

# Rectangular $W$ -algebras of types $so$ and $sp$ and dual coset CFTs

---

Takahiro Uetoko (Ritsumeikan Univ.)

Based on: [arXiv:1906.05872]

w/ Thomas Creutzig (Alberta Univ.), Yasuaki Hikida (YITP, Kyoto Univ.)

Aug. 22 (2019) @YITP “Strings and Fields 2019”

# Introduction

- Strings and Higher spins



# Introduction

- Strings and Higher spins

**String theory**

First Regge trajectory



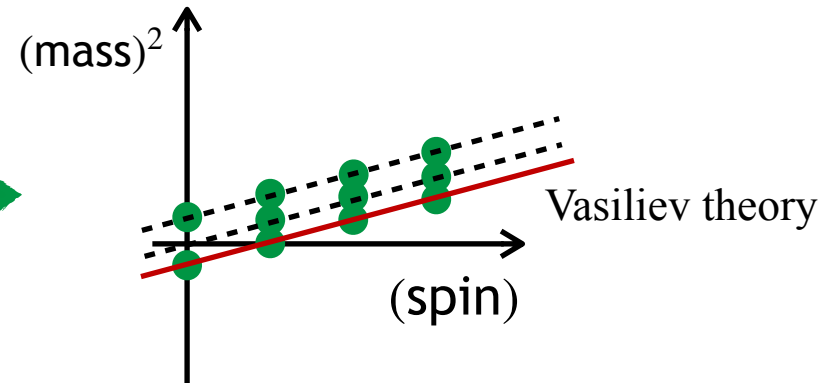
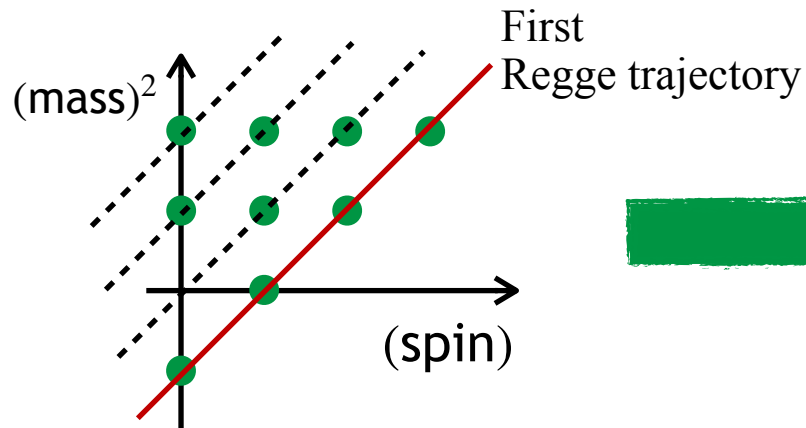
**Higher spin gravity**

Vasiliev theory

**Tensionless limit**

[Gross '88]

- String spectrum

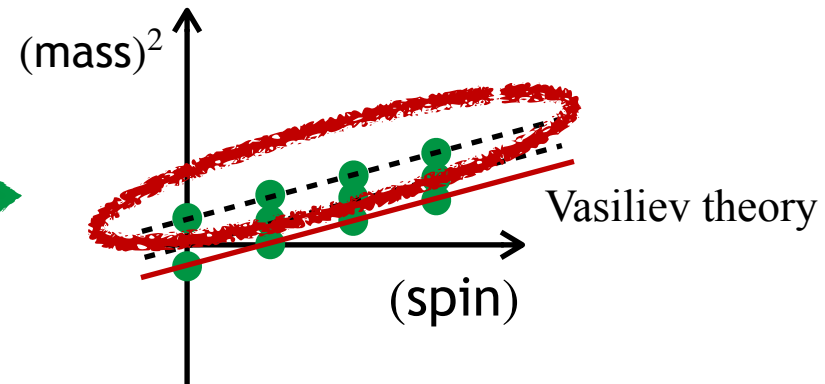
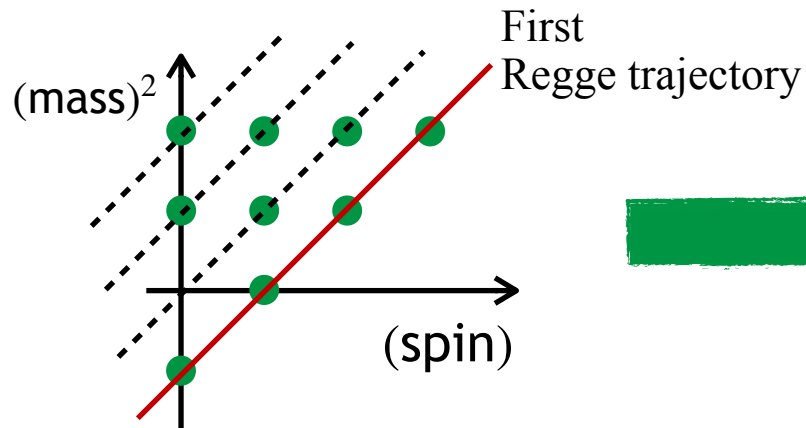


