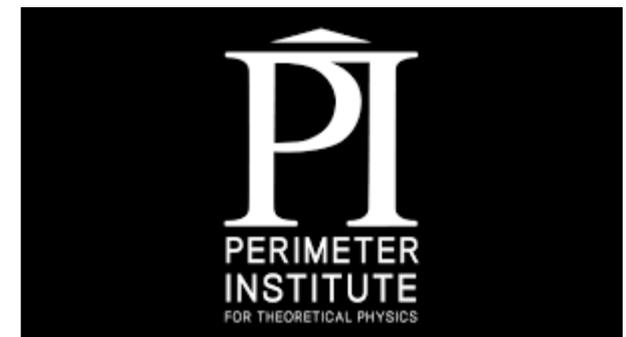


Entanglement branes, Modular flow, and Extended quantum field theory



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Based on collaboration with William Donnelly: [hep-th 1811.10785](https://arxiv.org/abs/hep-th/1811.10785)

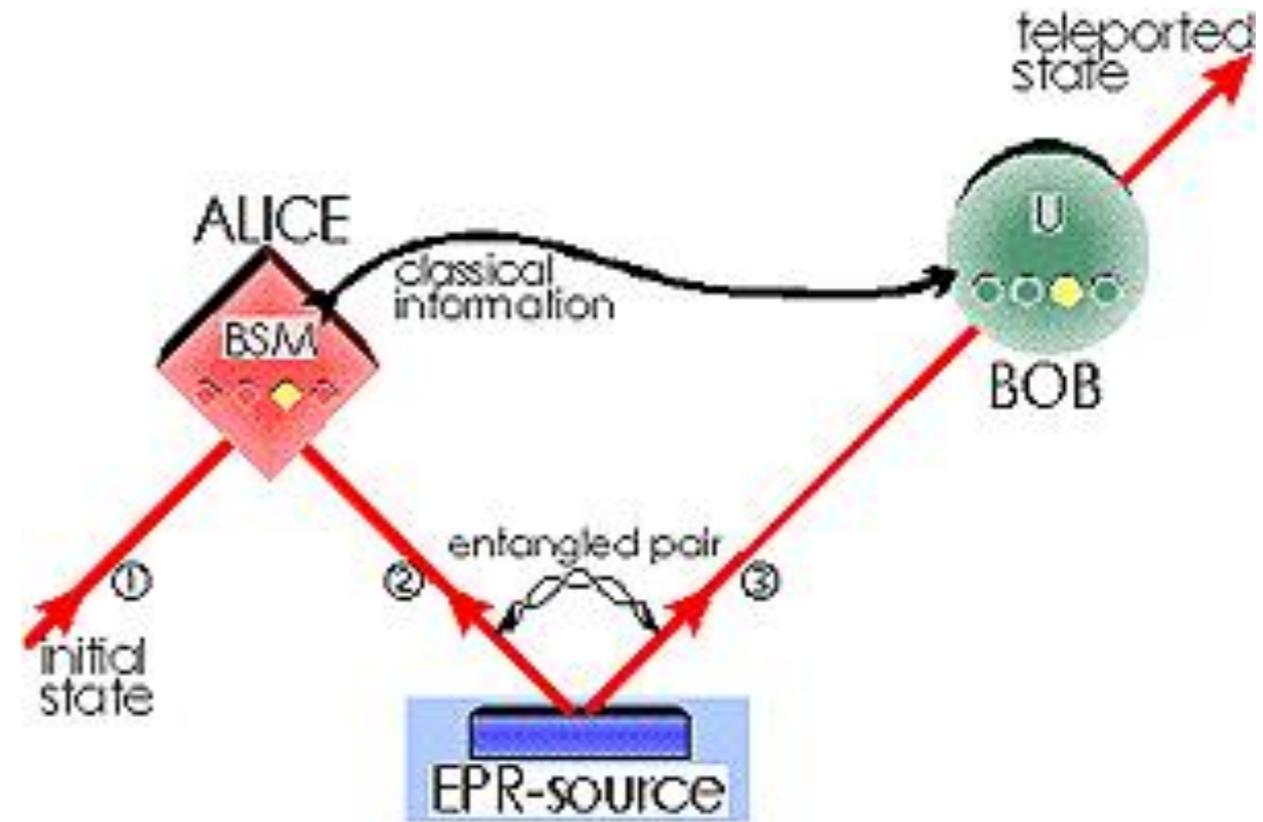
Quantum Entanglement is a **non local** feature of Quantum Mechanics

Unentangled

$$|0\rangle|0\rangle$$

Entangled EPR

$$\frac{1}{\sqrt{2}}(|1\rangle|0\rangle - |0\rangle|1\rangle)$$

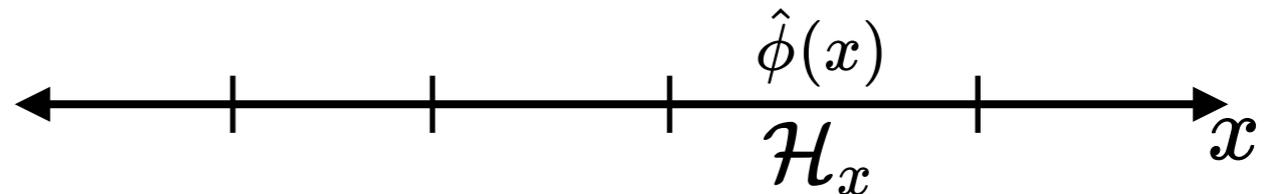


- Quantum Teleportation via EPR Pairs
- Non-local order parameter in Topological phases
- Many body localization
- Quantum gravity: Emergent smooth spacetimes from quantum entanglement

What do we mean by **locality** in quantum mechanics?

Hilbert Space factorization

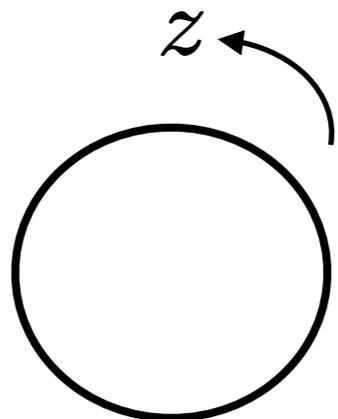
$$\mathcal{H} = \otimes_x \mathcal{H}_x$$



Different notions of locality can be assigned to the same Hilbert space

$$\mathcal{H} = \{ \text{Anti-symmetric wavefunctions } \psi(z_1, \dots, z_N), z_j = e^{i\theta_j} \}$$

N non-relativistic fermions in
on a spatial circle



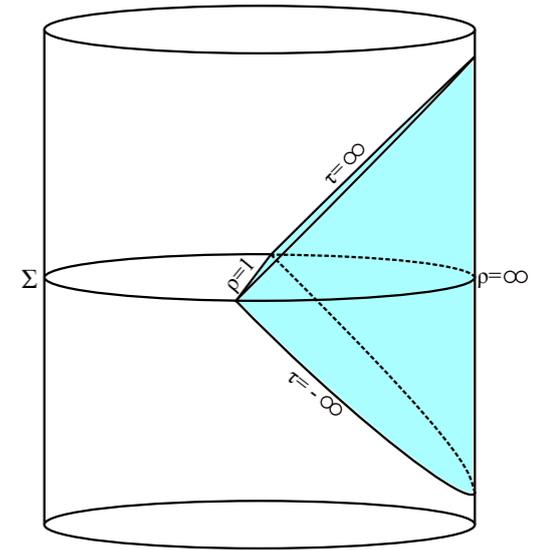
U(N) Yang Mills on a spatial
circle .

$$U = \text{P exp} \left(i \int A_x dx \right)$$

$$= \left(\begin{matrix} z_1 & & & \\ & \dots & & \\ & & & z_N \end{matrix} \right)$$

AdS/CFT and Bulk locality

- AdS/CFT provides a QG Hilbert space at asymptotic infinity
- How is the local bulk spacetime encoded in the CFT Hilbert space at infinity?



