

# *Quarter BPS classified by Brauer algebra*

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# The problem of AdS/CFT

Map between string states and gauge invariant operators

$$\left\langle O_{\alpha}(x)^{\dagger} O_{\beta}(y) \right\rangle = \frac{c(1/N) \delta_{\alpha\beta}}{(x-y)^{2\Delta_{\alpha}(g,1/N)}}$$

4D  $\mathcal{N}=4$  SYM (CFT)

Scaling dimension of local operator = Energy (in global time) of string state

$$\Delta(\lambda, N) = E(R/l_s, g_s)$$

$$4\pi\lambda / N = g_s$$

$$\sqrt{\lambda} = R^2 / \alpha'$$

The operator with the definite scaling dimension is a linear combination of the naive operators [operator mixing].

$$\hat{D} \tilde{O}_{\alpha} = \Delta_{\beta\alpha} \tilde{O}_{\beta}$$

$$\left\langle \tilde{O}_{\alpha}(x)^{\dagger} \tilde{O}_{\beta}(y) \right\rangle = \frac{1}{(x-y)^{2\Delta_0}} (S_{\alpha\beta} + T_{\alpha\beta} \log |\Lambda x|)$$

Definite scaling dimension = Eigenstate of the dilataton operator.

$$\Delta = TS^{-1}$$

## Simplification at large $N$ (planar theory)

Only single-traces (i.e. only the flavour structure should be taken care of. )

$$\text{tr}(XYYYXX) \leftrightarrow |\downarrow\uparrow\uparrow\uparrow\uparrow\downarrow\downarrow\rangle$$

Dilatation operator = Hamiltonian of integrable system

$$D_{\text{planar}}^{(1\text{-loop})} = \frac{\lambda}{8\pi^2} \sum_i (1 - P_{i,i+1})$$

Diagonalisation of the hamiltonian of integrable system solves the mixing problem.

$$\Delta(\lambda, N) = E(R/l_s, g_s)$$

**What to do if  $N$  is not big enough?**

## Operator mixing in the non-planar theory

holomorphic gauge inv. ops. built from two complex matrices,  $X, Y$ .

(This sector is closed in all order perturbation theory. )

$$\underline{D_{non-planar}^{(1-loop)} = -2tr \left( [X, Y] [\partial_X, \partial_Y] \right)}$$

$$X = \Phi_1 + i\Phi_2$$

$$Y = \Phi_3 + i\Phi_4$$

$$Z = \Phi_5 + i\Phi_6$$

--Example -----

$$O_1 = tr(XXYY)$$

$$O_2 = tr(XYXY)$$

$$\left[ \begin{array}{l} O_3 = tr(XX)tr(YY) \\ O_4 = tr(XY)tr(XY) \end{array} \right]$$

multi-trace

$$O'_1 = O_1 - O_2$$

$$O'_2 = O_1 + \frac{N}{2}O_4$$

$$O'_3 = O_2 - NO_4$$

$$O'_4 = O_3 + 2O_4$$

$$DO'_1 = 12NO'_1$$

$$DO'_a = 0 \quad (a = 2, 3, 4)$$

We will need a good method to organise gauge invariant operators of single traces and multi-traces.

Recall some important facts of the 1/2 BPS primary.

$$O_R(X) = \text{tr}_n \left( p_R X^{\otimes n} \right) \quad \text{Schur polynomial}$$

$$p_R = \frac{d_R}{n!} \sum_{\sigma \in S_n} \chi_R(\sigma) \sigma$$

$$X_{j_1}^{i_1} X_{j_2}^{i_2} \cdots X_{j_n}^{i_n}$$

The upper indices are transformed as the product of the fundamental rep. of  $GL(N)$ .

$$\langle O_R(X)^\dagger O_S(X) \rangle \propto \delta_{RS}$$

[Corley, Jevicki, Ramgoolam 01]

$$\langle X_j^i[x] X_l^{\dagger k}[y] \rangle_0 = \delta_l^i \delta_j^k \frac{1}{(x-y)^2}$$

$$\begin{array}{|c|c|} \hline \hline \hline \end{array} \quad p_{[2]} = \frac{1}{2}(1+s) \quad \frac{1}{2} \left( X_{j_1}^{i_1} X_{j_2}^{i_2} + X_{j_1}^{i_2} X_{j_2}^{i_1} \right) \Rightarrow \frac{1}{2} \left( \text{tr} X \text{tr} X + \text{tr} (XX) \right) = O_{[2]}$$

$$\begin{array}{|c|} \hline \hline \end{array} \quad p_{[1,1]} = \frac{1}{2}(1-s) \quad \frac{1}{2} \left( X_{j_1}^{i_1} X_{j_2}^{i_2} - X_{j_1}^{i_2} X_{j_2}^{i_1} \right) \Rightarrow \frac{1}{2} \left( \text{tr} X \text{tr} X - \text{tr} (XX) \right) = O_{[1,1]}$$

