

Operator analysis of magnetized T^2/Z_N orbifolds

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Collaborating with

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Mysteries of the Standard Model

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□ Chiral fermion

What is the origin of the chiral fermion ... ?

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Why so different the masses of the fermions are ... ?

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Why so different the masses of the fermions are ... ?

□ Flavor structure

What determine the flavor structure ... ?

⋮

Purpose

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We want to solve mysteries of the SM

- Chiral fermion Generations
- Mass hierarchy Flavor structure

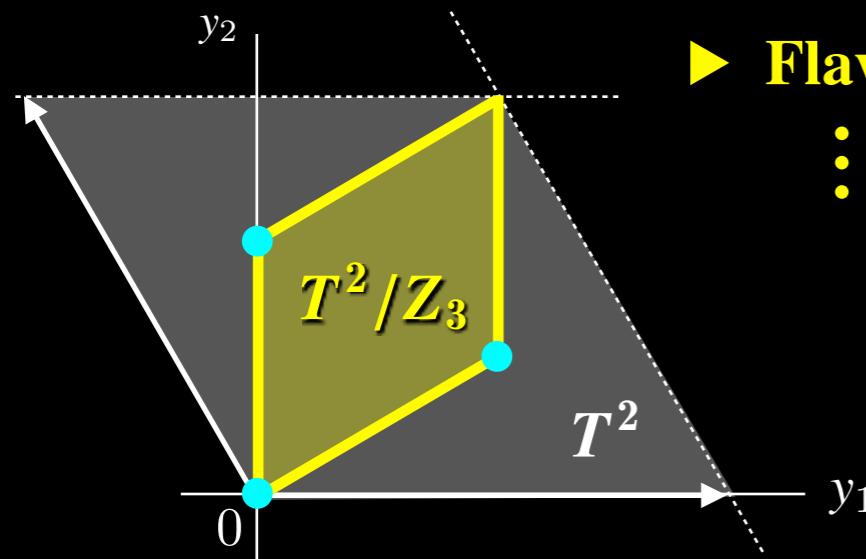
in the context of higher dimensional field theories.

Magnetic flux with Orbifold

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□ Orbifold

(e.g) T^2/Z_3

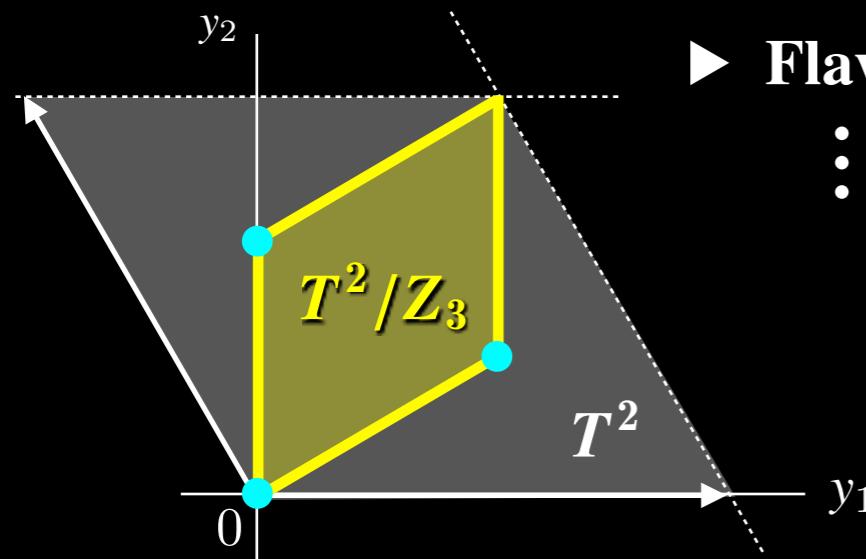


- ▶ Chiral fermion
- ▶ Mass hierarchy
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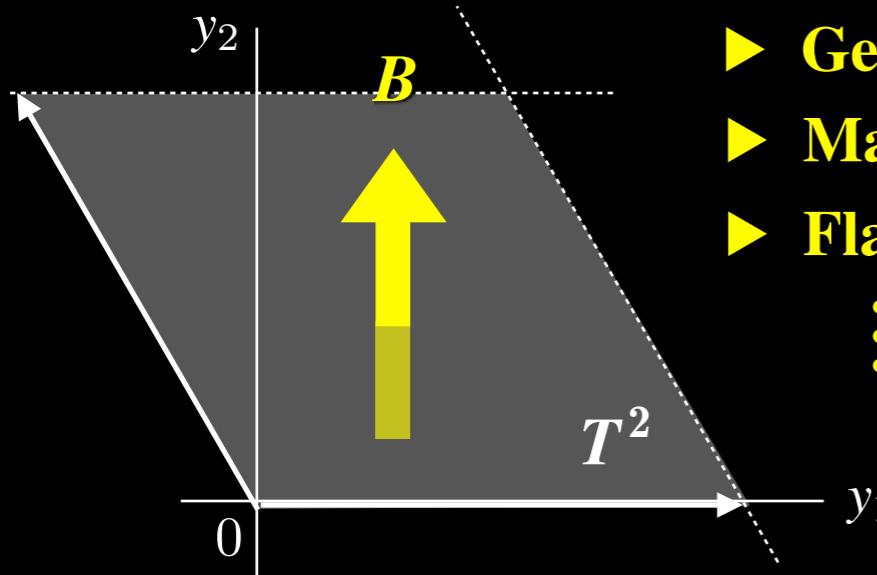
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□ Magnetic flux

