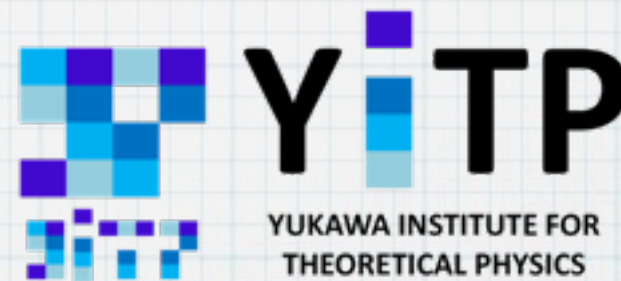


Conference @ YITP
22, July 2010

トポロジカルストリングとAGT予想

Topological Strings & AGT Relations

Masato Taki/ YITP, Kyoto University



based on [Maruyoshi-M.T. arXiv:1006.4505]
[M.T. arXiv:1007.2524]

1. AGT conjecture and Surface Operators

Today I will study “**AGT relation**” between gauge theory and 2-dim CFT. This mysterious relation has attracted attention recently.

2D Liouville CFT \longleftrightarrow $\mathcal{N} = 2$ superconformal quiver gauge theories

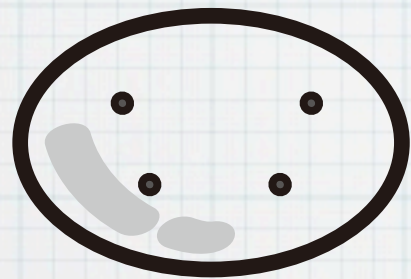
[Alday-Gaiotto-Tachikawa, '09]

$$\mathcal{B}_{\text{Liouville}}(q) = Z_{\text{Nekrasov}}(q)$$

Free-field CFT vs U(1)-theory

Let us consider the linear dilaton CFT with the background charge Q

$$\langle \phi(z_1) \phi(z_2) \rangle = -\ln(z_i - z_j)$$



$$\left\langle \prod_{i=1}^4 e^{\alpha_i \phi(z_i)} \right\rangle = \prod_{i < j} (z_i - z_j)^{-\alpha_i \alpha_j}$$

We choose $z_{1,\dots,4} = 0, q, 1, \infty$

$$\mathcal{B}_{\text{free CFT}} = (1 - q)^{-\alpha_2 \alpha_3}$$



$$m_1 m_2$$

Some algebra leads to the following expansion

$$(1 - q)^{-m_1 m_2} = \sum_Y q^{|Y|} \prod_{(i,j) \in Y} \frac{\prod_{f=1,2} \phi(m_f, (i, j))}{E(0, Y, (i, j))^2}$$

$$\phi(m, (i, j)) = m + i - j$$

$$E(a, Y, (i, j)) = a + (Y_j^T - i + 1) + (Y_i - j)$$

$$\left(\begin{array}{c} \text{where we use} \\ \sum_Y \chi_Y(x) \chi_Y(y) = \exp \sum_n n x_n y_n \end{array} \right)$$

This is the Nekrasov partition function for the “U(1) gauge theory with 2-flavors”!

$$\mathcal{B}_{\text{free CFT}} = Z_{U(1), N_f=2}$$

surface operator from Liouville theory [Alday-Gaiotto-Gukov-Tachikawa-Verlinde, '09]

In the chapter we give a rough sketch of an important extension of the AGT relation. The point is that the AGT relation, or Liouville/gauge-theory correspondence, also holds for extended objects, such as surface operators.

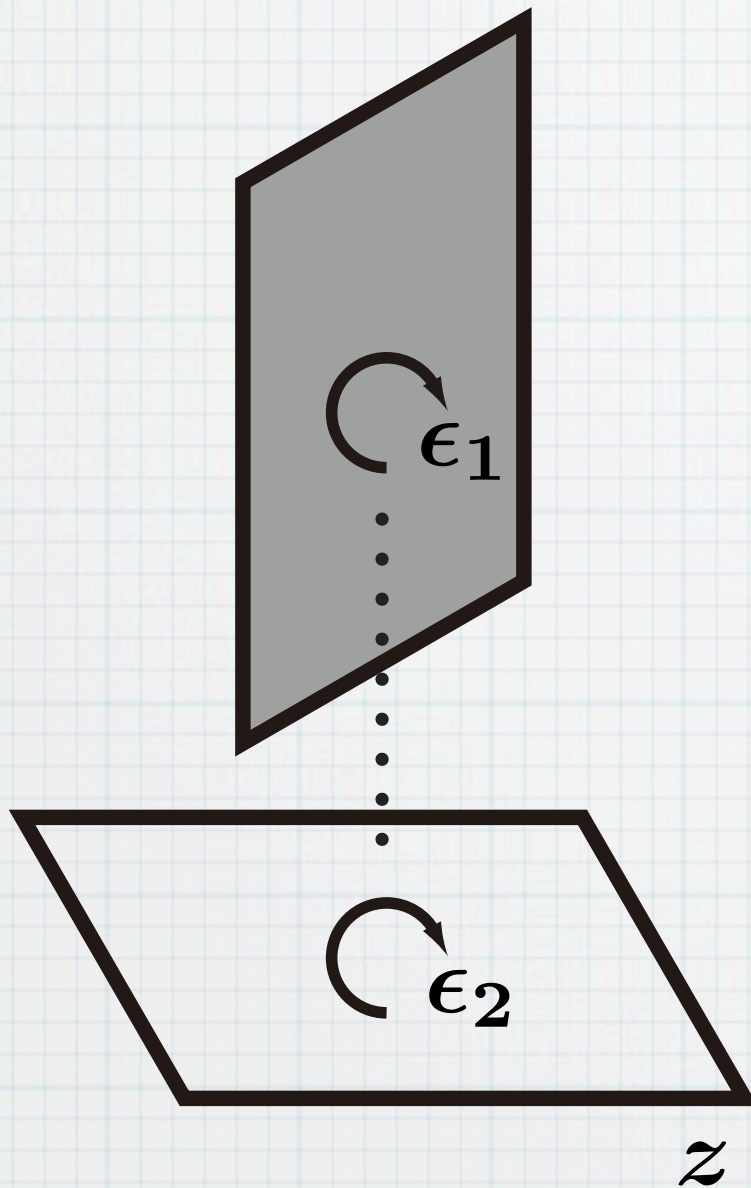
- * probe for the S-duality (geometric Langrange duality)
[Gukov-Witten, '06]**
- * ramified Donaldson-Witten inv.
[Kronheimer-Mrowka, '93,'95] [Braverman, '04]
[Alday-Tachikawa, '10]**
- * generator of 'tHooft-Wilson operator**
- * order parameter for phases of gauge theory**

