Fundamental exchange rate between coherence and asymmetry

H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

Hiroyasu Tajima @Kyoto university (YITP) H. Tajima, N. Shiraishi and K. Saito arXiv:1906.04076 (2019)

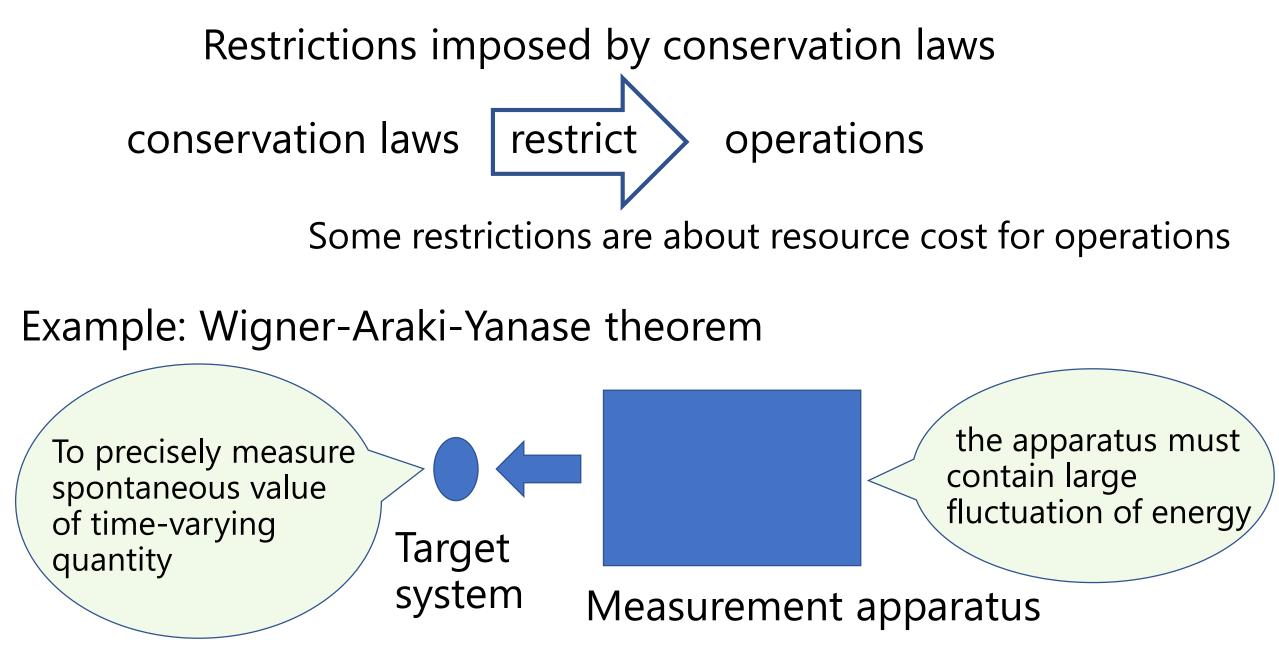
Collaborators:



N. Shiraishi @Gakushuin university



K. Saito @Keio university Topic: Resource cost for quantum operations under conservation laws



To perform precise measurement, we need large fluctuation of energy as a resource.

Restriction on unitary dynamics?

Is there any restriction similar to Wigner-Araki-Yanase theorem on implementing unitary dynamics under conservation laws?

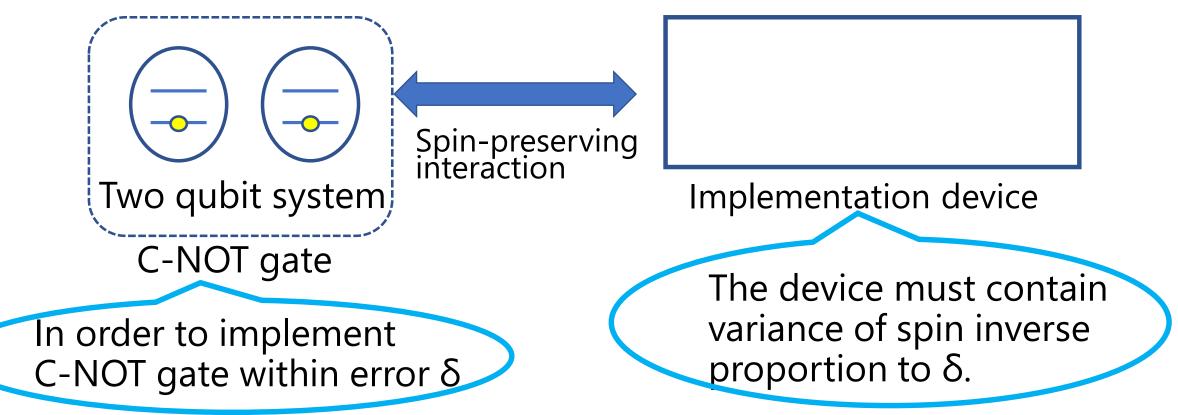
> initially proposed by Masanao Ozawa, about two decades ago: M. Ozawa, Phys. Rev. Lett. **89**, 057902 (2002).

The motivation is to clarify the restrictions on quantum computing imposed by conservation laws.

Restriction on C-NOT gate: Ozawa's result

Ozawa considers the implementation of Controlled-NOT gate under spinpreserving interaction. M. Ozawa, Phys. Rev. Lett. **89**, 057902 (2002).

Ozawa obtain a trade-off inequality between error and fluctuation for Controlled-NOT gate. (With using Wigner-Araki-Yanase theorem!)



