

Noise and decoherence induced by gravitons

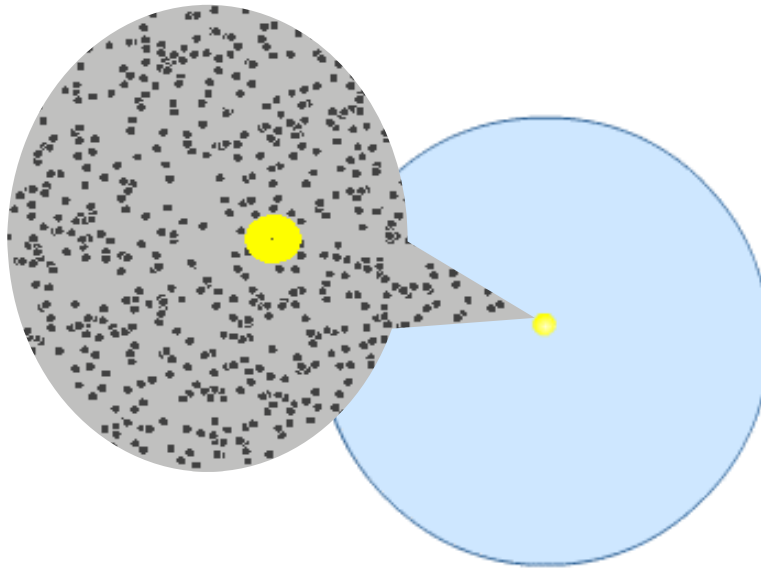
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Based on

S. Kanno, J. Soda (Kobe), J. Tokuda (Kobe),

Rhys. Rev. D 103 (2021), 044017 (selected as a PRD Editors' suggestion)

Brownian motion and gravitons



A phenomenon of random motion
of minute particles in water
Brown (1827)

The random motion occurs due to the random collisions of surrounding individual water molecules
Einstein (1905)

The size of the water molecules can be determined by the Brownian motion

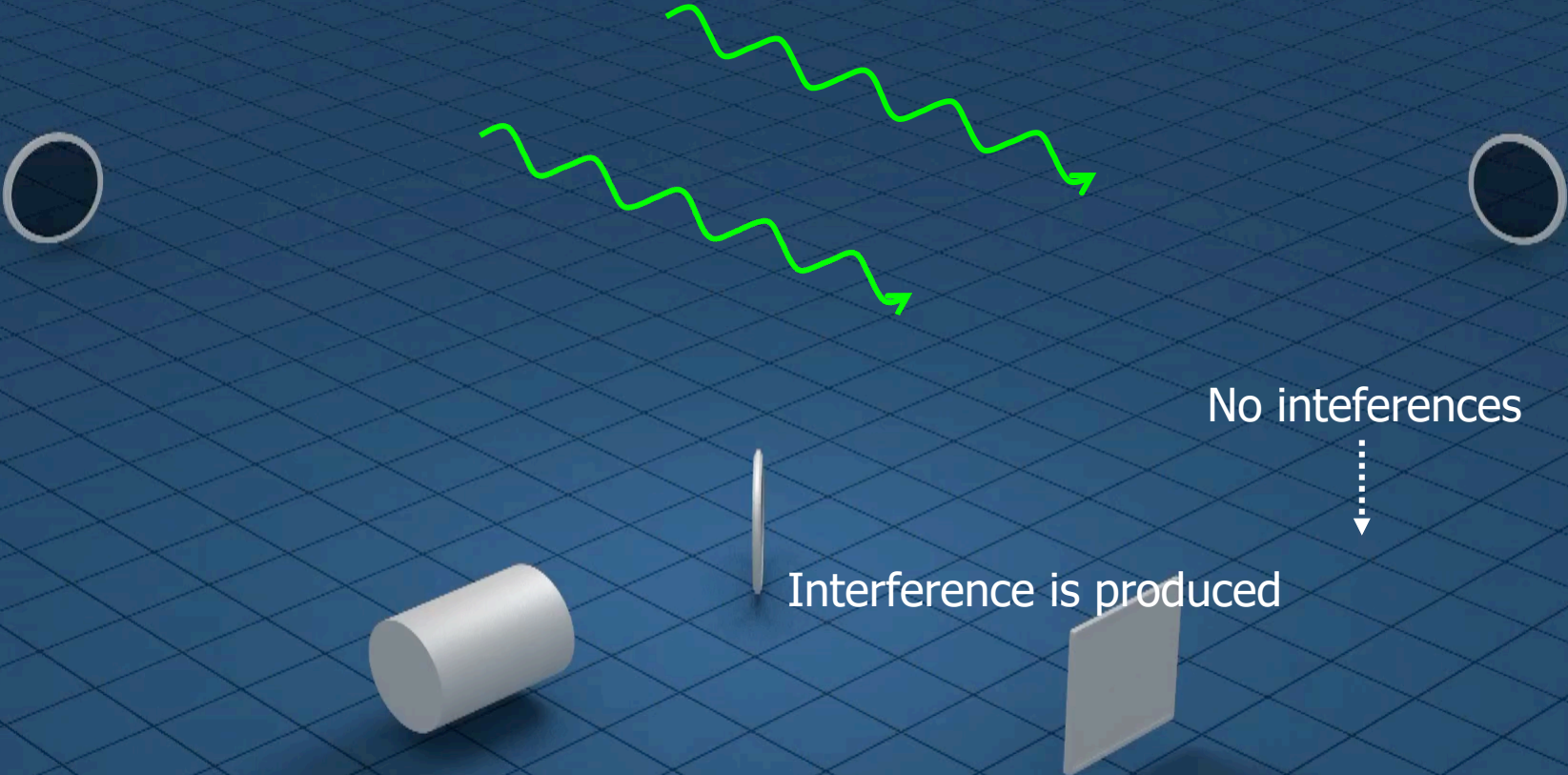
The size was determined by the experiment of the Brownian motion Perrin (1908)

