

# An Algebraic Model for the $su(2 | 2)$ Light-Cone String Field Theory

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森山翔文氏 (名大KMI) との共同研究に基づく :

I. Kishimoto and S. Moriyama, arXiv:1005.4719 (accepted for publication in JHEP)

# 21世紀の弦の場の理論の進展

- 開弦の場の理論

- 特にタキオン凝縮に応用され進展中...

[QFT02,05,06,07,08,09, FPA-QFT02,05]

- 閉弦の場の理論

- Non-polynomial型：Heterotic版...

- cubic型：冪等方程式, OSp covariantized版...

[QFT03,04]

- light-cone gauge: BMN対応, LCSFT/MST対応...

[FPA-QFT06]

# LCSFT on pp-wave

- AdS/CFT対応の観点で応用された：

- $AdS_5 \times S^5$  のPenrose limit: pp-wave時空

- BMN(Berenstein-Maldacena-Nastase)対応

4次元 $N=4$   $SU(N)$  SYMのalmost BPS operator



pp-wave上の超弦理論のstring state



弦の相互作用を含めて調べる

Spradlin-Volovich, Pankiewicz-Stefanski(2002), Pankiewicz(2003)による  
光円錐ゲージの弦の場の理論(LCSFT)の振動子表示の定式化が使われた

# Toward $su(2|2)$ LCSFT

- flat, pp-wave以外のより一般の背景でのSFTの構成を直接、具体的に振動子レベルでやるのは難しい。
- 代わりにflat, pp-waveの例からSFTのbuilding blockを抜き出し超対称代数で形を決めよう。→「模型」の提案
- pp-wave背景の対称性は $su(2|2)$ を部分代数として含んでいる。
- spin chain模型で $su(2|2)$ 代数が重要な役割を果たした。



$su(2|2)$ 対称性をもつ背景上のLCSFTの模型

pp-wave上のSFTをより簡潔に再現、  
一般化に向けた計算例

# Contents

- Introduction ✓
- Review of flat LCSFT
- Algebraic model for  $\mathfrak{su}(2|2)$  LCSFT
  - Building blocks, ansatz, assumptions
  - Solutions
- Summary and discussion

# Green-Schwarz形式の光円錐ゲージ

- Green-Schwarz形式光円錐ゲージの弦の座標、運動量

$$x^i(\sigma), \vartheta^a(\sigma) \quad p^i(\sigma), \lambda^a(\sigma) \quad : \text{モード展開される}$$

- 超電荷とハミルトニアン (free part)

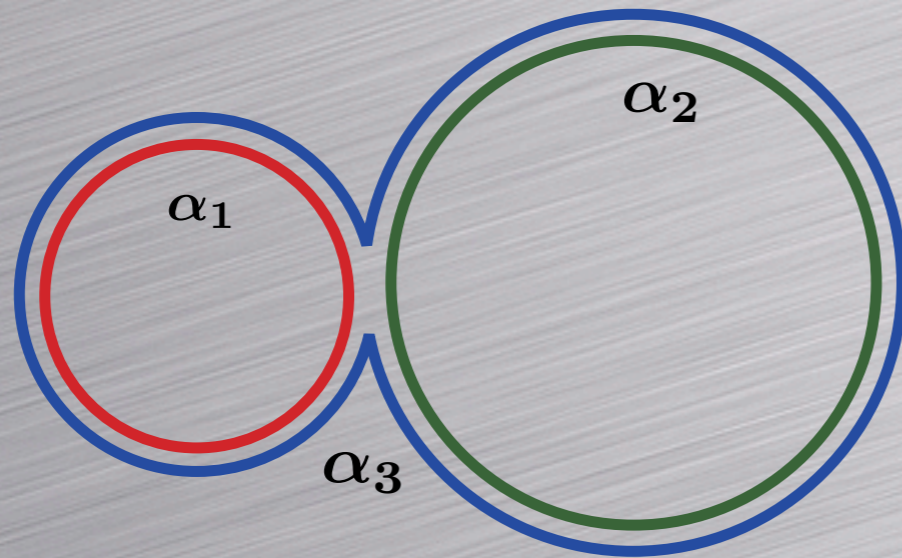
$$q^{\dot{a}} = \frac{\sqrt{2}}{\alpha} \sum_{n=-\infty}^{\infty} \gamma_{a\dot{a}}^i Q_{-n}^a \alpha_n^i,$$

$$\tilde{q}^{\dot{a}} = \frac{\sqrt{2}}{\alpha} \sum_{n=-\infty}^{\infty} \gamma_{a\dot{a}}^i \tilde{Q}_{-n}^a \tilde{\alpha}_n^i,$$

$$h = \frac{1}{\alpha} p^i p^i + \frac{1}{\alpha} \sum_{n=1}^{\infty} \left( \alpha_{-n}^i \alpha_n^i + \tilde{\alpha}_{-n}^i \tilde{\alpha}_n^i + \frac{n}{\alpha} (Q_{-n}^a Q_n^a + \tilde{Q}_{-n}^a \tilde{Q}_n^a) \right).$$

