Exotic Brane Junctions from F-theory

JHEP 05 (2016) 060 arXiv:1602.08606

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- What is exotic brane?
- How can we use exotic branes?

What is exotic brane?

Exotic brane is

- from standard brane via string dualities in lower dim
- ✓ vortex-like (codim 2)

Obers and Pioline: hep-th/9809039

Eyras and Lozano: hep-th/9908094

de Boer and Shigemori: arXiv:1209.6056

etc..

Exotic b_n^c -brane has a tension

$$rac{R_1 R_2 \cdots R_b \, (R_{b+1} \cdots R_{b+c})^2}{g_s^n \, \ell_s^{b+2c+1}}$$

Performing string dualities, the tension is transformed:

Example duality chain of 5-branes:

Exotic brane is

- from standard brane via string dualities in lower dim
- ✓ vortex-like (codim 2)

Obers and Pioline: hep-th/9809039

Eyras and Lozano: hep-th/9908094

de Boer and Shigemori: arXiv:1209.6056

etc..

NOTE

D7-brane is an object of codim 2 in 10D.

D7-brane physics has been studied for 20 years: F-theory

Vafa: hep-th/9602022

Sen: hep-th/9605150

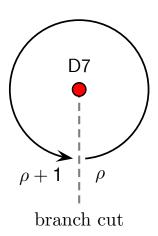
etc..

D7-brane:

$$ho(z) \; \equiv \; C^{(0)} + \mathrm{i}\mathrm{e}^{-\phi} \; = \; rac{ heta}{2\pi} + rac{\mathrm{i}}{2\pi} \log\left(rac{\Lambda}{r}
ight) \hspace{0.5cm} (z = x^8 + \mathrm{i} x^9 = r\,\mathrm{e}^{\mathrm{i} heta})$$

When ρ moves around D7-brane counterclockwise $\theta \to \theta + 2\pi$,

it receives a magnetic "charge" of D7-brane (monodromy) : $\rho \rightarrow \rho + 1$



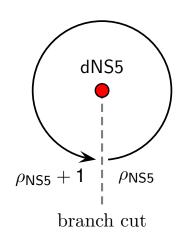
There exists a branch cut in z-plane.

defect NS5-brane: (transverse a, b directions are smeared)

$$ho_{ extsf{NS5}}(z) \; \equiv \; B_{ab}^{(2)} + \mathrm{i}\mathrm{e}^{+2\phi} \; = \; rac{ heta}{2\pi} + rac{\mathrm{i}}{2\pi} \log\left(rac{\Lambda}{r}
ight) \hspace{0.5cm} (z = x^8 + \mathrm{i} x^9 = r\,\mathrm{e}^{\mathrm{i} heta})$$

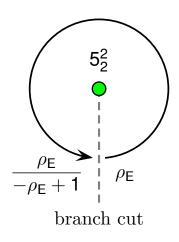
When ρ moves around NS5-brane counterclockwise $\theta \to \theta + 2\pi$,

it receives a magnetic "charge" of NS5-brane (monodromy) : $ho_{\rm NS5}
ightarrow
ho_{
m NS5} + 1$



In this case, we can identify $\rho_{\rm NS5}$ + 1 $\simeq \rho_{\rm NS5}$ by B-field gauge transformation.

Exotic 5_2^2 -brane : (T_{ab} -dualized from defect NS5-brane)



monodromy :
$$\frac{1}{\rho_{\rm E}} \to \frac{1}{\rho_{\rm E}} - 1$$
 where $\rho_{\rm E} = -\frac{1}{\rho_{\rm NS5}}$

We cannot identify this monodromy change by $\left\{ egin{array}{ll} B ext{-field gauge transformation} \\ ext{coordinate transformations} \end{array}
ight.$

52-brane's property is from that of defect NS5-brane

via
$$SL(2,\mathbb{Z}) \in SO(2,2;\mathbb{Z})$$
 \mathbf{T}_{ab} -duality

How can we use exotic branes?

Consider an NS5-brane crossing the branch cut of D7-brane from the left. If D7-brane goes across NS5-brane, a new 5-brane and a junction appear (Hanany-Witten effect).



Note: D7-brane(1234567), D5(1234X), NS5(1234Y), $X, Y \in z$ -plane

This is a brane junction in F-theory.

Gaberdiel and Zwiebach: hep-th/9709013 DeWolfe and Zwiebach: hep-th/9804210

etc..

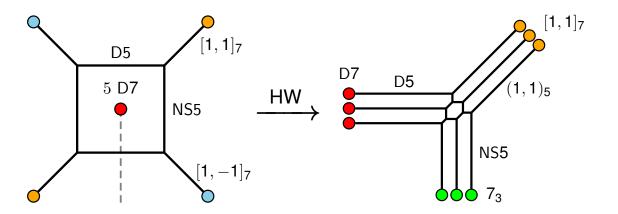
Perform the T_{34} - and S-duality of the system of NS5-brane with D7-brane.