surface operator

$\mathcal{N} = (2, 2)$ sigma model coupled the 4D bulk gauge field

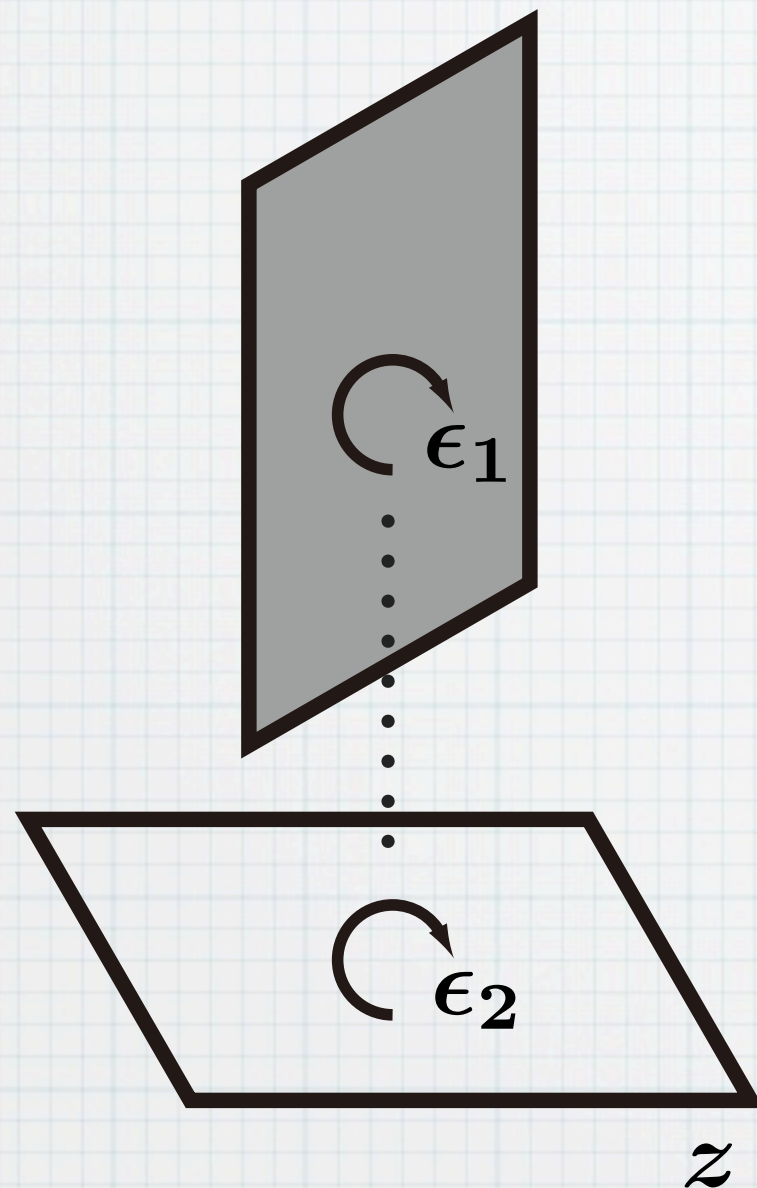
A surface operator is characterized by the subgroup of the gauge group which maintains the surface operator

$$\mathbb{L} \in G$$



surface operator

A surface operator creates the singularity near the locus of it's insertion:



$$A \sim \begin{pmatrix} a & 0 \\ 0 & -a \end{pmatrix} \frac{dz}{z}$$

$$\rightarrow \begin{cases} m = \int_{\mathbb{C}_{z=0}} F \\ k = \int_{\mathbb{C}^2} F \wedge F \end{cases}$$

Instantons in the presence of such a are operator labelled by these two topological numbers.

$$\Psi(z) = \sum_{k,m} q^k z^m \int_{\mathcal{M}_{N,k,m}} dV_{\epsilon_{1,2}}$$

AGGTV found that the partition function in the presence of the elementary surface operator is dual to the conformal block with the degenerate field insertion:

[Alday-Gaiotto-Gukov-Tachikawa-Verlinde, '09]

surface operator



insertion of degenerate field $\Phi_{2,1}$

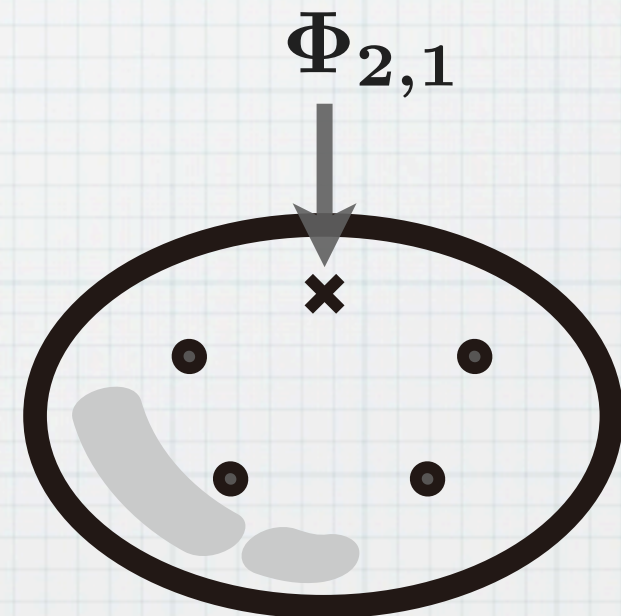
$$\Psi(z) = \exp \sqrt{\frac{\epsilon_1}{\epsilon_2}} \mathcal{G}(z) + \dots$$

$$\langle \Delta_1 | V_2(1) V_3(q) \Phi_{2,1}(z) | \Delta_4 \rangle$$

WKB

$$\mathcal{G}(z) = \int^z \sqrt{\phi_2(z')} dz'$$

$$x^2 = \phi_2(z) \quad \text{SW curve (Gaiotto form)}$$

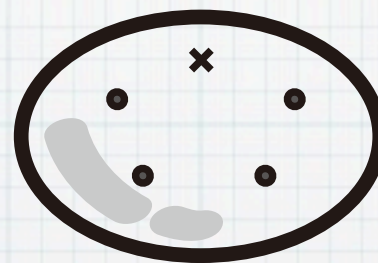


2. Geometric Engineering of Surface Operators

[M.T, arXiv:1007.2524]