- $i, a, \dot{a}$  はSO(8)の  $\mathfrak{8}_v, \mathfrak{8}_s, \mathfrak{8}_c$

# Kinematical overlap



:光円錐型の3弦相互作用

3弦の接続条件を表す $\delta$ -汎関数:

$$\delta(\alpha_1 + \alpha_2 + \alpha_3) \delta^8(x^{i(3)} - \Theta_1 x^{i(1)} - \Theta_2 x^{i(2)}) \delta^8(\vartheta^{i(3)} - \Theta_1 \vartheta^{i(1)} - \Theta_2 \vartheta^{i(2)}) \\ = \langle \alpha_1, x^{(1)}, \vartheta^{(1)} | \langle \alpha_2, x^{(2)}, \vartheta^{(2)} | \langle \alpha_3, x^{(3)}, \vartheta^{(3)} | V \rangle$$



振動子を用いてケット  $|V\rangle$  で表現できる

# Algebra and prefactors

- SUSY代数:  $\{Q^{\dot{a}}, Q^{\dot{b}}\} = \{\tilde{Q}^{\dot{a}}, \tilde{Q}^{\dot{b}}\} = 2H\delta^{\dot{a}\dot{b}}, \{Q^{\dot{a}}, \tilde{Q}^{\dot{b}}\} = 0$

- 非自明な最低次 (3弦相互作用項) の部分:

$$q^{\dot{a}}|Q^{\dot{b}}\rangle + q^{\dot{b}}|Q^{\dot{a}}\rangle = \tilde{q}^{\dot{a}}|\tilde{Q}^{\dot{b}}\rangle + \tilde{q}^{\dot{b}}|\tilde{Q}^{\dot{a}}\rangle = 2\delta^{\dot{a}\dot{b}}|H\rangle, \quad q^{\dot{a}}|\tilde{Q}^{\dot{b}}\rangle + \tilde{q}^{\dot{b}}|Q^{\dot{a}}\rangle = 0$$



up to level matching conditionで満たす解

Green-Schwarz-Brinkの公式 (の簡潔版 [I.K.-Moriyama(2006)])

$$|H\rangle = X^i \tilde{X}^j [\cosh \mathcal{Y}]^{ij} |V\rangle,$$

$$|Q^{\dot{a}}\rangle = \sqrt{-\alpha_{123}} \tilde{X}^i [\sinh \mathcal{Y}]^{\dot{a}i} |V\rangle,$$

$$|\tilde{Q}^{\dot{a}}\rangle = i\sqrt{-\alpha_{123}} X^i [\sinh \mathcal{Y}]^{i\dot{a}} |V\rangle$$

$$\mathcal{Y} = \sqrt{\frac{2}{-\alpha_{123}}} \eta^* Y^a \hat{\gamma}^a$$

$$\hat{\gamma}^a = \begin{pmatrix} 0 & \hat{\gamma}_{i\dot{a}}^a \\ \hat{\gamma}_{\dot{a}i}^a & 0 \end{pmatrix} = \begin{pmatrix} 0 & \gamma_{a\dot{a}}^i \\ \gamma_{\dot{a}a}^i & 0 \end{pmatrix}$$



# Building blocks (1)

- 3 弦相互作用項の構成要素:  $|V\rangle$   $Y^a$ ,  $X^i$ ,  $\widetilde{X}^i$

それぞれ振動子とノイマン係数を用いてあらわに表現されている:

$$|V\rangle = (2\pi)^9 \delta(\alpha_1 + \alpha_2 + \alpha_3) \delta^8(p_1^i + p_2^i + p_3^i) \delta^8(\lambda_1^a + \lambda_2^a + \lambda_3^a) \\ \times e^{\frac{1}{2} \sum \bar{N}_{nm}^{rs} (\alpha_{-n}^{(r)} \alpha_{-m}^{(s)} + \tilde{\alpha}_{-n}^{(r)} \tilde{\alpha}_{-m}^{(s)})} + \sum \bar{N}_n^r (\alpha_{-n}^{(r)} + \tilde{\alpha}_{-n}^{(r)}) P - \frac{\tau_0}{\alpha_{123}} P^2} \\ \times e^{\sum Q_{-n}^{\text{II}(r)} \alpha_r^{-1} n \bar{N}_{nm}^{rs} Q_{-m}^{\text{I}(s)} - \sqrt{2} \Lambda \sum \alpha_r^{-1} n \bar{N}_n^r Q_{-n}^{\text{II}(r)}} |0\rangle$$

$$P^i = \alpha_1 p_2^i - \alpha_2 p_1^i, \quad \Lambda^a = \alpha_1 \lambda_2^a - \alpha_2 \lambda_1^a, \quad Q_{-n}^{\text{I/II}a} = \frac{1}{\sqrt{2}} (\eta^{\pm 1} Q_{-n}^a + \eta^{\mp 1} \tilde{Q}_{-n}^a) \quad \alpha_{123} \equiv \alpha_1 \alpha_2 \alpha_3$$

$$\bar{N}_{nm}^{rs} = -\frac{\alpha_{123}}{\alpha_r/n + \alpha_s/m} \bar{N}_n^r \bar{N}_m^s, \quad \bar{N}_n^r = \frac{\Gamma(-n\alpha_{r+1}/\alpha_r) e^{n\tau_0/\alpha_r}}{\alpha_r n! \Gamma(1 - n(1 + \alpha_{r+1}/\alpha_r))}, \quad (\alpha_4 \equiv \alpha_1), \quad \tau_0 = \sum_{r=1}^3 \alpha_r \log |\alpha_r|$$