We obtain a config. that a new "3-brane" with exotic 5_2^2 -brane :

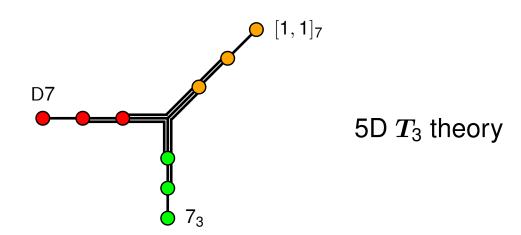


We find that D5(1234Y)-brane wrapped on T_{34}^2 (\equiv wD5) is ending on 5_2^2 -brane.

5D theory on 5-branes with 5 D7-branes:



IIB	0	1	2	3	4	5	6	7	8	9
5 D7	_	_	_	_	_	_	_	_		i i i
D5		_	_	_	_				_	i i i
NS5	_	_	_	_	_					
(1,1) ₅		_	_	_	_				an	gle



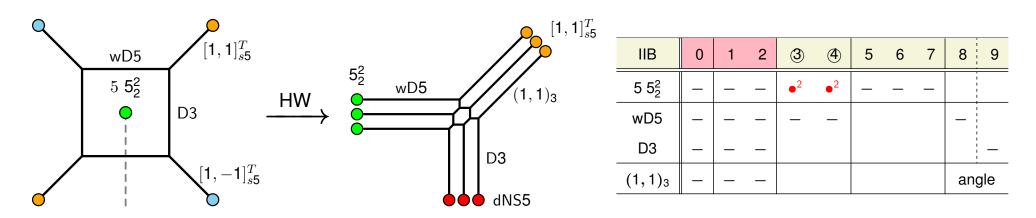
Aharony and Hanany: hep-th/9704170

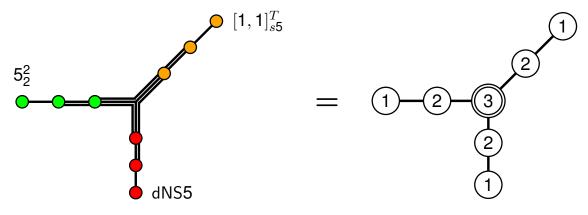
DeWolfe, Hanany, Iqbal and Katz: hep-th/9902179

Gaiotto and Witten: arXiv:0804.2902, 0807.3720

Benini, Benvenuti and Tachikawa: arXiv:0906.0359

3D theory on "3-branes" with 5 exotic 5^2_2 -branes :





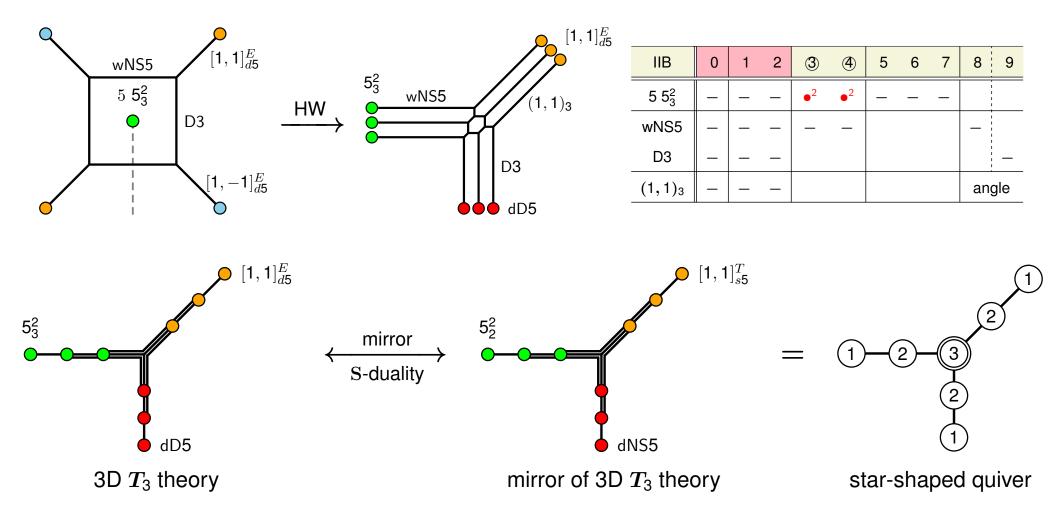
mirror of 3D T_3 theory

star-shaped quiver

 5_2^2 -brane yields mirror of 3D T_3 theory.

Benini, Tachikawa and Xie: arXiv:1007.0992

3D theory on "3-branes" with 5 exotic 5_3^2 -branes :



 5_3^2 -brane gives 3D T_3 theory.