Hard to observe water molecules directly → Proved the existence of water molecules indirectly through the minute particles

Hard to observe gravitons directly → Prove the existence of gravitons indirectly through something

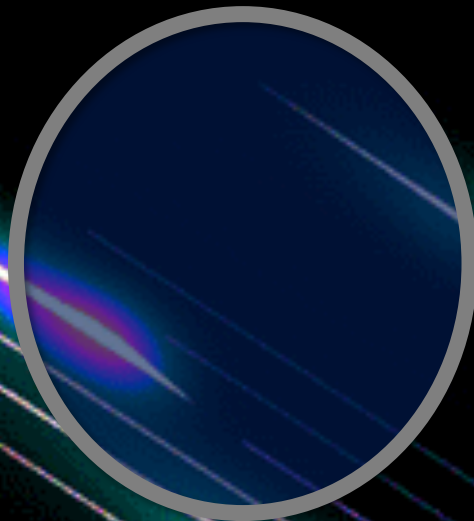
What should we target as the "something"

Laser interferometers



A suspended mirror

Gravitons : \leftarrow Environment molecules



System particle

The mirror

The same setup as the Brownian motion

The system of the mirror couples to an huge environment of gravitons



Noise due to gravitons should be induced on the mirror

Can we prove the existence of gravitons indirectly through this noise ?

What should we calculate?



Derive the Langevin equation
(stochastic differential equation)
of motion of this system

We first derive the total action of the mirror and the gravitational waves

How the mirror feels the effect of gravitational waves?

For simplicity, regard the mirror as a point particle

A single particle (mirror)? *Einstein's equivalence principle*

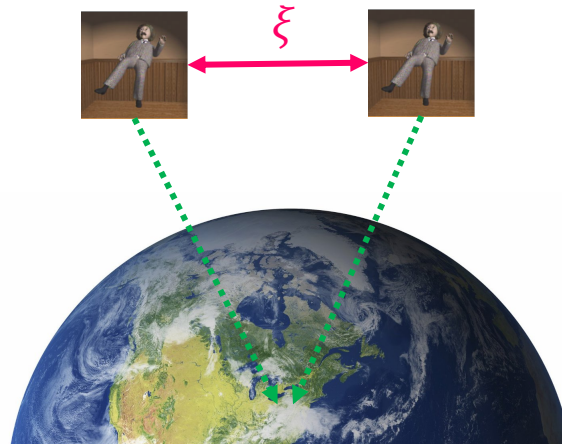
A free falling object feels no gravity

.....
local inertial frame



How to measure the gravitational force? *Geodesic deviation ξ*

Observe the geodesics of two free falling particles deviating relative to each other



The total action for the GWs and the mirrors

Total Action

$$S = \frac{M_{\text{pl}}^2}{2} \int d^4x \sqrt{-g} R - m_0 \int d\tau - m \int d\tau$$