# Introduction

- Strings and Higher spins

## How to explain the higher Regge trajectories?

- String spectrum



# Introduction

- Strings and **Extended** higher spins

**String theory**

**All** Regge trajectory



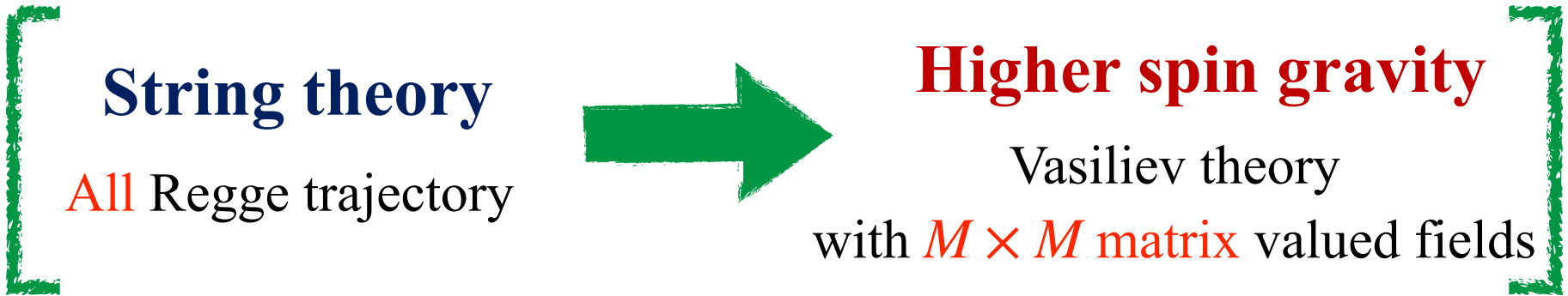
**Higher spin gravity**

Vasiliev theory

with  $M \times M$  **matrix** valued fields

# Introduction

- Strings and **Extended** higher spins



- Matrix extension of 3d Prokushkin-Vasiliev theory may be analyzed with the **infinite dimensional symmetry** of 2d CFT [Creutzig-Hikida '13]

# Introduction

- Strings and **Extended** higher spins

**String theory**

All Regge trajectory



**Higher spin gravity**

Vasiliev theory  
with  $M \times M$  matrix valued fields

- Matrix extension of 3d Prokushkin-Vasiliev theory may be analyzed with the **infinite dimensional symmetry** of 2d CFT [Creutzig-Hikida '13]
  - Dual model is 2d Grassmannian-like coset

$$\frac{su(N+M)_k}{su(N)_k \oplus u(1)_{kNM(N+M)}}$$

With  $M = 1$ , this reduce to  
the Gaberdiel-Gopakumar duality

[Gaberdiel-Gopakumar '10]

- Evidence: spectrum, **asymptotic symmetry**, ...

[Creutzig-Hikida-Rønne '13, Creutzig-Hikida '18]

# Introduction

- Strings and **Extended** higher spins

## Can we generalize this analysis to other models?

- Matrix extension of 3d Prokushkin-Vasiliev theory may be analyzed with the **infinite dimensional symmetry** of 2d CFT [Creutzig-Hikida '13]

- Dual model is 2d Grassmannian-like coset

$$\frac{su(N + M)_k}{su(N)_k \oplus u(1)_{kNM(N+M)}}$$

With  $M = 1$ , this reduce to  
the Gaberdiel-Gopakumar duality

[Gaberdiel-Gopakumar '10]

- Evidence: spectrum, **asymptotic symmetry**, ...

[Creutzig-Hikida-Rønne '13, Creutzig-Hikida '18]




# Introduction

- Our question and summary

## Can we generalize this analysis to other models?

An answer (my talk)

- 
- **We consider 2 ways to truncate the DOF**
    - Restricted matrix extensions;  $so(M)$ ,  $sp(M)$
    - Even spin truncation of  $hs[\lambda]$
  - **We propose the dual coset model and examine the asymptotic symmetry**

# Plan of talk

1. Introduction
2. HS gravity with  $sl(M)$  gauge sector
3. Some generalization for extended HS gravity
4. Summary

# Plan of talk

1. Introduction
2. HS gravity with  $sl(M)$  gauge sector
3. Some generalization for extended HS gravity
4. Summary

# HS gravity with $sl(M)$ gauge sector

- Chern-Simons description of HS gravity
  - 3d gravity:  $sl(2)$  CS theory [Witten '88]

# HS gravity with $sl(M)$ gauge sector

- Chern-Simons description of HS gravity
  - 3d gravity:  $sl(2)$  CS theory [Witten '88]



Include higher spin

- 3d HS gravity: [Prokushkin-Vasiliev '98]

$hs[\lambda]$

$$hs[\lambda] = B[\lambda] \ominus \mathbf{1}$$
$$B[\lambda] = \frac{U(sl(2))}{\langle C_2 - \frac{1}{4}(\lambda^2 - 1)\mathbf{1} \rangle}$$

# HS gravity with $sl(M)$ gauge sector

- Chern-Simons description of HS gravity

- 3d gravity:  $sl(2)$  CS theory [Witten '88]



Include higher spin

- 3d HS gravity: [Prokushkin-Vasiliev '98]

$hs[\lambda] \longrightarrow sl(n)$  CS theory with gravitational  $sl(2)$   
 [ $\lambda = n$  ( $n = 2, 3, \dots$ )]

$$hs[\lambda] = B[\lambda] \ominus \mathbf{1}$$

$$B[\lambda] = \frac{U(sl(2))}{\langle C_2 - \frac{1}{4}(\lambda^2 - 1)\mathbf{1} \rangle}$$

Ex) principal embedding

$$sl(n) = sl(2) \oplus \left( \bigoplus_{s=3}^n g^{(s)} \right)$$

spin-(s-1) representation

# HS gravity with $sl(M)$ gauge sector

- Chern-Simons description of HS gravity

- 3d gravity:  $sl(2)$  CS theory [Witten '88]



Include higher spin

$$hs[\lambda] = B[\lambda] \ominus \mathbf{1}$$

$$B[\lambda] = \frac{U(sl(2))}{\langle C_2 - \frac{1}{4}(\lambda^2 - 1)\mathbf{1} \rangle}$$

