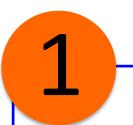
S-matrix Unitarity in Higher Derivative Gravity with Matter

Yugo Abe (Miyakonojo KOSEN, Miyazaki), Strings and Fields 2019 @ Yukawa Institute In collaboration with T. Inami (VAST), K. Izumi (Nagoya U.), T. Kitamura (Rikkyo U.), and T. Noumi (Kobe U.) Based on PTEP 2018 no.3, 031E01 and arXiv:1805.00262 to appear in PTEP

Discuss "unitarity" of \mathbb{R}^2 gravity and get some important suggestions about gravity theory



Motivation: Llewellyn Smith's conjecture

He suggests that there is an intimate connection between the high energy behavior of amplitudes and renormalizability.

The high energy behavior of amplitudes can be checked by the unitarity limit (= unitarity bound).

	Unitarity bound	UV renormalizability
QED	fulfilled	renormalizable
Yang-Mills theory	fulfilled	renormalizable
Weinberg-Salam model	fulfilled	renormalizable
4-Fermi theory	unfulfilled	non renormalizable
Massive vector theory	unfulfilled	non renormalizable



Question: Matter scattering and unitarity in \mathbb{R}^2 gravity (PTEP 2018 no.3, 031E01)

A higher derivative gravity theory (R^2 gravity theory) is known to be a renormalizable theory.

We are interested in the following two things:

What about unitarity in a higher derivative gravity theory?,

Is Llewellyn Smith's conjecture correct in a higher derivative gravity theory?

	Unitarity bound	UV renormalizability
$R^2gravity$?	renormalizable

Higher derivative gravity theory (R^2 gravity theory)

$$S = \int d^4x \sqrt{-g} \left(-\frac{1}{\kappa^2} R + \alpha R^2 + \beta R_{\mu\nu}^2 + \frac{1}{2} \nabla_{\mu} \phi \nabla^{\mu} \phi - \frac{1}{2} m^2 \phi^2 - \frac{1}{4!} \lambda \phi^4 \right)$$

graviton propagator

$$G_{\mu\nu,\alpha\beta} = \frac{2}{\beta p^4 + \kappa^{-2} p^2} P_{\mu\nu,\alpha\beta}^{(2)} + \frac{1}{2(3\alpha + \beta)p^4 - \kappa^{-2} p^2} P_{\mu\nu,\alpha\beta}^{(0)},$$

$$P_{\mu\nu,\alpha\beta}^{(2)} \equiv \frac{1}{2} (\theta_{\mu\alpha}\theta_{\beta\nu} + \theta_{\mu\beta}\theta_{\alpha\nu}) - \frac{1}{3}\theta_{\mu\nu}\theta_{\alpha\beta},$$

$$P_{\mu\nu,\alpha\beta}^{(2)} \equiv \frac{1}{2} \left(\theta_{\mu\alpha} \theta_{\beta\nu} + \theta_{\mu\beta} \theta_{\alpha\nu} \right) - \frac{1}{2} \theta_{\mu\nu} \theta_{\alpha\beta}$$

$$P_{\mu\nu,\alpha\beta}^{(0)} \equiv \frac{1}{3}\theta_{\mu\nu}\theta_{\alpha\beta}$$

where
$$heta_{\mu
u}\equiv\eta_{\mu
u}-rac{p_{\mu}p_{
u}}{p^2}$$

• scattering amplitude (tree graphs : contact and s-, t-, u-channels)

$$M(\phi\phi \to \phi\phi) = -\lambda + \frac{1}{(\beta s^2 + \kappa^{-2} s)} \left(2tu - \frac{1}{3} (s - 4m^2)^2 \right) - \frac{1}{3(2(3\alpha + \beta)s^2 - \kappa^{-2} s)} (s + 2m^2)^2$$

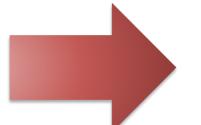
$$+ \frac{1}{(\beta t^2 + \kappa^{-2} t)} \left(2us - \frac{1}{3} (t - 4m^2)^2 \right) - \frac{1}{3(2(3\alpha + \beta)t^2 - \kappa^{-2} t)} (t + 2m^2)^2$$

$$+ \frac{1}{(\beta u^2 + \kappa^{-2} u)} \left(2st - \frac{1}{3} (u - 4m^2)^2 \right) - \frac{1}{3(2(3\alpha + \beta)u^2 - \kappa^{-2} u)} (u + 2m^2)^2$$

