Wall-Crossing of D4/D2/D0 on the Conifold

(arXiv: 1007.2731 [hep-th])

Takahiro Nishinaka (Osaka U.)

(In collaboration with Satoshi Yamaguchi)

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The d=4, N=2 string thoery have a special class of quantum states called 1/2 BPS states, whose "degeneracy" or index is piecewise constant in the moduli space.

$$\Omega(Q;t)=-rac{1}{2}{
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 $oldsymbol{Q}$: electro-magnetic charge, $oldsymbol{t}$: vacuum moduli

The trace is taken over the Hilbert space of charge $oldsymbol{Q}$, which depends on $oldsymbol{t}$.

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Wall-crossing phenomena

moduli space

$$\Omega(Q;t_1)$$
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Branes vs Black holes

wrapped D-branes in CY3





BPS black holes in 4dim (single - or multi-centered)



vacuum moduli = Calabi-Yau moduli



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The appearence/disappearence of BPS bound states is related to the existence of multi-centered BPS black holes. [F. Denef]

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<u>KS-formula</u>

Recently, Kontsevich and Soibelman have proposed a <u>wall-crossing formula</u> that tells us how the degeneracy changes at the walls of marginal stability.

 $\Omega(Q;t_1)$

discrete change

 $\Omega(Q;t_2)$

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Type IIA on Calabi-Yau

We study the wall-crossing of one non-compact D4-brane with arbitrary numbers of D2/DO on the resolved conifold.

The main topic of this talk

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Type IIA on Calabi-Yau

We study the wall-crossing of one non-compact D4-brane with arbitrary numbers of D2/DO on the resolved conifold.

The vacuum moduli are the Kahler moduli of the conifold.

We evaluate the partition function of D4/D2/D0 in various chambers in the moduli space by using the Kontsevich-Soibelman formula (KS-formula).

Definition

 $\mathbb{C}^4 \supset \mathcal{M}_{m{y}} := \{ (z_1, z_2, z_3, z_4) \, ; \ |z_1|^2 + |z_2|^2 - |z_3|^2 - |z_4|^2 = m{y} \}$ $U(1) \text{-action}: (z_1, z_2, z_3, z_4) \longrightarrow (e^{i\theta} Z_1, e^{i\theta} z_2, e^{-i\theta} z_3, e^{-i\theta} z_4)$

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(1) In the case of y > 0

The compact 2-cycle is $z_3 = z_4 = 0$,

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Two limits $y
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We put one D4-brane on a non-compact 4-cycle $z_3=0$



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The point y=0 is geometrically singular but the spectrum has no singularity if we tune the <u>B-field</u> for the compact 2-cycle.

<u>Kahler moduli</u>

 $z=x+iy\,$: Kahler parameter for the compact cycle $t = \mathbf{z}\mathcal{P} + \mathbf{\Lambda}e^{i\varphi}\mathcal{P}'$ $(\mathcal{P}, P' \in H^2(X))$ $(\mathcal{P}, e^{i\varphi} \in$

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In the final result, the local limit $\Lambda o \infty$ should be taken.

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 $oldsymbol{z} = x + i y$: Kahler parameter for the compact cycle $t = z \mathcal{P} + \Lambda e^{i arphi} \mathcal{P}' \; \; , \qquad (x: B-field, y: size)$

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 $(\mathcal{P},\,P'\in H^2(X))$ $iglace{} \Lambda e^{iarphi}:$ Kahler parameter for other non-compact cycles

In the final result, the local limit $\Lambda \to \infty$ should be taken.

Walls of marginal stability

For a decay channel $Q \rightarrow Q_1 + Q_2$, the walls are defined by $\arg[Z(Q)] = \arg[Z(Q_1)] = \arg[Z(Q_2)]$

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The relevant walls are

 $D4 + kD2 + lD0 \longrightarrow \begin{cases} D4 + (k \mp 1)D2 + (l - n)D0 \\ (\pm 1)D2 + nD0 \end{cases}$

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central charges

$$Z(\mathrm{D}4+ {k D}2+ {l D}0) \sim -rac{1}{2} \Lambda^2 e^{2iarphi}$$

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Walls of marginal stability are the subspace in the moduli space where these two central charges are aligned,

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These walls are labeled by $(\pm 1, n)$. So we denote them $W_n^{\pm 1}$.



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We will do this by using the Kontsevich–Soibelman's wall–crossing formula.

Consider an infinite dimensional Lie algebra with a commutation relation $[e_{Q_1}, e_{Q_2}] = (-1)^{\langle Q_1, Q_2 \rangle} \langle Q_1, Q_2 \rangle e_{Q_1+Q_2}$

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Consider an infinite dimensional Lie algebra with a commutation relation

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 $[e_{m{Q_1}},e_{m{Q_2}}]=(-1)^{\langle Q_1,Q_2
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and the following product

$$A = \prod_{oldsymbol{Q}_{ ext{BPS}}} K^{\Omega(oldsymbol{\Omega}_{ ext{BPS}};t)}_{oldsymbol{Q}_{ ext{BPS}}} \left(K_{oldsymbol{Q}} = \exp[\sum_{n=1}^{\infty} rac{1}{n^2} e_{noldsymbol{Q}}]
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The product A depends on the moduli t in two ways, the degeneracy $\Omega(Q;t)$ and the order in the product depend on t .