Multi-trace

Single-trace

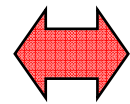
## Representation basis

The trace structure can be conveniently encoded in the Young diagram.

➤ Finite  $N$  constraint ( $\Rightarrow$  cut-off for angular momentum)

$$c_1(R) \leq N$$

$$\text{tr}(X^3) = \frac{3}{2} \text{tr} X \text{tr}(X^2) - \frac{1}{2} (\text{tr} X)^3 \quad (N=2)$$



$$\text{tr}_3(p_{[1,1,1]} X^{\otimes 3}) = 0$$

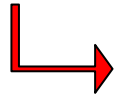
➤ Orthogonal (complete) at classical level

Representation basis seems to be a useful basis to organise gauge inv. operators *when  $N$  is not large enough*.

In this talk, I will study the mixing problem using a representation basis.

## Construction of a representation basis for the $X$ - $Y$ sector

$$tr_{m,n} \left( \underline{P^\gamma} X^{\otimes m} \otimes Y^{T \otimes n} \right)$$



Projector associated with an irreducible rep.  $\gamma$  of  $GL(N)$

$$\underbrace{X_{j_1}^{i_1} X_{j_2}^{i_2} \cdots X_{j_m}^{i_m} Y_{l_1}^{T k_1} Y_{l_2}^{T k_2} \cdots Y_{l_n}^{T k_n}}_{\text{blue oval}} \rightarrow \gamma$$

$3X_s \quad 2Y_s$

The irreducible representation of  $GL(N)$ :

$$\gamma = (\gamma_+, \gamma_-, k),$$

$$0 \leq k \leq \min(m, n), \quad \gamma_+ \mapsto (m - k), \quad \gamma_- \mapsto (n - k)$$

Roughly speaking, this  $k$  represents the mixing between  $X$  and  $Y$ .

$(3, 2)$	$\gamma_+$	$\gamma_-$
$k = 0$	$[3]$	$[2]$
	$[2, 1]$	$[2]$
	$[1, 1, 1]$	$[2]$
	$[3]$	$[1, 1]$
	$[2, 1]$	$[1, 1]$
	$[1, 1, 1]$	$[1, 1]$
$k = 1$	$[2]$	$[1]$
	$[1, 1]$	$[1]$
$k = 2$	$[1]$	$\emptyset$

See the example of the simplest case

$$N \otimes \bar{N} = (N^2 - 1) \oplus 1$$

$$X_{j_1}^{i_1} (Y^T)_{l_1}^{k_1} = \underbrace{X_{j_1}^{i_1} (Y^T)_{l_1}^{k_1} - \frac{1}{N} \delta^{i_1 k_1} X_{j_1}^m (Y^T)_{l_1}^m}_{O^{k=0}} + \underbrace{\frac{1}{N} \delta^{i_1 k_1} X_{j_1}^m (Y^T)_{l_1}^m}_{O^{k=1}}$$

$$O^{k=0} = \text{tr} X \text{tr} Y - \frac{1}{N} \text{tr}(XY)$$

$$k=0$$

$$O^{k=1} = \frac{1}{N} \text{tr}(XY)$$

$$k=1$$

$$\Rightarrow \langle O_{k=0}^\dagger O_{k=1} \rangle_0 = 0$$

In general, the  $k=0$  have the following structure:

$$O_R(X) O_S(Y) + O(1/N) \dots$$

$$O_R(X) = \text{tr}_m(p_R \cdot X^{\otimes m}), \quad O_S(Y) = \text{tr}_n(p_S \cdot Y^{\otimes n})$$