Restriction on general unitary gate: A long standing open problem

After Ozawa's result, similar trade-off relations were given for various (but specific) unitary gates:

Not gate and Fredkin gate: Hadamard gate:

T. Karasawa and M. Ozawa, M. Ozawa, Int. J. Quant. Phys. Rev. A **75**, 032324 (2007). Inf. **1**, 569 (2003).

Question :

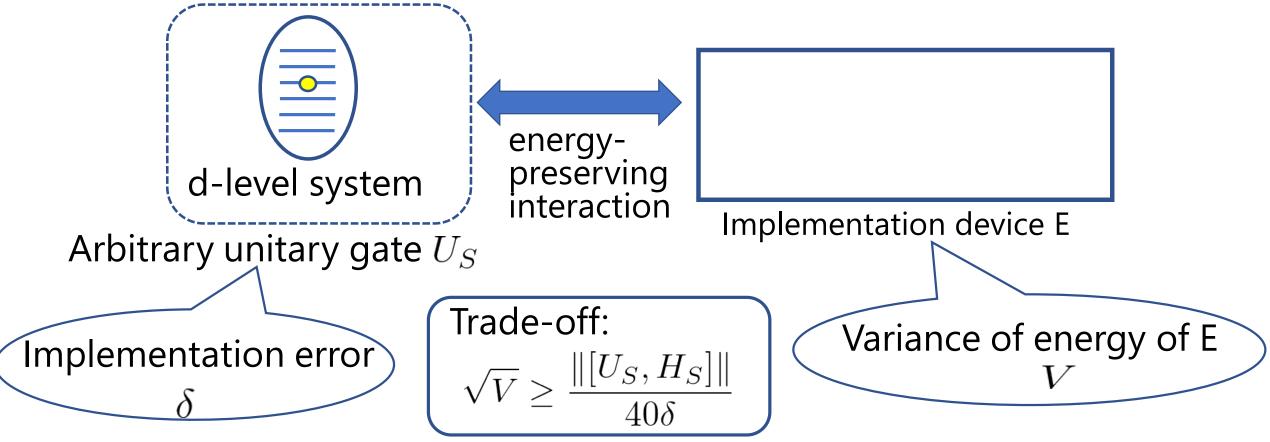
Is there any universal trade-off between fluctuation and error for general unitary, other than qubit gates ?

Although the above strong circumstantial evidence, the trade-off was never given.

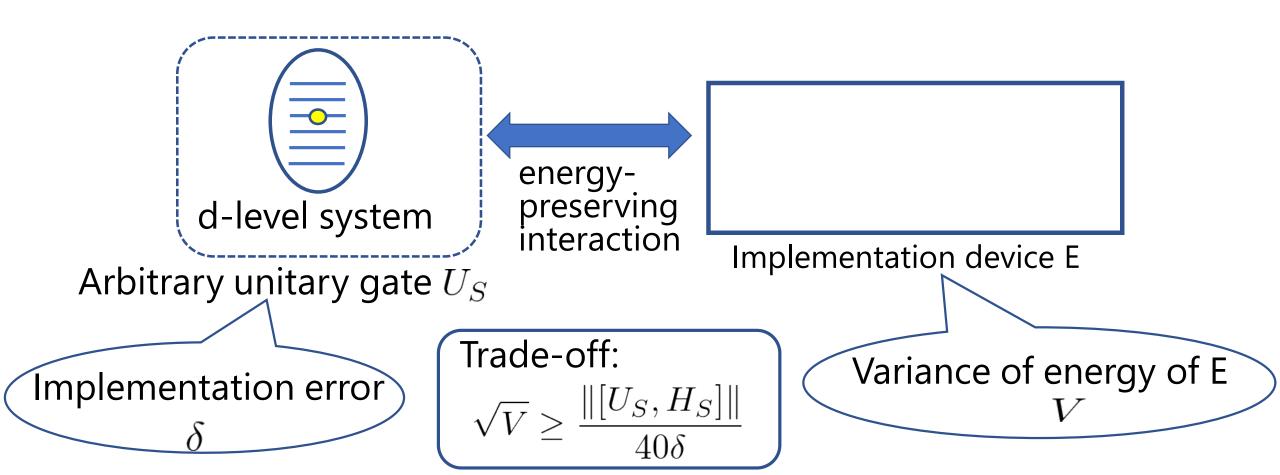
H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

We consider the implementation of arbitrary unitary gate under conservation law of energy.

We derive a universal trade-off inequality between fluctuation of energy and implementation error of unitary operations. (without using Wigner-Araki-Yanase theorem)

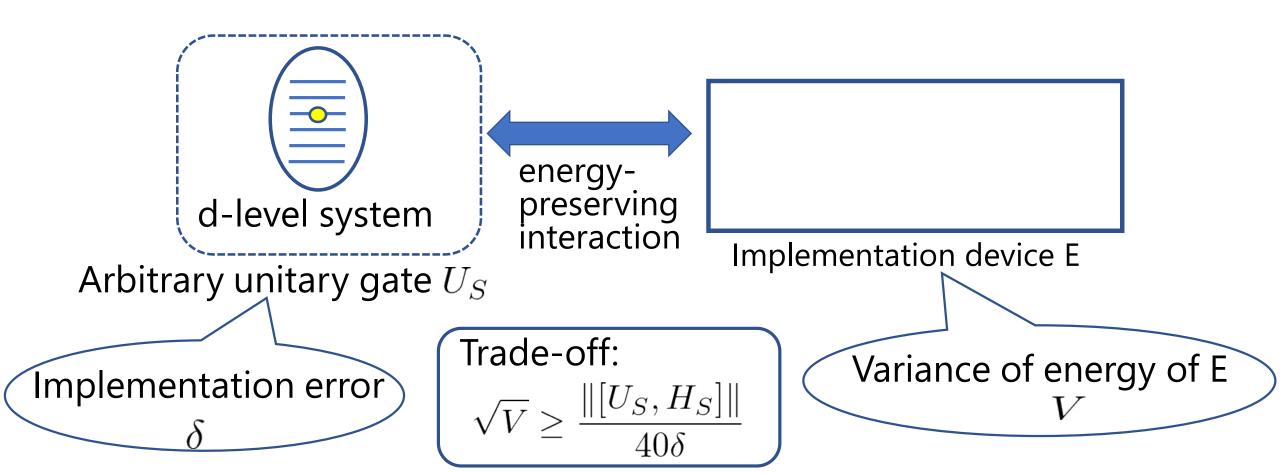


H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. 121, 110403 (2018)