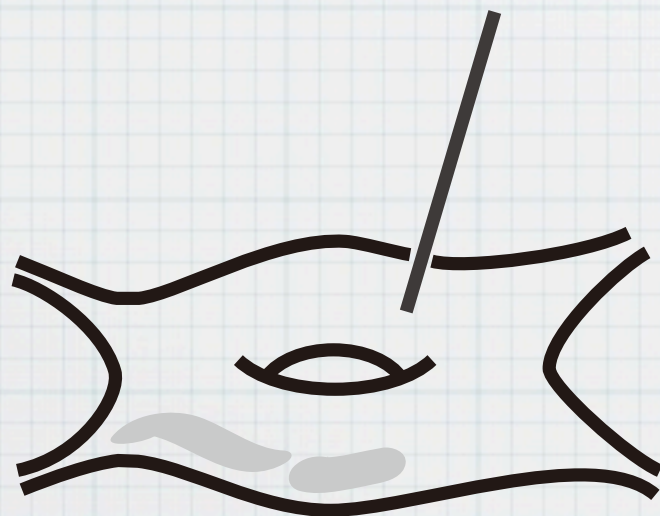
In the chapter we report our work in progress. Let us consider the the following proposals for stringy derivations of surface operators

insertion of degenerate field $\Phi_{2,1}$

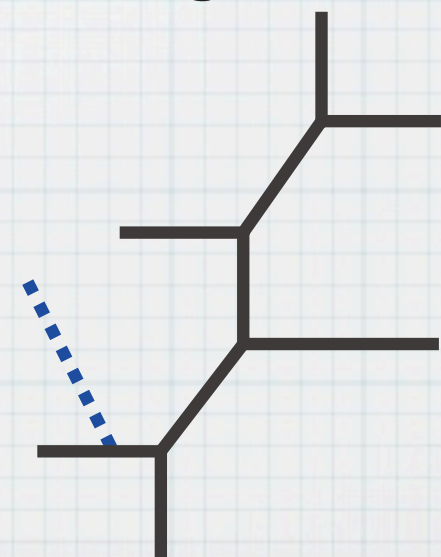


AGGTV

M2-brane construction



engineering via toric brane



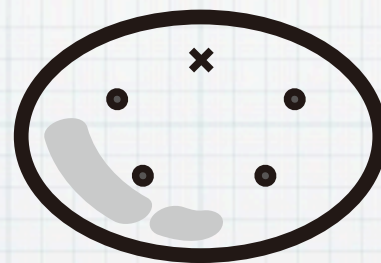
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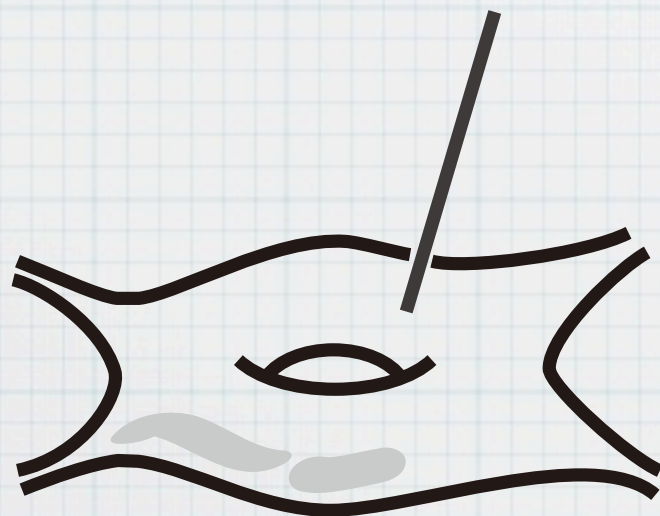
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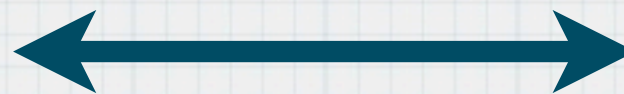
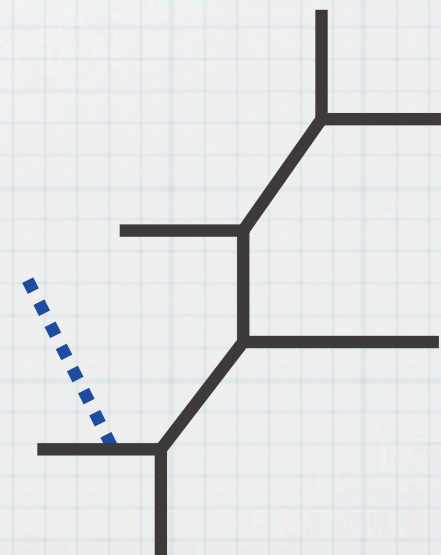
AGGTV