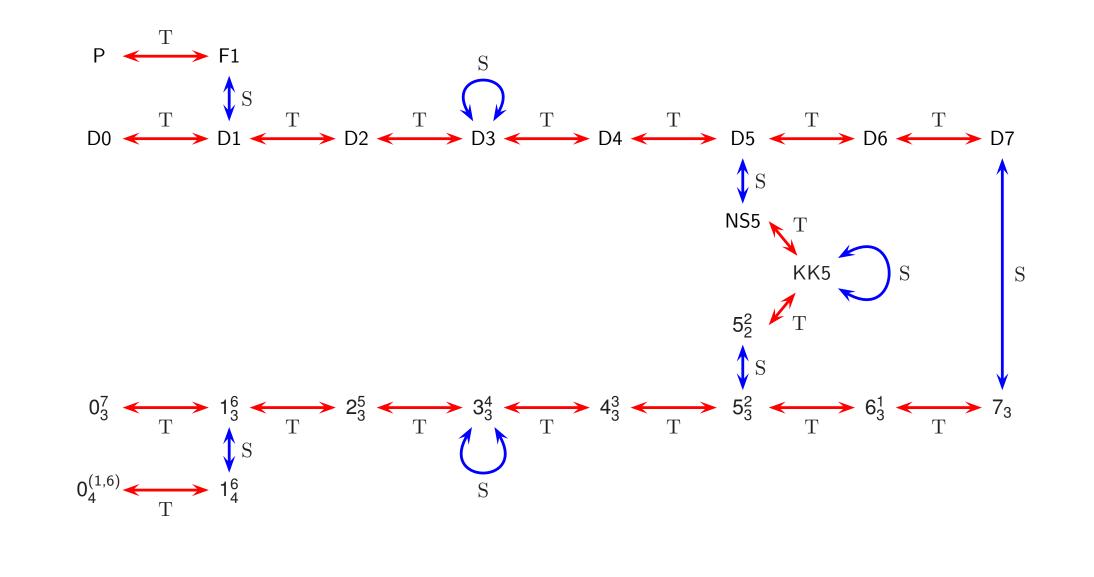
Exotic branes also play a role in generating (non-)Lagrangian theories.

Summary and Discussions

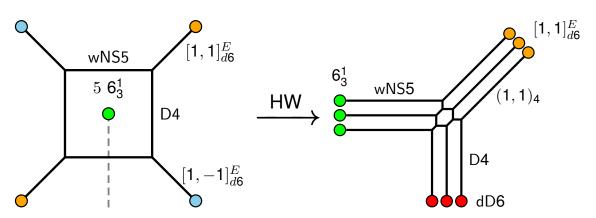
- Exotic brane has a monodromy (charge) with branch cut.
- When D3-brane cross the branch cut of 5_2^2 -brane, it jumps to D3 + "D5 wrapped on two-torus".

- Exotic 5-branes are the building blocks of 3D T_3 -theory and its mirror theory.
- They correctly provide 3D theory even when spacetime is compactified.

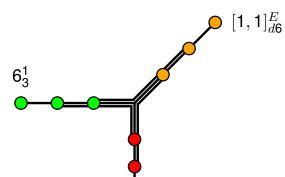
ullet Non-Lagrangian theory from non-geometric background



4D theory on "4-branes" with 5 exotic 6^1_3 -branes :



IIA	0	1	2	3	4	5	6	7	8	9
5 6 ₃ ¹	_	_	_	_	•2	_	_	_		
wNS5	_	_	_	_	_				_	
D4	_	_	_	_						_
$(1,1)_4$	_	_	_	_					an	gle



4D T_3 theory can be realized.

Uplift to M-theory : dD6
$$\to$$
 KK6, $6^1_3 \to$ KK6; D4 \to M5, wNS5 \to M5 $g_s\ell_s=R_{\natural},\ g_s\ell_s^3=\ell_{\mathtt{P}}^3$