H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

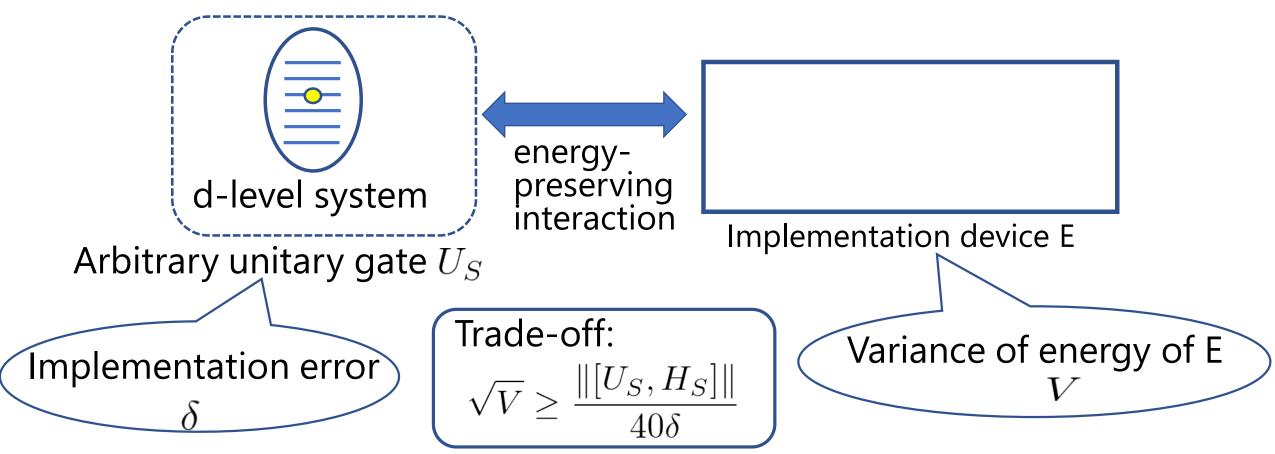
We also show that the required fluctuation must have quantum origin.



H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

We also show that the required fluctuation must have quantum origin.

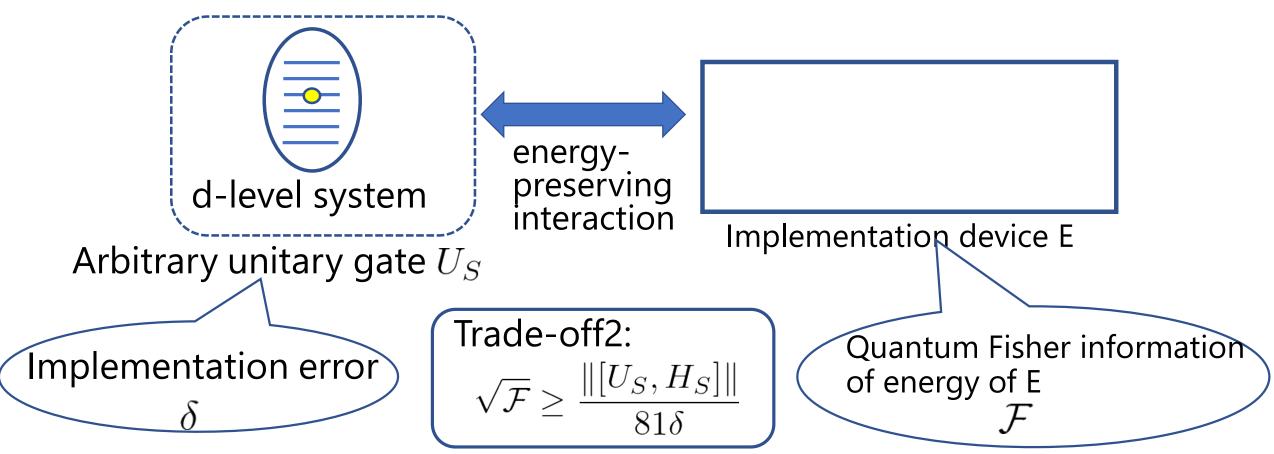
We derive another trade-off between implementation error and quantum Fisher information, which is a well-known measure of quantum fluctuation.



H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

We also show that the required fluctuation must have quantum origin.

We derive another trade-off between implementation error and quantum Fisher information, which is a well-known measure of quantum fluctuation.



Quantum Fisher information: A measure of coherence

$$\begin{aligned} \mathcal{F}_A(\rho) &:= 2 \sum_{a,b} \frac{(p_a - p_b)^2}{p_a + p_b} |A_{ab}|^2 \\ A_{ab} &:= \langle \psi_a | A | \psi_b \rangle \\ \{p_a, \psi_a\} \text{ is eivenvalues and eivenvectors of } \rho \end{aligned}$$

Important feature :
$$\mathcal{F}_A(\rho) = 4 \min_{\{q_j, \phi_j\}: \rho = \sum_j q_j \phi_j} \sum_j q_j V_A(\phi_j).$$