M2-brane construction



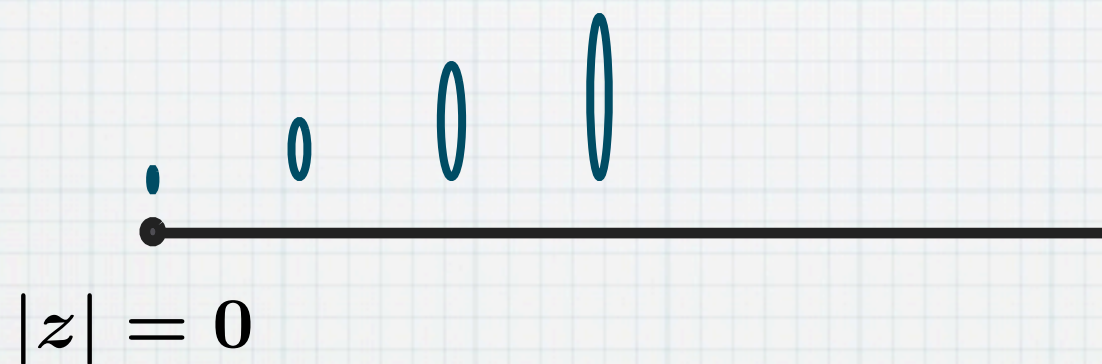
engineering via toric brane



toric Calabi-Yau three-folds

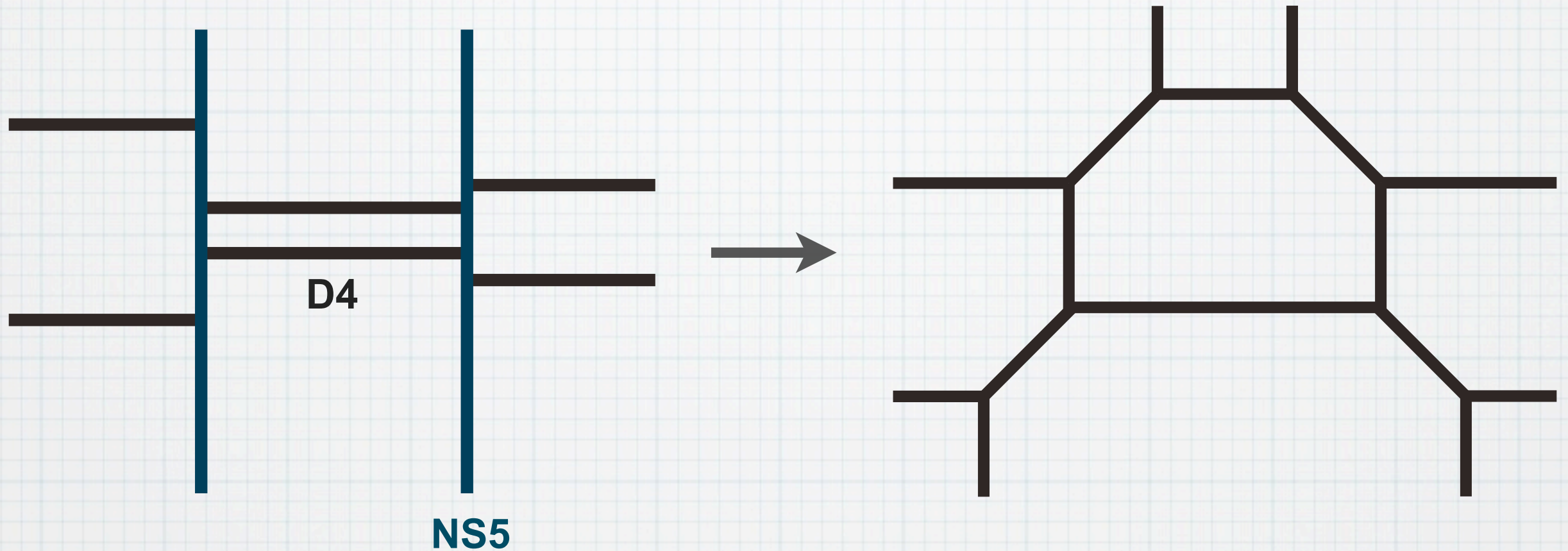
* \mathbb{C}^1

$$z = |z|e^{i\theta} \quad S^1(\theta) \longrightarrow \mathbb{R}_{\geq 0}(|z|)$$



Any toric manifold is labelled by such a skeleton diagram which encodes the degeneration locus of the torus fiber.

* $SU(N)$ geometry



The low-energy effective action for Type IIA superstring on the Calabi-Yau is equal to the Seiberg-Witten solution for $SU(2)$ gauge theory with 4-flavors. Thus the Calabi-Yau engineers the gauge theory.

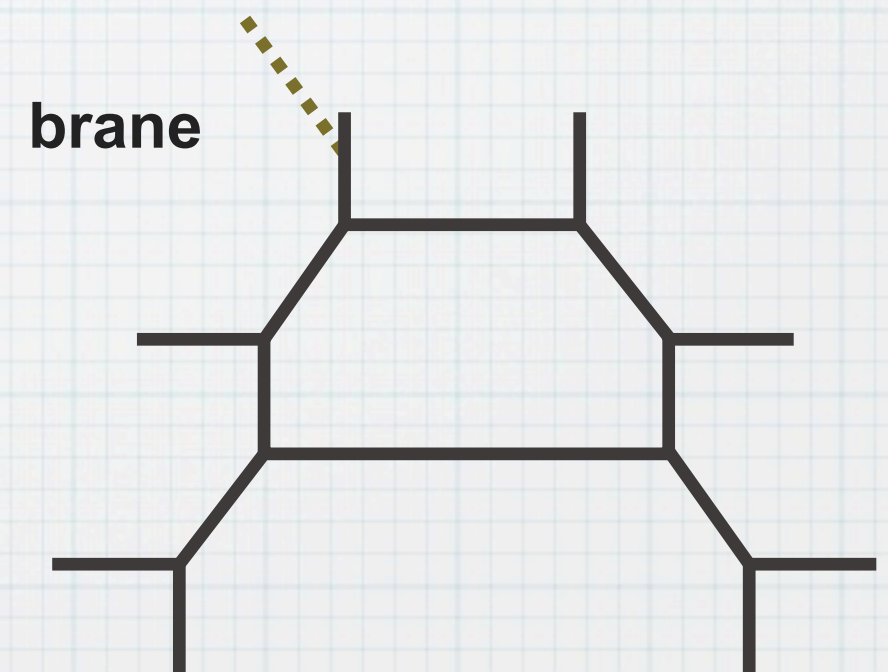
geometric engineering : The closed string topological string partition function of a certain toric CY is equal to Nekrasov partition function.



[Kozcaz-Pasquetti-Wyllard, '10]

insertion of degenerate field $\Phi_{2,1}$ \longleftrightarrow open topological string amplitude on certain toric CY

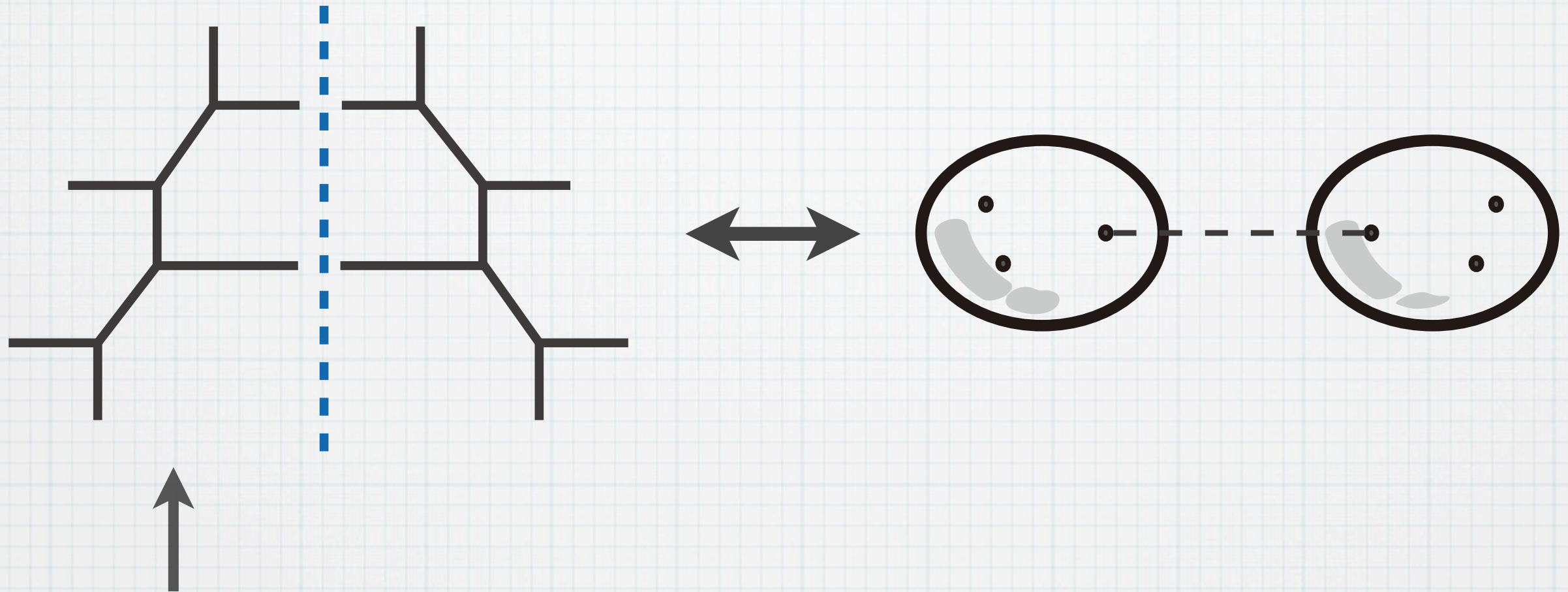
$$\Psi(z) = \exp \sqrt{\frac{\epsilon_1}{\epsilon_2}} \mathcal{G}(z) + \dots$$