- 3d HS gravity: [Prokushkin-Vasiliev '98]

$$hs[\lambda] \longrightarrow sl(n) \text{ CS theory with gravitational } sl(2)$$

$[\lambda = n (n = 2, 3, \dots)]$



Matrix extension

[Gaberdiel-Gopakumar '13, Creutzig-Hikida-Rønne '13]

- 3d HS gravity **with  $M \times M$  fields**:

$$hs_M[\lambda] \simeq gl(M) \otimes B[\lambda] \ominus \mathbf{1}_M \otimes \mathbf{1}$$

$$\simeq sl(M) \otimes \mathbf{1} \oplus \mathbf{1}_M \otimes hs[\lambda] \oplus sl(M) \otimes hs[\lambda]$$

Ex) principal embedding

$$sl(n) = sl(2) \oplus \left( \bigoplus_{s=3}^n g^{(s)} \right)$$

spin-(s-1) representation

# HS gravity with $sl(M)$ gauge sector

- Chern-Simons description of HS gravity

- 3d gravity:  $sl(2)$  CS theory [Witten '88]



Include higher spin

$$hs[\lambda] = B[\lambda] \ominus \mathbf{1}$$

$$B[\lambda] = \frac{U(sl(2))}{\langle C_2 - \frac{1}{4}(\lambda^2 - 1)\mathbf{1} \rangle}$$

- 3d HS gravity: [Prokushkin-Vasiliev '98]

$$hs[\lambda] \longrightarrow sl(n) \text{ CS theory with gravitational } sl(2)$$

$[\lambda = n \ (n = 2, 3, \dots)]$



Matrix extension

[Gaberdiel-Gopakumar '13, Creutzig-Hikida-Rønne '13]

- 3d HS gravity **with  $M \times M$  fields**:

$$hs_M[\lambda] \simeq gl(M) \otimes B[\lambda] \ominus \mathbf{1}_M \otimes \mathbf{1}$$

$$\simeq sl(M) \otimes \mathbf{1} \oplus \mathbf{1}_M \otimes hs[\lambda] \oplus sl(M) \otimes hs[\lambda]$$

$$\longrightarrow sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n) \simeq sl(Mn)$$

$[\lambda = n]$

Include gravitational  $sl(2)$

Ex) principal embedding

$$sl(n) = sl(2) \oplus \left( \bigoplus_{s=3}^n g^{(s)} \right)$$

spin-(s-1) representation



# HS gravity with $sl(M)$ gauge sector

- Chern-Simons description of HS gravity
  - 3d gravity:  $sl(2)$  CS theory [Witten '88]

**We examine  $sl(Mn)$  CS theory decomposed as**

$$sl(Mn) \simeq sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n)$$

- 3d HS gravity with  $M \times M$  fields:

$$hs_M[\lambda] \simeq gl(M) \otimes B[\lambda] \ominus \mathbf{1}_M \otimes \mathbf{1}$$

spin-(s-1) representation

$$\simeq sl(M) \otimes \mathbf{1} \oplus \mathbf{1}_M \otimes hs[\lambda] \oplus sl(M) \otimes hs[\lambda]$$

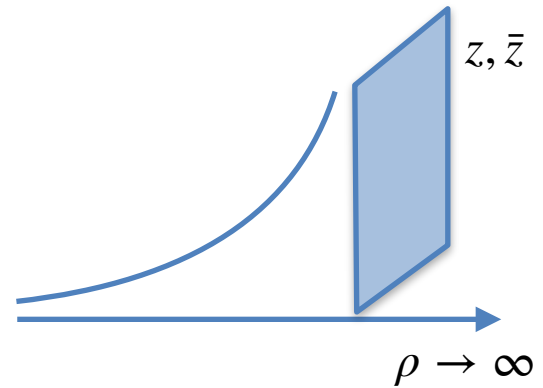
$$\longrightarrow sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n) \simeq sl(Mn)$$

$[\lambda = n]$

Include gravitational  $sl(2)$

# HS gravity with $sl(M)$ gauge sector

- Gauge field (Solution of EOM)
  - 3d gravity:  $A = e^{-\rho L_0} a(z) e^{\rho L_0} dz + L_0 d\rho$



# HS gravity with $sl(M)$ gauge sector

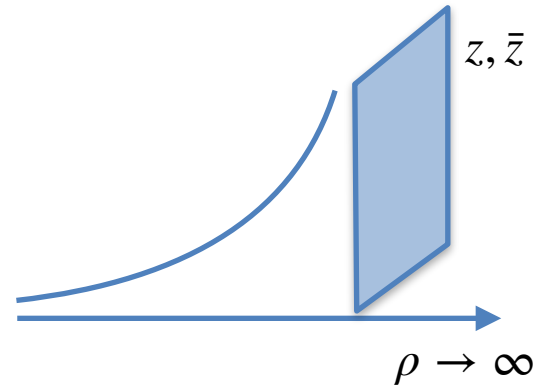
- Gauge field (Solution of EOM)
  - 3d gravity:  $A = e^{-\rho L_0} a(z) e^{\rho L_0} dz + L_0 d\rho$



Include higher spin

- 3d HS gravity:

$$A = e^{-\rho V_0^2} a(z) e^{\rho V_0^2} dz + V_0^2 d\rho$$



Ex)  $s = 2$  case of  $V_{-s+1, \dots, s-1}^s$   
 $V_{0, \pm 1}^2 \equiv L_{0, \pm 1}$

# HS gravity with $sl(M)$ gauge sector

- Gauge field (Solution of EOM)

- 3d gravity:  $A = e^{-\rho L_0} a(z) e^{\rho L_0} dz + L_0 d\rho$



Include higher spin

- 3d HS gravity:

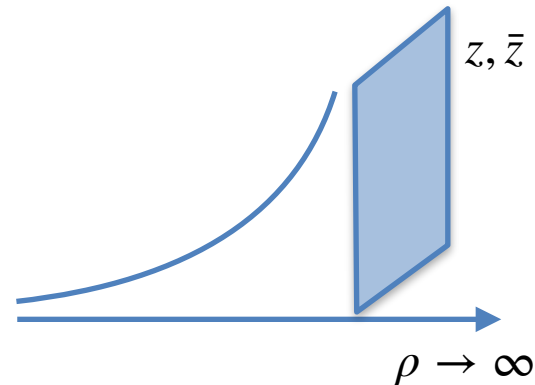
$$A = e^{-\rho V_0^2} a(z) e^{\rho V_0^2} dz + V_0^2 d\rho$$



Matrix extension

- 3d HS gravity **with  $M \times M$  fields**:

$$A = e^{-\rho(\mathbf{1}_M \otimes V_0^2)} a(z) e^{\rho(\mathbf{1}_M \otimes V_0^2)} dz + (\mathbf{1}_M \otimes V_0^2) d\rho$$



Ex)  $s = 2$  case of  $V_{-s+1, \dots, s-1}^s$   
 $V_{0, \pm 1}^2 \equiv L_{0, \pm 1}$

# HS gravity with $sl(M)$ gauge sector

- Asymptotically AdS condition

- 3d gravity:  $A = e^{-\rho L_0} a(z) e^{\rho L_0} dz + L_0 d\rho$



- 3d HS gravity:

$$A = e^{-\rho V_0^2} a(z) e^{\rho V_0^2} dz + V_0^2 d\rho$$



- 3d HS gravity **with  $M \times M$  fields**:

$$A = e^{-\rho(\mathbf{1}_M \otimes V_0^2)} a(z) e^{\rho(\mathbf{1}_M \otimes V_0^2)} dz + (\mathbf{1}_M \otimes V_0^2) d\rho$$

Asymptotically AdS

$$(A - A_{AdS})|_{\rho \rightarrow \infty} = \mathcal{O}((e^\rho)^0)$$

# HS gravity with $sl(M)$ gauge sector

- Asymptotically AdS condition

- 3d gravity:  $A = e^{-\rho L_0} a(z) e^{\rho L_0} dz + L_0 d\rho$

Asymptotically AdS  
 $(A - A_{AdS})|_{\rho \rightarrow \infty} = \mathcal{O}((e^\rho)^0)$



$$a(z) = L_1 + \frac{1}{k_{CS}} T(z) L_{-1}$$

Virasoro generator

- 3d HS gravity:

$$A = e^{-\rho V_0^2} a(z) e^{\rho V_0^2} dz + V_0^2 d\rho$$



- 3d HS gravity **with  $M \times M$  fields**:

$$A = e^{-\rho(\mathbf{1}_M \otimes V_0^2)} a(z) e^{\rho(\mathbf{1}_M \otimes V_0^2)} dz + (\mathbf{1}_M \otimes V_0^2) d\rho$$

# HS gravity with $sl(M)$ gauge sector

- Asymptotically AdS condition

- 3d gravity:  $A = e^{-\rho L_0} \underline{a(z)} e^{\rho L_0} dz + L_0 d\rho$

Asymptotically AdS  
 $(A - A_{AdS})|_{\rho \rightarrow \infty} = \mathcal{O}((e^\rho)^0)$



$a(z) = L_1 + \frac{1}{k_{CS}} T(z) L_{-1}$

- 3d HS gravity:

$A = e^{-\rho V_0^2} \underline{a(z)} e^{\rho V_0^2} dz + V_0^2 d\rho$



$a(z) = V_1^2 + \sum_{s=2}^n W^{(s)}(z) V_{-s+1}^s$

$W_N$  generators

Virasoro generator

- 3d HS gravity **with  $M \times M$  fields**:

$A = e^{-\rho(\mathbf{1}_M \otimes V_0^2)} a(z) e^{\rho(\mathbf{1}_M \otimes V_0^2)} dz + (\mathbf{1}_M \otimes V_0^2) d\rho$

# HS gravity with $sl(M)$ gauge sector

- Asymptotically AdS condition

Asymptotically AdS  
 $(A - A_{AdS})|_{\rho \rightarrow \infty} = \mathcal{O}((e^\rho)^0)$

- 3d gravity:  $A = e^{-\rho L_0} \underline{a(z)} e^{\rho L_0} dz + L_0 d\rho$



$$a(z) = L_1 + \frac{1}{k_{CS}} T(z) L_{-1}$$

Virasoro generator

- 3d HS gravity:

$$A = e^{-\rho V_0^2} \underline{a(z)} e^{\rho V_0^2} dz + V_0^2 d\rho$$



$$a(z) = V_1^2 + \sum_{s=2}^n W^{(s)}(z) V_{-s+1}^s$$

$W_N$  generators

- 3d HS gravity **with  $M \times M$  fields**:

$sl(Mn) \simeq sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n)$

$$A = e^{-\rho(\mathbf{1}_M \otimes V_0^2)} \underline{a(z)} e^{\rho(\mathbf{1}_M \otimes V_0^2)} dz + (\mathbf{1}_M \otimes V_0^2) d\rho$$

$$a(z) = J_a(z) (t^a \otimes \mathbf{1}_n) + \mathbf{1}_M \otimes V_1^2$$

$$+ \sum_{s=2}^n W^{(s)}(z) (\mathbf{1}_M \otimes V_{-s+1}^s) + \sum_{s=2}^n Q_a^{(s)}(z) (t^a \otimes V_{-s+1}^s)$$