$$Y^a = \Lambda^a - \frac{\alpha_{123}}{2} \sum_{r=1}^3 \sum_{n=1}^{\infty} \alpha_r^{-1} n \bar{N}_n^r (\eta Q_{-n}^{(r)a} + \eta^* \tilde{Q}_{-n}^{(r)a})$$

$$X^i = P^i - \alpha_{123} \sum_{r=1}^3 \sum_{n=1}^{\infty} \alpha_r^{-1} n \bar{N}_n^r \alpha_{-n}^{i(r)}, \quad \widetilde{X}^i = P^i - \alpha_{123} \sum_{r=1}^3 \sum_{n=1}^{\infty} \alpha_r^{-1} n \bar{N}_n^r \tilde{\alpha}_{-n}^{i(r)}.$$

# Building blocks (2)

- 3 弦相互作用項の構成要素:  $|V\rangle$ ,  $Y^a$ ,  $X^i$ ,  $\widetilde{X}^i$

実は... 
$$\lambda^a(\sigma)|V\rangle \sim \frac{1}{4\pi|\alpha_{123}|^{1/2}} \frac{1}{|\sigma - \sigma_{\text{int}}|^{1/2}} Y^a |V\rangle$$

$$\{q^{\dot{a}}, Y^a\} = \frac{1}{\sqrt{2}} \eta \gamma_{a\dot{a}}^i X^i, \quad \{\tilde{q}^{\dot{a}}, Y^a\} = \frac{1}{\sqrt{2}} \eta^* \gamma_{a\dot{a}}^i \widetilde{X}^i$$

またハミルトニアンには  $[q, X] \sim \partial Y$ ,  $[\tilde{q}, \widetilde{X}] \sim \bar{\partial} Y$

を用いた不定性も入りうる: [Lee-Russo(2004), Dobashi-I.K.-Moriyama(unpublished)]

$$|Q^{\dot{a}}\rangle = \left( \tilde{f}_0 \widetilde{X}^i [\sinh Y]^{i\dot{a}} + X^i \left[ f_1 Y + \frac{1}{7!} f_7 Y^7 \right]^{i\dot{a}} \right) |V\rangle = \left( \tilde{f}_0 \widetilde{X}^i [\sinh Y]^{i\dot{a}} + q^{\dot{a}i} \sqrt{-\alpha_{123}} (f_1 + f_7 y_0^8 \delta^8(Y)) \right) |V\rangle$$

$$|\tilde{Q}^{\dot{a}}\rangle = \left( g_0 [\sinh Y]_{\dot{a}i} X^i + \left[ \tilde{g}_1 Y + \frac{1}{7!} \tilde{g}_7 Y^7 \right]^{\dot{a}i} \widetilde{X}^i \right) |V\rangle = \left( g_0 [\sinh Y]_{\dot{a}i} X^i + \tilde{q}^{\dot{a}i} \sqrt{-\alpha_{123}} (\tilde{g}_1 + \tilde{g}_7 y_0^8 \delta^8(Y)) \right) |V\rangle$$

$$|H\rangle = \left( \frac{1}{\sqrt{-\alpha_{123}}} g_0 \widetilde{X}^i X^j [\cosh Y]^{ij} + h \sqrt{-\alpha_{123}} (\tilde{g}_1 + y_0^8 \tilde{g}_7 \delta^8(Y)) \right) |V\rangle$$



# Algebraic model

- ここで扱う超代数：  $\mathfrak{psu}(2|2) \times \mathbb{R}^3$

$$\begin{aligned}\{Q^\alpha_a, Q^\beta_b\} &= \epsilon^{\alpha\beta} \epsilon_{ab} \mathcal{H}, \\ \{\tilde{Q}^\alpha_a, \tilde{Q}^\beta_b\} &= \epsilon^{\alpha\beta} \epsilon_{ab} \tilde{\mathcal{H}}, \\ \{Q^\alpha_a, \tilde{Q}^\beta_b\} &= \epsilon^{\alpha\beta} \epsilon_{ac} \mathcal{R}^c_b + \epsilon_{ab} \mathcal{L}^\alpha_\gamma \epsilon^{\gamma\beta} + \epsilon^{\alpha\beta} \epsilon_{ab} \mathcal{N}.\end{aligned}$$

$\mathcal{J} = j + g_s |J\rangle + \dots$  のように展開し free 部分は知っているとし、

$$\begin{aligned}q^\alpha_a |Q^\beta_b\rangle + q^\beta_b |Q^\alpha_a\rangle &= \epsilon^{\alpha\beta} \epsilon_{ab} |H\rangle, \\ \tilde{q}^\alpha_a |\tilde{Q}^\beta_b\rangle + \tilde{q}^\beta_b |\tilde{Q}^\alpha_a\rangle &= \epsilon^{\alpha\beta} \epsilon_{ab} |\tilde{H}\rangle, \\ q^\alpha_a |\tilde{Q}^\beta_b\rangle + \tilde{q}^\beta_b |Q^\alpha_a\rangle &= \epsilon^{\alpha\beta} \epsilon_{ab} |N\rangle.\end{aligned}$$