We will derive the correspondence by utilizing the string dualities!

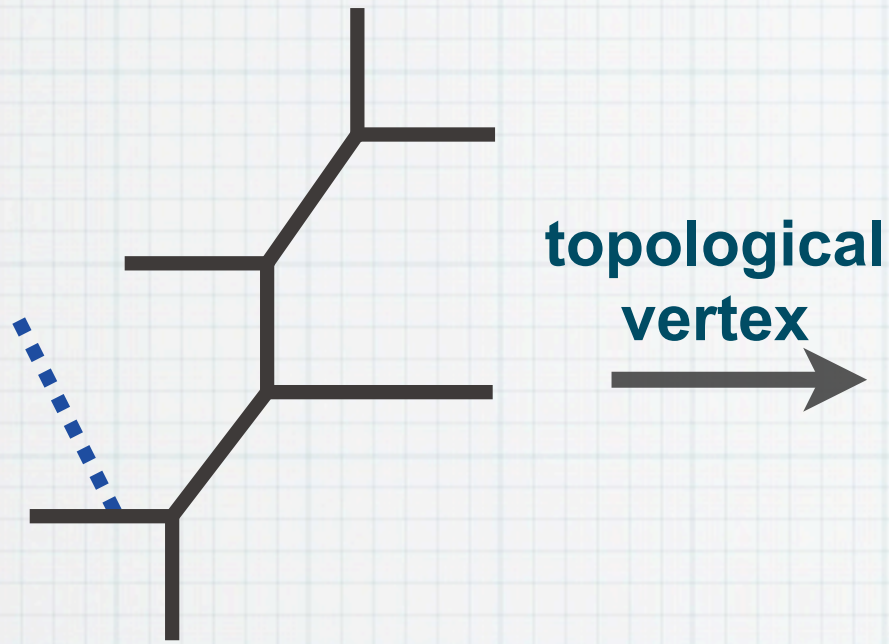
simple example

The degeneration of the Riemann surface in CFT side corresponds to the following optical-theorem-like computation methods.



4 hypermultiplets theory : $\mathcal{T}_{3,0}(A_1)$

Recall the Gaiotto's theory, which is that the 3-punctured sphere describes 4-hypermultiplets of gauge theory. So, the partition function for the half-toric-geometry will give the partition function of it. Let us study this half-geometry for example.



$$\Phi(z) = \sum_n z^n \prod_{k=1}^n \frac{(k\hbar + m_1)(k\hbar + m_1 + m_2 + m_3)}{k\hbar(k\hbar + m_1 + m_3)}$$

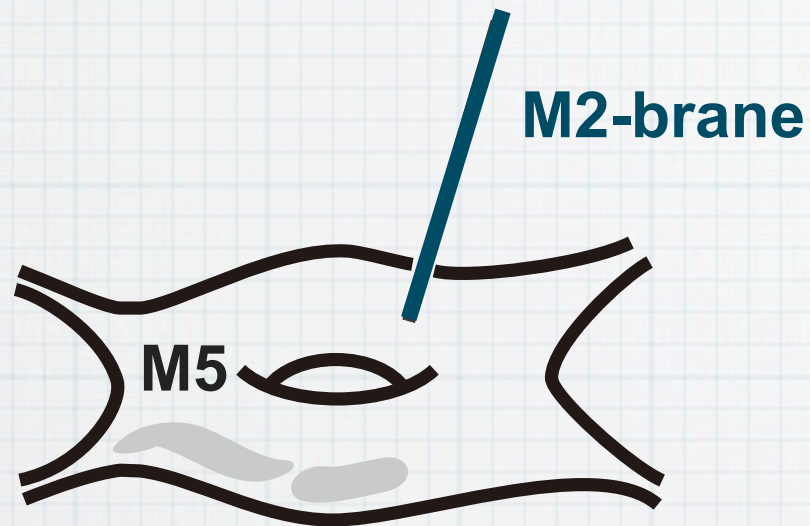


$$G(z) = \int^z \frac{\sqrt{\alpha_1^2 z(z-1) + \alpha_2^2(1-z) + \alpha_3^2 z}}{z(1-z)}$$

* This result reproduces the correct Seiberg-Witten curve $x^2 = \phi_2(z)$
c.f. [Schiappa-Wyllard, '09]