A free fermion Hilbert space in N=4 SYM

Ten- Dimensional Geometry for the IIB string
(Lin, Lunin, Maldacena)

1/2 BPS sector of N=4 SYM

Free Fermions in a Magnetic field in the LLL

$$ds^2 = -h^{-2}(dt + V_i dx^i)^2 + h^2(dy^2 + dx^i dx^i) + ye^G d\Omega_3^2 + ye^{-G} d\tilde{\Omega}_3^2$$

$$h^{-2} = 2y \cosh G,$$

$$y\partial_y V_i = \epsilon_{ij} \partial_j z, \quad y(\partial_i V_j - \partial_j V_i) = \epsilon_{ij} \partial_y z$$

$$z = \frac{1}{2} \tanh G$$

$$F = dB_t \wedge (dt + V) + B_t dV + d\hat{B},$$

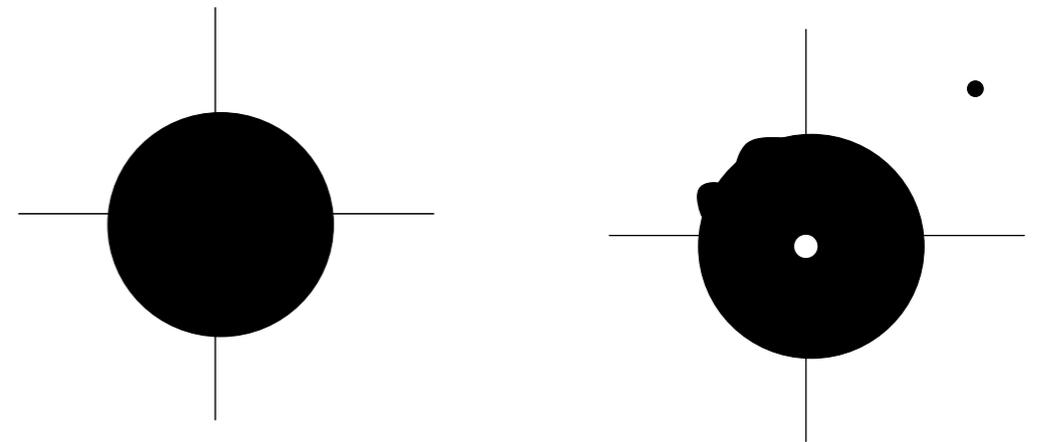
$$\tilde{F} = d\tilde{B}_t \wedge (dt + V) + \tilde{B}_t dV + d\hat{\tilde{B}}$$

$$B_t = -\frac{1}{4} y^2 e^{2G},$$

$$\tilde{B}_t = -\frac{1}{4} y^2 e^{-2G}$$

$$d\hat{B} = -\frac{1}{4} y^3 *_3 d\left(\frac{z + \frac{1}{2}}{y^2}\right),$$

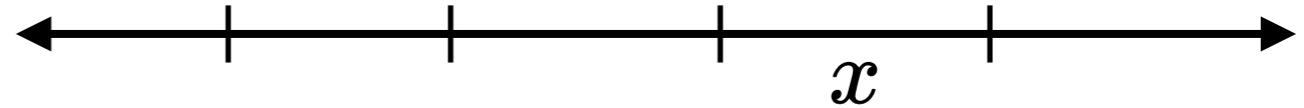
$$d\hat{\tilde{B}} = -\frac{1}{4} y^3 *_3 d\left(\frac{z - \frac{1}{2}}{y^2}\right)$$



The extended Hilbert space construction

Hilbert Space factorization

$$\mathcal{H} = \bigotimes_x \mathcal{H}_x$$

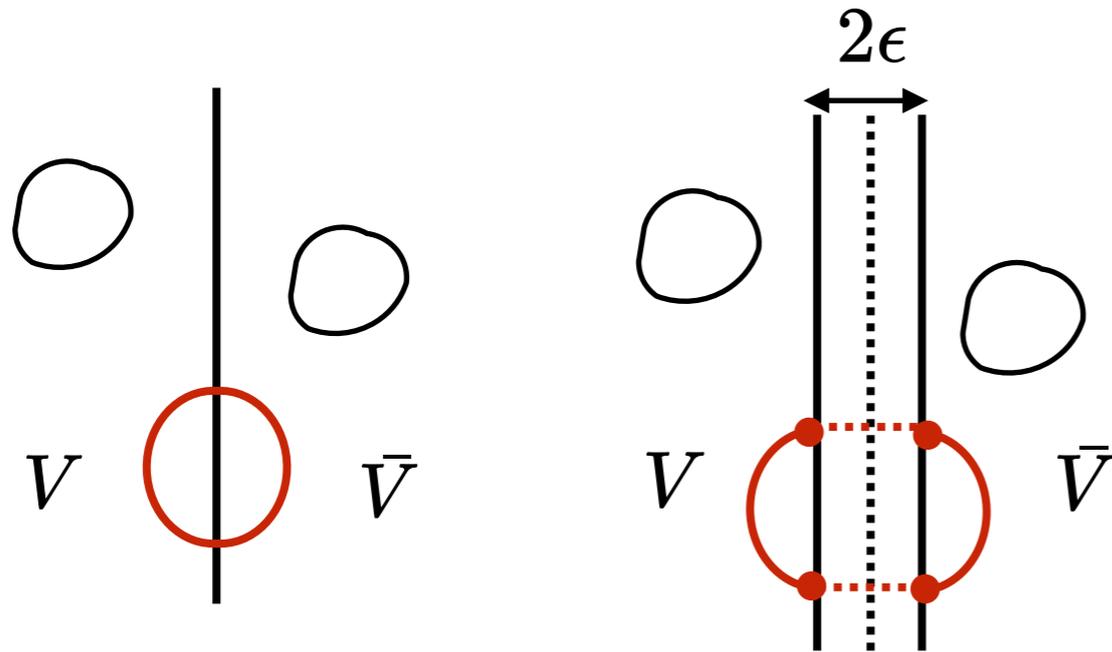


Two obstructions

- 1 A subregion has a boundary and therefore edge modes, even for a scalar field! (Agon, Headrick, Jefferis, Kasko) (Campagnia, Freidel, et al)
- 2 Degrees of freedom in subregions are not independent
 - continuity in a quantum field theory
 - Gauss Law constraint in gauge theory. Even on a lattice !

The extended Hilbert space construction provides a solution by combining 1 and 2 (Donnelly, Freidel, Buividovich, ...)

Extended Hilbert space for gauge theories



$$\mathcal{H}_{\text{physical}} \subset \mathcal{H}_V \otimes \mathcal{H}_{\bar{V}}$$

Contains edge modes transforming under boundary symmetry group $G_L \times G_R$

Gauss law

$$Q|\psi\rangle = 0 \quad \text{for} \quad |\psi\rangle \in \mathcal{H}_{\text{Physical}}$$

Invariance under $G = \text{Diag} (G_L \times G_R)$

Entangling product

$$\mathcal{H}_{\text{physical}} = \mathcal{H}_V \otimes_G \mathcal{H}_{\bar{V}}$$

Reduced density matrix

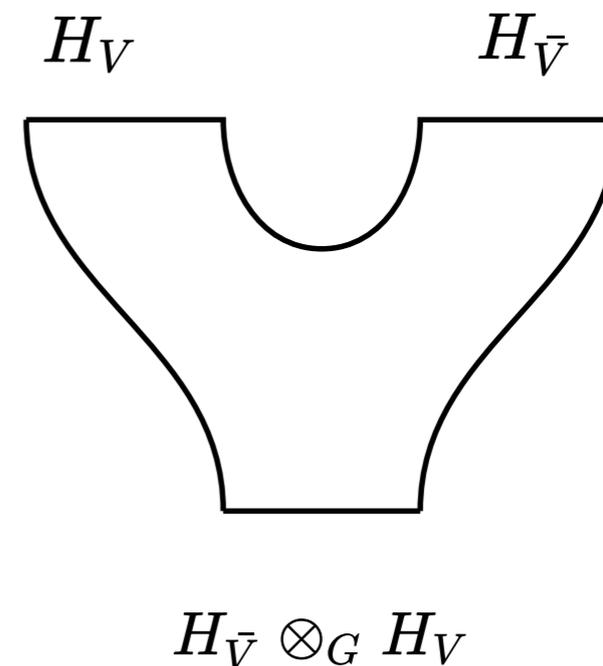
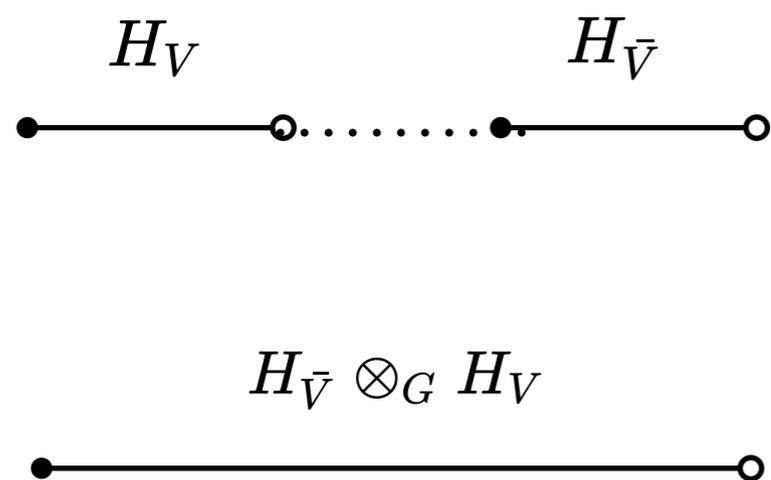
$$\rho_V = \text{tr}_{\bar{V}} |\psi\rangle \langle \psi|$$