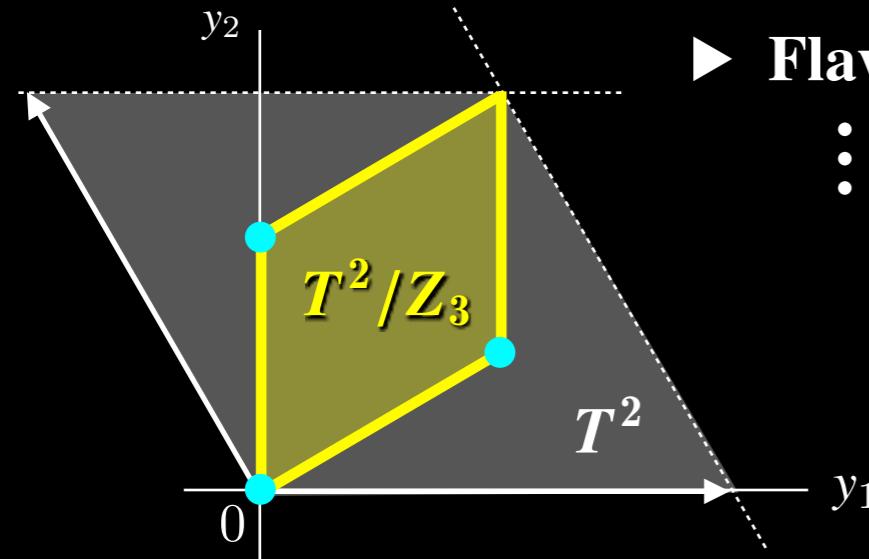


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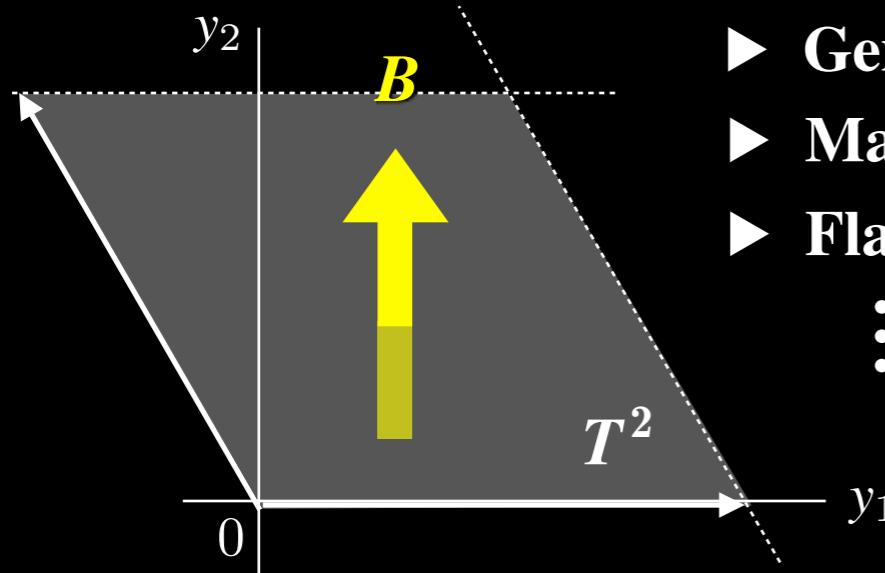
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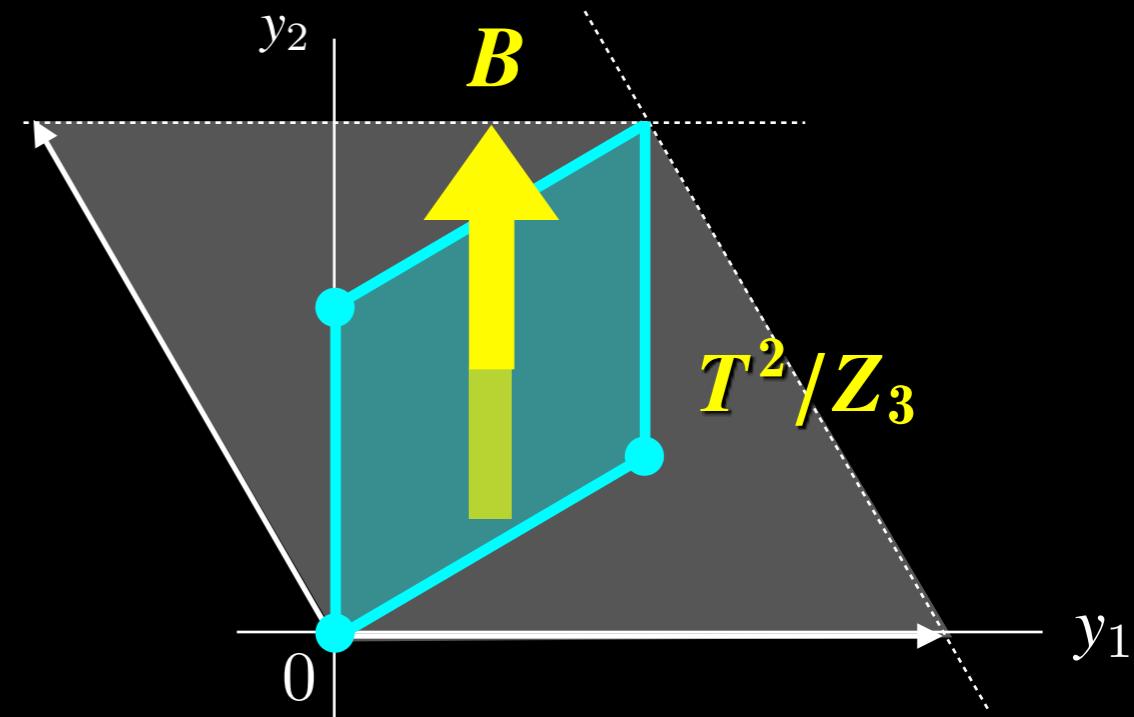
□ Magnetic flux