 ρ is pure $\mathcal{F}_A(\rho) = 4V_A(\rho)$

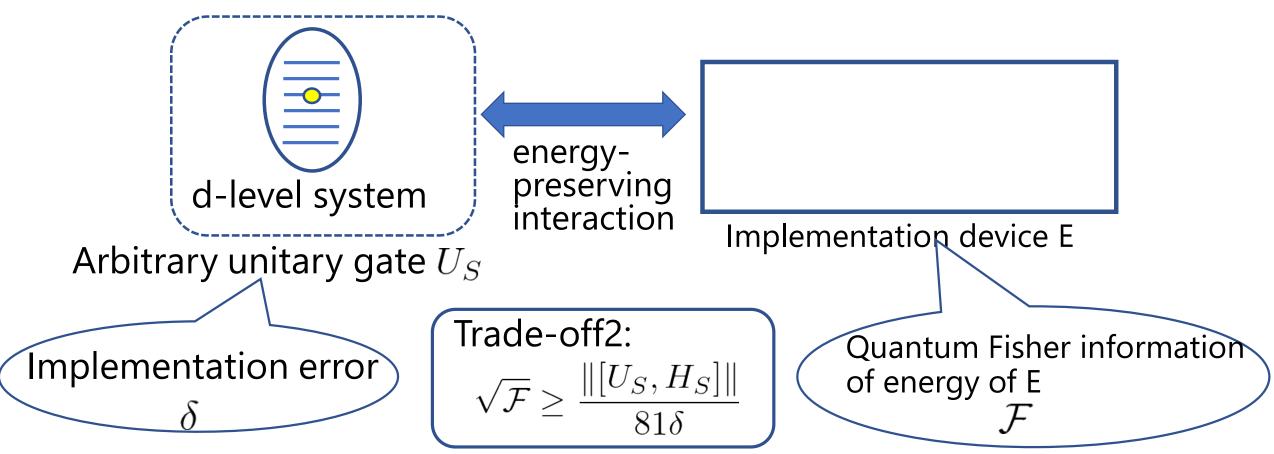
$$[\rho, A] = 0 \quad \qquad \mathcal{F}_A(\rho) = 0$$

Namely, QFI is "quantum part" of fluctuation of the physical quantity A.

H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

We also show that the required fluctuation must have quantum origin.

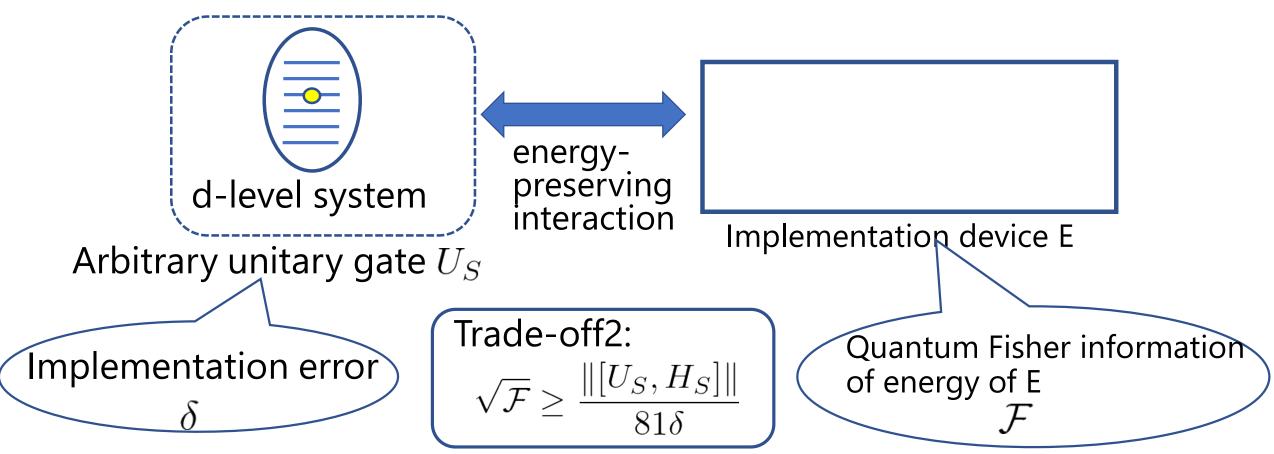
We derive another trade-off between implementation error and quantum Fisher information, which is a well-known measure of quantum fluctuation.



H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

We also show that the required fluctuation must have quantum origin.

We derive another trade-off between implementation error and quantum Fisher information, which is a well-known measure of quantum fluctuation.



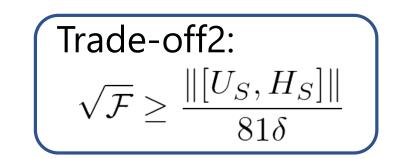
The remaining question: a generalization of Ozawa's question H. Tajima, N. Shiraishi and K. Saito Phys. Rev. Lett. **121**, 110403 (2018)

Quantum Fisher information \mathcal{F} is a measure of coherence.

So, Trade-off 2 is a lower bound for coherence necessary to implement unitary dynamics under conservation law.

Question': How much coherence is "necessary and sufficient" to implement unitary dynamics under conservation law?

This is a generalization of Ozawa's question.

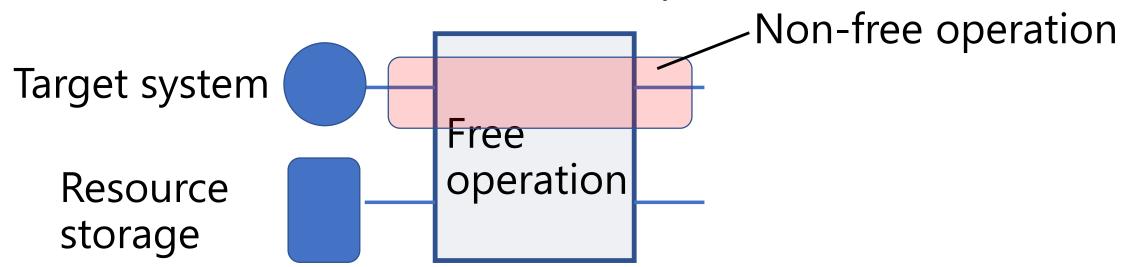


Approach from quantum information -resource theory of quantum channels

Free operations and free states: we can use freely

Resource states:

the states we cannot create from free operations and free states



Key question of resource theory of quantum channels: How much resource do we need to implement the desired operations? Approach from quantum information -resource theory of quantum channels

Partially solved in various cases

 \rightarrow

Quantum thermodynamics:

P. Faist and R. Renner. Phys. Rev. X, 8 021011, (2018).
P. Faist, M. Berta and F. Brandao, Phys. Rev. Lett. 122, 200601 (2019).