Entanglement Entropy

$$S_V = -\text{tr} \rho_V \log \rho_V$$

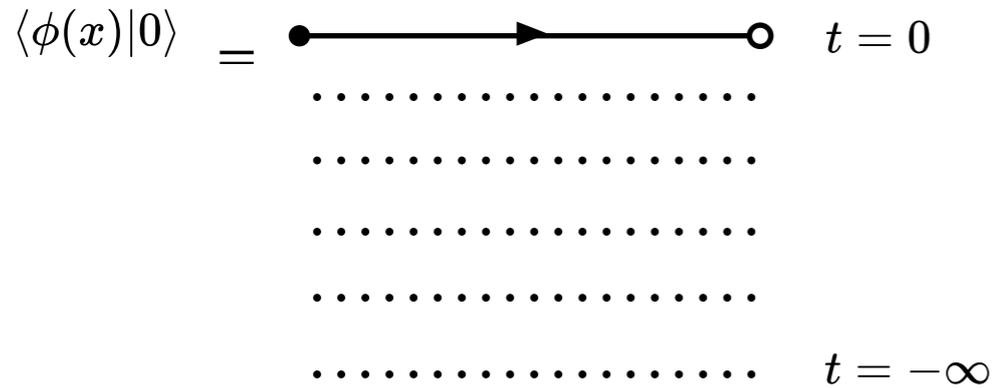
Extended Hilbert space and extended TQFT

- Edge modes are not unique e.g. in quantum hall states (Cano, Cheng, Mulligan, ...et. al) (Fliss, Wen, Parrikar, ..et. al.)
- What are the rules for determining the “correct” edge modes and their gluings?
- In 2D, we provide constraints on the Hilbert space extension using the frame work of **extended topological quantum field theory**
- Key insight: View the entangling product as a spacetime process=cobordism

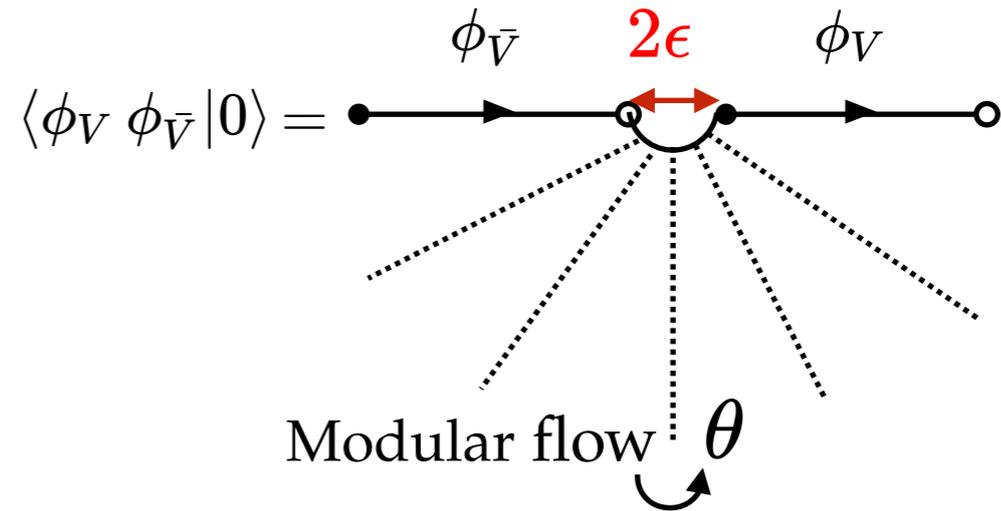


Entangling product from the path integral

Euclidean path integral prepares the (unnormalized) vacuum



Angular quantization

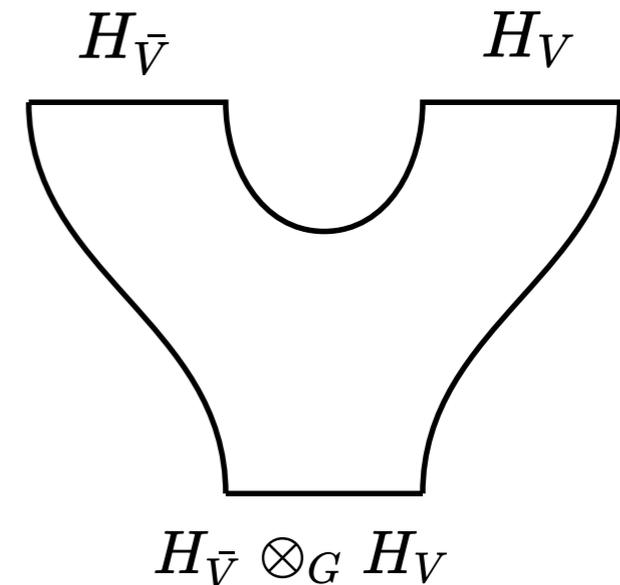
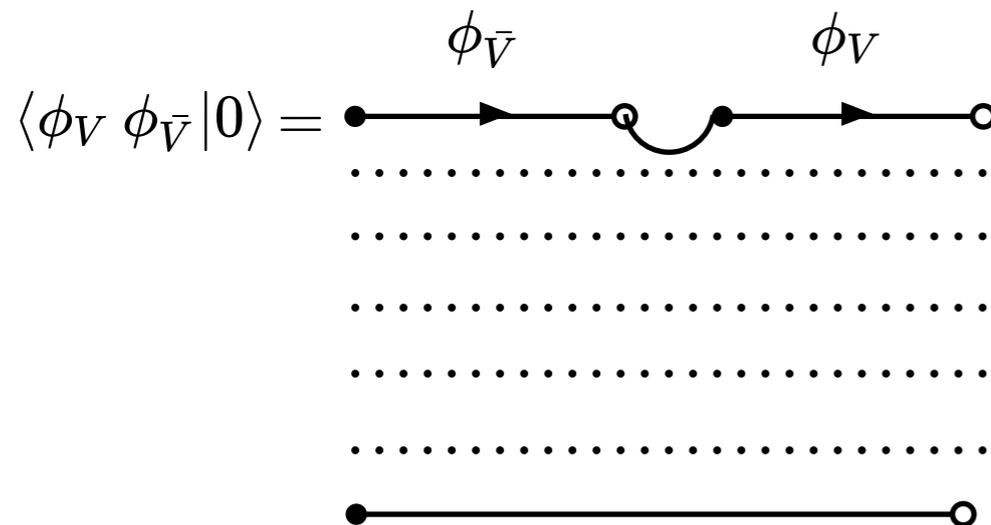


State-Channel duality

Modular Hamiltonian

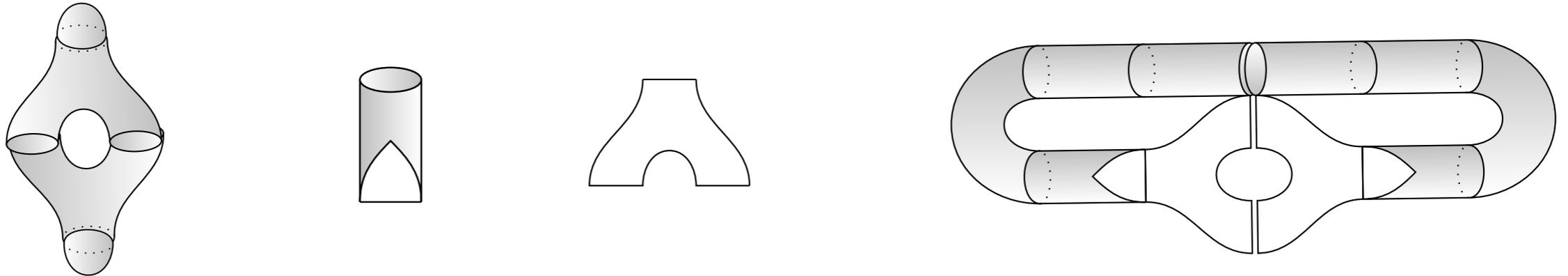
$$\langle \phi_V \phi_{\bar{V}} | 0 \rangle = \langle \phi_V | \exp\left(\frac{-H_V}{2}\right) \mathcal{J} | \phi_{\bar{V}} \rangle \stackrel{\text{CPT}}{=} \sum_n \exp\left(-\frac{E_n}{2}\right) \langle \phi_V | n \rangle \langle \phi_{\bar{V}} | \bar{n} \rangle$$

Entangling product in CS gauge theory (Wong 2018)



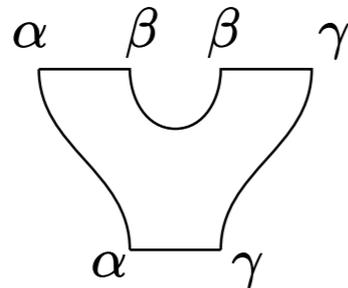
Locality in Extended TQFT

Cut path integral along surfaces of increasing codimension (Atiyah, Segal, Freed, Baez,...)

