/2)^{-1}$$

$$[q, p] = i \cancel{2\pi k}$$

- Introducing Grand Potential

$$e^{J(\mu)} = \sum_{N=0}^{\infty} Z(N) e^{\mu N} = \det (1 + e^{\mu} e^{-H})$$

WKB (small k) expansion

$$\begin{aligned} J(\mu) &= k^{-1} J_0(\mu) + k^1 J_1(\mu) + k^3 J_2(\mu) + k^5 J_3(\mu) + \dots \\ &= k^{-1} [J_0^{\text{pert}}(\mu) + (\#\mu^2 + \#\mu + \#) e^{-2\mu} + (\dots) e^{-4\mu} + \dots] \\ &\quad + k [J_1^{\text{pert}}(\mu) + (\#\mu^2 + \#\mu + \#) e^{-2\mu} + (\dots) e^{-4\mu} + \dots] \\ &\quad + \dots \dots \\ &= (\#\mu^3 + \#\mu + \#) + (\#\mu^2 + \#\mu + \#) e^{-2\mu} + (\dots) e^{-4\mu} + \dots \end{aligned}$$

Airy Function Reproduced

[Marino-Putrov]

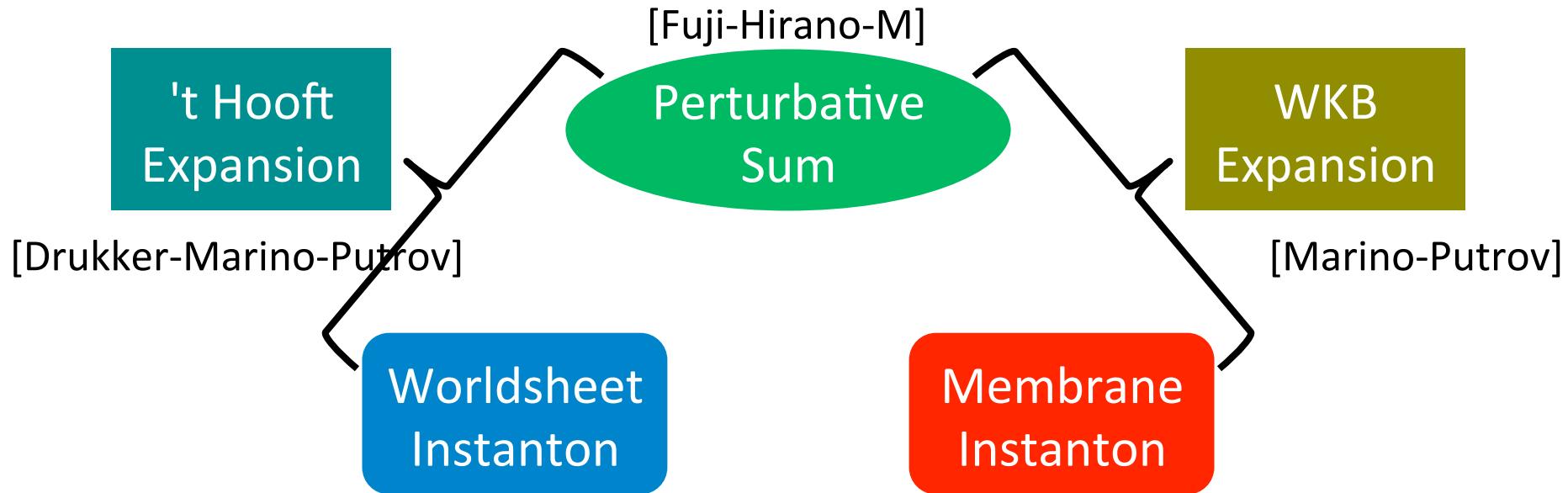
Membrane Instanton

$$e^{-2\mu} \approx e^{-\pi\sqrt{2Nk}} = \exp[-T_{D2} \text{Area}(\mathbf{RP}^3)]$$