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□ Magnetized orbifold

(e.g) T^2/Z_3



**T-H.Abe, YF, T.Kobayashi,
T.Miura, K.Nishiwaki, M.Sakamoto,
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Features

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$$\psi_{\eta}^{(j)}(z)_{T^2/Z_N} = \mathcal{N} \sum_{k=0}^{N-1} \bar{\eta}^k \psi^{(j)}(\omega^k z)_{T^2}$$

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Complex coordinate
Zero-mode solution
on T^2 with a magnetic flux
Eigenstate of T^2/Z_N

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index for degenerated sol. : $j \in \{0, 1, 2, \dots, |M| - 1\}$
 $\omega \equiv e^{i \frac{2\pi}{N}}$

Z_N eigenvalue

$$\eta \in \{1, \omega, \omega^2, \dots, \omega^{N-1}\}$$

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(e.g) T^2/Z_3	M	2	4	6	8	\dots
η	1	1	1	3	3	\dots
ω	0	2	2	2	2	\dots
ω^2	1	1	1	1	3	\dots

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$$\psi_{\eta}^{(j)}(z)_{T^2/Z_N} = N \sum_{k=0}^{N-1} \bar{\eta}^k \psi^{(j)}(\omega^k z)_{T^2}$$
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$$\psi(z)_{T^2} \supset \vartheta \begin{bmatrix} a \\ b \end{bmatrix}(Mz, M\tau) = \sum_{l \in \mathbb{Z}} e^{i\pi(a+l)^2 M\tau} e^{2\pi i(a+l)(Mz+b)}$$

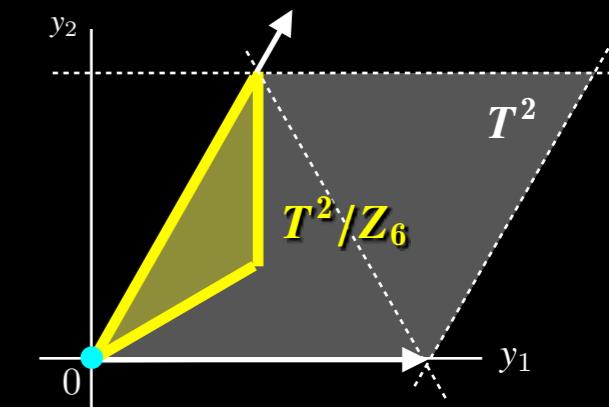
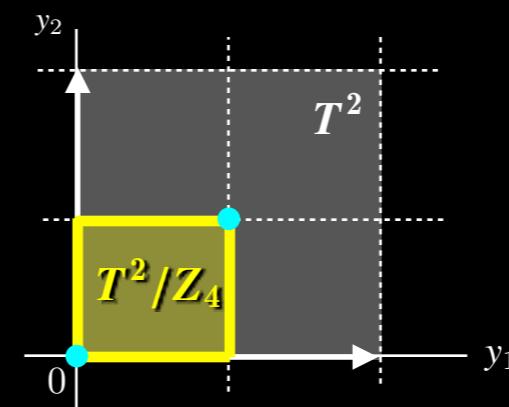
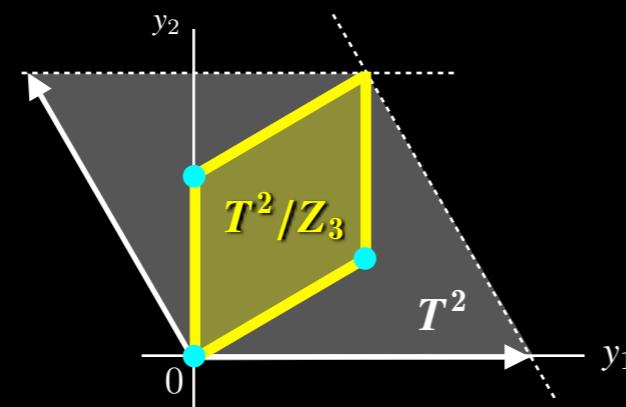
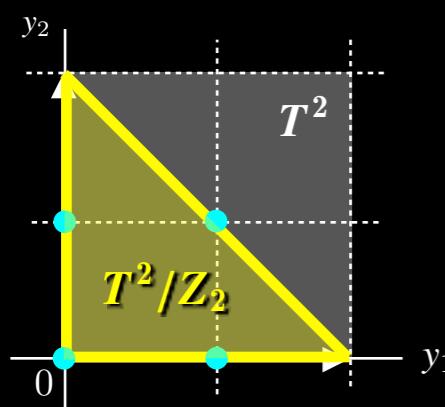
Our result

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□ Exact analytic results for

- The number of generations
- Expand coefficient
etc.

from 2d Quantum Mechanics analysis.