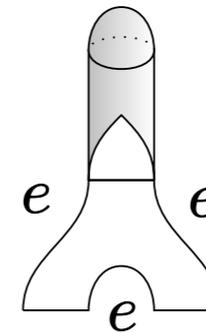


Moore-Segal gave sewing constraints for cutting and gluing path integrals. These rules determine allowed **boundary conditions=D branes**.

D branes



E branes



What we did:

- Introduce the **Entanglement brane** boundary condition
- Interpret Moore-Segal as constraints for extended Hilbert space and edge modes
- Formulate 2D Yang Mills as an extended TQFT a la Moore-Segal
- Compute multi-interval modular flows, EE, negativity

Outline

- Open-closed extended TQFT (Moore-Segal)
- Entanglement brane
- Multi-interval Modular flows, EE
- 2DYM as an open-closed TQFT
- Future works: CFT, higher dimensions, holography

Atiyah's formulation of Axiomatic TQFT

In 2D, a TQFT is a rule assigning

1-dim closed manifolds = Hilbert space over \mathbb{C}



\mathcal{H}_{S^1}



$\mathcal{H}_{-S^1} = \mathcal{H}_{S^1}^*$

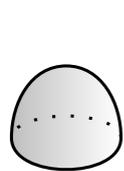


$\mathcal{H}_{S^1} \otimes \mathcal{H}_{S^1}$

\emptyset

\mathbb{C}

Cobordism between circles = Linear maps (quantum evolution)



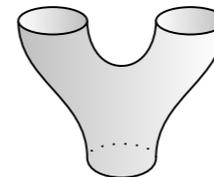
\mathbb{C}
↓
 \mathcal{H}_{S^1}

Wavefunction $\Psi[\phi(x)]$



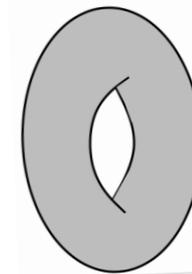
\mathcal{H}_{S^1}
↓ e^{-tH}
 \mathcal{H}_{S^1}

Propagator=Identity



$\mathcal{H}_{S^1} \otimes \mathcal{H}_{S^1}$
↓
 \mathcal{H}_{S^1}

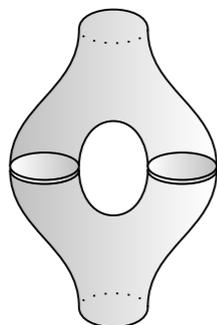
Fusion/Multiplication



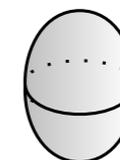
\mathbb{C}
↓
 \mathbb{C}

Partition function

Gluing Cobordisms = Composing linear maps



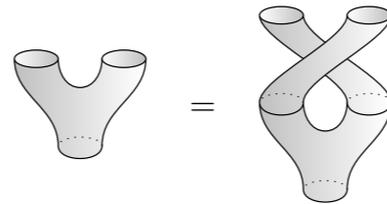
\mathcal{H}_{S^1}
↓
 $\mathcal{H}_{S^1} \otimes \mathcal{H}_{S^1}$
↓
 \mathcal{H}_{S^1}



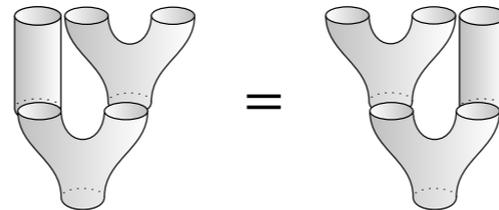
$$\int D\phi(x) \Psi[\phi(x)] \Psi^*[\phi(x)] = \langle \Psi | \Psi \rangle$$

A 2D Closed TQFT is a commutative Frobenius algebra

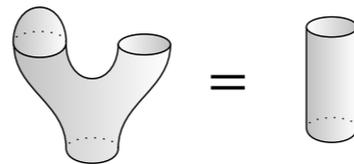
Commutative



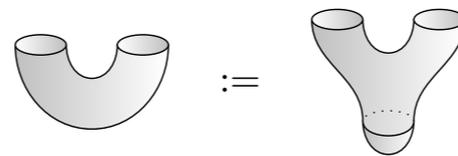
Associative



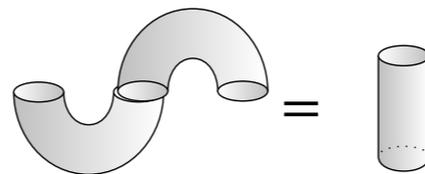
Unit



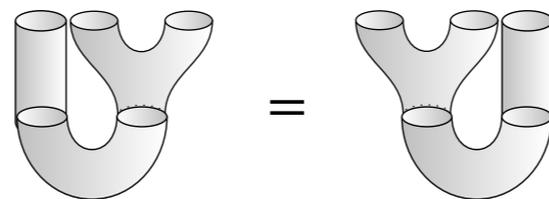
Symmetric bilinear form



Invertible



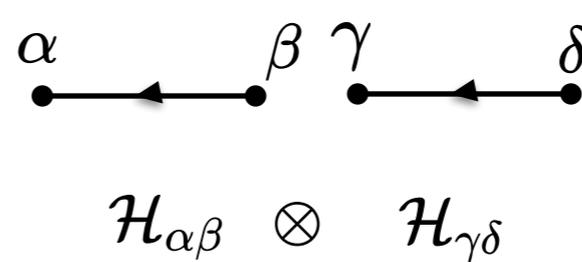
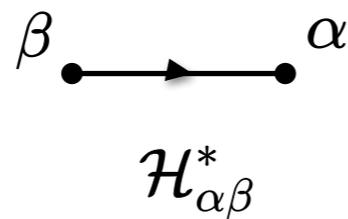
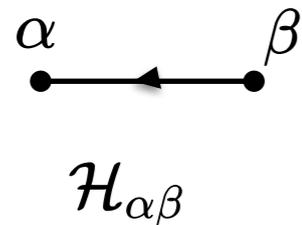
Invariance



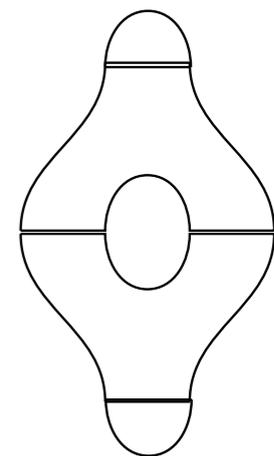
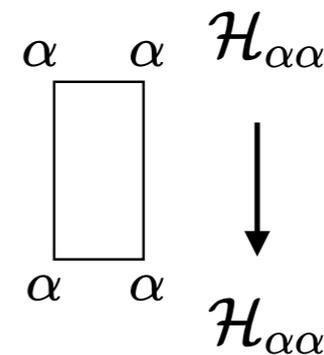
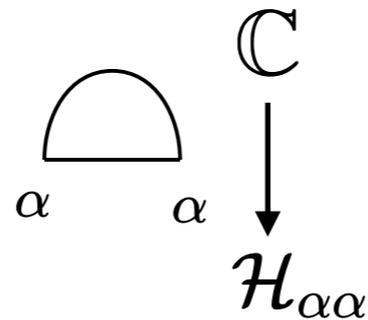
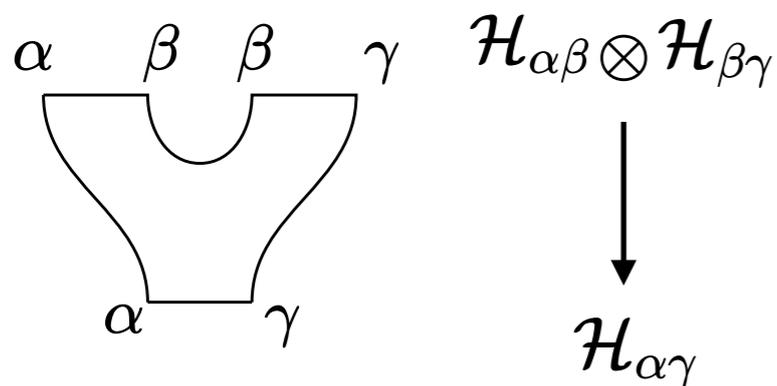
Open TQFT is a symmetric Frobenius Algebra

An open TQFT assigns

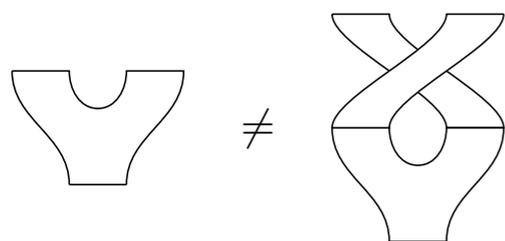
Hilbert space to oriented intervals with boundary conditions :



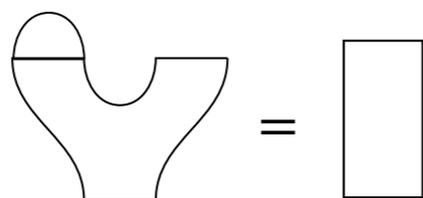
Open cobordisms to linear maps



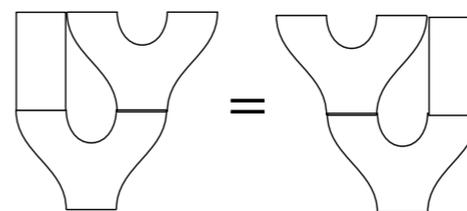
Non-commutative mult.