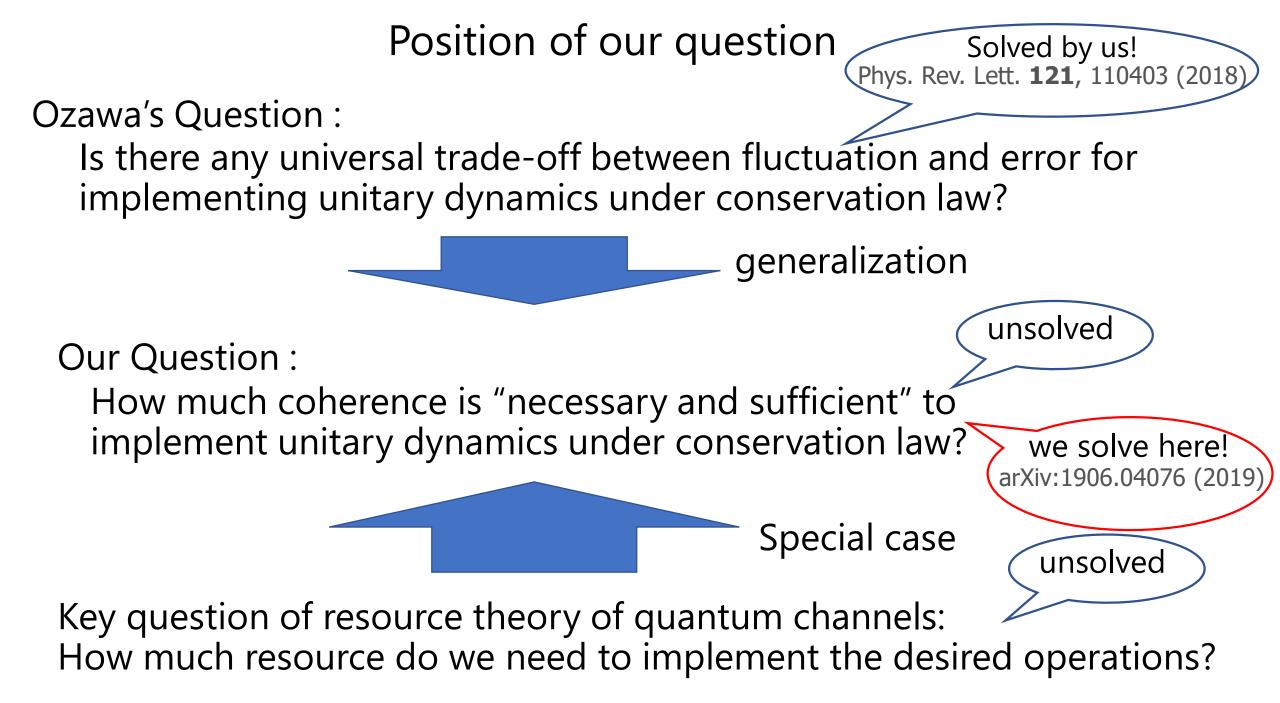
Resource erasure:

Z.-W. Liu and A. Winter, arXiv:1904.04201 (2019).

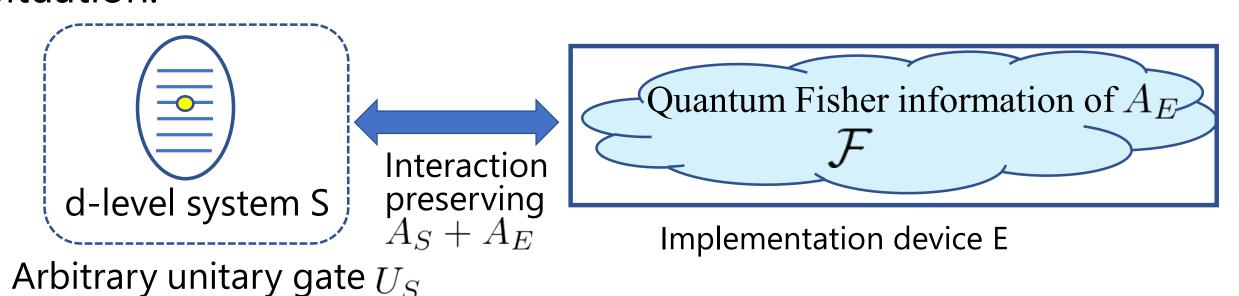
Incoherent operations:

M. G. Diaz, K. Fang, X. Wang, M. Rosati, M. Skotiniotis, J. Calsamiglia and A. Winter, Quantum **2**, 100 (2018).

Upper and lower bounds for "necessary and sufficient" resource to implement the desired operations Key question of resource theory of quantum channels: How much resource do we need to implement the desired operations?