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<u>KS-formula says that "nevertheless, the product A is independent of t !"</u>

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We can read off the change in $\ \Omega(Q_{ ext{BPS}};t)$ from the invariance of $oldsymbol{A}$.

Partition function

$$\mathcal{Z}(u,v) = \sum_{oldsymbol{m},oldsymbol{n}} \Omega(\mathrm{D4} + oldsymbol{m}\mathrm{D2} + oldsymbol{n}\mathrm{D0}) \,\, v^{oldsymbol{m}} u^{oldsymbol{n}}$$

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Partition function



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 \mathcal{Z}

$$T(u,v) = \sum_{m,n} \Omega(\mathrm{D}4 + m\mathrm{D}2 + n\mathrm{D}0) \ v^m u^n$$

Suppose that we move the moduli from $\operatorname{Im} z = \infty$ to $\operatorname{Im} z = -\infty$ as above.

(1) For $\operatorname{Im} z > 0$, there are the walls of $\{W_{\infty}^{-1}, \dots, W_{2}^{-1}, W_{1}^{-1}\}$. When the moduli is in the chamber between W_{n}^{-1} and W_{n-1}^{-1} , $\mathcal{Z}(u,v) = \mathcal{Z}_{+\infty}(u,v) \times \prod_{r=n}^{\infty} \frac{1}{(1-u^{r}v^{-1})}$ \boldsymbol{z}

→ Rez

 W_{-2}^{+1}

 W_{-1}^{+1}

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2 arphi

 $W_0^{\pm 1}$

 W_{1}^{+1}

0

(2) For Im z < 0, there are the walls of $\{W_0^{+1}, W_1^{+1}, \cdots, W_{\infty}^{+1}\}$.



$$W_{-1}^{-1}$$

$$-2 -1 0 1 2^{\varphi}$$

$$W_{-2}^{+1}$$

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(2) For $\operatorname{Im} z < 0$, there are the walls of $\{W_0^{+1}, W_1^{+1}, \cdots, W_{\infty}^{+1}\}$. When the moduli is in the chamber between W_n^{+1} and W_{n+1}^{+1} ,

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In particular, we obtain

$$\mathcal{Z}_{-\infty}(u,v) = \mathcal{Z}_{+\infty}(u,v) imes \prod_{r=1}^{\infty} rac{1}{(1-u^rv^{-1})} imes \prod_{r=0}^{\infty} rac{1}{(1-u^rv)}$$

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Moreover, these two limits coinside with the attractor moduli of the MSW black holes, where BPS microstates are counted in the field theory on D4-brane.



In particular, we obtain

$$\mathcal{Z}_{-\infty}(u,v) = \mathcal{Z}_{+\infty}(u,v) imes \prod_{r=1}^{\infty} rac{1}{(1-u^r v^{-1})} imes \prod_{r=0}^{\infty} rac{1}{(1-u^r v)}$$

where $\mathcal{Z}_{\pm\infty}(u, v)$ denotes the partition function in the limit $\operatorname{Im} z = \pm \infty$. These two limits correspond to the large \mathbb{P}^1 limit in left and right hand side of the following picture:



Moreover, these two limits coinside with the attractor moduli of the MSW black holes, where BPS microstates are counted in the field theory on D4-brane. Actually, $\mathcal{Z}_{+\infty}(u,v)$ was already evaluated in a literature: [Aganagic-Ooguri-Saulina-Vafa'04]

$$Z_{+\infty}(u,v) = f(u)(1-v)\prod_{r=1}^{\infty}(1-u^r)(1-u^rv)(1-u^rv^{-1})$$

In particular, we obtain

$$\mathcal{Z}_{-\infty}(u,v) = \mathcal{Z}_{+\infty}(u,v) imes \prod_{r=1}^{\infty} rac{1}{(1-u^rv^{-1})} imes \prod_{r=0}^{\infty} rac{1}{(1-u^rv)}$$

where $\mathcal{Z}_{\pm\infty}(u, \sigma)$ denotes the partition function in the limit $\operatorname{Im} z = \pm \infty$. These two limits correspond to the large \mathbb{P}^1 limit in left and right hand side of the following picture:



Moreover, these two limits coinside with the attractor moduli of the MSW black holes, where BPS microstates are counted in the field theory on D4-brane.

Actually, \mathcal{Z}_{+} \gtrsim (u,v) was already evaluated in a literature: [Aganagic-Ooguri-Saulina-Vafa '04]

$$Z_{+\infty}(u,v) = f(u)(1-v) \prod_{r=1}^{\infty} (1-u^r)(1-u^rv)(1-u^rv^{-1})$$



$$Z_{+\infty}(u,v) = f(u)(1-v) \prod_{r=1}^{\infty} (1-u^r)(1-u^rv)(1-u^rv^{-1})$$







Summary

- We have discussed the wall-crossing phenomena of D4/D2/D0 bound states on the resolved conifold.
- We considered one non-compact D4-brane and various numbers of D2/D0 on it.
- We identified all walls of marginal stability.
- By moving the Kahler moduli, we can consider the flop transition of the conifold through which the topology of the conifold is changed.
- We evaluate the partition function in all chambers by using the Kontsevich-Soibelman wall-crossing formula.
- The result is completely consistent with the known facts about the field theory on D4-branes and the flop transition.

That's all for my presentation. Thank you very much.

3