を満たす「3弦相互作用項」を構成しよう！

$\mathcal{L}^\alpha_\beta = l^\alpha_\beta$ ,  $\mathcal{R}^a_b = r^a_b$  は kinematical だと仮定している。

# Building blocks

- 3弦相互作用項  $|H\rangle, |\widetilde{H}\rangle, |N\rangle, |Q^\alpha_a\rangle, |\widetilde{Q}^\alpha_a\rangle$

を構成するもの (仮定)

$|V\rangle$  : kinematical overlap  $\sim$  弦の接続条件を表す $\delta$ -汎函数

$$Y^\alpha_{\dot{b}}, Y'^a_{\dot{\beta}}, X^a_{\dot{b}}, \widetilde{X}^a_{\dot{b}}, X'^\alpha_{\dot{\beta}}, \widetilde{X}'^\alpha_{\dot{\beta}}, W^\alpha_{\dot{b}}, \widetilde{W}^\alpha_{\dot{b}}, W'^a_{\dot{\beta}}, \widetilde{W}'^a_{\dot{\beta}}$$

: prefactorの構成要素  $(\alpha, \dot{\alpha}, a, \dot{a}, \dots = 1, 2)$

$Y^\alpha_{\dot{b}}, Y'^a_{\dot{\beta}}$   $\sim$  fermion運動量を $|V\rangle$ にかけ相互作用点に近づけ発散を取り除いた部分

交換関係:

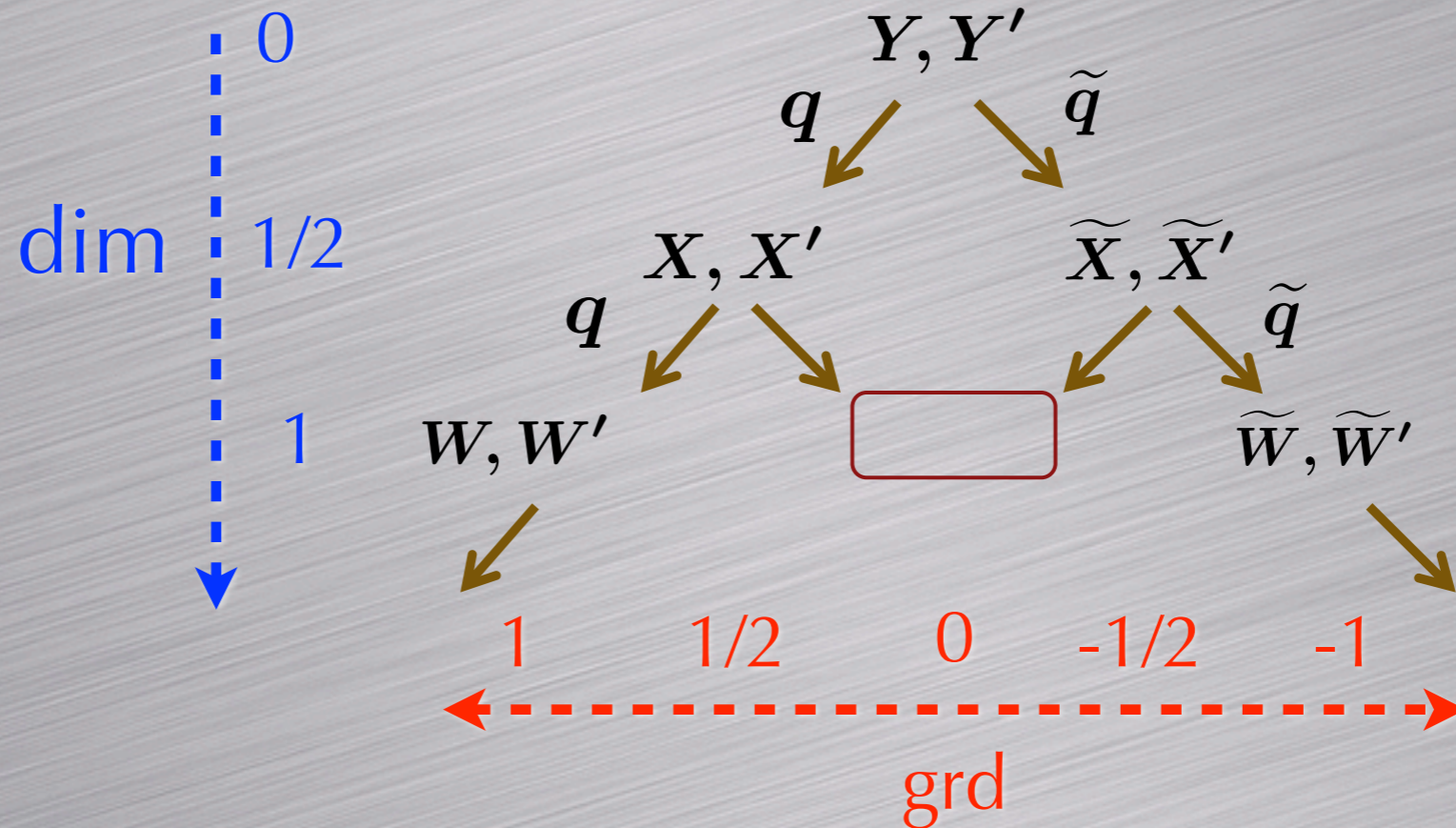
$$\begin{aligned} \{q^\alpha_a, Y^\beta_{\dot{b}}\} &= -\epsilon^{\beta\alpha} (\epsilon X)_{a\dot{b}}, & \{q^\alpha_a, Y'^b_{\dot{\beta}}\} &= \delta_a^b X'^\alpha_{\dot{\beta}}, & \sim X, X' \text{の定義} \\ [q^\alpha_a, X^b_{\dot{b}}] &= \delta_a^b W^\alpha_{\dot{b}}, & [q^\alpha_a, X'^\beta_{\dot{\beta}}] &= -\epsilon^{\beta\alpha} (\epsilon W')_{a\dot{\beta}}. & \sim W, W' \text{の定義} \end{aligned}$$

