$A_{\infty}$  structure from the Berkovits formulation of open superstring field theory Tomoyuki Takezaki, The University of Tokyo, Komaba Based on arXiv: 1505.01659 with Theodore Erler and Yuji Okawa See also arXiv: 1505.02069 and 1510.00364 by Theodore Erler

## 1. Introduction

For the Neveu-Schwarz sector of open superstring field theory, we now have two formulations: the Berkovits formulation [1] and the  $A_{\infty}$  formulation [2].

Berkovits  $A_{\infty}$  In the  $A_{\infty}$  formulation, the singularity in Witten's original theory is regularized.  $\{M_n\}_{n\in\mathbb{N}}$  are constructed from Q, star product \* and  $\xi$  [2].

# 4. Main results of our paper

In our paper, we transform the action of the  $A_{\infty}$  formulation into the same form as the Berkovits formulation.

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Hilbert space	large HS	small HS
gauge fixing	difficult	straightforward
closed form expression	yes	no

We find the same structure in these two formulations.

# 2. The Berkovits formulation

The Berkovits formulation [1] is based on the large Hilbert space of the superconformal ghosts. The equation of motion and the gauge transformation in the free theory are

 $Q\eta\Phi_B=0, \quad \delta\Phi_B=Q\Lambda+\eta\Omega$ 

where Q is the BRST operator of the superstring, and  $\eta$  is the zero mode of the superconformal ghost  $\eta(z)$ . The full action takes the Wess-Zumino-Witten-like form

 $S_{
m WZW}[\Phi_B] = - \int_0^1 dt \, \langle \, B_t(t), Q B_\eta(t) \, 
angle \,$ 

with

 $B_n(t) = (\eta e^{\Phi_B(t)}) e^{-\Phi_B(t)}, \ \ B_t(t) = (\partial_t e^{\Phi_B(t)}) e^{-\Phi_B(t)}$ where  $\Phi_B(0) = 0$  and  $\Phi_B(1) = \Phi_B$ . The *t* dependence of

 $S_{\mathrm{A}_{\infty}}[\Psi_A] = - \int_0^1 dt \left\langle \, A_t(t), Q A_\eta(t) \, 
ight
angle$ where  $A_{\eta}(t)$  and  $A_{t}(t)$  satisfies  $\eta A_{\eta}(t) = A_{\eta}(t)A_{\eta}(t),$  $\partial_t A_n(t) = \eta A_t(t) - A_n(t)A_t(t) + A_t(t)A_n(t)$ . The expression of  $A_{\eta}(t)$  and  $A_{t}(t)$  to the second order is  $A_\eta(t)=\Psi_A(t)+rac{1}{2}\xi(\Psi_A(t)\Psi_A(t))$  $-rac{1}{3}(\xi\Psi_A(t))\Psi_A(t)+rac{1}{3}\Psi_A(t)(\xi\Psi_A(t))+\cdots$  $A_t(t) = \xi \partial_t \Psi_A(t) + rac{1}{3} \xi \Big( \Psi_A(t) (\xi \partial_t \Psi_A(t)) - (\xi \partial_t \Psi_A(t)) \Psi_A(t) \Big)$  $-rac{1}{3}(\xi\partial_t\Psi_A(t))(\xi\Psi_A(t))-rac{1}{3}(\xi\Psi_A(t))(\xi\partial_t\Psi_A(t))+\cdots$ where  $\Psi_A(0) = 0$  and  $\Psi_A(1) = \Psi_A$ . Equating  $A_n$  and  $B_n$ , we obtain a field redefinition between the reduced Berkovits formulation and the  $A_{\infty}$  formulation. We also uplift the  $A_{\infty}$  formulation to the large Hilbert space, and find its relation to the Berkovits formulation.

the action is topological, and the action is a functional of  $\Phi_B$ . The gauge invariance and the topological t dependence follow from

> $\eta B_{\eta}(t) = B_{\eta}(t)B_{\eta}(t),$  $\partial_t B_\eta(t) = \eta B_t(t) - B_\eta(t) B_t(t) + B_t(t) B_\eta(t)$ .

A regular formulation in small Hilbert space (reduced Berkovits formulation) can be obtained by fixing gauge freedom generated by  $\eta$ . The point is that we can use  $\xi$ ; a line integral of superconformal ghost  $\xi(z)$  [3].

## 3. The $A_{\infty}$ formulation

The  $A_{\infty}$  formulation [2] is based on the small Hilbert space of the superconformal ghosts. The action and the gauge transformation are given by the same set of multi-string products  $\{M_n\}_{n\in\mathbb{N}}$ :

$$S_{A_\infty}[\Psi_A] = rac{1}{2}\,\omega(\Psi_A,Q\Psi_A) + rac{1}{3}\,\omega(\Psi_A,M_2(\Psi_A,\Psi_A)) 
onumber \ + rac{1}{4}\,\omega(\Psi_A,M_3(\Psi_A,\Psi_A,\Psi_A)) + \cdots$$

### 5. Conclusion

We can understand the Berkovits formulation and the  $A_{\infty}$ formulation as different parametrizations of  $A_n, A_t$ .



#### 6. Future directions

Recentry, an action of open superstring field theory including the Ramond sector is constructed [4]. They started with the Berkovits action, and coupled it to the Ramond string field. In our next paper [5], we construct an  $A_{\infty}$  action with Ramond sector, and show its equivalence with [4].

 $\delta_{\Lambda}\Psi_A = Q\Lambda + M_2(\Lambda,\Psi_A) + M_2(\Psi_A,\Lambda)$ 

 $+M_3(\Lambda,\Psi_A,\Psi_A)+M_3(\Psi_A,\Lambda,\Psi_A)+M_3(\Psi_A,\Psi_A,\Lambda)+\cdots$ where  $M_1 \equiv Q$ . The gauge invariance of the action follows from  $A_{\infty}$  relations

 $0 = Q^2 a$ 

 $0 = QM_2(a, b) + M_2(Qa, b) + M_2(a, Qb)$  $0 = QM_3(a, b, c) + M_2(M_2(a, b), c) + M_2(a, M_2(b, c))$  $+ M_3(Qa, b, c) + M_3(a, Qb, c) + M_3(a, b, Qc)$  $0 = QM_4(a, b, c, d) + \cdots$ 

where  $a, b, c, d, \ldots$  are arbitrary string fields. If  $M_2$  is associative, we can set  $M_{n>3} = 0$ . However, such a choice leads to divergent results (Witten's original theory).

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