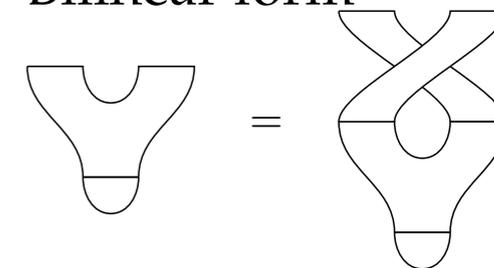
Unit



Associative

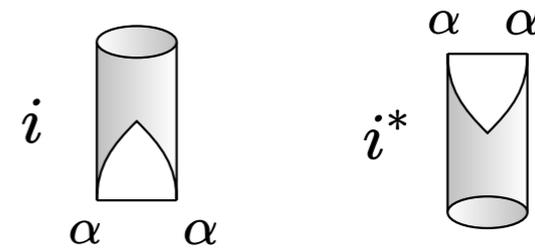


Invariant symmetric
Bilinear form

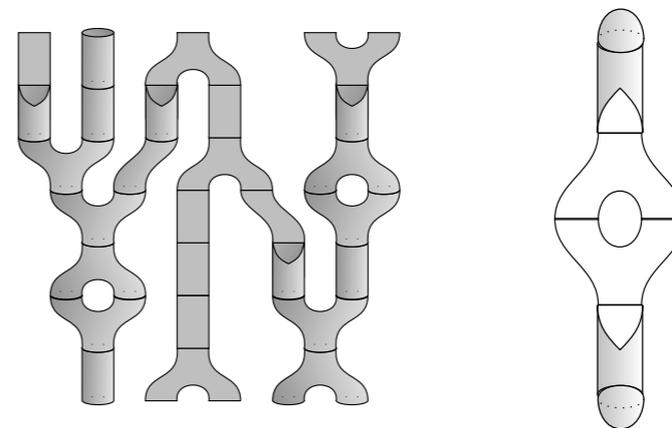
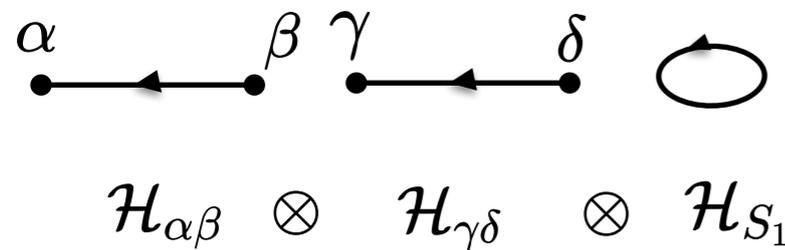


Open closed TQFT

The zipper relates the closed and open algebra...

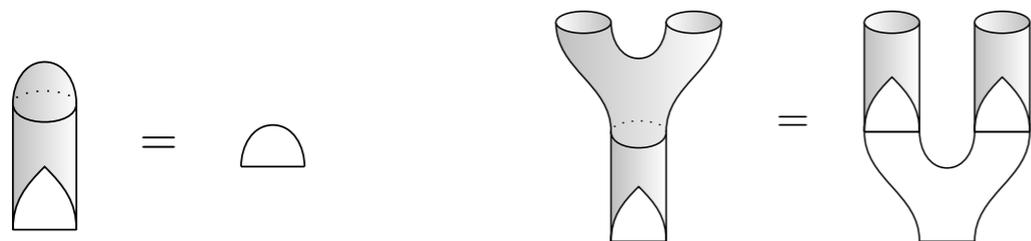


Open-closed Hilbert spaces and cobordisms:

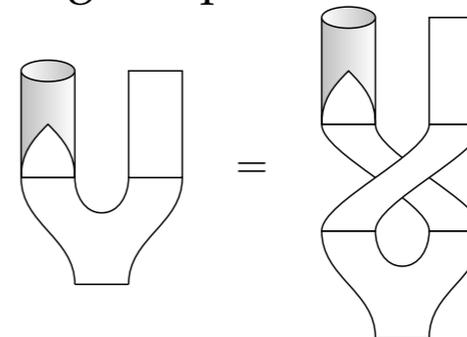


Moore-Segal Sewing rules : Ensures compatibility of gluing

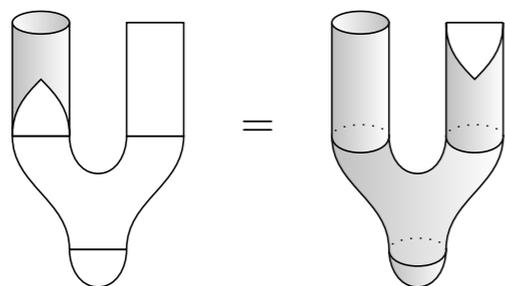
i is an algebra homomorphism



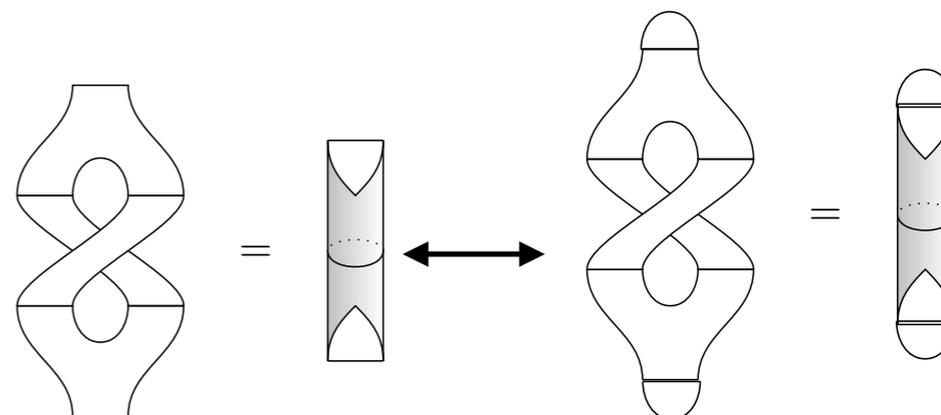
Closed strings maps to the center of open strings



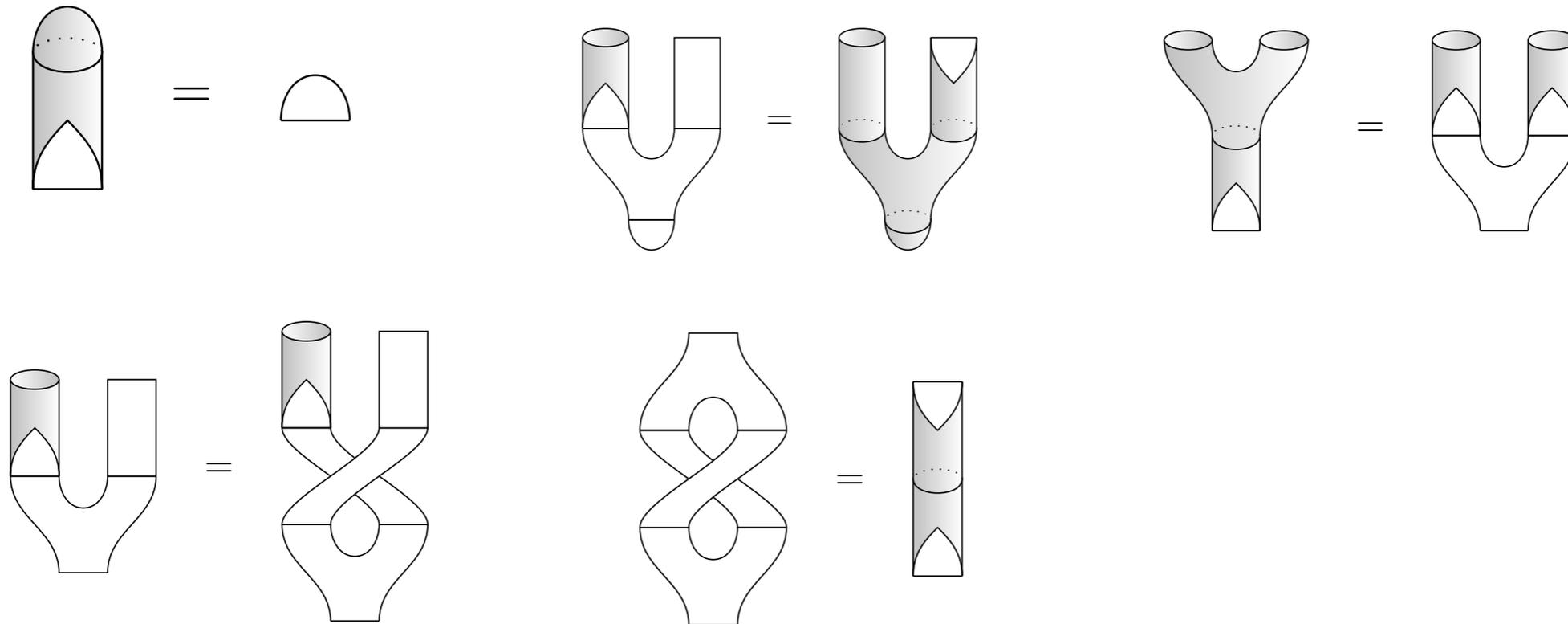
i is the adjoint of i^*



Cardy (Spacetime covariance)