# HS gravity with $sl(M)$ gauge sector

- Asymptotically AdS condition

Asymptotically AdS  
 $(A - A_{AdS})|_{\rho \rightarrow \infty} = \mathcal{O}((e^\rho)^0)$

- 3d gravity:  $A = e^{-\rho L_0} a(z) e^{\rho L_0} dz + L_0 d\rho$



$$a(z) = L_1 + \frac{1}{k_{CS}} T(z) L_{-1}$$

Virasoro generator

- 3d HS gravity:

$$A = e^{-\rho V_0^2} a(z) e^{\rho V_0^2} dz + V_0^2 d\rho$$



$$a(z) = V_1^2 + \sum_{s=2}^n W^{(s)}(z) V_{-s+1}^s$$

$W_N$  generators

- 3d HS gravity with  $M \times M$  fields:

$sl(Mn) \simeq sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n)$

$$A = e^{-\rho(\mathbf{1}_M \otimes V_0^2)} a(z) e^{\rho(\mathbf{1}_M \otimes V_0^2)} dz + (\mathbf{1}_M \otimes V_0^2) d\rho$$

Charged higher spin generators

$$a(z) = \underbrace{J_a(z)}_{\text{generators of } sl(M)} (t^a \otimes \mathbf{1}_n) + \mathbf{1}_M \otimes V_1^2 + \sum_{s=2}^n W^{(s)}(z) (\mathbf{1}_M \otimes V_{-s+1}^s) + \sum_{s=2}^n Q_a^{(s)}(z) (t^a \otimes V_{-s+1}^s)$$

$W_N$  generators

# HS gravity with $sl(M)$ gauge sector

- Asymptotically AdS condition

Asymptotically AdS  
 $(A - A_{AdS})|_{\rho \rightarrow \infty} = \mathcal{O}((e^\rho)^0)$

**They are the generators of asymptotic symmetries**



$$a(z) = V_1^2 + \sum_{s=2}^n W^{(s)}(z) V_{-s+1}^s$$

- 3d HS gravity **with  $M \times M$  fields**:

$$A = e^{-\rho(\mathbf{1}_M \otimes V_0^2)} \underline{a(z)} e^{\rho(\mathbf{1}_M \otimes V_0^2)} dz + (\mathbf{1}_M \otimes V_0^2) d\rho$$

Charged higher spin generators

$$a(z) = \underbrace{J_a(z)}_{\text{generators of } sl(M)} (t^a \otimes \mathbf{1}_n) + \mathbf{1}_M \otimes V_1^2 + \sum_{s=2}^n W^{(s)}(z) (\mathbf{1}_M \otimes V_{-s+1}^s) + \sum_{s=2}^n \underbrace{Q_a^{(s)}(z)}_{W_N \text{ generators}} (t^a \otimes V_{-s+1}^s)$$

# HS gravity with $sl(M)$ gauge sector

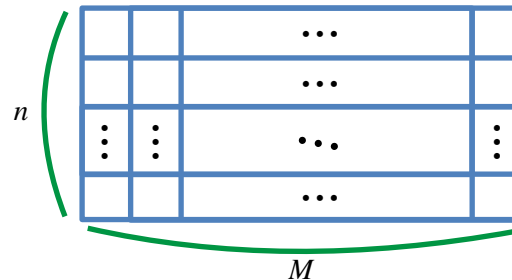
- Asymptotic symmetry

$$sl(Mn) \simeq sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n)$$

- Hamiltonian reduction of  $sl(Mn)$  with the  $sl(2)$  embedding

→ We obtain the **rectangular W-algebra**

Principal embedding correspond to the partition



→ Young diagram  
of rectangular type

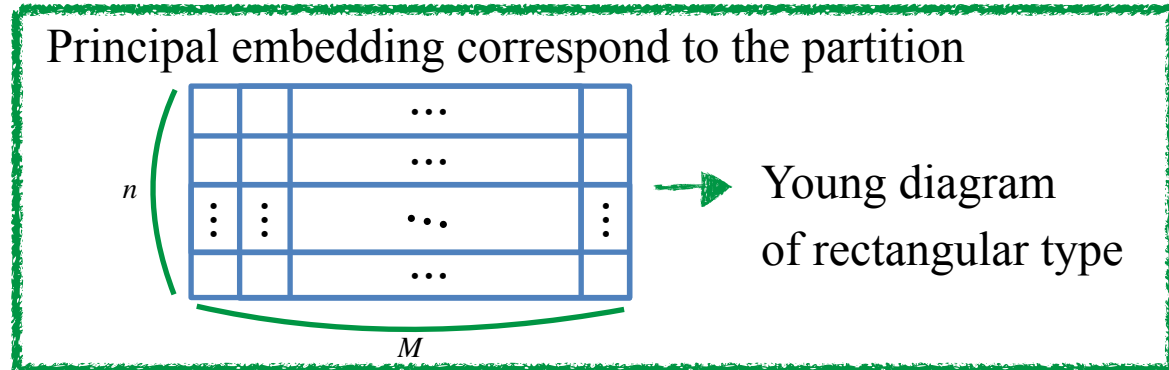
# HS gravity with $sl(M)$ gauge sector

- Asymptotic symmetry

$$sl(Mn) \simeq sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n)$$

- Hamiltonian reduction of  $sl(Mn)$  with the  $sl(2)$  embedding

→ We obtain the **rectangular W-algebra**



- The central charge and level of  $sl(M)$  are evaluated (e.g. [Kac-Wakimoto '03])

- These are consistent with the dual coset model with **finite  $N$**

[Creutzig-Hikida '18]