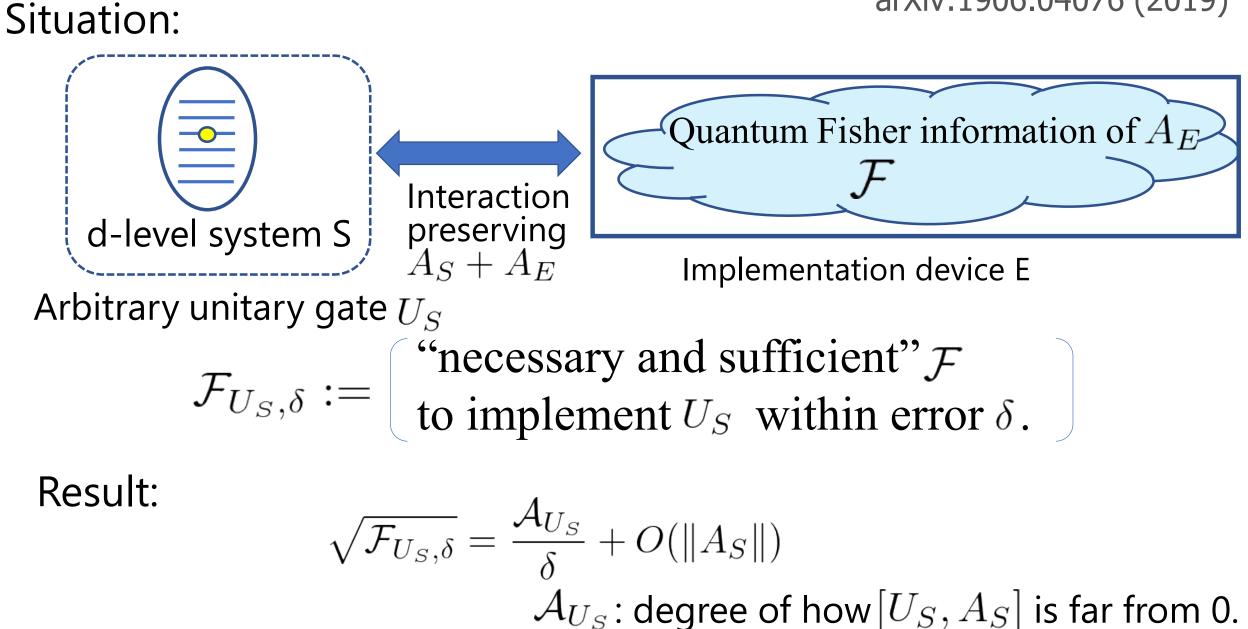
Situation that we treat (detail is shown later) arXiv:1906.04076 (2019)



We derive "necessary and sufficient" \mathcal{F} to implement U_S within error δ .

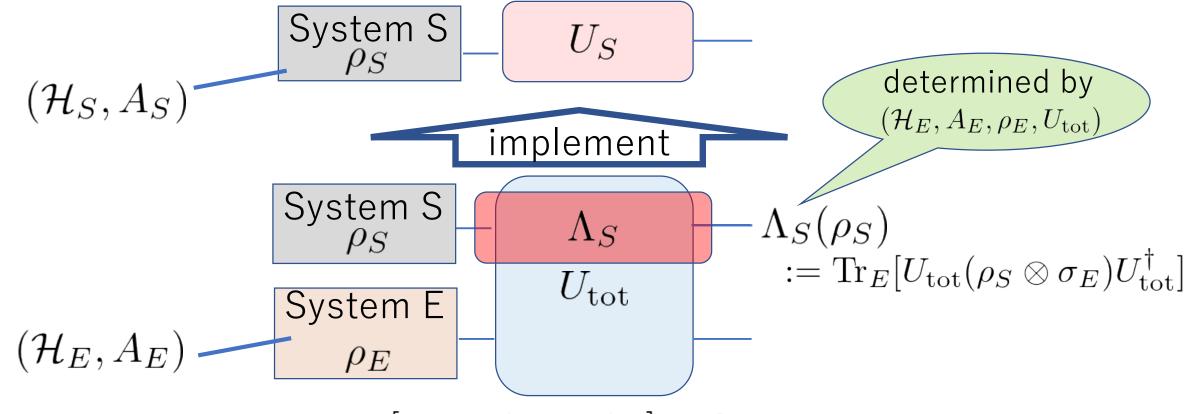
Our result

arXiv:1906.04076 (2019)



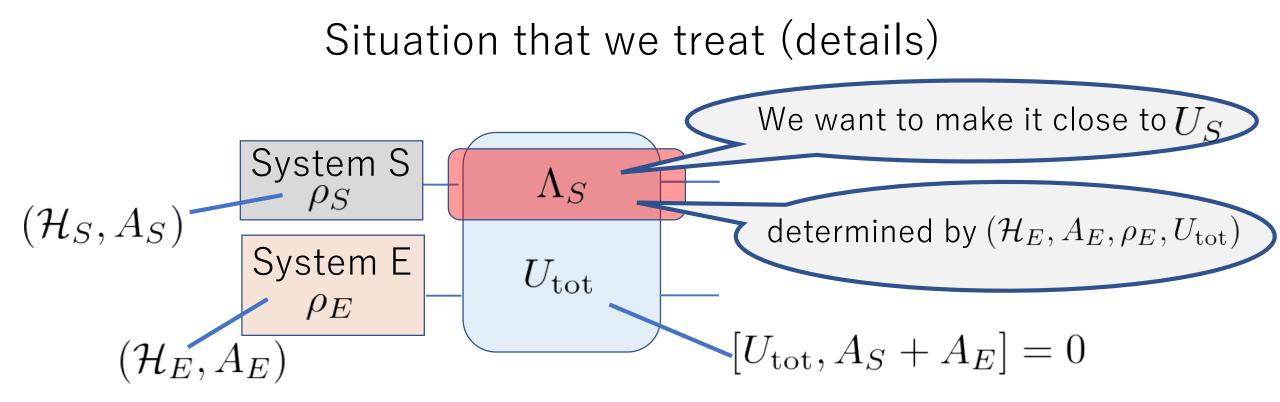
Situation that we treat (details)

We approximately implement U_S on the target system S by the interaction with an external system E.