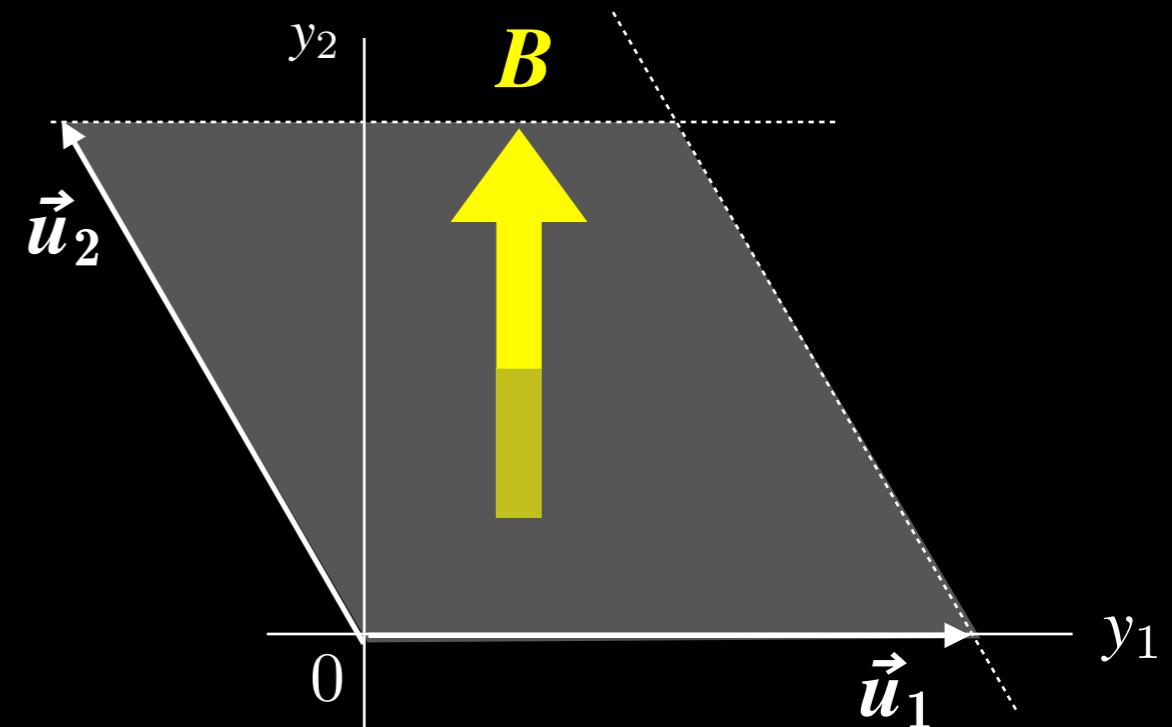


2d QM on magnetized orbifold



2d QM on magnetized orbifold

- Wave-function form ($a_w = 0$ basis)
 - └ Wilson line phase