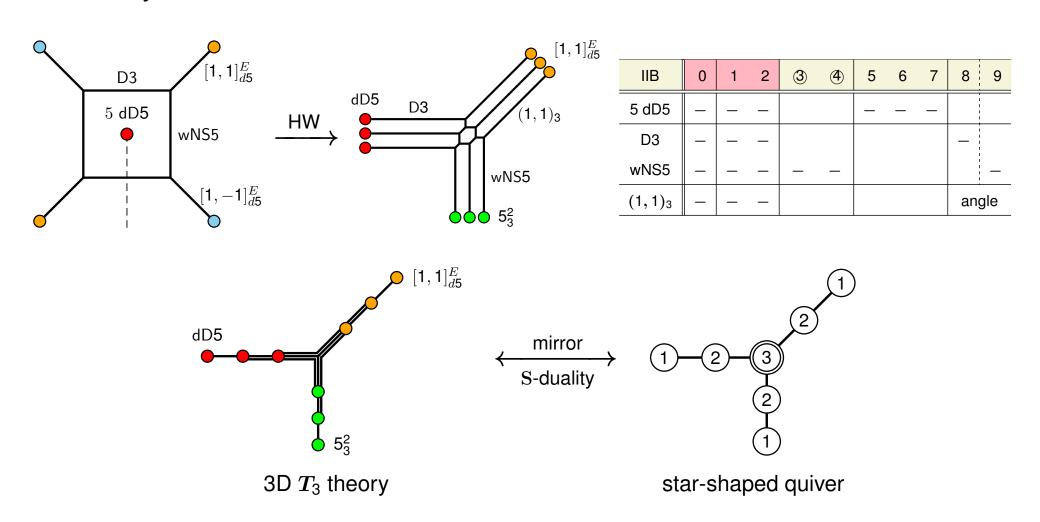
dD6

Gaiotto: arXiv:0904.2715

Thanks

Appendix

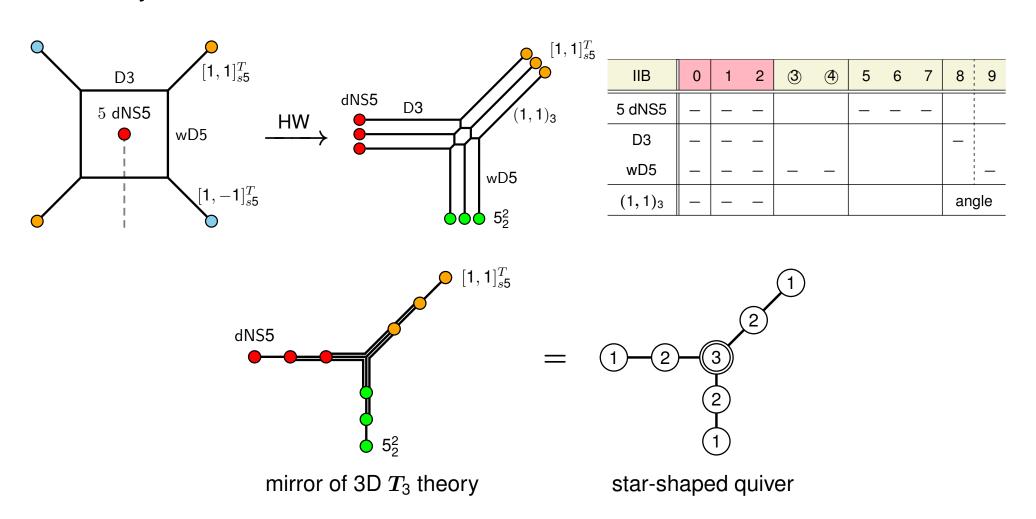
3D theory on "3-branes" with 5 defect D5-branes:



3D T_3 theory is realized.

Benini, Tachikawa and Xie: arXiv:1007.0992

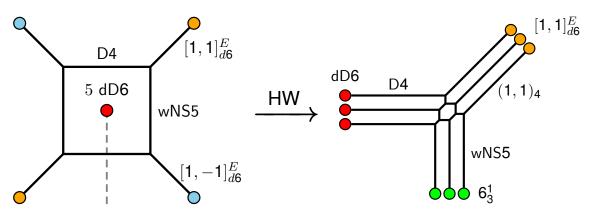
3D theory on "3-branes" with 5 defect NS5-branes:



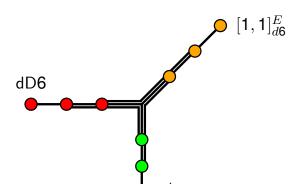
Mirror of 3D T_3 theory is realized.

Benini, Tachikawa and Xie: arXiv:1007.0992

4D theory on "4-branes" with 5 defect D6-branes:



IIA	0	1	2	3	4	5	6	7	8	9
5 dD6	_	_	_	_		_	_	_		
D4	_	_	_	_					_	
wNS5	_	_	_	_	_					_
$(1,1)_4$	_	_	_	_					an	gle

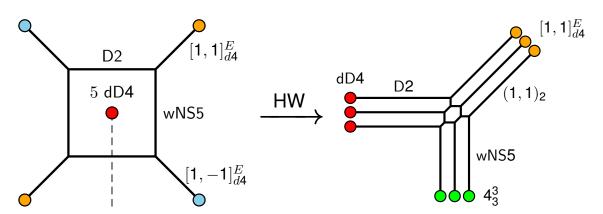


4D T_3 theory can be realized.

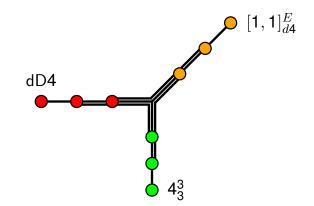
Uplift to M-theory : dD6
$$o$$
 KK6, 6^1_3 o KK6; D4 o M5, wNS5 o M5 $g_s\ell_s=R_{\natural},\ g_s\ell_s^3=\ell_{\mathtt{P}}^3$

Gaiotto: arXiv:0904.2715

2D theory on "2-branes" with 5 defect D4-branes:



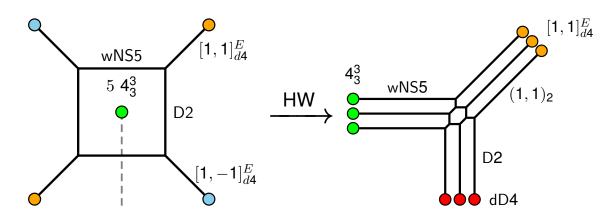
IIA	0	1	2	3	4	5	6	7	8	9
5 dD4	_	-				_	_	_		
D2	_	_							_	
wNS5	_	_	_	_	_					_
$(1,1)_2$	_	_							an	gle



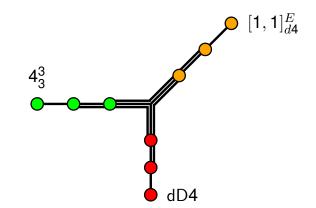
2D T_3 theory can be realized.

Uplift to M-theory : $dD4 \rightarrow M5$, $4^3_3 \rightarrow 5^3$; $wNS5 \rightarrow M5$, $D2 \rightarrow M2$

2D theory on "2-branes" with 5 exotic 4_3^3 -branes :



IIA	0	1	2	3	4	5	6	7	8	9
5 4 ³ ₃	_	_	•2	•2	•2	_	_	_		
wNS5	_	_	_	_	_				_	
D2	_	_								_
(1,1)2	_	_							an	gle



2D T_3 theory can be realized.

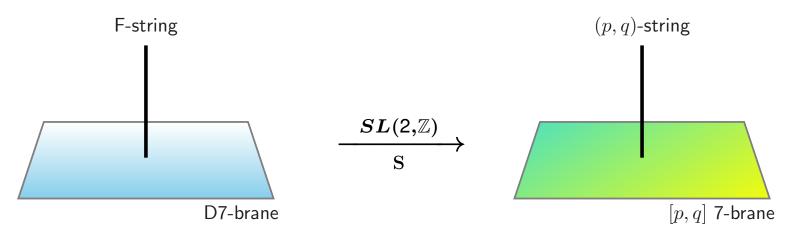
Uplift to M-theory : $dD4 \rightarrow M5$, $4^3_3 \rightarrow 5^3$; $wNS5 \rightarrow M5$, $D2 \rightarrow M2$

F-string: couple to $B_{(2)}$

D-string : couple to $C_{(2)}$

$$\rho(z) = \frac{\mathrm{i}}{2\pi} \log\left(\frac{\Lambda}{z}\right) = \frac{\theta}{2\pi} + \frac{\mathrm{i}}{2\pi} \log\left(\frac{\Lambda}{r}\right)$$

D7(1234567): couple to $ho(z)=C+\mathrm{i}\mathrm{e}^{-\phi}$ $(z=x^8+\mathrm{i}x^9=r\,\mathrm{e}^{\mathrm{i}\theta})$



$$(1,0)$$
-string = F1

$$(1,0)$$
-string = F1 $[1,0]$ 7-brane = D7(1234567)

$$(0,1)$$
-string = D1

$$(0,1)$$
-string = D1 $[0,1]$ 7-brane = $7_3(1234567)$

Open D-string is ending on $7_3(1234567)$.

This is a setup in F-theory. We perform ST_{67} -duality and reduce 67-directions.

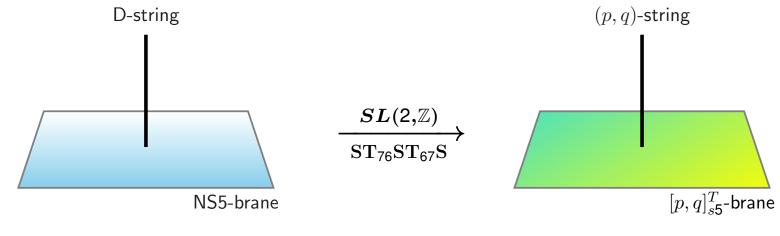
D-string: couple to $C_{(2)}$

D3-brane : couple to $C_{(4)}$

$$ho(z) = rac{\mathrm{i}}{2\pi}\log\left(rac{\Lambda}{z}
ight) = rac{ heta}{2\pi} + rac{\mathrm{i}}{2\pi}\log\left(rac{\Lambda}{r}
ight)$$

NS5(12345): couple to
$$ho(z) = B_{67}^{(2)} + \mathrm{i} \mathrm{e}^{+2\phi}$$
 $(z = x^8 + \mathrm{i} x^9 = r\,\mathrm{e}^{\mathrm{i}\theta})$

$$(z = x^8 + ix^9 = r e^{i\theta})$$



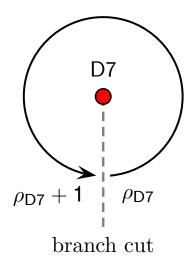
$$(1,0)$$
-string = D1

$$[1,0]_{s5}^{T}$$
-brane = NS5(12345)

$$(0,1)$$
-string = D3 wrapped on T_{67}^2 $[0,1]_{s5}^T$ -brane = $5_2^2(12345,67)$

$$[0, 1]_{s5}^{T}$$
-brane = $5_{2}^{2}(12345,67)$

Open D3-brane wrapped on T_{67}^2 is ending on exotic 5_2^2 (12345,67).