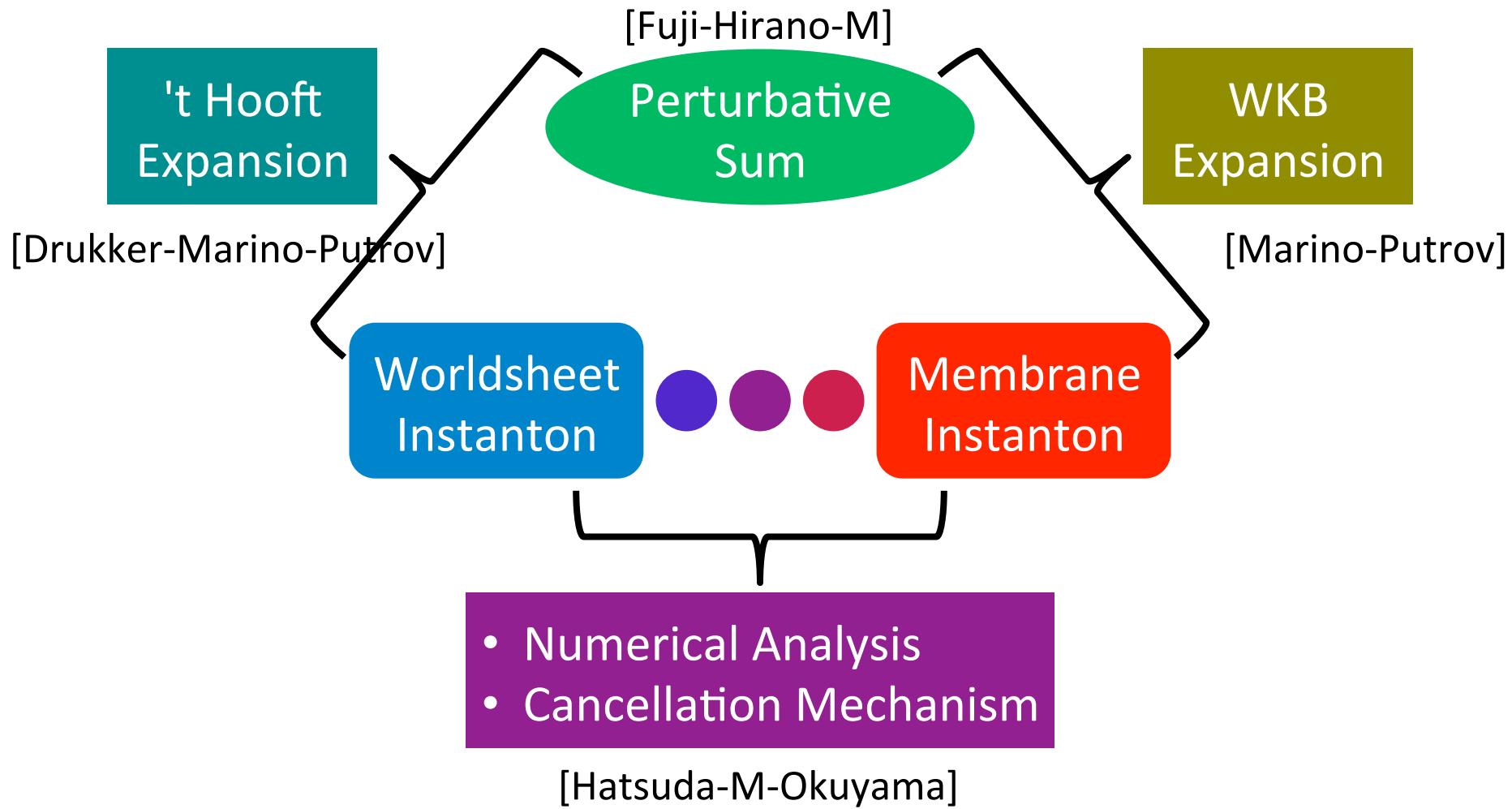
D2-brane Wrapping $\mathbf{RP}^3 \subset \mathbb{CP}^3$

[Drukker-Marino-Putrov]

Short Summary



Furthermore



Bound States [Hatsuda-M-Okuyama 1301]