- We need the proper t' Hooft parameter

$$\lambda = \frac{k}{k + N}$$

# HS gravity with $sl(M)$ gauge sector

- Asymptotic symmetry

$$sl(Mn) \simeq sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n)$$

- Hamiltonian reduction of  $sl(Mn)$  with the  $sl(2)$  embedding

→ We obtain the **rectangular W-algebra**

## Let's turn to the other matrix extension!

- The central charge and level of  $sl(M)$  are evaluated (e.g. [Kac-Wakimoto '03])

- These are consistent with the dual coset model with **finite  $N$**

[Creutzig-Hikida '18]

- We need the proper t' Hooft parameter

$$\lambda = \frac{k}{k+N} \text{ or } \lambda' = -\frac{k}{k+N+M}$$

# Plan of talk

1. Introduction
2. HS gravity with  $sl(M)$  gauge sector
3. Some generalization for extended HS gravity
4. Summary

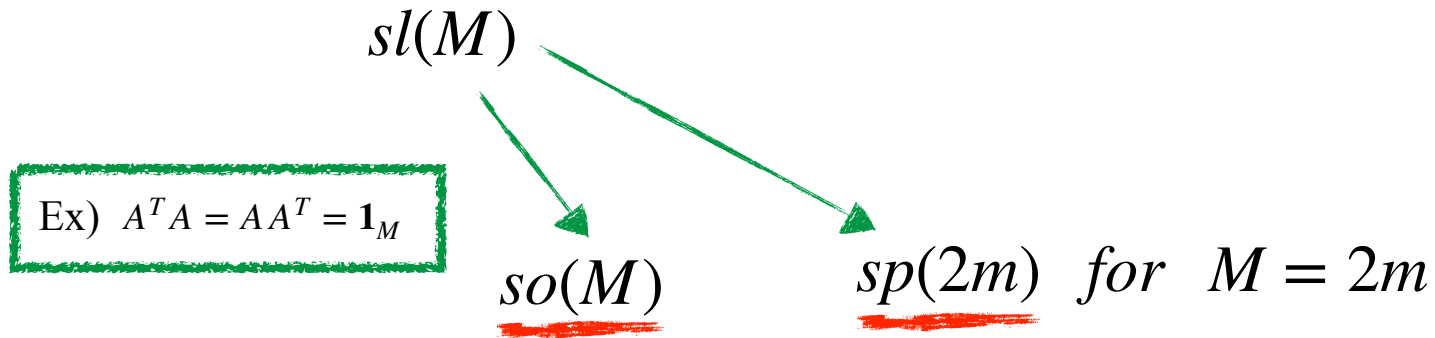
# Some generalization

- Restricted  $M \times M$  matrix extension
  - 2 restrictions on the extra matrix DOF

$$sl(M)$$

# Some generalization

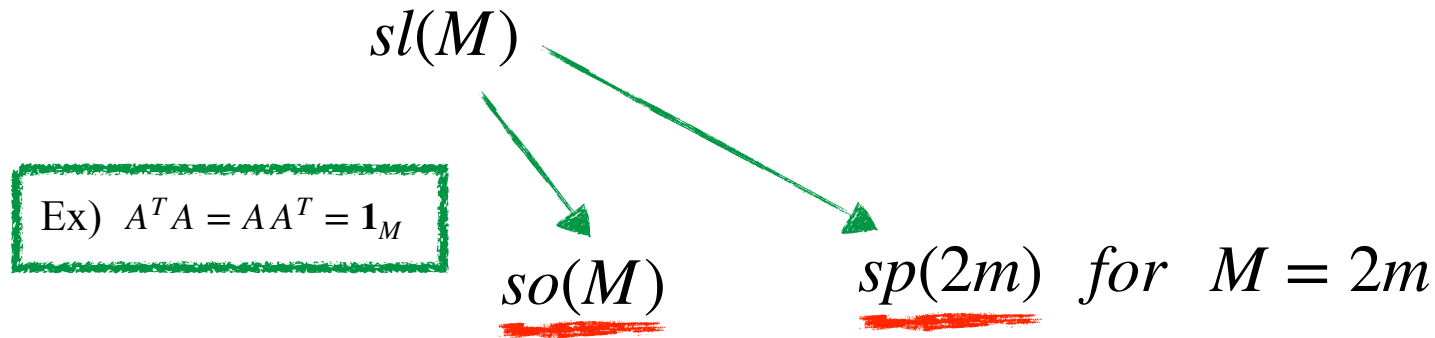
- Restricted  $M \times M$  matrix extension
  - 2 restrictions on the extra matrix DOF





# Some generalization

- Restricted  $M \times M$  matrix extension
  - 2 restrictions on the extra matrix DOF