“Proof” of the AGGV-conjecture

* M2-brane construction



	0	1	2	3	4	5	6	7	8	9	10
M5	○	○	○	○			Σ		×	×	×
M2	○	○	×	×	×	×	×	×	○	×	×

$$\Sigma : H(x, y) = 0.$$

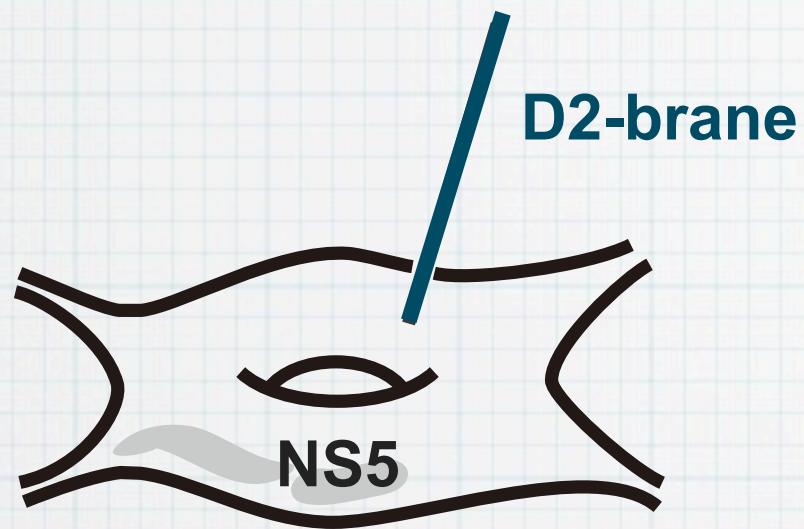
dim. reduction along x^{10}

* IIA superstring with N NS5-branes on Σ and a transverse D2-brane

Next, let us use the following duality:

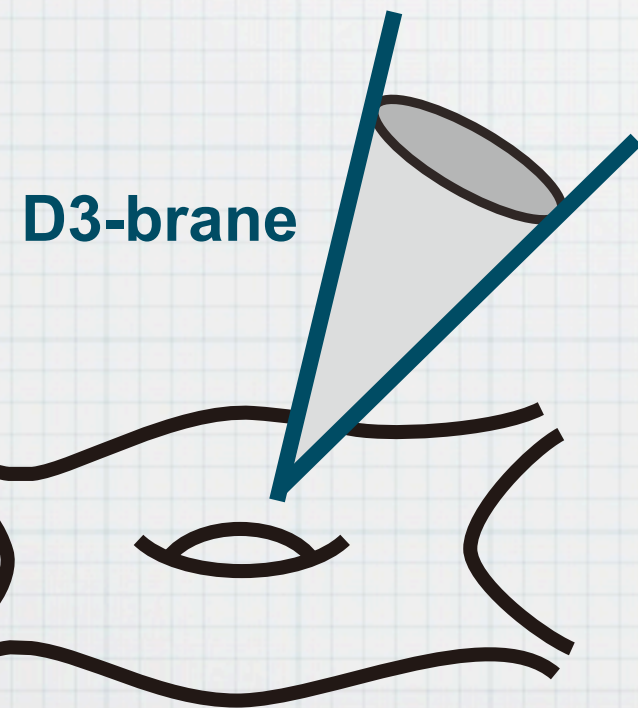
N NS5-branes (IIA) are T-dual to A_{N-1} -singularity

* IIA superstring with N NS5-branes on Σ and a transverse D2-brane



↓
T-duality along a circle transverse to Σ

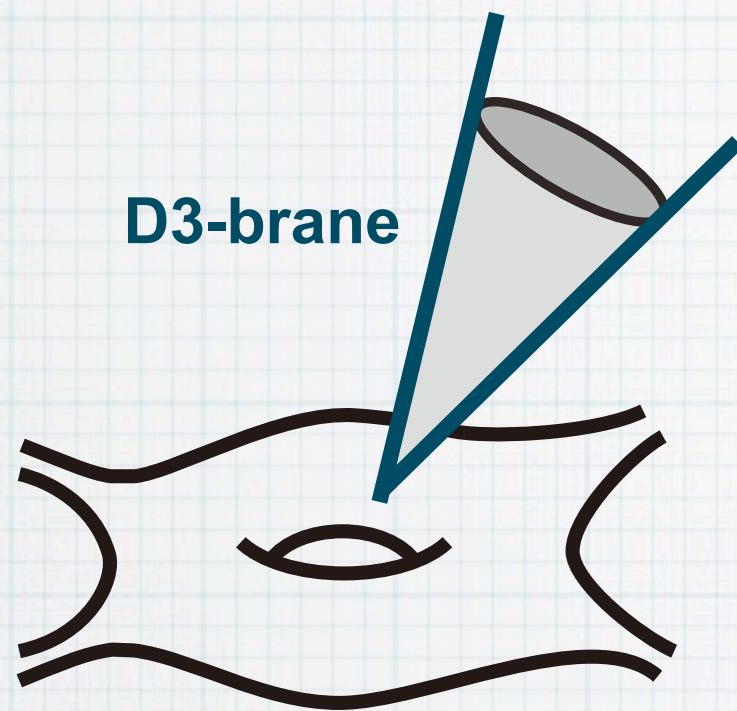
* IIB superstring on the Calabi-Yau $uv + H(x, y) = 0$



and D3 brane on

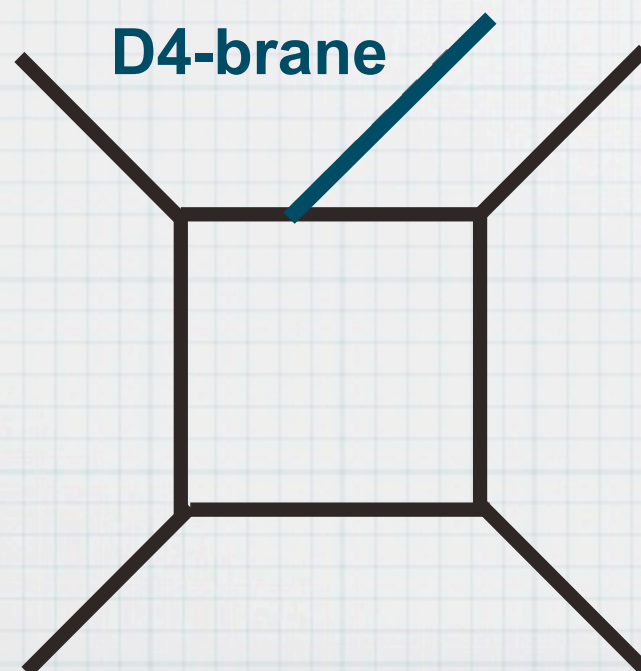
$$u = 0, \quad v = \text{arbitrary}, \quad H(x_0, y_0) = 0.$$

* IIB superstring on the Calabi-Yau with a D3-brane



mirror symmetry [Aganagic-Vafa, '01]