$\widetilde{q}^\alpha_a$  との交換関係も同様

※さらに  $q^\alpha_a |V\rangle, [q^\alpha_a, \widetilde{X}^b_{\dot{b}}], [q^\alpha_a, \widetilde{X}'^\beta_{\dot{\beta}}]$  の形を決めておく必要がある!

# Dimension and grade

- dim と grd を assign する:



- consistency

$$\{[q^\alpha_a, \tilde{X}^c_c], q^\beta_b\} + \{q^\alpha_a, [q^\beta_b, \tilde{X}^c_c]\} = \epsilon^{\alpha\beta} \epsilon_{ab} [h, \tilde{X}^c_c],$$

$$\{q^\alpha_a, q^\beta_b\} |V\rangle = \epsilon^{\alpha\beta} \epsilon_{ab} h |V\rangle, \quad \{q^\alpha_a, \tilde{q}^\beta_b\} |V\rangle = \epsilon^{\alpha\beta} \epsilon_{ab} n |V\rangle, \quad \text{等は満たすように決める}$$

$$r^a_b |V\rangle = l^a_\beta |V\rangle = 0 \quad \text{は仮定}$$

# Ansatz

$$q^{\alpha}_a |V\rangle = \sum_{n,m} \left\{ v_{nm} (Y^n X Y'^m)^{\alpha}_a + v'_{mn} (Y^m X' Y'^n)^{\alpha}_a \right\} |V\rangle$$

up to dim 1/2,  
grd=1/2 (12項)