$$ho \, o \,
ho + \mathsf{1} \, = \, rac{a
ho + b}{c
ho + d} \, \equiv \, M_{[\mathsf{1}, \mathsf{0}]} \cdot
ho \, , \quad \, M_{[\mathsf{1}, \mathsf{0}]} \, = \, \left(egin{array}{c} \mathsf{1} & \mathsf{1} \ \mathsf{0} & \mathsf{1} \end{array}
ight) \, \in \, SL(\mathsf{2}, \mathbb{Z})$$

or

$$K_{[1,0]}\cdot(\rho+1) = \rho\,, \quad K_{[1,0]} = (M_{[1,0]})^{-1} = \left(\begin{array}{cc} 1 & -1 \\ 0 & 1 \end{array}\right)$$

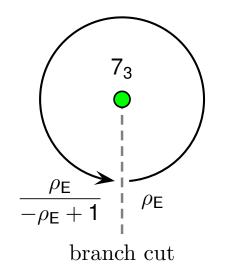
 $M_{[p,q]}$: moving around the 7-brane

 $K_{[p,q]}$: going across the branch cut

By $SL(2,\mathbb{Z})$, the monodromy matrix for general [p,q] 7-brane is given as

$$K_{[p,q]} \ = \ g \, K_{[1,0]} \, g^{-1} \ = \ \left(egin{array}{ccc} 1 + pq & -p^2 \ q^2 & 1 - pq \end{array}
ight) \ g \in SL(2,\mathbb{Z})$$

ex) monodromy $K_{[0,1]}$ for 7_3 -brane : $K_{[0,1]}=\begin{pmatrix}1&0\\1&1\end{pmatrix}$



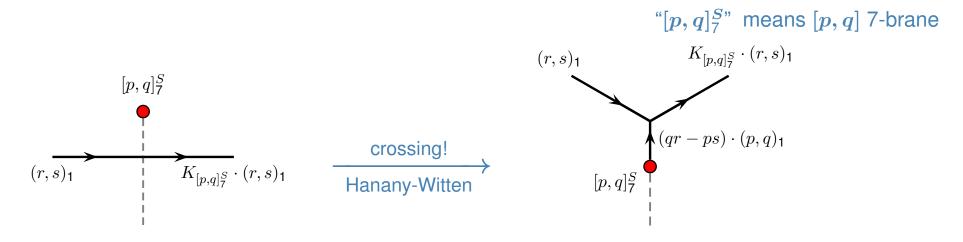
$$ho_{\mathsf{E}} \, = \, -rac{\mathsf{1}}{
ho_{\mathsf{D7}}}$$

Consider an (r, s)-string crossing the branch cut of [p, q] 7-brane from the left.

The string charge is jumped by monodromy.

If the 7-brane goes across the string, the string is no longer crossing the branch cut.

Further, a new string and a junction appear (Hanany-Witten effect).



Note: 7-brane is stretched in 1234567-directions.

This is a string junction in F-theory.

Consider a D-string crossing the branch cut of D7-brane from the left.

A new string and a junction appear by Hanany-Witten effect.



Note: D7-brane is stretched in 1234567-directions.

This is a string junction in F-theory.

Gaberdiel and Zwiebach: hep-th/9709013

DeWolfe and Zwiebach: hep-th/9804210

etc..

Consider a D3-brane wrapped on T_{ab}^2 (wD3) and defect NS5-brane (dNS5). If dNS5 goes across wD3,

a new D-string and a junction appear by Hanany-Witten effect.



Consider a D-string and 5_2^2 -brane.

D-string charge is jumped by monodromy.

If 5₂²-brane goes across D-string,

a new wD3 and a junction appear (Hanany-Witten effect).

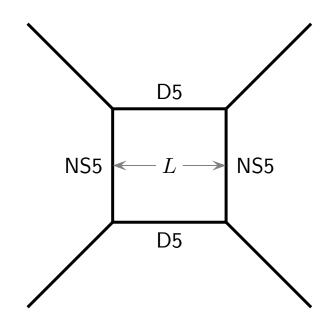


• defect NS5-brane and D5-brane wrapped on T_{ab}^2 :



• 5²₂-brane and D3-brane :



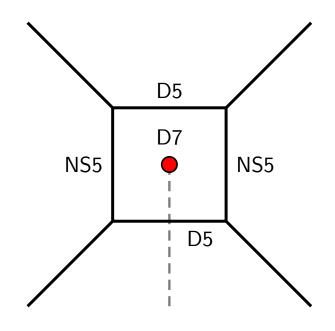


IIB	0	1	2	3	4	5	6	7	8	9
<i>n</i> D7	_	_	_	_	_	_	_	_		;
D5	_	_	_	_	_				_	
NS5	_	_	_	_	_					<u> </u>
$(1,1)_5$	_	_	_	_	_				an	gle

color : N_c D5 between 2 NS5 = $SU(N_c)$ gauge symmetry

flavor : N_f D5 outside 2 NS5 = N_f flavors

coupling : $rac{1}{g_{
m YM}^2} \simeq rac{L}{g_s \, \ell_s^2}$ (This can be derived from Dirac-Born-Infeld action of D5-brane.)