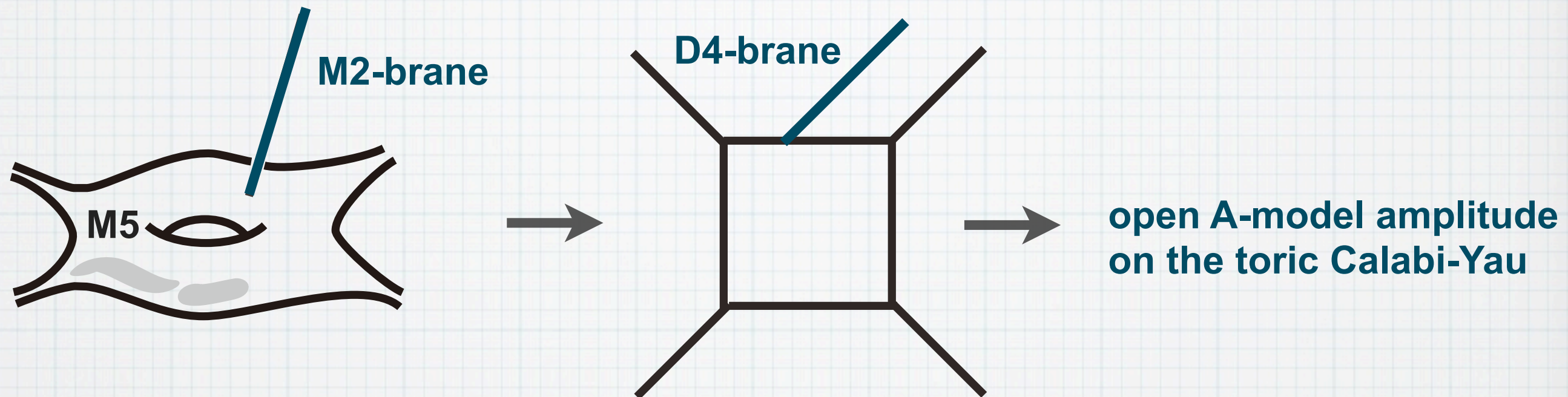
* IIA superstring on the toric Calabi-Yau and the toric D4 brane



This is precisely the situation which open topological string amplitude engineers!!

[Ooguri-Vafa, '99]

In this way we have proven the proposal of [Kozcaz-Pasquetti-Wyllard, '10]



Moreover, by employing the derivation of [Dijkgraaf-Vafa, '09], we can show that such a brane is described by the following degenerate field!

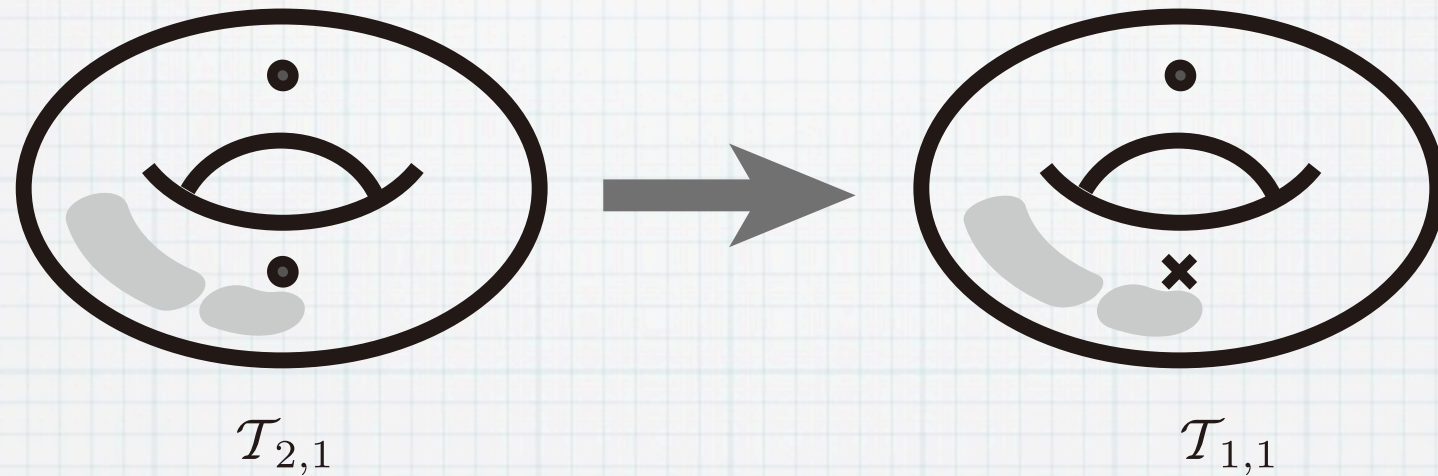
$$V(q) = e^{-b\phi(q)}$$

Thus string dualities also verify the AGTV conjecture!

3. AGTV relation via Bubbling Calabi-Yau

[Maruyoshi-M.T. arXiv:1006.4505][M.T, arXiv:1007.2524]

[Alday-Gaiotto-Gukov-Tachikawa-Verlinde, '09] implies



\hat{A}_1 quiver gauge theory

a, a', m, m'

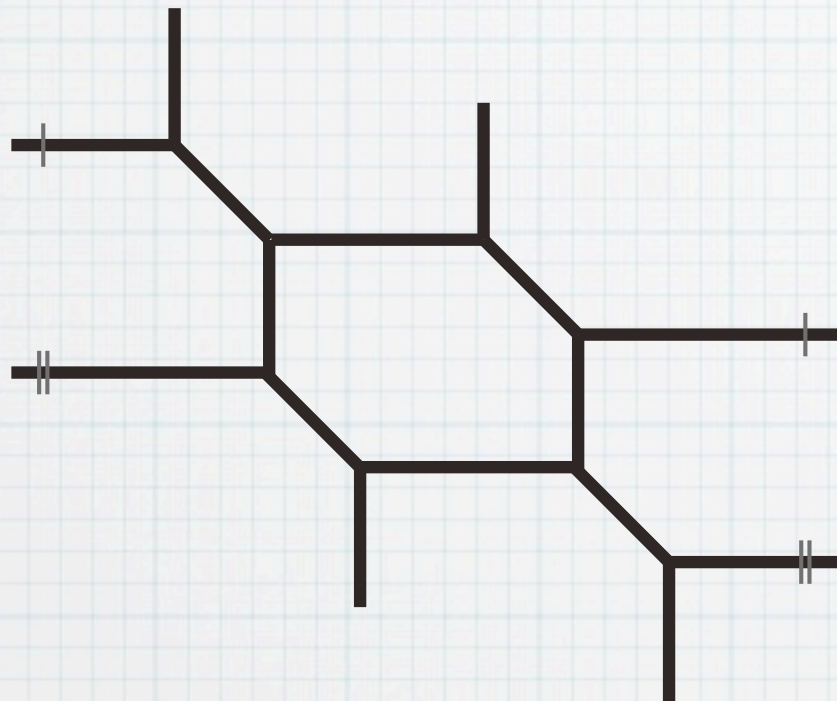
$\mathcal{N} = 2^*$ gauge theory

a, m

$$\begin{aligned} a' &= a + \frac{\epsilon_2}{2} \\ m' &= -\frac{\epsilon_2}{2} \end{aligned}$$

In the stringy realization, what will happen in the degeneration of parameter?