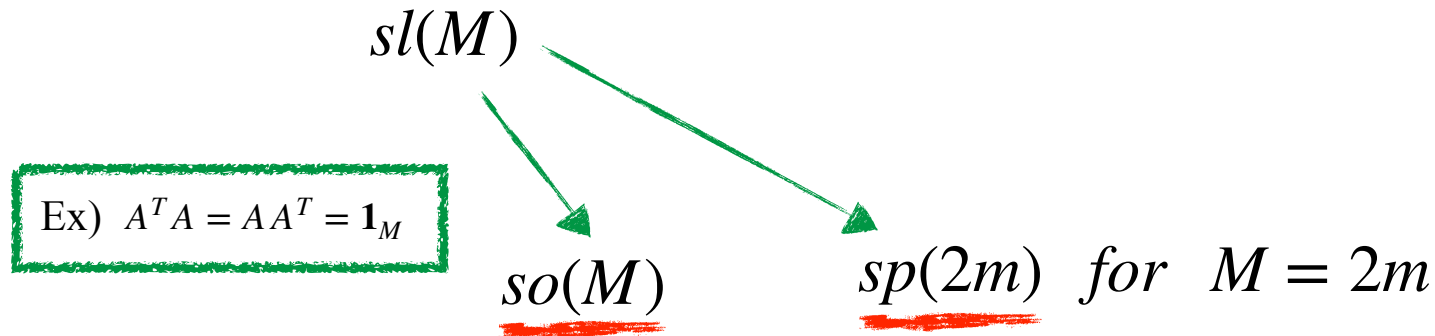


- Even spin truncation
  - $hs[\lambda]$  can be truncated to  $hs^e[\lambda]$

$$hs[\lambda] \simeq hs^e[\lambda] \oplus hs^o[\lambda]$$

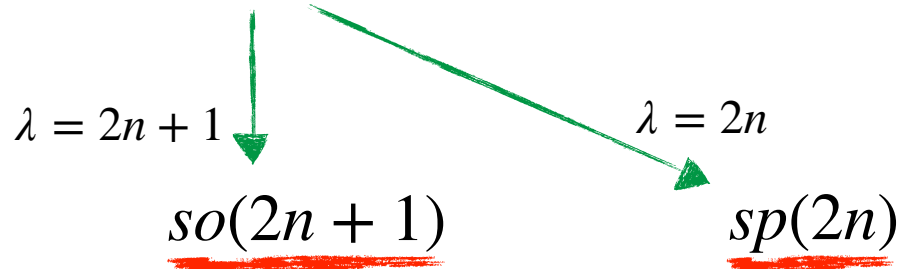
# Some generalization

- Restricted  $M \times M$  matrix extension
  - 2 restrictions on the extra matrix DOF



- Even spin truncation
  - $hs[\lambda]$  can be truncated to  $hs^e[\lambda]$

$$hs[\lambda] \simeq hs^e[\lambda] \oplus hs^o[\lambda]$$



# Some generalization

- Restricted  $M \times M$  matrix extension

## We consider 4 CS gravities

$\begin{array}{c} M \times M \\ \hline hs[\lambda] \end{array}$	$so(M)$	$sp(2m)$
$so(2n + 1)$	$so(M(2n + 1))$	$sp(2m(2n + 1))$
$sp(2n)$	$sp(2Mn)$	$so(4mn)$

Ex) Decomposition of previous model

$$sl(Mn) \simeq sl(M) \otimes \mathbf{1}_n \oplus \mathbf{1}_M \otimes sl(n) \oplus sl(M) \otimes sl(n)$$

$$\underline{so(2n + 1)}$$

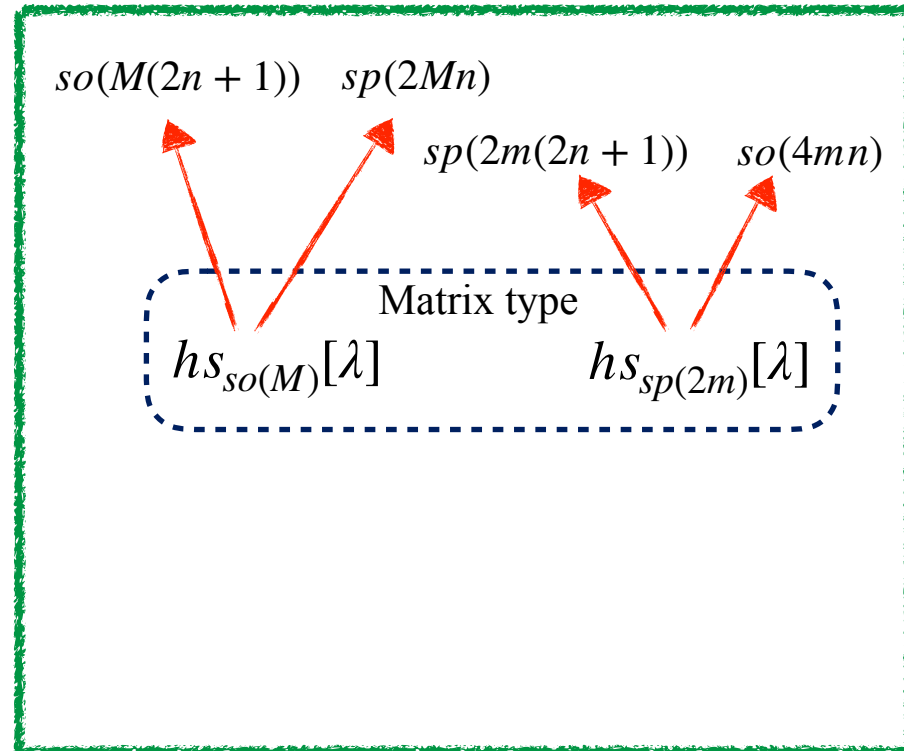
$$\underline{sp(2n)}$$

# Some generalization

- Asymptotic symmetry
  - We obtain 4 types of rectangular W-algebras by Hamiltonian reduction



Their **central charge** and **level** are evaluated



# Some generalization

- Asymptotic symmetry
  - We obtain 4 types of rectangular W-algebras by Hamiltonian reduction