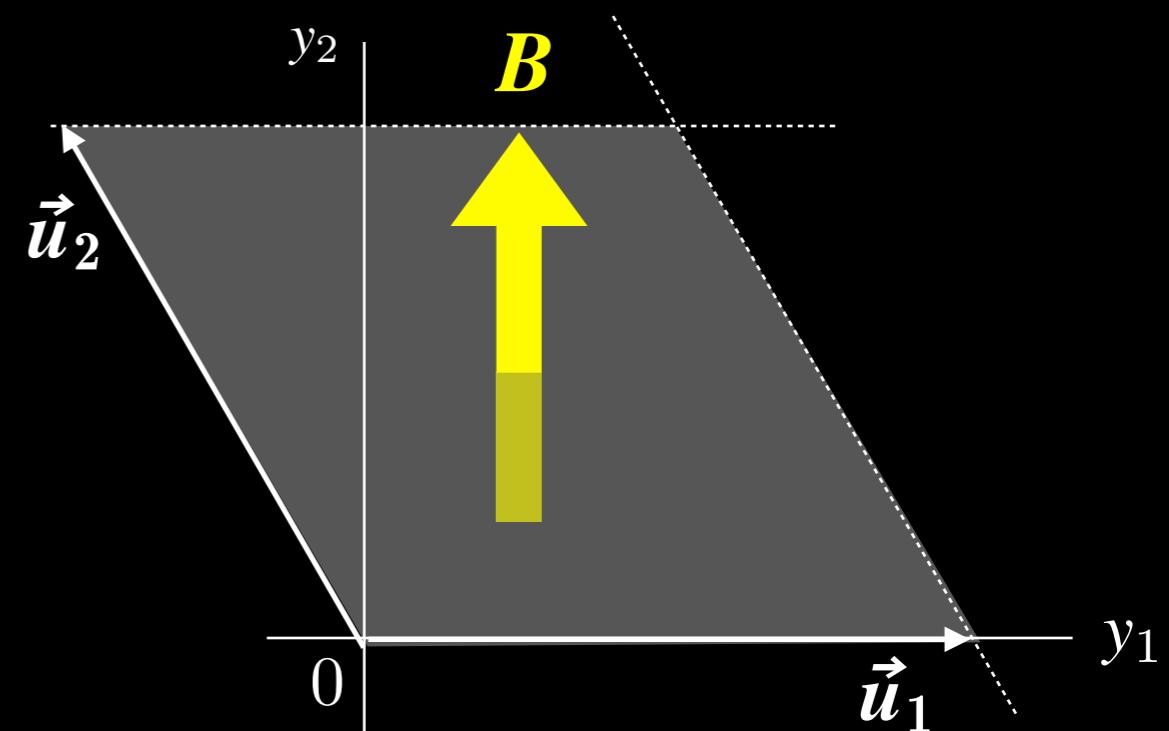


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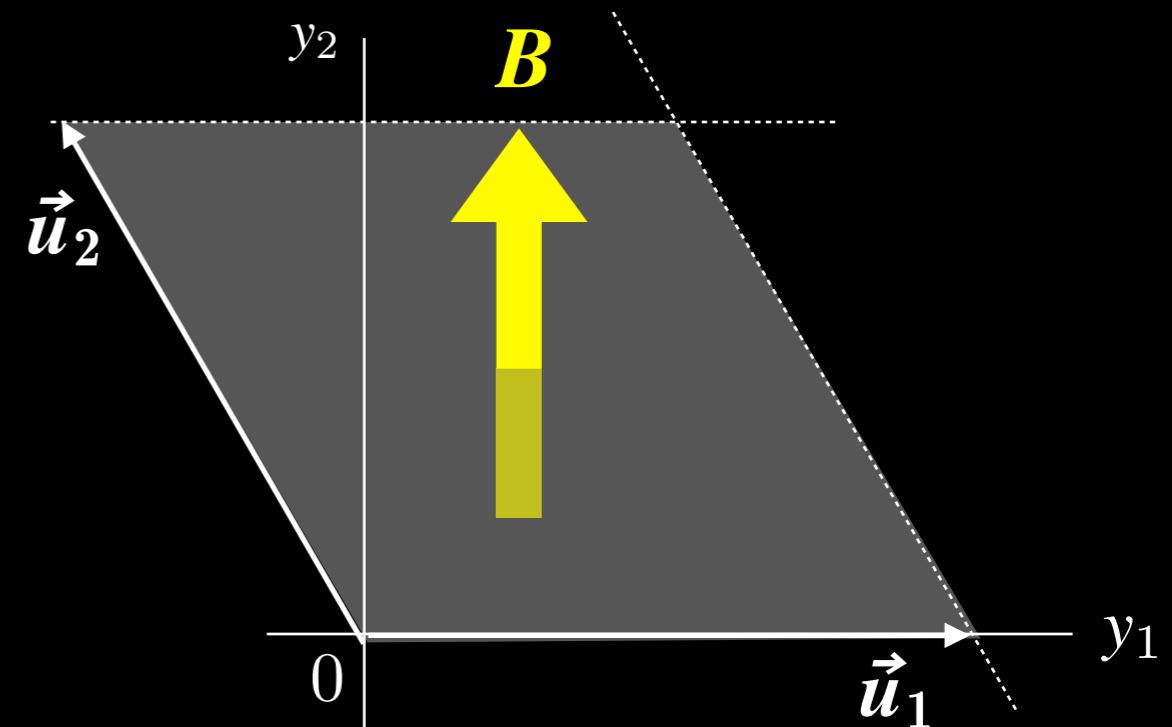
► Hamiltonian $H = (-i\vec{\nabla} - q\vec{A}(y))^2$