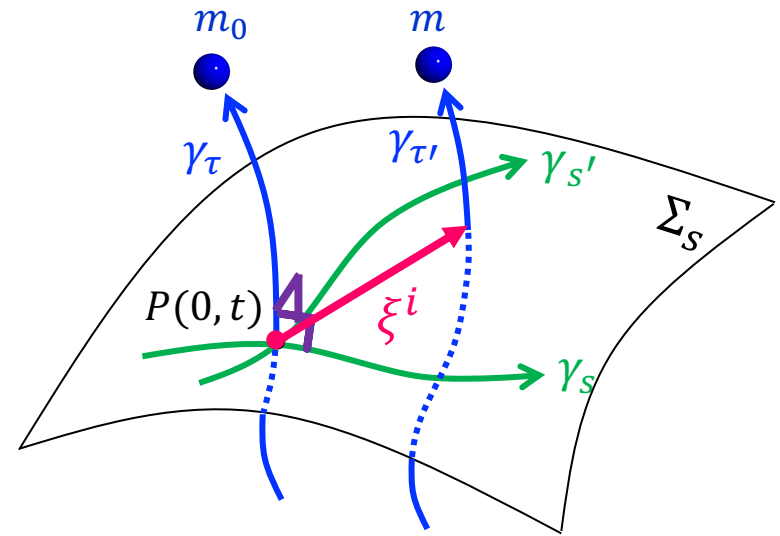
No dynamical variables

$\left[\begin{array}{l} h_k^A : \text{Fourier modes of metric field (GWs)} \\ \xi^i : \text{Geodesic deviation between the mirrors} \end{array} \right.$

$$S = \int dt \sum_{\mathbf{k}} \sum_{A=+, \times} \left[\frac{1}{2} \dot{h}_{\mathbf{k}}^A \dot{h}_{\mathbf{k}}^A - \frac{1}{2} k^2 h_{\mathbf{k}}^A h_{\mathbf{k}}^A \right]$$

$$+ \int dt \left[\frac{m}{2} (\dot{\xi}^i)^2 + \frac{1}{2} \frac{m}{M_{\text{pl}}} \frac{1}{\sqrt{V}} \sum_{\mathbf{k}, A} \left[e_{ij}^A(\mathbf{k}) \ddot{h}_{\mathbf{k}}^A \xi^i \xi^j \right] \right]$$

cubic derivative interaction



Fermi normal coordinates

Quantize the GWs and the deviation between mirrors

For simplicity, work in the interaction picture below

We promote the free metric field $h_k^A(t)$ to the operator $\hat{h}_k^A(t)$

$$\hat{h}_k^A(t) = \hat{a}_k^A \overset{\substack{\text{Positive freq. mode in the Minkowski space} \\ \downarrow}}{v_k(t)} + \hat{a}_{-k}^{A\dagger} v_k^*(t) \quad \hat{a}_k^A |0\rangle = 0 \quad [\hat{a}_k^A, \hat{a}_p^{B\dagger}] = \delta^{AB} \delta_{kp}$$

(Note: A dotted arrow points from the text "The Minkowski vacuum" to the state $|0\rangle$ in the second equation.)

Since $\xi^i(t)$ is just a position, we promote $\hat{\xi}^i(t)$ to the operator as well

Solve the EOMs of the system of \hat{h}_k^A and $\hat{\xi}^i$ to derive the Langevin equation of geodesic deviation $\hat{\xi}^i$

Langevin equation of geodesic deviation

Kanno, Soda & Tokuda (2020)
Parikh, Wilczek & Zahariade (2020)

We obtained

$$\ddot{\xi}^i + \frac{m}{40\pi M_{\text{pl}}^2} \left(\delta_{ik}\delta_{j\ell} + \delta_{i\ell}\delta_{jk} - \frac{2}{3}\delta_{ij}\delta_{k\ell} \right) \xi^i \frac{d^5}{dt^5} (\xi^k \xi^\ell)$$

frictional force (radiation reaction force)

$$= -\hat{N}_{ij}(0, t) \hat{\xi}^j(t)$$

random force (noise)

$$\langle \hat{h}_{\mathbf{k}}^A(0) \rangle = 0$$

classical free GWs

where

$$\hat{N}_{ij}(t) = \frac{1}{M_{\text{pl}}\sqrt{V}} \sum_{\mathbf{k}, A} k^2 e_{ij}^A(\mathbf{k}) \left\{ \hat{h}_{\mathbf{k}}^A(0) \cos kt + \dot{\hat{h}}_{\mathbf{k}}^A(0) \frac{\sin kt}{k} \right\}$$

The noise consists of quantized fluctuations by gravitational fields (gravitons)

Noise due to gravitons

Detectability is discussed by effective strain in the frequency domain

$$h_{\text{eff}} = \frac{N(f)}{(2\pi f)^2} \equiv \frac{1}{(2\pi f)^2} \left(\int_{-\infty}^{\infty} dt \langle \{ \hat{N}^{ij}(t), \hat{N}_{ij}(0) \} \rangle e^{2\pi i f t} \right)^{\frac{1}{2}}$$