$$\begin{aligned} [q^{\alpha}_a, \tilde{X}^b_b] = & v_1 \delta_a^b Y^{\alpha}_b + v_2 \delta_a^b (Y^3)^{\alpha}_b + v_3 Y^{\alpha}_b (Y'^2)^b_a + v_4 \delta_a^b Y^{\alpha}_b (Y'^4)^c_c + v_5 (Y^3)^{\alpha}_b (Y'^2)^b_a + v_6 \delta_a^b (Y^3)^{\alpha}_b (Y'^4)^c_c + v_7 Y^{\alpha}_b (\tilde{X} X)^b_a \\ & + v_8 Y^{\alpha}_b (X \tilde{X})^b_a + v_9 (Y X \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{10} (Y X)^{\alpha}_a \tilde{X}^b_b + v_{11} (Y \tilde{X} \epsilon)^{ab} (\epsilon X)_{ab} + v_{12} (Y \tilde{X})^{\alpha}_a X^b_b + v_{13} (Y^3)^{\alpha}_b (\tilde{X} X)^b_a \\ & + v_{14} (Y^3)^{\alpha}_b (X \tilde{X})^b_a + v_{15} (Y^3 X \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{16} (Y^3 X)^{\alpha}_a \tilde{X}^b_b + v_{17} (Y^3 \tilde{X} \epsilon)^{ab} (\epsilon X)_{ab} + v_{18} (Y^3 \tilde{X})^{\alpha}_a X^b_b \\ & + v_{19} Y^{\alpha}_b (Y'^2)^b_a (X \tilde{X})^c_c + v_{20} Y^{\alpha}_b (\tilde{X} X Y'^2)^b_a + v_{21} Y^{\alpha}_b (Y'^2 X \tilde{X})^b_a + v_{22} Y^{\alpha}_b (X \tilde{X} Y'^2)^b_a + v_{23} Y^{\alpha}_b (Y'^2 \tilde{X} X)^b_a \\ & + v_{24} (Y X \tilde{X})^{\alpha}_b (Y'^2)^b_a + v_{25} (Y X Y'^2)^{\alpha}_a \tilde{X}^b_b + v_{26} (Y X Y'^2 \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{27} (Y \tilde{X} X)^{\alpha}_b (Y'^2)^b_a + v_{28} (Y \tilde{X} Y'^2)^{\alpha}_a X^b_b \\ & + v_{29} (Y \tilde{X} Y'^2 \epsilon)^{ab} (\epsilon X)_{ab} + v_{30} (Y X)^{\alpha}_a (Y'^2 \tilde{X})^b_b + v_{31} (Y X \epsilon)^{ab} (\epsilon Y'^2 \tilde{X})_{ab} + v_{32} (Y \tilde{X})^{\alpha}_a (Y'^2 X)^b_b + v_{33} (Y \tilde{X} \epsilon)^{ab} (\epsilon Y'^2 X)_{ab} \\ & + v_{34} Y^{\alpha}_b (\tilde{X} X)^b_a (Y'^4)^c_c + v_{35} Y^{\alpha}_b (X \tilde{X})^b_a (Y'^4)^c_c + v_{36} (Y X \epsilon)^{ab} (\epsilon \tilde{X})_{ab} (Y'^4)^c_c + v_{37} (Y X)^{\alpha}_a \tilde{X}^b_b (Y'^4)^c_c \\ & + v_{38} (Y \tilde{X} \epsilon)^{ab} (\epsilon X)_{ab} (Y'^4)^c_c + v_{39} (Y \tilde{X})^{\alpha}_a X^b_b (Y'^4)^c_c + v_{40} (Y^3)^{\alpha}_b (Y'^2)^b_a (X \tilde{X})^c_c + v_{41} (Y^3)^{\alpha}_b (\tilde{X} X Y'^2)^b_a \\ & + v_{42} (Y^3)^{\alpha}_b (Y'^2 X \tilde{X})^b_a + v_{43} (Y^3)^{\alpha}_b (X \tilde{X} Y'^2)^b_a + v_{44} (Y^3)^{\alpha}_b (Y'^2 \tilde{X} X)^b_a + v_{45} (Y^3 X \tilde{X})^{\alpha}_b (Y'^2)^b_a + v_{46} (Y^3 X Y'^2)^{\alpha}_a \tilde{X}^b_b \\ & + v_{47} (Y^3 X Y'^2 \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{48} (Y^3 \tilde{X} X)^{\alpha}_b (Y'^2)^b_a + v_{49} (Y^3 \tilde{X} Y'^2)^{\alpha}_a X^b_b + v_{50} (Y^3 \tilde{X} Y'^2 \epsilon)^{ab} (\epsilon X)_{ab} + v_{51} (Y^3 X)^{\alpha}_a (Y'^2 \tilde{X})^b_b \\ & + v_{52} (Y^3 X \epsilon)^{ab} (\epsilon Y'^2 \tilde{X})_{ab} + v_{53} (Y^3 \tilde{X})^{\alpha}_a (Y'^2 X)^b_b + v_{54} (Y^3 \tilde{X} \epsilon)^{ab} (\epsilon Y'^2 X)_{ab} + v_{55} (Y^3)^{\alpha}_b (\tilde{X} X)^b_a (Y'^4)^c_c \\ & + v_{56} (Y^3)^{\alpha}_b (X \tilde{X})^b_a (Y'^4)^c_c + v_{57} (Y^3 X \epsilon)^{ab} (\epsilon \tilde{X})_{ab} (Y'^4)^c_c + v_{58} (Y^3 X)^{\alpha}_a \tilde{X}^b_b (Y'^4)^c_c + v_{59} (Y^3 \tilde{X} \epsilon)^{ab} (\epsilon X)_{ab} (Y'^4)^c_c \\ & + v_{60} (Y^3 \tilde{X})^{\alpha}_a X^b_b (Y'^4)^c_c + v_{61} Y^{\alpha}_b \delta_a^b (X' \tilde{X}')^{\beta}_{\beta} + v_{62} (X' \tilde{X}' Y)^{\alpha}_b \delta_a^b + v_{63} (\tilde{X}' X' Y)^{\alpha}_b \delta_a^b + v_{64} (Y^3)^{\alpha}_b \delta_a^b (X' \tilde{X}')^{\beta}_{\beta} \\ & + v_{65} (X' \tilde{X}' Y^3)^{\alpha}_b \delta_a^b + v_{66} (\tilde{X}' X' Y^3)^{\alpha}_b \delta_a^b + v_{67} Y^{\alpha}_b (Y'^2)^b_a (X' \tilde{X}')^{\beta}_{\beta} + v_{68} Y^{\alpha}_b \delta_a^b (X' \tilde{X}')^{\beta}_{\beta} (Y'^4)^c_c + v_{69} (X' \tilde{X}' Y)^{\alpha}_b \delta_a^b (Y'^4)^c_c \\ & + v_{70} (\tilde{X}' X' Y)^{\alpha}_b \delta_a^b (Y'^4)^c_c + v_{71} (Y^3)^{\alpha}_b (Y'^2)^b_a (X' \tilde{X}')^{\beta}_{\beta} + v_{72} (Y^3)^{\alpha}_b \delta_a^b (X' \tilde{X}')^{\beta}_{\beta} (Y'^4)^c_c + v_{73} (X' \tilde{X}' Y^3)^{\alpha}_b \delta_a^b (Y'^4)^c_c \\ & + v_{74} (\tilde{X}' X' Y^3)^{\alpha}_b \delta_a^b (Y'^4)^c_c + v_{75} (\tilde{X}' Y' \epsilon)^{ab} (\epsilon X)_{ab} + v_{76} (\tilde{X}' Y')^{\alpha}_a X^b_b + v_{77} (X' Y' \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{78} (X' Y')^{\alpha}_a \tilde{X}^b_b \\ & + v_{79} (\tilde{X}' Y'^3 \epsilon)^{ab} (\epsilon X)_{ab} + v_{80} (\tilde{X}' Y'^3)^{\alpha}_a X^b_b + v_{81} (X' Y'^3 \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{82} (X' Y'^3)^{\alpha}_a \tilde{X}^b_b + v_{83} (\tilde{X}' Y' \epsilon)^{ab} (\epsilon X)_{ab} (Y^4)^{\beta}_{\beta} \\ & + v_{84} (\tilde{X}' Y')^{\alpha}_a X^b_b (Y^4)^{\beta}_{\beta} + v_{85} (X' Y' \epsilon)^{ab} (\epsilon \tilde{X})_{ab} (Y^4)^{\beta}_{\beta} + v_{86} (X' Y')^{\alpha}_a \tilde{X}^b_b (Y^4)^{\beta}_{\beta} + v_{87} (\tilde{X}' Y'^3 \epsilon)^{ab} (\epsilon X)_{ab} (Y^4)^{\beta}_{\beta} \\ & + v_{88} (\tilde{X}' Y'^3)^{\alpha}_a X^b_b (Y^4)^{\beta}_{\beta} + v_{89} (X' Y'^3 \epsilon)^{ab} (\epsilon \tilde{X})_{ab} (Y^4)^{\beta}_{\beta} + v_{90} (X' Y'^3)^{\alpha}_a \tilde{X}^b_b (Y^4)^{\beta}_{\beta} + v_{91} (Y^2 \tilde{X}' Y' \epsilon)^{ab} (\epsilon X)_{ab} \\ & + v_{92} (Y^2 \tilde{X}' Y')^{\alpha}_a X^b_b + v_{93} (\tilde{X}' Y' \epsilon)^{ab} (\epsilon X Y^2)_{ab} + v_{94} (\tilde{X}' Y')^{\alpha}_a (X Y^2)^b_b + v_{95} (Y^2 X' Y' \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{96} (Y^2 X' Y')^{\alpha}_a \tilde{X}^b_b \\ & + v_{97} (X' Y' \epsilon)^{ab} (\epsilon \tilde{X} Y^2)_{ab} + v_{98} (X' Y')^{\alpha}_a (\tilde{X} Y^2)^b_b + v_{99} (Y^2 \tilde{X}' Y'^3 \epsilon)^{ab} (\epsilon X)_{ab} + v_{100} (Y^2 \tilde{X}' Y'^3)^{\alpha}_a X^b_b \\ & + v_{101} (\tilde{X}' Y'^3 \epsilon)^{ab} (\epsilon X Y^2)_{ab} + v_{102} (\tilde{X}' Y'^3)^{\alpha}_a (X Y^2)^b_b + v_{103} (Y^2 X' Y'^3 \epsilon)^{ab} (\epsilon \tilde{X})_{ab} + v_{104} (Y^2 X' Y'^3)^{\alpha}_a \tilde{X}^b_b \\ & + v_{105} (X' Y'^3 \epsilon)^{ab} (\epsilon \tilde{X} Y^2)_{ab} + v_{106} (X' Y'^3)^{\alpha}_a (\tilde{X} Y^2)^b_b \end{aligned}$$