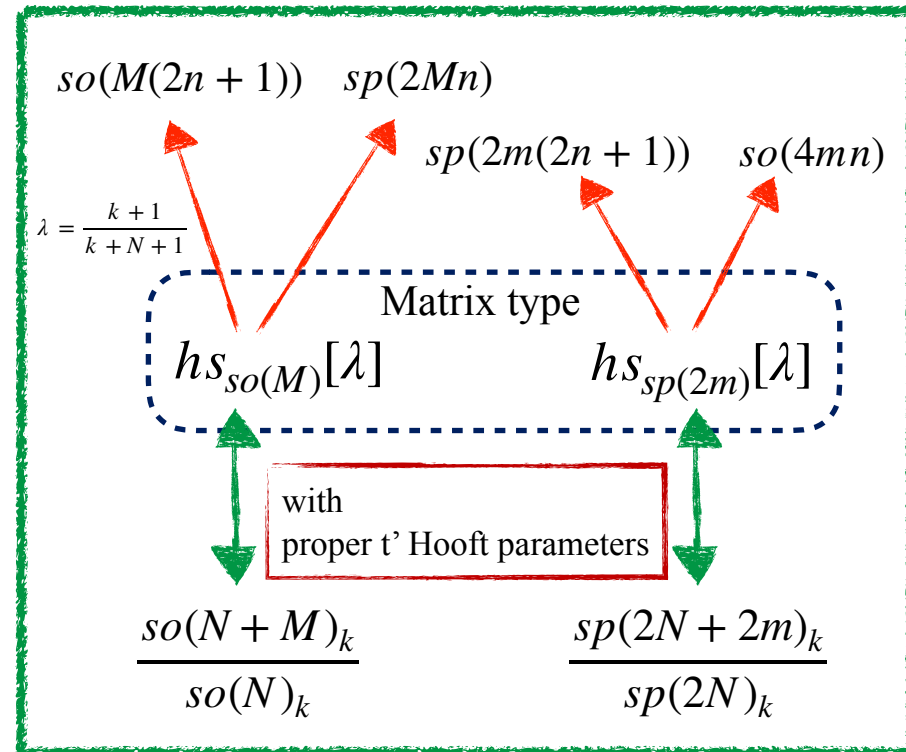


Their **central charge** and **level** are evaluated

- Dual coset model
  - We propose 2 dual coset models



Their **central charge** and **level** are consistent with above algebras with **finite  $N$**



# Some generalization

- Asymptotic symmetry
  - We obtain 4 types of rectangular W-algebras by Hamiltonian reduction

$so(M(2n+1))$   $sp(2Mn)$

## Other checks of the duality



## We focus on the OPEs

$so(N)_k$

$sp(2N)_k$

Their **central charge** and **level** are consistent with above algebras with **finite  $N$**

# Some generalization

- OPEs (for  $\lambda = 2$ )

*sl(2)* with  $M \times M$  matrix

# Some generalization

- OPEs (for  $\lambda = 2$ )

*sl(2)* with  $M \times M$  matrix

- We compute the OPEs among generators of each **algebras**

by requiring **associatively**



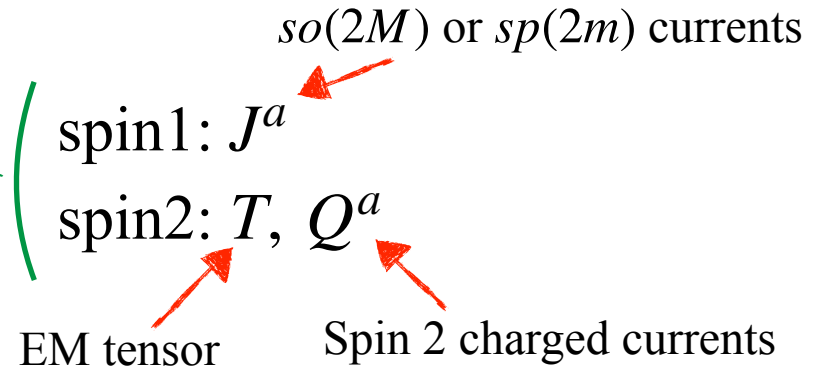
# Some generalization

- OPEs (for  $\lambda = 2$ )

$sl(2)$  with  $M \times M$  matrix

- We compute the OPEs among generators of each **algebras**

by requiring **associatively**



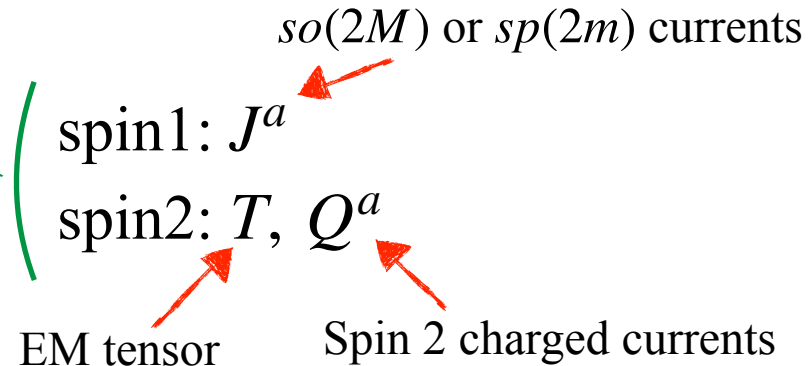
# Some generalization

- OPEs (for  $\lambda = 2$ )

$sl(2)$  with  $M \times M$  matrix

- We compute the OPEs among generators of each **algebras**

by requiring **associatively**



- We compute OPEs among generators of each **coset algebras**

**Above OPEs are reproduced !!**

# Plan of talk


1. Introduction
2. HS gravity with  $sl(M)$  gauge sector
3. Some generalization for extended HS gravity
4. Summary

# Summary

- Our question and summary

## Can we generalize this analysis to other models?

An answer (my talk)

- 
- **We consider 2 ways to truncate the DOF**
    - Restricted matrix extensions;  $so(M)$ ,  $sp(M)$
    - Even spin truncation of  $hs[\lambda]$
  - **We propose the dual coset model and examine the asymptotic symmetry**

## That's all for my presentation

Back up slides

# Some generalization

- $\mathcal{N} = 1$  super symmetric models
  - We analyze 4 types algebras and propose 2 coset models