MB

WS

Bound States $e^{-2/\mu - 4m\mu/k}$ Taken Care Of

By WS instantons with Chemical Potential Redef

$$\mu \rightarrow \mu_{\text{eff}}$$

$$\mu_{\text{eff}} = \mu + \# e^{-2\mu} + \# e^{-4\mu} + \dots$$

Cancelation Mechanism [Hatsuda-M-Okuyama 1211]

$$J_{k=1}(\mu) = [\# \mu^2 + \# \mu + \#] e^{-4\mu} + [\# \mu^2 + \# \mu + \#] e^{-8\mu} + [\# \mu^2 + \# \mu + \#] e^{-12\mu} + \dots$$

$$J_{k=2}(\mu) = [\# \mu^2 + \# \mu + \#] e^{-2\mu} + [\# \mu^2 + \# \mu + \#] e^{-4\mu} + [\# \mu^2 + \# \mu + \#] e^{-6\mu} + \dots$$

$$J_{k=3}(\mu) = [\#] e^{-4\mu/3} + [\#] e^{-8\mu/3} + [\# \mu^2 + \# \mu + \#] e^{-4\mu} + \dots$$

$$J_{k=4}(\mu) = [\#] e^{-\mu} + [\# \mu^2 + \# \mu + \#] e^{-2\mu} + [\#] e^{-3\mu} + \dots$$

$$J_{k=6}(\mu) = [\#] e^{-2\mu/3} + [\#] e^{-4\mu/3} + [\# \mu^2 + \# \mu + \#] e^{-2\mu} + \dots$$

WS(1)

WS(2)

WS(3)

↑
Free Energy of Topological String

Cancelation Mechanism [Hatsuda-M-Okuyama 1211]

(Raison d'Etre for M-theory!?)

$$\begin{aligned} J_{k=1}(\mu) &= [\# \mu^2 + \# \mu + \#] e^{-4\mu} + [\# \mu^2 + \# \mu + \#] e^{-8\mu} + [\# \mu^2 + \# \mu + \#] e^{-12\mu} + \dots \\ J_{k=2}(\mu) &= [\# \mu^2 + \# \mu + \#] e^{-2\mu} + [\# \mu^2 + \# \mu + \#] e^{-4\mu} + [\# \mu^2 + \# \mu + \#] e^{-6\mu} + \dots \\ J_{k=3}(\mu) &= [\#] e^{-4\mu/3} + [\#] e^{-8\mu/3} + [\# \mu^2 + \# \mu + \#] e^{-4\mu} + \dots \\ J_{k=4}(\mu) &= [\#] e^{-\mu} + [\# \mu^2 + \# \mu + \#] e^{-2\mu} + [\#] e^{-3\mu} + \dots \quad \text{MB(2)} \\ J_{k=6}(\mu) &= [\#] e^{-2\mu/3} + [\#] e^{-4\mu/3} + [\# \mu^2 + \# \mu + \#] e^{-2\mu} + \dots \end{aligned}$$

WS(1)

WS(2)

WS(3) MB(1)

↑
Free Energy of Topological String

All Explicitly In Topological Strings

[..., Hatsuda-M-Marino-Okuyama]

$$J(\mu) = J^{\text{pert}}(\mu^{\text{eff}}) + J^{\text{WS}}(\mu^{\text{eff}}) + J^{\text{MB}}(\mu^{\text{eff}})$$

$$J^{\text{pert}}(\mu) = C\mu^3/3 + B\mu + A$$

$$J^{\text{WS}}(\mu^{\text{eff}}) = F_{\text{top}}(T_1^{\text{eff}}, T_2^{\text{eff}}, \lambda)$$

$$J^{\text{MB}}(\mu^{\text{eff}}) = (2\pi i)^{-1} \partial_\lambda [\lambda F_{\text{NS}}(T_1^{\text{eff}}/\lambda, T_2^{\text{eff}}/\lambda, 1/\lambda)]$$

$$F_{\text{top}}(T_1, T_2, \tau) = \dots$$

$$F_{\text{NS}}(T_1, T_2, \tau) = \dots$$

$$C=2/\pi^2 k, B=\dots, A=\dots$$

$$T_1^{\text{eff}} = 4\mu^{\text{eff}}/k - i\pi$$

$$T_2^{\text{eff}} = 4\mu^{\text{eff}}/k + i\pi$$

$$\lambda = 2/k$$

$$\mu^{\text{eff}} = \begin{cases} \mu - (-1)^{k/2} 2e^{-2\mu} {}_4F_3(1, 1, 3/2, 3/2; 2, 2, 2; (-1)^{k/2} 16e^{-2\mu}) & k = \text{even} \\ \mu + e^{-4\mu} {}_4F_3(1, 1, 3/2, 3/2; 2, 2, 2; -16e^{-4\mu}) & k = \text{odd} \end{cases}$$