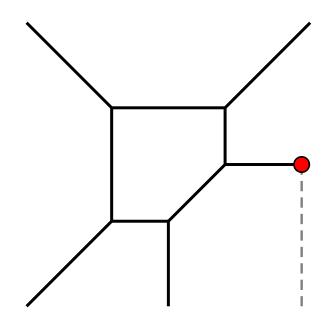


IIB	0	1	2	3	4	5	6	7	8	9
<i>n</i> D7										1
										! ! !
D5	_	_	_	_	_				_	
NS5	_	_	_	_	_					<u> </u>
$(1,1)_5$	_	_	_	_	_				an	gle

color : N_c D5 between 2 NS5 = $SU(N_c)$ gauge symmetry

flavor : N_f D5 outside 2 NS5 = N_f flavors

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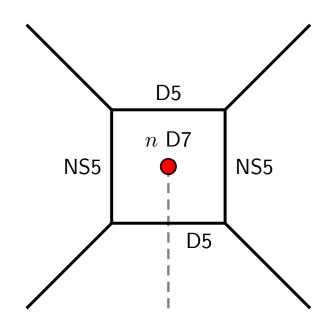


IIB	0	1	2	3	4	5	6	7	8	9
<i>n</i> D7	_	_	_	_	_	_	_	_		;
D5	_	_	_	_	_				_	
NS5		_	_	_	_					: -
$(1,1)_5$	_	_	_	_	_				an	gle

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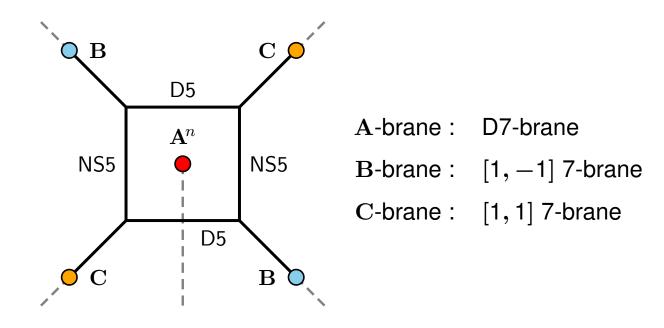
IIB	0	1	2	3	4	5	6	7	8	9
<i>n</i> D7	_	_	_	_	_	_	_	_		
D5	_	_	_	_	_				_	
NS5	_	_	_	_	_					_
$(1,1)_5$	_	_		_	_				an	gle

This config. indicates 5D SU(2) gauge symmetry with n flavors on 01234-directions.

In the $L \propto 1/g_{
m YM}^2
ightarrow 0$ limit, this gauge theory flows to SCFT with E_{n+1} symmetry.

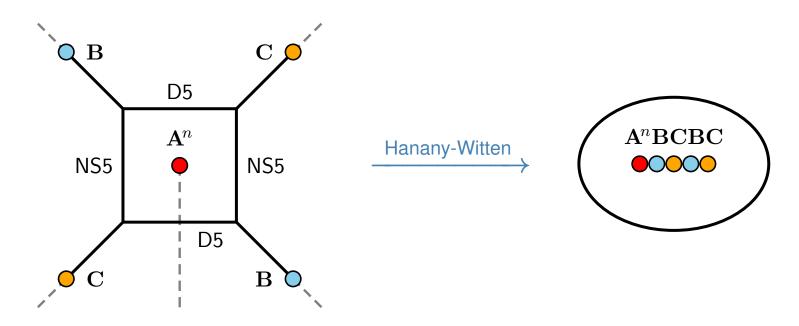
(UV fixed point)

5D $\mathcal{N}=1$ SU(2) gauge theory with n flavors on D5-brane can be illustrated as



Without changing the 5D gauge theory on D5-brane, semi-infinite (p, q) 5-branes are terminated by [p, q] 7-branes.

5D $\mathcal{N}=1$ SU(2) gauge theory with n flavors on D5-brane can be illustrated as



Moving B- and C-branes along the (p,q) 5-branes and going inside the "box", the (p,q) 5-branes are annihilated by the Hanany-Witten effect.

Further, the "box" becomes "loop" by back reaction of A^n -, B-, C-, B-, and C-branes.

(skipped drawing the branch cuts.)