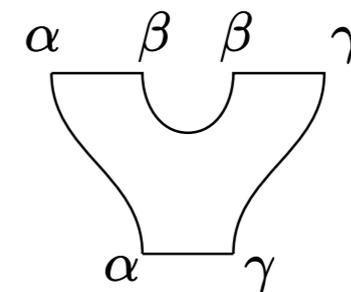
Moore-Segal Sewing relations



Moore-Segal:

Q: Given a closed string theory, what are the possible boundaries, i.e. D Branes?

A: D branes correspond to extensions to an open string algebra satisfying these constraints.



For us: Open string algebra ~ choice of Hilbert space extension i.e. edge modes

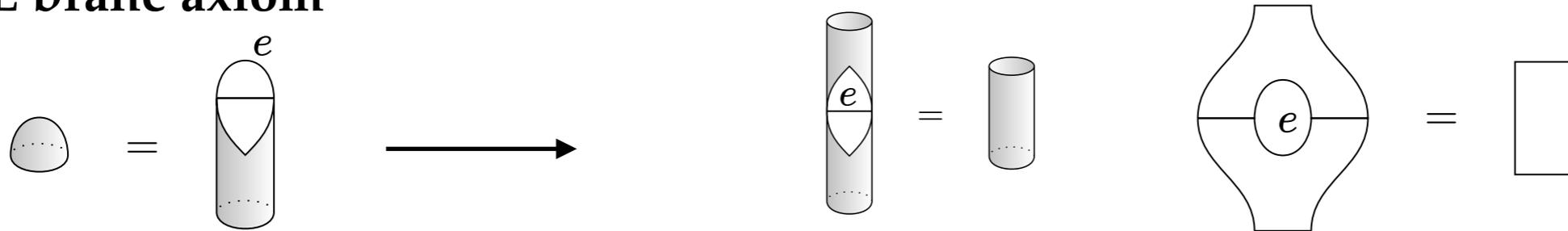
Outline

- Open-closed extended TQFT (Moore-Segal)
- Entanglement brane
- Multi-interval Modular flows, EE
- 2DYM as an open closed TQFT
- Future works: CFT, higher dimensions, holography

The Entanglement Brane axiom

Holes originating from splitting the Hilbert space can be sewed up

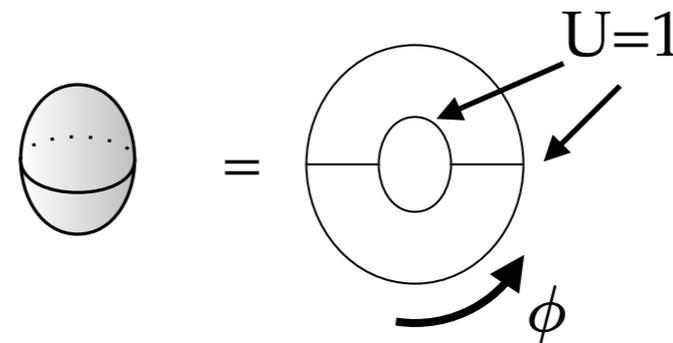
E brane axiom



e= choice of (possibly nonlocal) boundary conditions

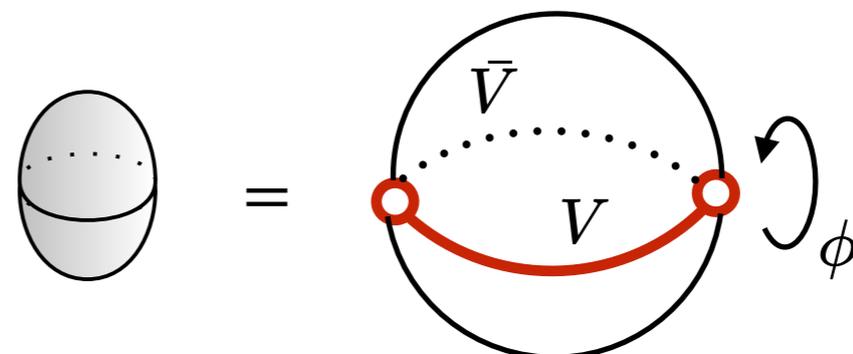
In 2D Yang Mills: e = trivial holonomy along boundary circles
 \sim sum over electric boundary conditions.

$$U = P \exp(i \int A_x dx)$$



Implies correlations are preserved under reduction to V :

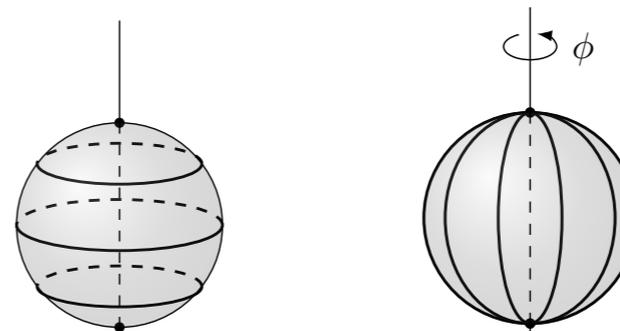
$$\langle O_V \rangle = \text{tr}(\rho_V O_V)$$



The Entanglement Brane in a toy string theory

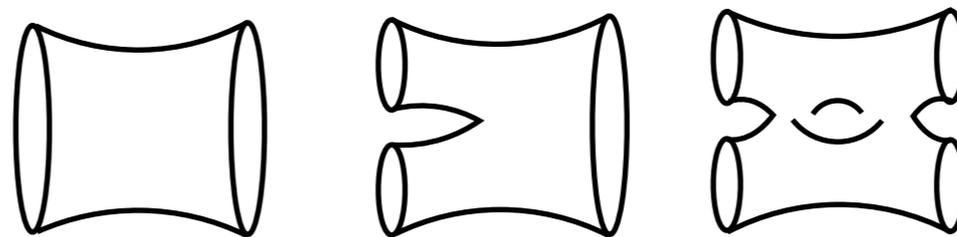
2D Yang Mills = Closed String theory (Gross-Taylor)

Entanglement brane axiom
relates open and closed sector:

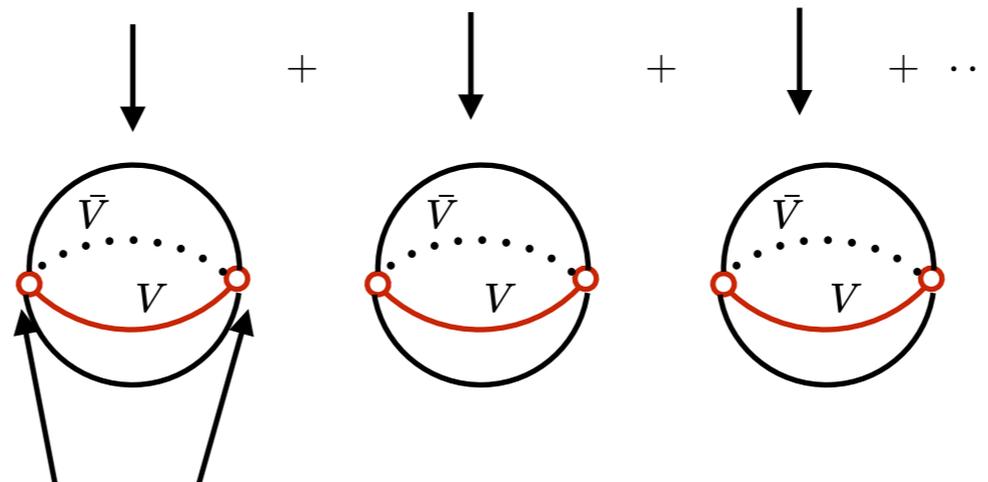


c.f. Susskind-Uglum

Cutting a closed string results in $N = \frac{1}{g_{\text{string}}}$ Chan-Paton factors.