Under the restriction $[U_{tot}, A_S + A_E] = 0$,

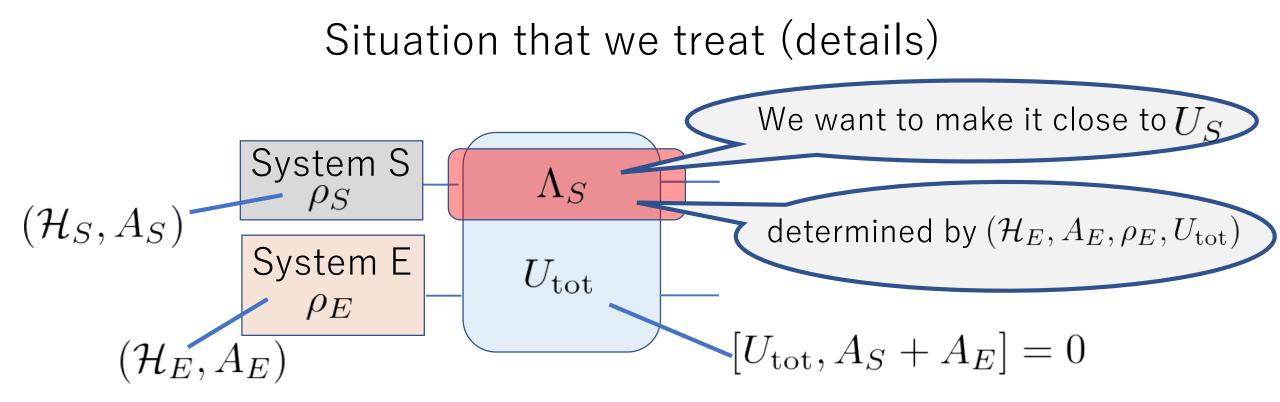
we take $(\mathcal{H}_E, A_E, \rho_E, U_{\mathrm{tot}})$ freely, and try to make Λ_S close to U_S .



Under this setup, we define the following three quantities: degree of how $[U_S, A_S]$ is far from 0 $\sqrt{\mathcal{F}_{U_S,\delta}} = \frac{\mathcal{A}_{U_S}}{\delta} + O(||A_S||)$ "necessary and sufficient" amount

of Coherence to implement U_S

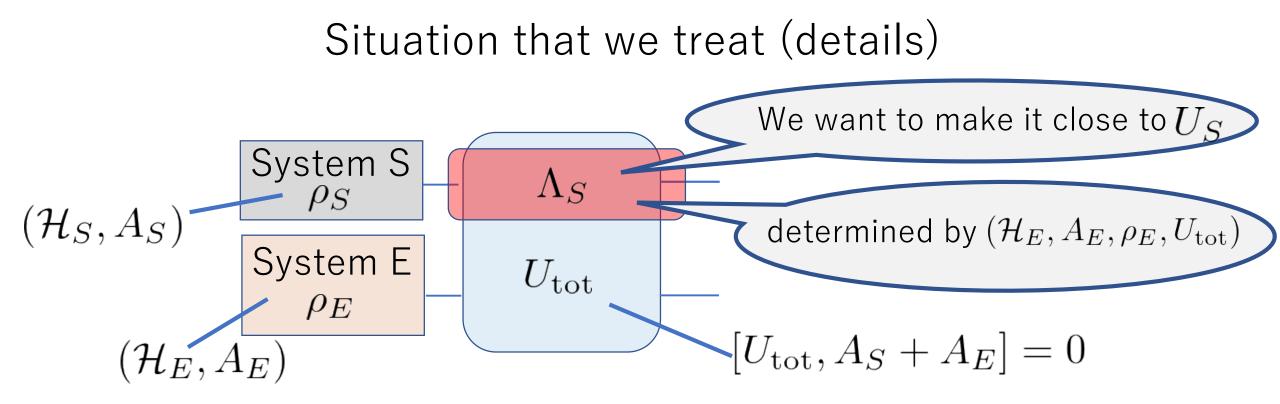
" implementation error



Under this setup, we define the following three quantities: degree of how $[U_S, A_S]$ is far from 0 $\sqrt{\mathcal{F}_{U_S,\delta}} = \frac{\mathcal{A}_{U_S}}{\delta} + O(||A_S||)$ "necessary and sufficient" amount

of Coherence to implement U_S

" implementation error

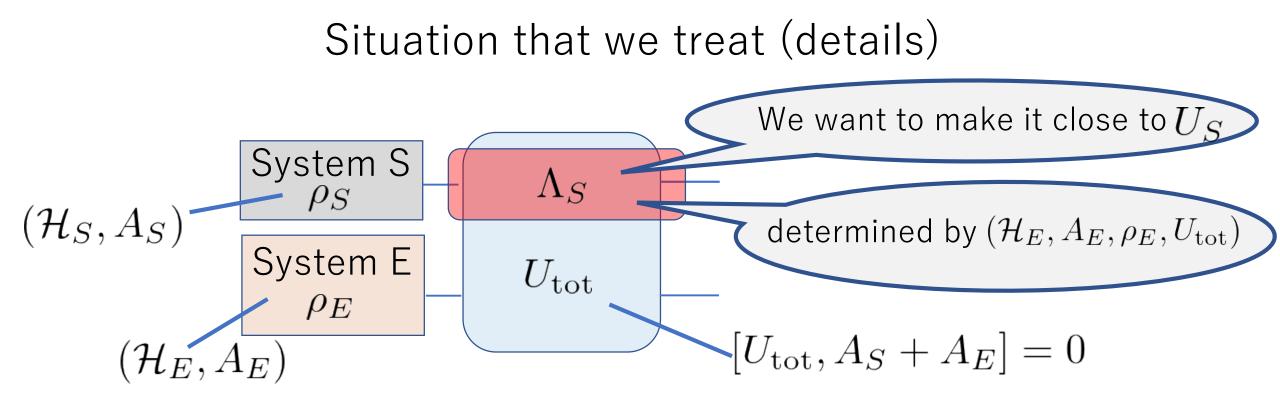


We define δ as maximal entanglement Bures distance between $U_S \rho_S U_S^{\dagger}$ and $\Lambda_S (\rho_S)$:

$$\mathcal{I} = (\mathcal{H}_E, A_E, \rho_E, U_{\text{tot}}) \text{ implements } U_S \text{ within error } \delta$$

$$\stackrel{\longleftarrow}{\underset{\text{def}}{\longleftarrow}} \delta \ge \max_{\rho_S} L_e(\rho_S, \Lambda_{U_S^{\dagger}} \circ \Lambda_S)$$