► Vector potential $\vec{A}(y) = -\frac{1}{2}\Omega\vec{y}$



2d QM on magnetized orbifold

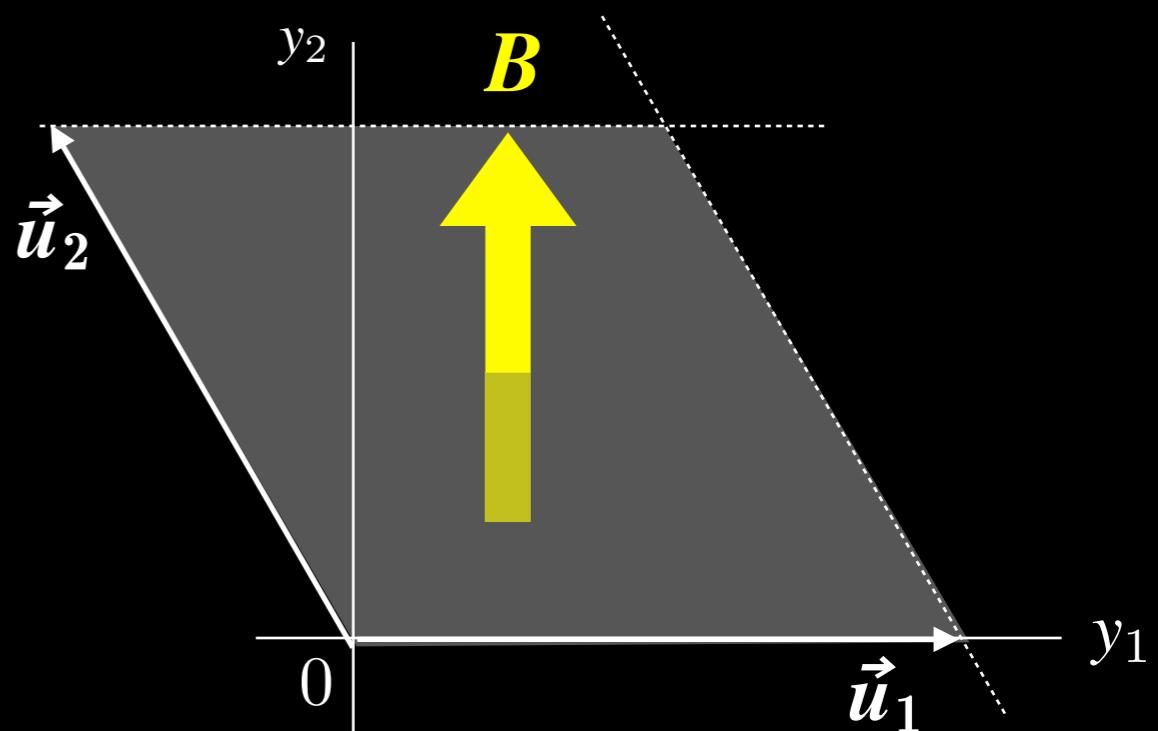
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- Pseudo-periodic $\psi(\vec{y} + \vec{u}_a) = e^{-i\frac{q}{2}\vec{y}^T\Omega\vec{u}_a + 2\pi i\alpha_a}\psi(\vec{y})$
($a = 1, 2$)



2d QM on magnetized orbifold

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- └ Basis of the torus



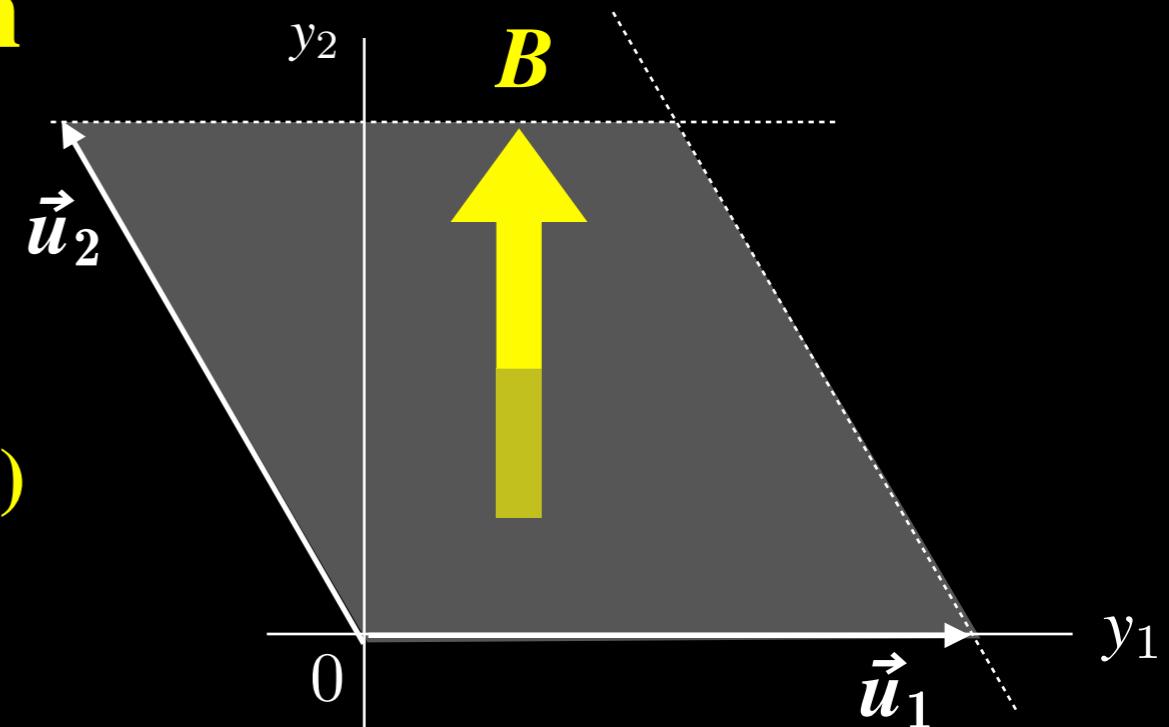
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($a = 1, 2$)
- ▶ Magnetic flux quantization

$$\frac{q(\vec{u}_1^T B \vec{u}_2)}{2\pi} = M \in \mathbb{Z}$$

where $B = \frac{1}{2}(\Omega - \Omega^T)$



2d QM on magnetized orbifold

- Operator formalism ($a_w = 0$ basis)

$$\psi(y) = \langle y | \psi \rangle, \quad \vec{p} \equiv -i\vec{\nabla}, \quad [\hat{y}_i, \hat{p}_j] = i\delta_{i,j}$$

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Canonical transformation

$$\{\hat{y}_1, \hat{y}_2, \hat{p}_1, \hat{p}_2\} \mapsto \{\hat{Y}, \hat{P}, \hat{\tilde{Y}}, \hat{\tilde{P}}\} \quad (\quad [\hat{Y}, \hat{P}] = i, \\ [\hat{\tilde{Y}}, \hat{\tilde{P}}] = i \quad)$$