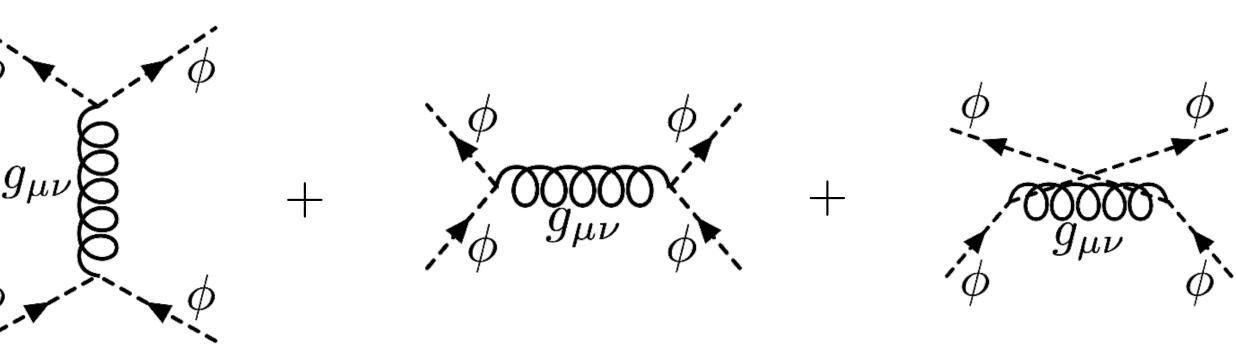
scattering amplitude in high-energy limit

$$M(\phi\phi \to \phi\phi) \simeq -\lambda + \frac{2}{\beta} \left(\frac{tu}{s^2} + \frac{us}{t^2} + \frac{st}{u^2} \right) - \frac{3(2\alpha + \beta)}{2\beta(3\alpha + \beta)}$$
$$= -\lambda + \frac{2}{\beta} g(\cos\theta) - \frac{3(2\alpha + \beta)}{2\beta(3\alpha + \beta)} = \mathcal{O}(s^0)$$

where
$$g(x) = \frac{1}{4}(1-x)(1+x) - 2(1+x)/(1-x)^2 - 2(1-x)/(1+x)^2$$



We calculate s-, t-, u-channel graviton exchanges in ϕ - ϕ scattering.



Einstein gravity theory

$$S = \int d^4x \sqrt{-g} \left(-\frac{1}{\kappa^2} R + \frac{1}{2} \nabla_{\mu} \phi \nabla^{\mu} \phi - \frac{1}{2} m^2 \phi^2 - \frac{1}{4!} \lambda \phi^4 \right)$$

graviton propagator

$$G^{\mu\nu\rho\sigma} = \frac{\kappa^2}{p^2} \left(\eta^{\mu\rho} \eta^{\nu\sigma} + \eta^{\mu\sigma} \eta^{\nu\rho} - \eta^{\mu\nu} \eta^{\rho\sigma} \right)$$

scattering amplitude (tree graphs : contact and s-, t-, u-channels)

$$M(\phi\phi \to \phi\phi) = -\lambda + 2\kappa^2 \left[\left(\frac{tu}{s} + \frac{su}{t} + \frac{st}{u} \right) + 6m^2 - 2m^4 \left(\frac{1}{s} + \frac{1}{t} + \frac{1}{u} \right) \right]$$

scattering amplitude in high-energy limit

$$M \simeq -\lambda + \kappa^2 [2sf(\cos\theta) + \mathcal{O}(s^0)] = \mathcal{O}(s^1)$$

where
$$f(x) = \frac{1}{4}(1-x^2) + 2(1+x^2)/(1-x^2)$$



We examine the unitarity bound in the UV limit :

$$s \to \infty$$
, $t \simeq -\frac{1}{2}s(1-\cos\theta)$, $u \simeq -\frac{1}{2}s(1+\cos\theta)$

where θ is the scattering angle in the center-of-mass frame

	Unitarity bound	UV renormalizability
Einstein gravity + matter	unfulfilled	non renormalizable
R^2 gravity + matter	<u>fulfilled</u>	renormalizable

Consideration: S-matrix Unitarity (arXiv:1805.00262 to appear in PTEP)

A higher derivative gravity theory does not have (Commonly used) unitarity because of the negative norm. A positive norm theory Therefore, there is no guarantee that the unitarity bound will hold in a higher derivative gravity theory. What should we evaluate to the high energy behavior of amplitudes in a higher derivative gravity theory?

We analyzed a higher derivative scalar theory as a simple toy model.

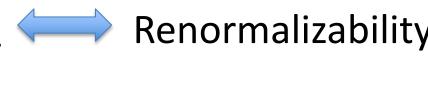
From our investigation of a higher derivative scalar theory, we have studied that the high energy behavior of amplitudes can be evaluated by S-matrix unitarity ($S^{\dagger}S=1$), regardless of the norm positivity.

We expected that there is a connection between S-matrix unitarity and renormalizability. We are going to be able to evaluate the high energy behavior of amplitudes by S-matrix unitarity $(S^{\dagger}S=1)$ in comparing a higher derivative gravity theory and Einstein gravity theory.

In future work, we will calculate s-, t-, u-channel graviton exchanges in two-graviton elastic scattering.

bound holds.

Unitarity (= $S^{\dagger}S = 1$ & positive norm) \longleftrightarrow Renormalizability There is a guarantee that Unitarity





A theory, regardless of the norm positivity S-matrix unitarity Renormalizability