up to dim 1,  
grd=0 (106項)

$[q^{\alpha}_a, \tilde{X}'^{\beta}_{\beta}]$  も同様に106項

※consistencyから係数に関係がつく

# Linearized version

- $Y, Y'$ について線形な場合：
 
$$q^\alpha_a |V\rangle = \frac{i}{2} [(YX)^\alpha_a + (X'Y')^\alpha_a] |V\rangle$$

$$[q^\alpha_a, \widetilde{X}^b_b] = \frac{i}{2} \delta_a^b Y^\alpha_b, \quad [q^\alpha_a, \widetilde{X}'^\beta_\beta] = \frac{i}{2} \epsilon^{\beta\alpha} (\epsilon Y')_{a\beta}$$

- このとき  $su(2|2)$  代数を満たすもの：

$$|Q^\alpha_a\rangle = \frac{1}{2} q_1 \left\{ \eta^* [(\sinh \bar{Y}) \widetilde{X} (\cosh \bar{Y}')]^\alpha_a + \eta [(\cosh \bar{Y}) \widetilde{X}' (\sinh \bar{Y}')]^\alpha_a \right\} |V\rangle + q^\alpha_a |W\rangle$$

$$|H\rangle = \frac{1}{2} q_1 \left\{ \frac{1}{12} [\text{Tr} Y^4 - \text{Tr} Y'^4] \right. \\ \left. + \text{Tr} X \cosh \bar{Y} \widetilde{X} \cosh \bar{Y}' - i \text{Tr} \sinh \bar{Y} \widetilde{X} \sinh \bar{Y}' X' \right. \\ \left. + \text{Tr} \cosh \bar{Y} \widetilde{X}' \cosh \bar{Y}' X' + i \text{Tr} X \sinh \bar{Y} \widetilde{X}' \sinh \bar{Y}' \right\} |V\rangle + h |W\rangle$$

$$\bar{Y} = Y\eta, \quad \bar{Y}' = Y'\eta^*, \quad (\eta \equiv e^{i\pi/4})$$

pp-wave の場合の LCSFT の 3 弦相互作用項

[Pankiewicz(2003)] + [Di Vecchia-Petersen-Petrini-Russo-Tanzini(2003)]  
を再現している！