\hat{A}_1 quiver gauge theory

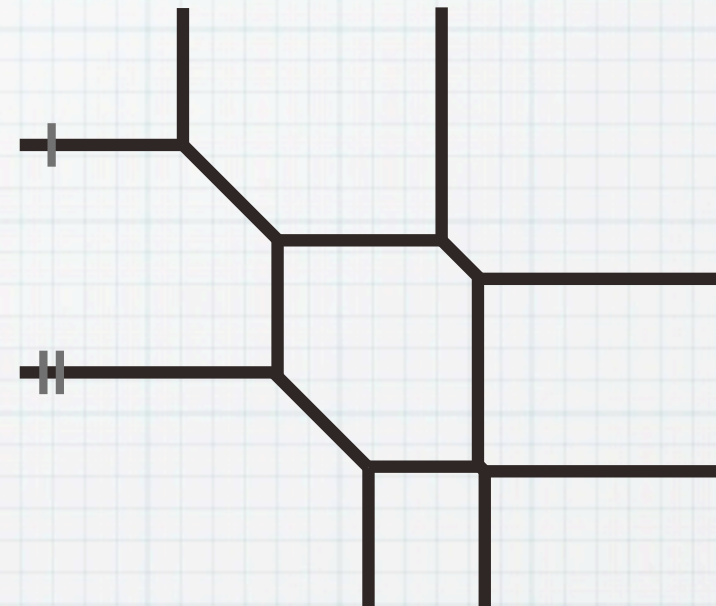


$$a' = a + \frac{\epsilon_2}{2}$$

$$m' = -\frac{\epsilon_2}{2}$$



$\mathcal{N} = 2^*$ gauge theory
with surface operator

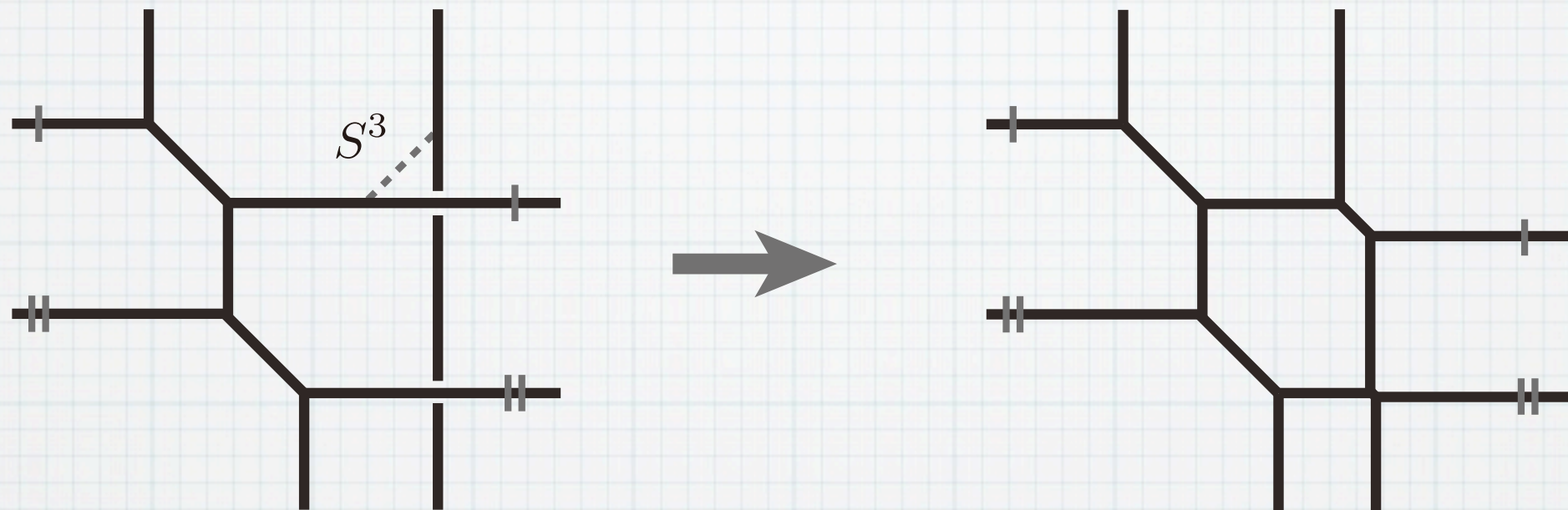


Relation to the toric brane construction??



Bubbling of Calabi-Yau !! [Okuda-Gomiz, '07]

**Toric barnes melts into the background geometry via geometric transition
(open/closed duality)**



**In this way the AGGTV conjecture includes a gauge theory version of open/
closed duality**

decoupling limit

By taking the decoupling limit of the hypermultiplet $m \rightarrow \infty$, we obtain the ramified instanton partition function for the SU(2) pure SYM

$$\Psi_{\text{pure SYM}}(\Lambda, \lambda) = \sum_{\vec{Y}, \vec{W}} \lambda^{|\vec{Y}|} (\Lambda^4 \lambda^{-1})^{|\vec{W}|} z_{\text{vect}}(a, \vec{Y}) z_{\text{vect}}(b = a - \epsilon_2/2, \vec{W}) \\ \times z_{\text{bifund}}(b = a + \epsilon_2/2, \vec{W}; a, \vec{Y}; m_2 = -\epsilon_2/2)$$

Meanwhile, Maruyoshi and M.T. proposed the following CFT expression for the partition function by using irregular conformal block

[Maruyoshi-M.T., '10]

$$\Psi_{\text{pure SYM}} = \sum_{n=0}^{\infty} \sum_Y \Lambda^{4|Y|} \lambda^{n-|Y|} \beta_{1^n}^Y Q_{\Delta}^{-1}([1^{|Y|}]; Y)$$

where $\Phi_{2,1}(z)|\Delta, Y\rangle = \sum_{Y'} z^{|Y'|-|Y|+\delta} \beta_{Y'}^Y |\Delta', Y'\rangle$

Do they match? It is a ramified version of **Gaiotto-Marshakov-Mironov-Morozov conjecture**.

SUMMARY

- We develop AGT relation in order to describes surface operators
- We can prove the AGT relation in the presence of the surface operators by employing string dualities.
- Topological vertex is useful to study AGT relation.
- Understanding of the AGGTV relation involves geometric transition.

Fin