All Explicitly In Topological Strings

[..., Hatsuda-M-Marino-Okuyama]

$$J(\mu) = J^{\text{pert}}(\mu^{\text{eff}}) + J^{\text{WS}}(\mu^{\text{eff}}) + J^{\text{MB}}(\mu^{\text{eff}})$$

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$F(T_1, T_2, \tau_1, \tau_2)$: Free Energy

of Refined Top Strings

T_1, T_2 : Kahler Moduli

τ_1, τ_2 : Coupling Constants

Topological Limit $F_{\text{top}}(T_1, T_2, \tau) = \lim_{\tau_1 \rightarrow \tau, \tau_2 \rightarrow -\tau} F(T_1, T_2, \tau_1, \tau_2)$

NS Limit $F_{\text{NS}}(T_1, T_2, \tau) = \lim_{\tau_1 \rightarrow \tau, \tau_2 \rightarrow 0} 2\pi i \tau_2 F(T_1, T_2, \tau_1, \tau_2)$

All Explicitly In Topological Strings

[..., Hatsuda-M-Marino-Okuyama]

$$J(\mu) = J^{\text{pert}}(\mu^{\text{eff}}) + J^{\text{WS}}(\mu^{\text{eff}}) + J^{\text{MB}}(\mu^{\text{eff}})$$

$$J^{\text{pert}}(\mu) = C\mu^3/3 + B\mu + A$$

$$J^{\text{WS}}(\mu^{\text{eff}}) = F_{\text{top}}(T_1^{\text{eff}}, T_2^{\text{eff}}, \lambda)$$

$$J^{\text{MB}}(\mu^{\text{eff}}) = (2\pi i)^{-1} \partial_\lambda [\lambda F_{\text{NS}}(T_1^{\text{eff}}/\lambda, T_2^{\text{eff}}/\lambda, 1/\lambda)]$$

$$\begin{aligned} F(T_1, T_2, \tau_1, \tau_2) = & \sum_{j_L, j_R} \sum_n \sum_{d_1, d_2} N_{j_L, j_R}^{d_1, d_2} \\ & \chi_{j_L}(q_L) \chi_{j_R}(q_R) e^{-n(d_1 T_1 + d_2 T_2)} \\ & / [n(q_1^{n/2} - q_1^{-n/2})(q_2^{n/2} - q_2^{-n/2})] \end{aligned}$$

$$q_1 = e^{2\pi i \tau_1} \quad q_2 = e^{2\pi i \tau_2} \quad q_L = e^{\pi i (\tau_1 - \tau_2)} \quad q_R = e^{\pi i (\tau_1 + \tau_2)}$$

$N_{j_L, j_R}^{d_1, d_2}$: BPS Index of local $\mathbb{P}^1 \times \mathbb{P}^1$
(Gopakumar-Vafa or Gromov-Witten invariants)

Looking Back, What Makes ABJM Solvable?

- Max SUSY ($N=6$) !?
- Relations to Topological Strings?

In Fact, ABJ ($N=6$ & Topological Strings)

Solved Similarly.

[Awata-Hirano-Shigemori, Honda, Matsumoto-M, Honda-Okuyama, Kallen]

Generalizations

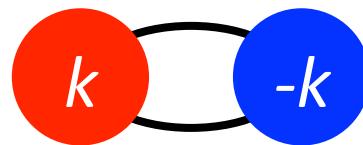
- ABJM only Two Types of Instantons
→ More Non-Trivial Cancellations
To Understand Membrane Moduli Space
- What If Relation To Topological Strings Is Unclear?