$$Z_{\text{YM}}(S^2) = \text{tr}_V \rho_V =$$



N D0-branes

$$N = \frac{1}{g_{\text{string}}}$$

Outline

- Open-closed extended TQFT (Moore-Segal)
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Single interval Modular flow

Tensor product factorization

$$|\psi\rangle = \begin{array}{c} \text{[Diagram: A vertical tube with a shaded top section and a horizontal line through the middle, connected to a base with two lobes. Labels } \bar{V} \text{ and } V \text{ are below the lobes.]} \\ \bar{V} \quad V \end{array} = \begin{array}{c} \text{[Diagram: A shape with a rounded top and a central indentation, resembling a wide 'U' with a hump.]} \\ \end{array} = \begin{array}{c} \text{[Diagram: A simple semi-circular arc.]} \\ \end{array}$$

State-Channel duality

$$\hat{\psi} = \begin{array}{c} \text{[Diagram: A shape with a central hump and two side lobes. Labels } \bar{V} \text{ and } V \text{ are above and below the lobes.]} \\ \bar{V} \quad V \end{array} = \begin{array}{c} \text{[Diagram: A vertical rectangle. Labels } \bar{V} \text{ and } V \text{ are above and below the rectangle.]} \\ \bar{V} \\ V \end{array}$$

Unnormalized reduced density matrix

$$\rho_V = \begin{array}{c} \text{[Diagram: A circular shape with a central hole. Labels } \bar{V} \text{ and } V \text{ are at the top and bottom of the hole.]} \\ \bar{V} \quad V \\ \bar{V} \quad V \end{array} = \hat{\psi}\hat{\psi}^\dagger = \begin{array}{c} \text{[Diagram: A vertical rectangle divided into two equal horizontal sections. Labels } V \text{ are above and below the rectangle.]} \\ V \\ V \end{array}$$

Effective partition function

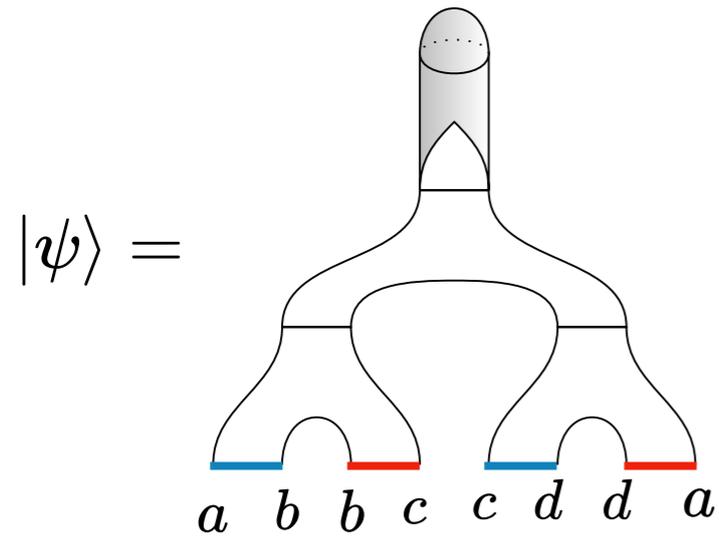
$$Z = \text{tr}_V \rho_V = \begin{array}{c} \text{[Diagram: A torus (donut shape).]} \\ \end{array} = \begin{array}{c} \text{[Diagram: A shaded sphere with a horizontal line through the middle.]} \\ \end{array}$$

Entanglement entropy

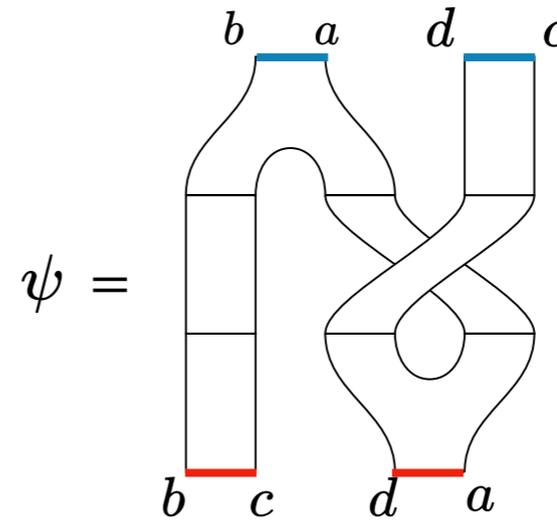
$$S = -\text{tr}_V \frac{\rho_V}{Z} \log \frac{\rho_V}{Z}$$

Multi-interval Modular flow

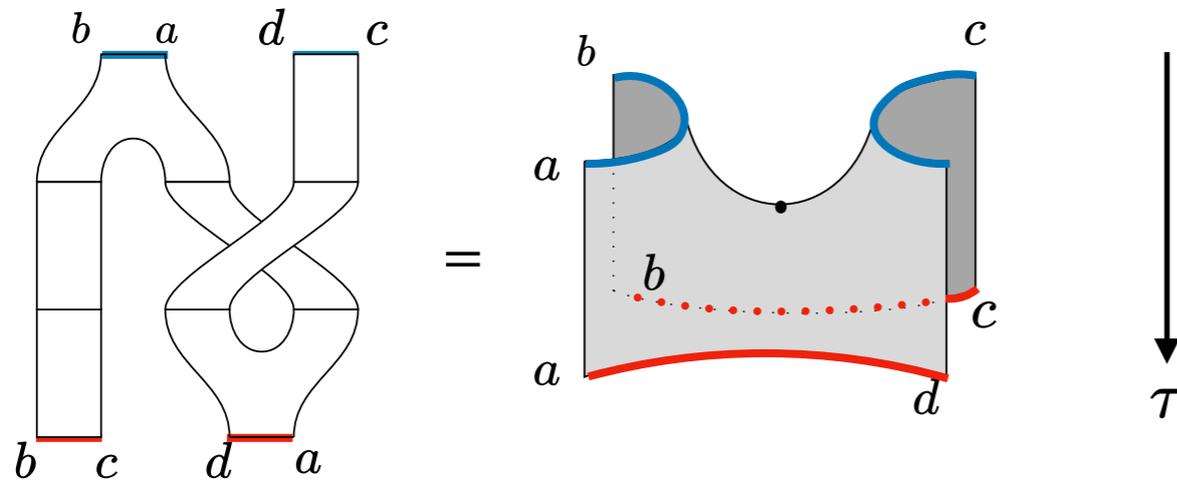
Tensor product factorization



State-Channel duality



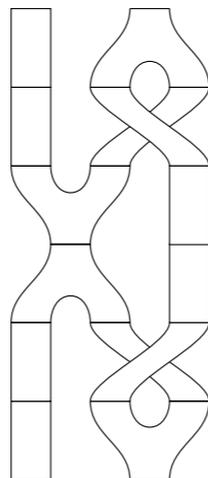
Saddle point



Modular time τ is a morse function !

Reduced density matrix

$$\rho = \psi\psi^\dagger =$$



$$Z = \text{tr}\rho =$$

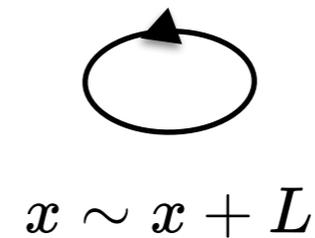
Outline

- Open-closed extended TQFT (Moore-Segal)
- Entanglement brane
- Multi-interval Modular flows, EE, and Negativity
- 2DYM as an open-closed TQFT
- Future works: CFT, higher dimensions, holography

2DYM as a closed TQFT

Configuration space

$$U = \mathcal{P} \exp \left(i \int_0^L dx A_x(x) \right) \in G$$



Hilbert space on a circle = **Class functions on G**

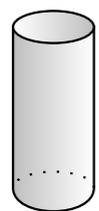
Representation Basis

$$\Psi[U] = \Psi[gUg^{-1}]$$

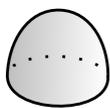
$$\langle U | R \rangle = \text{Tr}_R(U)$$