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- Hamiltonian (same as the harmonic oscillator)

$$\hat{H} = \frac{1}{2}\hat{P}^2 + \frac{\omega^2}{2}\hat{Y}$$

- Constraint condition (BC's)

$$e^{i\hat{P}} |\psi\rangle = |\psi\rangle$$

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**Simultaneously
diagonalizable.**

2d QM on magnetized orbifold

- Eigenstates on T^2

$$\left| n, \frac{j}{M} \right\rangle_{T^2}$$

2d QM on magnetized orbifold

□ Eigenstates on T^2

Index for the degeneracy
(Eigenvalue of the operator \hat{Y})

$$\left| n, \frac{j}{M} \right\rangle_{T^2} \quad j \in \{0, 1, 2, \dots, |M| - 1\}$$

KK-index (from Hamiltonian)

$$n = 0, 1, 2, \dots$$

2d QM on magnetized orbifold

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$$\left(\leftrightarrow \psi^{(j)}(z)_{T^2} \right)$$

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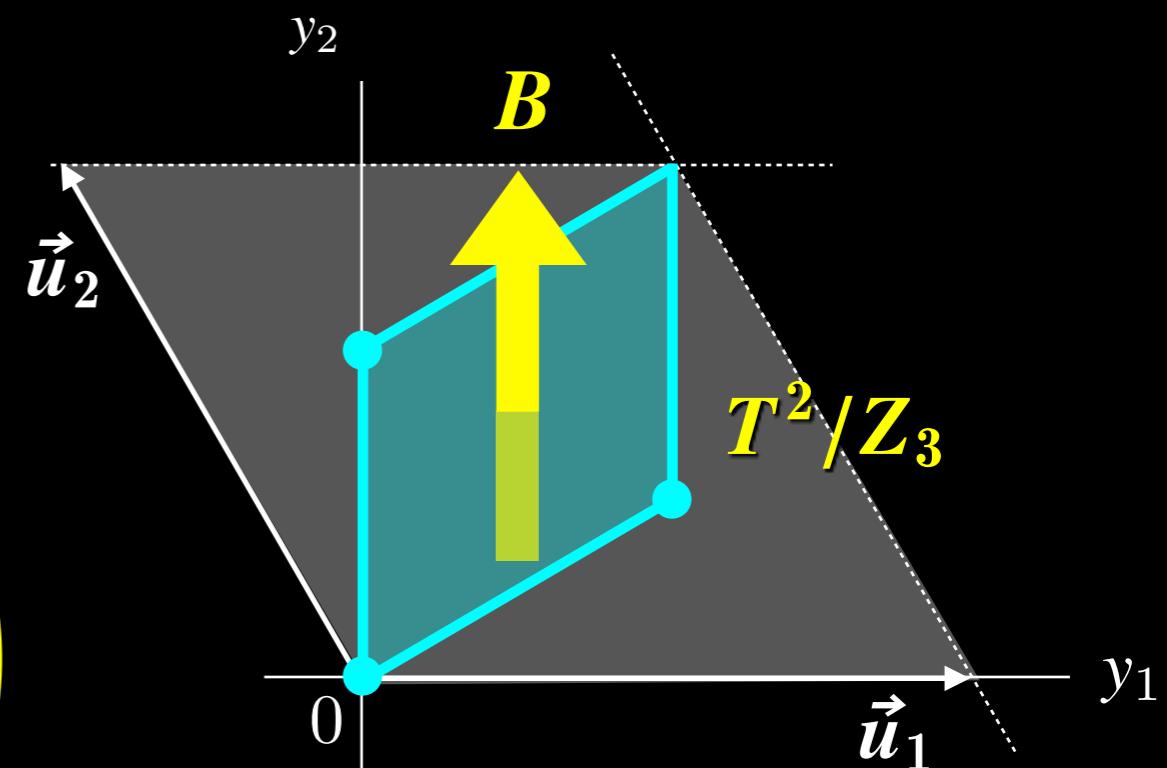
$$U \hat{\vec{y}} U^\dagger = R_\omega \hat{\vec{y}}$$

$$U \hat{\vec{p}} U^\dagger = R_\omega \hat{\vec{p}}$$

where

$$R_\omega = \begin{pmatrix} \cos \omega & -\sin \omega \\ \sin \omega & \cos \omega \end{pmatrix}$$

$$\omega \equiv e^{i \frac{2\pi}{N}}$$



2d QM on magnetized orbifold

- Expand coefficient

2d QM on magnetized orbifold

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► Operator formalism:

$$U_{Z_3} \left| n, \frac{j}{M} \right\rangle_{T^2} = \sum_k \tilde{C}_{jk} \left| n, \frac{k}{M} \right\rangle_{T^2}$$

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↑ theta functions ...