Contents

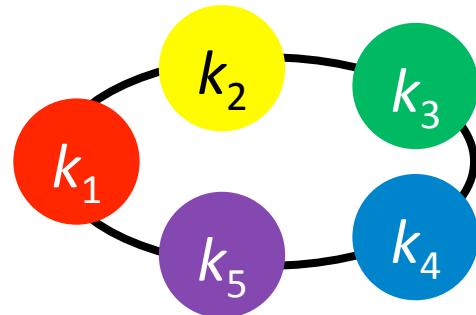
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Generalizations

- ABJM in Quiver Diagram



- More General N=3 Circular Quivers ($\sum_i k_i = 0$)



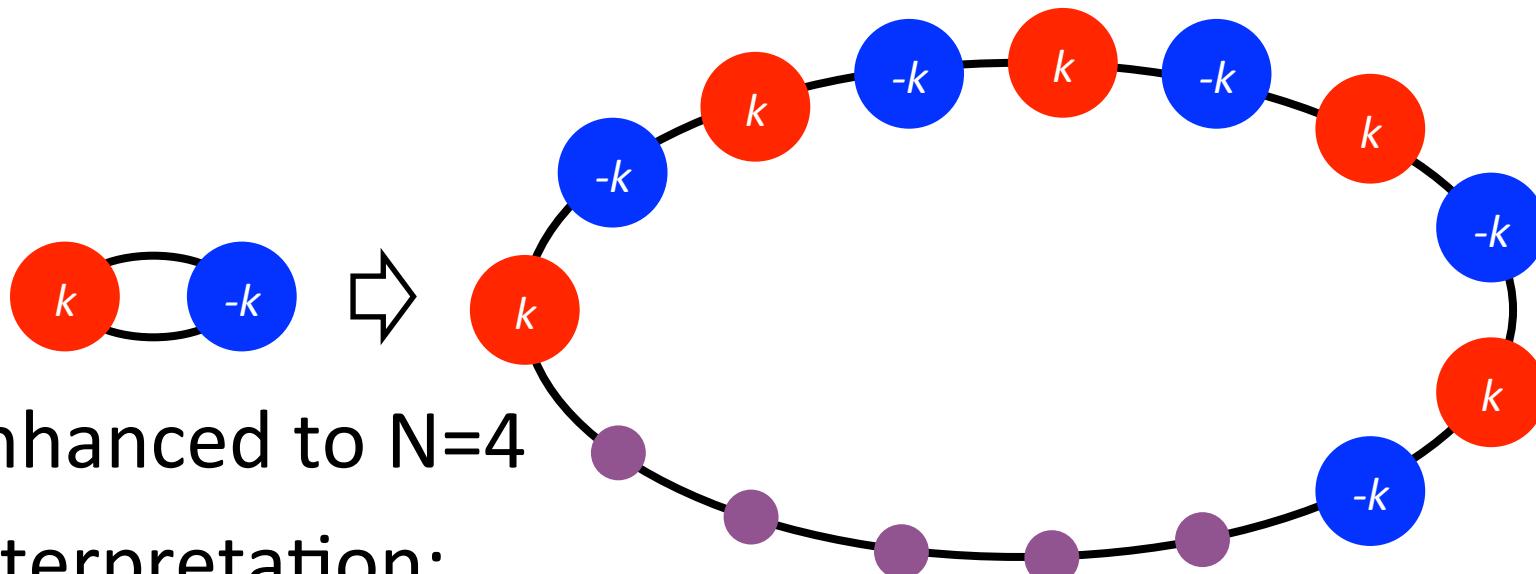
(Expected to be M2 on General Background)

Generalizations

Especially,

[Gaiotto-Witten, Hosomichi-Lee³-Park]

$$[U(N)_k \times U(N)_{-k}]^r$$



Enhanced to $N=4$

Interpretation:

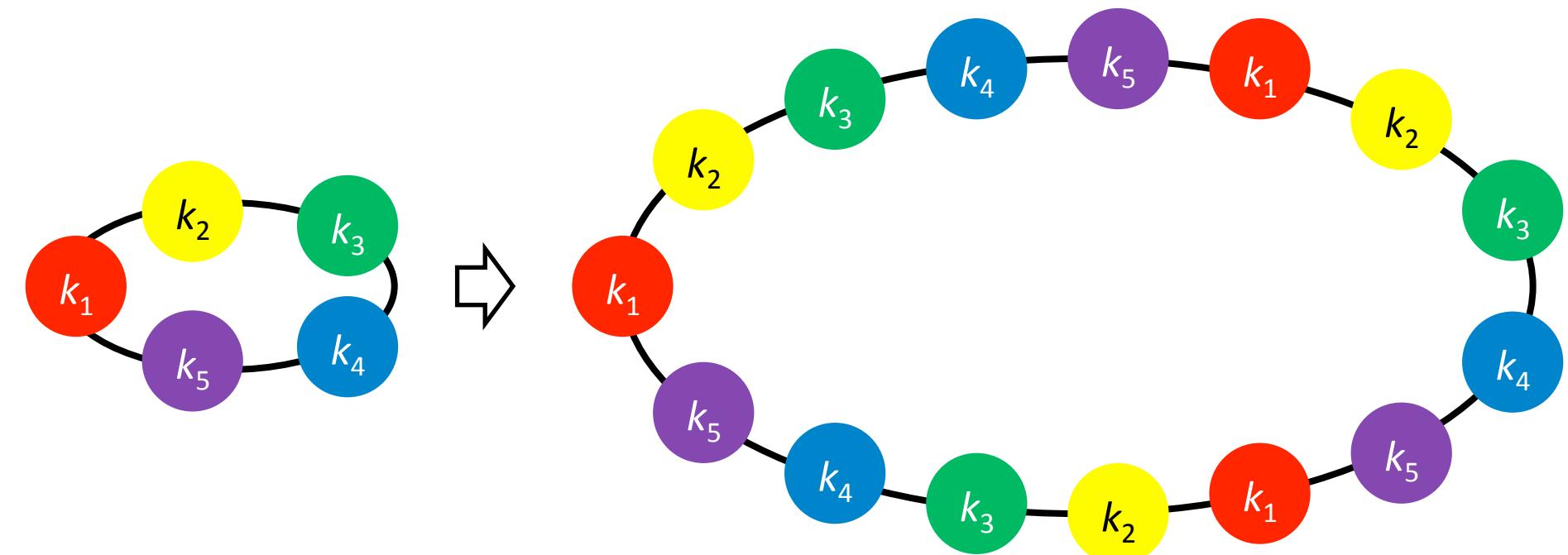
$$\text{Orbifold } C^4 / (Z_{kr} \times Z_r)$$

[Benna-

Klebanov-Klose-Smedback, Imamura-Kimura, Terashima-Yagi]

In Fact, More Generally

As Long As Multiple Repetition



Results [Honda-M]

- Grand Potential for Multiple Repetition

$$J_r(\mu) = \dots \text{ in terms of } J_1(\mu)$$

- Especially, Perturbative Coefficients

$$C_r = C_1 / r^2 \quad B_r = B_1 - \pi^2 C_1 (1 - r^{-2}) / 3 \quad A_r = r A_1$$

Results [Honda-M]

- Orbifold ABJM Partition Function Solved
 - A New **Instanton** Effect Originates From Pert
$$\exp(-2\pi^2 \textcolor{green}{C}\mu/k) = \exp(-2\mu/r)$$
 - Interpretation:
D2 wrapping A New Lagragian Submfld RP^3/\mathbb{Z}_r

More N=4 Superconformal Theories

N=4 Enhancement, Not Restricted To

$$\{k_a\} = \{k, -k, k, -k, \dots, k, -k\}$$

But

[Imamura-Kimura]

$$k_a = (k/2) (s_a - s_{a-1}) \quad s_a = \pm 1$$

Many Properties Similar to ABJM

Airy Functions, Cancellation Mechanism, ...

[M-Nosaka, see Nosaka's poster]

Summary

- ABJM Partition Function Solved
- Generalizations to General $N=3$
Superconformal Circular Quiver Theories?
- Start with $N=4$
 - Orbifold Case: Solved AND Beyond
 - General Cases: Just Starting
- Hope: Solve All Superconformal Theories &
Understand **Membrane** Moduli Space

Thank you for your attention.