Hamiltonian $\sim \text{tr}(E^2) =$ **Casimir**

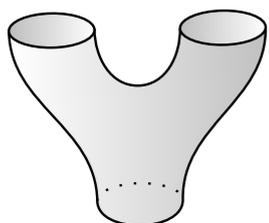
$$H |R\rangle = \frac{g_{\text{YM}}^2 L}{2} C_2(R) |R\rangle$$



$$= \sum_R e^{-AC_2(R)} |R\rangle \langle R|$$



$$= \sum_R \dim R e^{-AC_2(R)} |R\rangle$$



$$= \sum_R \frac{1}{\dim R} e^{-AC_2(R)} |R\rangle \langle R| \langle R|$$



Euler characteristic

$$Z(M) = \sum_R (\dim R)^{\chi(M)} e^{-AC_2(R)}$$



2DYM as an open TQFT

Configuration space

$$U = \text{P exp} \int_a^b A_x dx$$



Hilbert space on an interval

General functions on gauge group G

Basis

$$\langle U | Rab \rangle = \sqrt{\dim R} R_{ab}(U)$$

Edge modes $a, b = 1, \dots, \dim R$

Boundary symmetry :

$$U \rightarrow g(a)Ug^{-1}(b)$$

$$R \rightarrow GRG^{-1}$$

Entangling product = Matrix Multiplication

$$R_{ac}(U_{\bar{V}}U_V) = \sum_b R_{ab}(U_{\bar{V}})R_{bc}(U_V)$$



$$= \sum_{R,a,b,c} \frac{e^{-AC_2(R)}}{\sqrt{\dim R}} |Rab\rangle |Rbc\rangle \langle Rac|$$

$$= \sum_{R,a} \frac{e^{-AC_2(R)}}{\sqrt{\dim(R)}} |Raa\rangle \langle R|$$

$$= \text{[Cylinder with pair of pants]} = \sum_{R,a} \sqrt{\dim R} e^{-AC_2(R)} |Raa\rangle$$

Single interval Modular flow and EE

Tensor product factorization

$$|\psi\rangle = \text{[Diagram 1]} = \text{[Diagram 2]} = \text{[Diagram 3]} = \sum_{R,a,b} e^{\frac{-AC_2(R)}{2}} |Rab\rangle |Rba\rangle$$

Effective partition function

$$Z = \text{tr}_V \rho_V = \text{[Diagram 4]} = \sum_{R,a,b} e^{-AC_2(R)} = \sum_R (\dim R)^2 e^{-AC_2(R)} = \text{[Diagram 5]}$$

Entanglement entropy in terms of $P(R) = \frac{(\dim R)^2 e^{-AC_2(R)}}{Z}$

Edge modes

$$S = -\text{tr}_V \frac{\rho_V}{Z} \log \frac{\rho_V}{Z} = \sum_R -P(R) \log P(R) + 2P(R) \log \dim R$$

Multi-interval Modular flow

State-Channel duality

$$|\psi\rangle = \text{[Diagram: A genus-2 surface with a vertical tube on top and four horizontal intervals at the bottom labeled a, b, b, c, c, d, d, a.]} \longrightarrow \psi = \text{[Diagram: A genus-2 surface with four horizontal intervals at the top labeled b, a, d, c and four at the bottom labeled b, c, d, a.]} = \sum_{R,a,b,c,d} \frac{e^{\frac{-AC_2(R)}{2}}}{\dim R} |Rbc\rangle |Rcd\rangle \langle Rba| \langle Rdc|$$

Density matrix

$$\rho_V = \psi\psi^\dagger = \text{[Diagram: A genus-2 surface with four horizontal intervals at the top labeled b, a, d, c and four at the bottom labeled b, c, d, a.]} = \sum_{R,a,b,c,d} \frac{e^{-AC_2(R)}}{(\dim R)^2} |Rbc\rangle |Rda\rangle \langle Rbc| \langle Rda|$$

Entropy

$$P(R) = \frac{(\dim R)^2 e^{-AC_2(R)}}{Z}$$

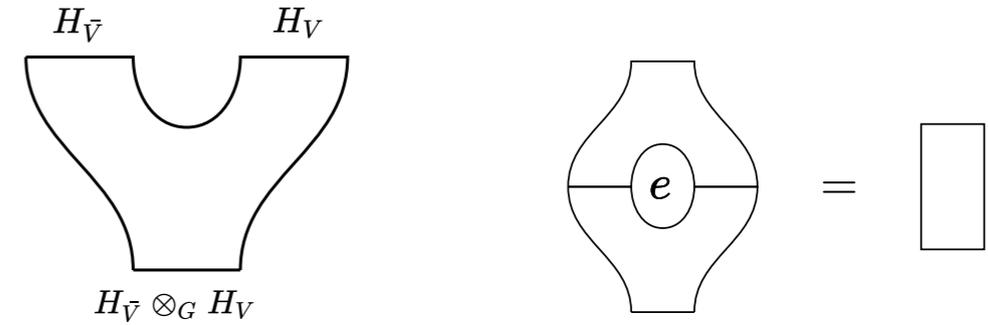
$$S = \sum_R -P(R) \log P(R) + 4 \sum_R P(R) \log \dim R$$

Number of entangling surfaces



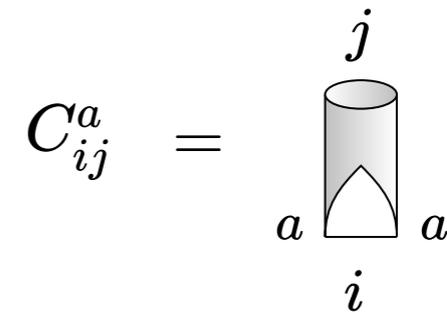
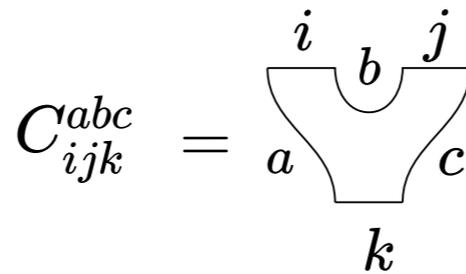
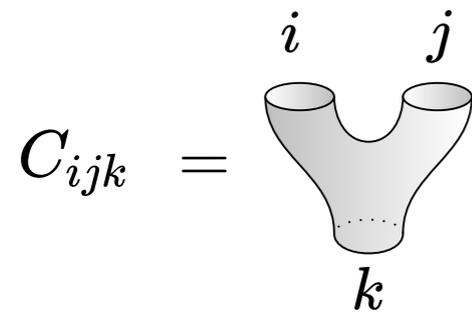
Summary

- Entanglement probes the structure of **extended QFT**
e.g. extension defines an open string algebra
- The extension satisfies the **E-brane axiom**



In Progress: Entanglement and Extended CFT

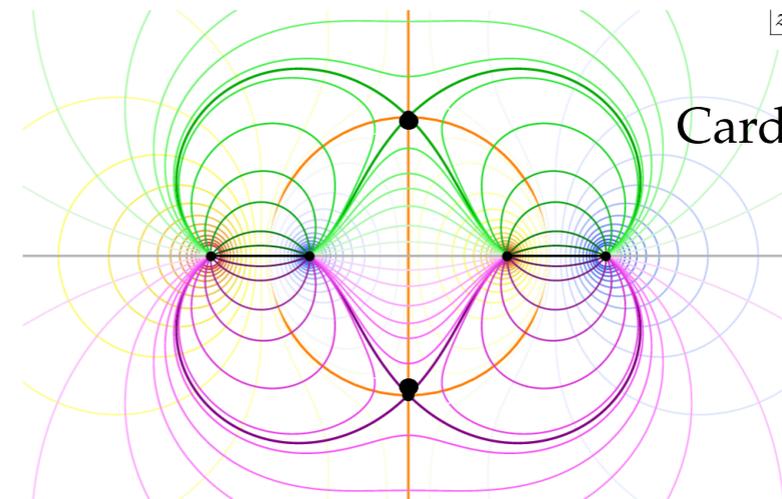
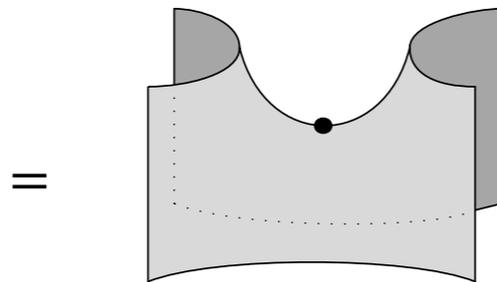
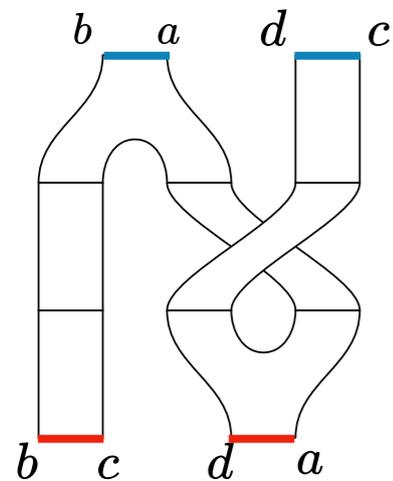
OPE's of a BCFT



E brane boundary condition ~ Conformally Inv. BC

Fusion Rule ~ Entangling product ?

A hint from free fermions



Cardy-Tonni 2016

Extra slides