2d QM on magnetized orbifold

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(e.g.) **T^2/Z_3 orbifold**

2d QM on magnetized orbifold

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(e.g.) T^2/Z_3 orbifold

► Eigenstate of \hat{Y} from Z_3 -rotated state $U_{Z_3} \left| n, \frac{j}{M} \right\rangle_{T^2}$

2d QM on magnetized orbifold

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$$\left| n, 0 \right\rangle_{T^2} \propto \sum_{k=1}^{M-1} e^{2\pi i \eta_k} U_{Z_3} \left| n, \frac{k}{M} \right\rangle_{T^2}$$

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$$\downarrow$$

$$U_{Z_3} \left| n, \frac{j}{M} \right\rangle_{T^2} = \sum_k \tilde{C}_{jk} \left| n, \frac{k}{M} \right\rangle_{T^2}$$

$$\boxed{\tilde{C}_{jk} = \frac{1}{\sqrt{M}} e^{-i \frac{\pi}{12} + \frac{3\pi\alpha^2}{M} + i\pi \frac{k(k+6\alpha)}{M} + 2\pi i \frac{j \cdot k}{M}}} : \text{Exact form !!}$$

2d QM on magnetized orbifold

□ Physical states

2d QM on magnetized orbifold

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$$\left| n, \frac{j}{M} \right\rangle_{T^2/Z_N, \eta} = \frac{1}{N} \sum_{l=0}^{N-1} \bar{\eta}^l \left(\hat{U}_{Z_N}^l \left| n, \frac{j}{M} \right\rangle_{T^2} \right)$$

2d QM on magnetized orbifold

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Eigenvalue of Z_N orbifold

($\eta = \omega^k$; $\omega \equiv e^{i\frac{2\pi}{N}}$, $k = 0, 1, 2, \dots, N-1$)

2d QM on magnetized orbifold

□ Physical states

$$\begin{aligned} \left| n, \frac{j}{M} \right\rangle_{T^2/Z_N, \eta} &= \frac{1}{N} \sum_{l=0}^{N-1} \bar{\eta}^l \left(\hat{U}_{Z_N}^l \left| n, \frac{j}{M} \right\rangle_{T^2} \right) \\ &= \sum_{k=0}^{M-1} M_{jk} \left| n, \frac{k}{M} \right\rangle_{T^2} \end{aligned}$$

2d QM on magnetized orbifold

- Physical states

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- # of physical states = Rank [M_{jk}]

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$$= \sum_{k=0}^{M-1} M_{jk} \left| n, \frac{k}{M} \right\rangle_{T^2}$$

□ # of physical states = Rank [M_{jk}]

(e.g) T^2/Z_3

η	2	4	6	8	\cdots
1	1	1	3	3	\cdots
ω	0	2	2	2	\cdots
ω^2	1	1	1	3	\cdots

Consistent with the numerical analysis !!

2d QM on magnetized orbifold

New formula...?

2d QM on magnetized orbifold

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(e.g.) T^2/Z_3 case ($\tau = e^{i\frac{2\pi}{3}}$)

2d QM on magnetized orbifold

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$$C_{jk} = \int_{T^2} dz d\bar{z} \psi^{(j)}(\omega z)_{T^2} \psi^{(k)*}(z)_{T^2}$$

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$$\psi^{(j)}(z)_{T^2} = \mathcal{N} e^{i\pi M z \frac{\text{Im } z}{\text{Im } \tau}} \vartheta \begin{bmatrix} \frac{j+\alpha_1}{M} \\ -\alpha_\tau \end{bmatrix} (Mz, M\tau)$$

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$$C_{jk} = \int_{T^2} dz d\bar{z} \psi^{(j)}(\omega z)_{T^2} \psi^{(k)*}(z)_{T^2}$$

$$\psi^{(j)}(z)_{T^2} = \underbrace{\mathcal{N} e^{i\pi M z \frac{\text{Im } z}{\text{Im } \tau}} \vartheta \begin{bmatrix} \frac{j+\alpha_1}{M} \\ -\alpha_\tau \end{bmatrix} (Mz, M\tau)}$$

$$\vartheta \begin{bmatrix} a \\ b \end{bmatrix} (\nu, \tau) = \sum_{l \in \mathbb{Z}} e^{i\pi(a+l)^2\tau + 2\pi i(a+l)(\nu+b)}$$

2d QM on magnetized orbifold

□ New formula...?

(e.g.) T^2/Z_3 case ($\tau = e^{i\frac{2\pi}{3}}$)

$$C_{jk} = \int_{T^2} dz d\bar{z} \psi^{(j)}(\omega z)_{T^2} \psi^{(k)*}(z)_{T^2}$$

$$\psi^{(j)}(z)_{T^2} = \mathcal{N} e^{i\pi M z \frac{\text{Im } z}{\text{Im } \tau}} \vartheta \left[\begin{array}{c} j+\alpha_1 \\ -\alpha_\tau \end{array} \right] (Mz, M\tau)$$

$$\vartheta \left[\begin{array}{c} a \\ b \end{array} \right] (\nu, \tau) = \sum_{l \in \mathbb{Z}} e^{i\pi(a+l)^2\tau + 2\pi i(a+l)(\nu+b)}$$

Operator formalism said

$$\downarrow$$

$$= \frac{1}{\sqrt{M}} e^{-i\frac{\pi}{12} + \frac{3\pi\alpha^2}{M} + i\pi \frac{k(k+6\alpha)}{M} + 2\pi i \frac{j\cdot k}{M}} \quad (\alpha \equiv \alpha_1 = \alpha_2)$$

Conclusion and Discussion

Conclusion and Discussion

- We obtained exact analytic results for

- The number of generations
- Expand coefficient
etc.

from 2d Quantum Mechanics analysis.