## Finite $N$ constraint (stringy exclusion principle)

[YK-Ramgoolam 0709.2158]

$$tr_{m,n} \left( P^\gamma X^{\otimes m} \otimes Y^{T \otimes n} \right)$$

$$c_1(\gamma_+) + c_1(\gamma_-) \leq N$$

$$\gamma_+ \mapsto (m-k), \quad \gamma_- \mapsto (n-k)$$

$$tr_{m,n} \left( P^{\gamma(k=0,R,S)} X^{\otimes m} \otimes Y^{T \otimes n} \right) = tr_m \left( p_R X^{\otimes m} \right) tr_n \left( p_S Y^{\otimes m} \right) + \dots$$

$$c_1(R) + c_1(S) \leq N$$

$$\gamma_+ = R \mapsto m, \quad \gamma_- = S \mapsto n$$

This is stronger than the naively expected one :  $c_1(R) \leq N, \quad c_1(S) \leq N$

$$tr(X^2 Y) = tr X tr XY - \frac{1}{2} (tr X)^2 tr Y + tr(X^2) tr Y \quad (N=2)$$

$$tr_{2,1} \left( P_{[1,1][1]} X^{\otimes 2} \otimes Y^T \right) = 0$$

This would give a cut-off for the angular momentum of the composite system.

## One-loop analysis

[YK 1002.2424]

The one-loop mixing problem : to look for eigenstates of  $H$ .

$$\hat{H} \equiv tr \left( [X, Y] [\partial_X, \partial_Y] \right)$$

Our goal will be to understand the mixing pattern in terms of Young diagrams.

It is not so easy in general, but we can find some eigenstates easily *based on the new language*.

Easy to find that the  $k=0$  ops are annihilated by  $H$ :

$$(\partial_X \partial_Y)_{pq} X_j^i Y_l^{Tk} = \underline{\delta_{ik}} \delta_{pj} \delta_{ql}$$

$$(\partial_Y \partial_X)_{pq} X_j^i Y_k^{Tk} = \underline{\delta_{jl}} \delta_{pk} \delta_{qi}$$

$$C X_j^i Y_l^{Tk} = \underline{\delta_{ik}} X_j^s Y_l^{Ts}$$

$$X_j^i Y_l^{Tk} \tilde{C} = \underline{\delta_{jl}} X_s^i Y_s^{Tk}$$

$$\hat{H} \cdot tr_{m,n} \left( P^{\gamma(k=0)} X^{\otimes m} \otimes Y^{T \otimes n} \right) = 0 \quad \leftarrow C \cdot P^{\gamma(k=0)} = 0$$

The other eigenstates

$$\begin{aligned}
& \hat{H} \cdot \text{tr}_{m,n} \left( P^\gamma X^{\otimes m} \otimes Y^{T \otimes n} \right) \\
&= \sum_{r,s} \text{tr}_{m,n} \left( \left[ P^\gamma, C_{r,s} \right] X^{\otimes r-1} \otimes [X, Y] \otimes X^{\otimes m-r} \otimes Y^{T \otimes s-1} \otimes 1 \otimes Y^{T \otimes n-s} \right) \\
&= 0
\end{aligned}$$

$$C_{r,s} P^\gamma = P^\gamma C_{r,s} \quad (\text{Schur's lemma})$$

This is valid for any  $m, n, N$ .

$$\gamma = (\gamma_+, \gamma_-, k),$$

$$0 \leq k \leq \min(m, n), \quad \gamma_+ \mapsto (m - k), \quad \gamma_- \mapsto (n - k)$$

$(3, 2)$	$\gamma_+$	$\gamma_-$
$k = 0$	$[3]$	$[2]$
	$[2, 1]$	$[2]$
	$[1, 1, 1]$	$[2]$
	$[3]$	$[1, 1]$
	$[2, 1]$	$[1, 1]$
	$[1, 1, 1]$	$[1, 1]$
$k = 1$	$[2]$	$[1]$
	$[1, 1]$	$[1]$
$k = 2$	$[1]$	$\emptyset$

$$3X_s \quad 2Y_s$$

## On the complete set

[YK Ramgoolam 0709.2158, 0807.3696]

$$tr_{m,n} \left( P^\gamma X^{\otimes m} \otimes Y^{T \otimes n} \right)$$

The number of the operators is not enough to provide a complete set. A complete basis is given by

$$O_{A,ij}^\gamma(X, Y) \equiv tr_{m,n} \left( Q_{A,ij}^\gamma X^{\otimes m} \otimes (Y^T)^{\otimes n} \right)$$