# Toy model I

仮定：
$$q^\alpha_a |V\rangle = \frac{i}{2} [(YX)^\alpha_a + (X'Y')^\alpha_a] |V\rangle + (v_{30}Y^3X + v_{34}Y^3XY'^4 + v'_{03}X'Y'^3 - v_{34}Y^4X'Y'^3)^\alpha_a |V\rangle$$

$$[q^\alpha_a, \widetilde{X}^b_b] = \frac{i}{2} \delta_a^b Y^\alpha_b, \quad [q^\alpha_a, \widetilde{X}'^\beta_\beta] = \frac{i}{2} \epsilon^{\beta\alpha} (\epsilon Y')_{a\beta} \longleftarrow \text{pp-waveのときと同じ}$$

このとき  $\text{su}(2|2)$  代数を満たすもの：

$$|Q^\alpha_a\rangle = \frac{1}{2} q_1 \left\{ \eta^* [(\sinh \bar{Y}) \widetilde{X} (\cosh \bar{Y}' - (v'_{03}/2) Y'^4)]^\alpha_a + \eta [(\cosh \bar{Y} + (v_{30}/2) Y^4) \widetilde{X}' \sinh \bar{Y}']^\alpha_a \right\} |V\rangle + q^\alpha_a |W\rangle$$

$$|H\rangle = \frac{1}{2} q_1 \left\{ \frac{1}{12} [\text{Tr } Y^4 (1 - (v'_{03}/4) \text{Tr } Y'^4) - \text{Tr } Y'^4 (1 + (v_{30}/4) \text{Tr } Y^4)] \right. \\ + \text{Tr } X (\cosh \bar{Y} - (v_{30}/2) Y^4) \widetilde{X} (\cosh \bar{Y}' - (v'_{03}/2) Y'^4) - i \text{Tr } \sinh \bar{Y} \widetilde{X} \sinh \bar{Y}' X' \\ + \text{Tr} (\cosh \bar{Y} + (v_{30}/2) Y^4) \widetilde{X}' (\cosh \bar{Y}' + (v'_{03}/2) Y'^4) X' + i \text{Tr } X \sinh \bar{Y} \widetilde{X}' \sinh \bar{Y}' \\ \left. - (v_{34}/4) (\text{Tr } XY^4 \widetilde{X} Y'^4 + \text{Tr } Y^4 \widetilde{X}' Y'^4 X') \right\} |V\rangle + h |W\rangle$$



# Toy model II

仮定：  $q^\alpha_a |V\rangle = \frac{i}{2} [(YX)^\alpha_a + (X'Y')^\alpha_a] |V\rangle \quad \longleftarrow \text{pp-waveのときと同じ}$

$$[q^\alpha_a, \widetilde{X}^b_b] = -\frac{i}{2} (Y\epsilon)^{b\alpha} \epsilon_{ab} + y_{03} [(Y^3\epsilon)^{b\alpha} \epsilon_{ab} + 2i(Y\epsilon)^{b\alpha} (\epsilon X \widetilde{X})_{ab} + 4i(YX)^\alpha_a \widetilde{X}^b_b],$$

$$[q^\alpha_a, \widetilde{X}'^\beta_\beta] = \frac{i}{2} \epsilon^{\beta\alpha} (\epsilon Y')_{a\beta} - y'_{30} [\epsilon^{\beta\alpha} (\epsilon Y'^3)_{a\beta} + 2i(\widetilde{X}' X' \epsilon)^{\beta\alpha} (\epsilon Y')_{a\beta} + 4i(X'Y')^\alpha_a \widetilde{X}'^\beta_\beta]$$

up to dim 1 でconsistentに変形

このとき up to dim 1 で  $su(2|2)$  代数を満たすものを求めた。



pp-waveからのハミルトニアンズレ：

$$\frac{1}{2} q_1 \left[ y_{03} \text{Tr} Y^4 \left( 1 - \frac{1}{48} \text{Tr} Y'^4 \right) + y'_{30} \text{Tr} Y'^4 \left( 1 - \frac{1}{48} \text{Tr} Y^4 \right) \right] |V\rangle.$$

超電荷はpp-waveのときと同じ。

# Summary and discussion

- $su(2|2)$ 対称性をもつ背景でのLCSFTの構成に向けた「模型」を提案。
- 従来知られているpp-waveの場合の形を簡潔に再現した。
- “SUGRA極限”は  $\tilde{X}^a_{\dot{\alpha}} = X^a_{\dot{\alpha}}, \tilde{X}'^{\alpha}_{\dot{\alpha}} = X'^{\alpha}_{\dot{\alpha}} \quad W^{\alpha}_{\dot{b}} = \tilde{W}^{\alpha}_{\dot{b}} = W'^{\alpha}_{\dot{\beta}} = \tilde{W}'^{\alpha}_{\dot{\beta}} = 0$   
でとれてSUGRAの式と比較することにより係数を制限できる。
- LCSFTの3弦相互作用項をdim.に関する展開で系統的に構成できる。
- pp-waveから少し一般化した具体例(toy model I,II)の計算をした。
- pp-waveより一般の背景の場合の正当化?あるいは対応する背景は?  
(bubbling geometry?) また、高次の相互作用項は?
- AdS/CFT (BMN) 対応の一般化への応用?



3rd International Conference:

# String Field Theory and Related Aspects

*Yukawa Institute for Theoretical Physics,  
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October 18-22, 2010*

弦の場の理論 10

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