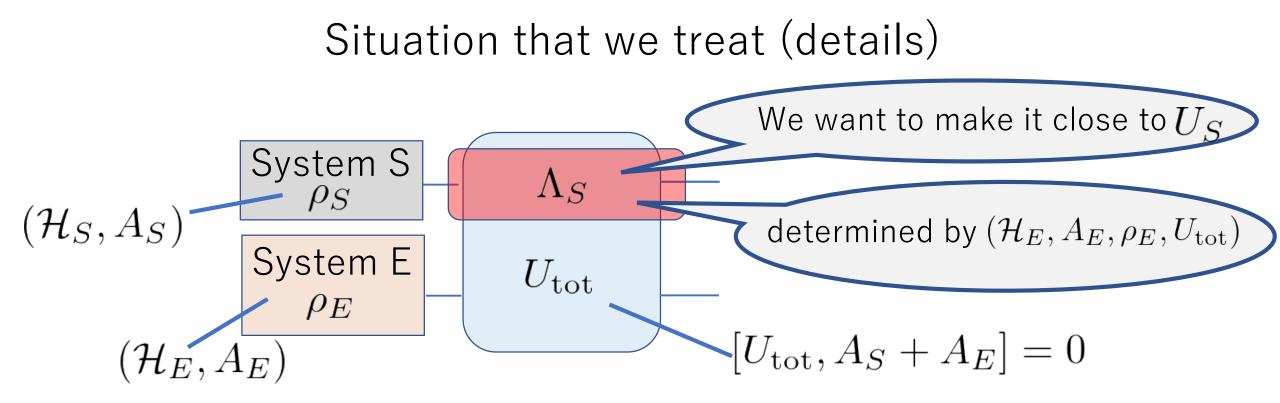
$$\Lambda_{U_S^{\dagger}}(\rho) := U_S^{\dagger} \rho_S U_S \quad L_e(\rho_S, \Lambda_S) := \sqrt{2(1 - F_e(\rho, \Lambda_S))}$$



We define $\mathcal{F}_{U_s,\delta}$ as the minimal sufficient amount of QFI to implement U_S within error δ .

$$\mathcal{F}_{U_S,\delta} := \min_{\mathcal{I} \models_{\delta} U_S} \mathcal{F}_{A_E}(\rho_E)$$

$$\begin{array}{c} \mathcal{I} \models_{\delta} U_S \ \text{means} \\ ``\mathcal{I} \ \text{implements} \ U_S \ \text{within} \ \text{error} \ \delta \ `` \end{array}$$



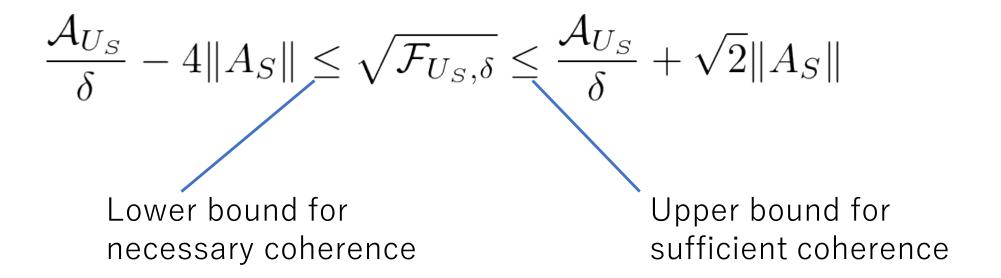
We define \mathcal{A}_{U_S} as degree of how U_S changes the conserved quantity A_S

Maximum and minimum eigenvalues

$$\mathcal{A}_{U_S} := \frac{\lambda_{\max}(U_S^{\dagger}A_S U_S - A_S) - \lambda_{\min}(U_S^{\dagger}A_S U_S - A_S)}{2},$$

Results

The following two bounds hold :



Results

Combining two bounds, we obtain an asymptotic equality:

$$\sqrt{\mathcal{F}_{U_S,\delta}} = \frac{\mathcal{A}_{U_S}}{\delta} + O(\|A_S\|) \qquad \delta \to 0$$

Simple equality between degree of asymmetry (degree of violation of conservation law) and amount of coherence!

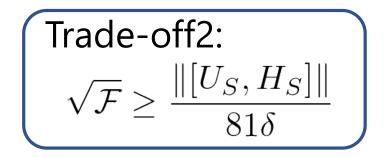
Summary

Phys. Rev. Lett. **121**, 110403 (2018) arXiv:1906.04076 (2019)

We derived two inequalities and one equality.

Trade-off:
$$\sqrt{V} \ge \frac{\|[U_S, H_S]\|}{40\delta}$$

Fundamental trade-off between error and fluctuation Answer to Ozawa's question



A lower bound for necessary coherence for implementing unitary dynamics under energy-conservation law

Trade-off3:
$$\sqrt{\mathcal{F}_{U_S,\delta}} = \frac{\mathcal{A}_{U_S}}{\delta} + O(||A_S||)$$

Asymptotic equality for "necessary and sufficient" coherence for implementing unitary dynamics under conservation laws



Answer to the key question

of resource theory of quantum channels in a special case.