$$P^\gamma = \sum_{A,i} Q_{A,ii}^\gamma$$

$$\begin{aligned} \gamma_+ &\mapsto (m-k), \quad \gamma_- \mapsto (n-k) & R &\mapsto m, \quad S \mapsto n \\ A &= (R, S) \end{aligned}$$

$$M_{\gamma \rightarrow A} = \sum_{\delta \mapsto k} g(\gamma_+, \delta; R) g(\gamma_-, \delta; S)$$

$$\left\langle O_{A_1, i_1 j_1}^{\gamma_1} (X, Y)^\dagger O_{A_2, i_2 j_2}^{\gamma_2} (X, Y) \right\rangle_0 \propto \delta^{\gamma_1 \gamma_2} \delta_{A_1 A_2} \delta_{i_1 i_2} \delta_{j_1 j_2}$$

[Orthogonality]

When  $k=0$ ,  $i, j=1$  and  $\gamma=A=(R, S)$ .

$$Q_{A,ij}^{\gamma(k=0)} = P^{\gamma(k=0, \gamma_+, \gamma_-)} = P_{RS}$$

## On the mixing pattern

An  $X$  and a  $Y$  are always combined after the action of the dilatation operator.

This means the  $k=0$  operators can not appear as the image of the dilatation operator.

$$\begin{aligned}\hat{H} \cdot O^{k=0} &= 0 \\ \hat{H} \cdot O^{k \neq 0} &= \sum_{k' \neq 0} O^{k'}\end{aligned}$$

In this sense, the  $k=0$  operators are not mixed with the other sectors ( $k \neq 0$ ).

$$\begin{aligned}\left\langle O^{\gamma(k=0)}(X, Y)^\dagger O_{A_2, i_2 j_2}^{\gamma_2(k \neq 0)}(X, Y) \right\rangle_1 &= 0 \\ \left\langle O^{\gamma_1(k=0)}(X, Y)^\dagger O^{\gamma_2(k=0)}(X, Y) \right\rangle_1 &\propto \delta^{\gamma_1 \gamma_2}\end{aligned}$$

## Summary

➤ Proposed to use *the representation basis* at finite  $N$

Young diagrams  $\Leftrightarrow$  multi-trace structure, flavour structure

Finite  $N$  constraint (stringy exclusion principle)

Orthogonal at classical level

➤ Will be useful to solve the mixing problem

The operator labelled by an irreducible rep of  $GL(N)$  is annihilated by  $D$ .

$$\hat{H} \cdot \text{tr}_{m,n} \left( P^\gamma X^{\otimes m} \otimes Y^{T \otimes n} \right) = 0$$

Looks like we moved to a proper language, with the help of *Brauer algebra*.

## Examples of the basis

1  $X$  1  $Y$

$$O^{\gamma(k=0)} = \text{tr} X \text{tr} Y - \frac{1}{N} \text{tr}(XY) \quad O^{\gamma(k=1)} = \frac{1}{N} \text{tr}(XY)$$

$(1, 1)$	$\gamma_+$	$\gamma_-$
$k = 0$	$[1]$	$[1]$
$k = 1$	$\emptyset$	$\emptyset$

2  $X$ s 1  $Y$

$$O^{\gamma(k=0, [2], [1])} = \frac{1}{2} \left( (\text{tr} X)^2 + \text{tr}(X)^2 \right) \text{tr} Y - \frac{1}{N+1} \left( \text{tr} X \text{tr}(XY) + \text{tr}(X^2 Y) \right)$$

$$O^{\gamma(k=0, [1, 1], [1])} = \frac{1}{2} \left( (\text{tr} X)^2 - \text{tr}(X)^2 \right) \text{tr} Y - \frac{1}{N+1} \left( \text{tr} X \text{tr}(XY) - \text{tr}(X^2 Y) \right)$$

$$O^{\gamma(k=1, [1], [0])} = \frac{2}{N^2 - 1} \left( N \text{tr} X \text{tr}(XY) - \text{tr}(X^2 Y) \right)$$

$(2, 1)$	$\gamma_+$	$\gamma_-$
$k = 0$	$[2]$	$[1]$
	$[1, 1]$	$[1]$
$k = 1$	$[1]$	$\emptyset$