5D $\mathcal{N}=1$ SU(2) gauge theory with n flavors on D5-brane can be illustrated as

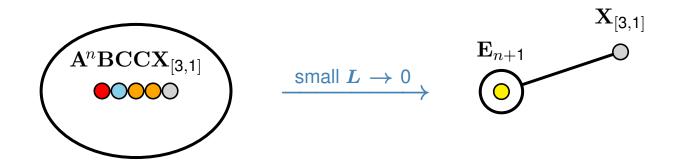


Perform the re-ordering of 7-branes with branch cuts.

When a 7-brane goes across another's branch cut, its monodromy is modified.

$$A^nBCBC \rightarrow A^nBCCX_{[3,1]}$$
 with $X_{[3,1]}$ -brane $\equiv [3,1]$ -brane

5D $\mathcal{N}=1$ SU(2) gauge theory with n flavors on D5-brane can be illustrated as



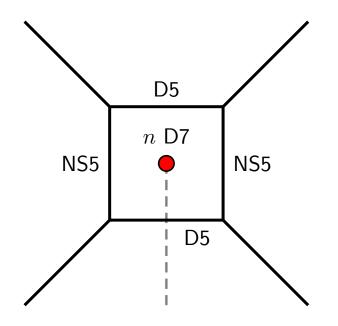
Perform the small loop limit $L \to 0$. This implies the strong coupling limit $g_{YM} \to \infty$. There exists a non-trivial UV fixed point of 5D gauge theory \to CFT.

 ${\bf A}^n$ -, ${\bf B}$ -, and ${\bf C}^2$ -branes are collapsed to ${\bf E}_{n+1}$ -brane.

 $\mathbf{X}_{[3,1]}$ -brane is gone far away from \mathbf{E}_{n+1} -brane.

Open string ending on 5-brane loop and \mathbf{E}_{n+1} -brane provides \mathbf{E}_{n+1} symmetry.

We argued 5D SUSY gauge theory and its strong coupling limit by brane construction.



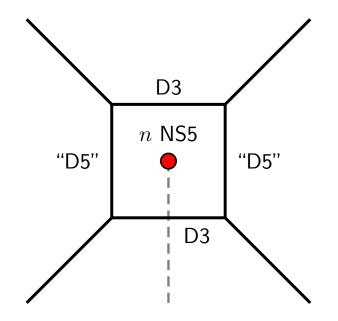
IIB	0	1	2	3	4	5	6	7	8	9
<i>n</i> D7	_	_	_	_	_	_	_	_		i i i
D5	_	_	_	_	_				_	
NS5	_	_	_	_	_					<u> </u>
$(1,1)_5$	_	_	_	_	_				an	gle

 ST_{67} -dual \rightarrow 5D

Perform string dualities : ST_{47} -dual ightarrow 4D

 ST_{34} -dual \rightarrow 3D

${ m ST}_{ m 34} ext{-dualized system}$:

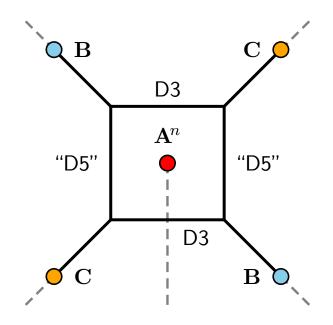


IIB	0	1	2	3	4	5	6	7	8	9
n NS5	_	_	_			_	_	_		
D3	_	_	_						_	
"D5"		_	_	1	_					_
$(1,1)_3$	_	_	_						an	gle

We can see 3D SU(2) gauge symmetry with n flavors on 012-directions.

There exits IR fixed point.

${ m ST}_{ m 34} ext{-dualized system}$:



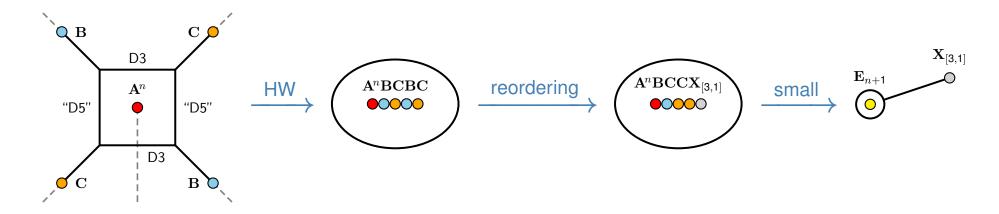
IIB	0	1	2	3	4	5	6	7	8	9
n NS5	_	_	_			_	_	_		
D3		_	_						_	
"D5"	_	_	_	_	_					<u> </u>
$(1,1)_3$	_	_	_						an	gle

A-brane: $[1,0]_{s5}^T$ -brane = NS5

B-brane: $[1,-1]_{s5}^T$ -brane

C-brane: $[1,1]_{s5}^T$ -brane

3D $\mathcal{N}=4$ SU(2) gauge with n flavors \rightarrow SCFT with E_{n+1} symmetry



3D gauge coupling is given by Dirac-Born-Infeld action of D3-brane:

$$rac{1}{g_{
m YM}^2} \simeq rac{L}{g_s}$$

Then the strong coupling limit is given by