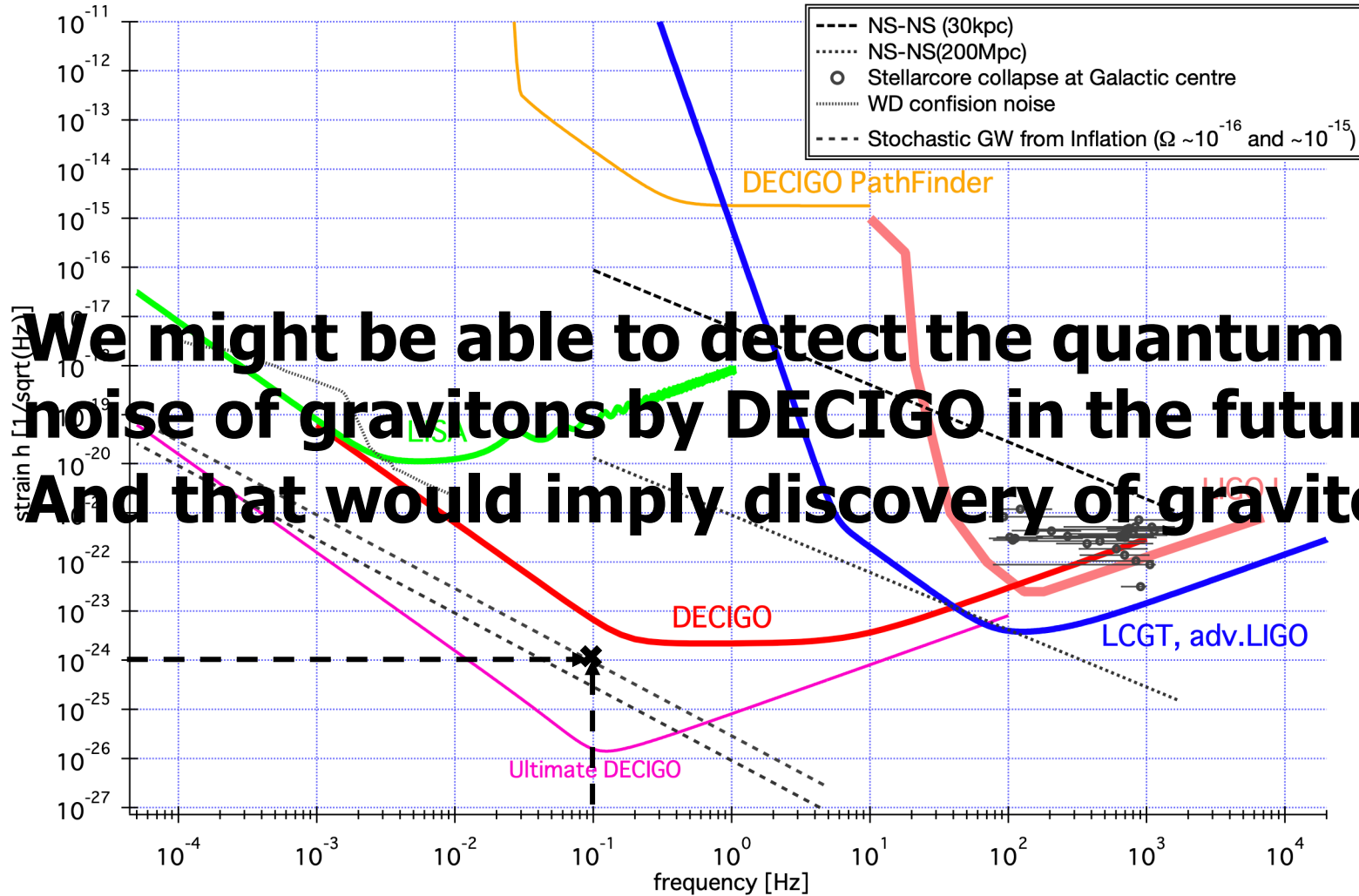
For primordial GWs

$$h_{\text{eff}} \approx 2 \times 10^{-42} \left(\frac{f}{1 \text{ Hz}} \right) \left(\frac{f_c}{f} \right)^2 \text{ Hz}^{-\frac{1}{2}}$$


, f_c : cutoff frequency $\sim 2 \times 10^8$ Hz

The latest sensitivity curves

Kanno, Soda & Tokuda (2020)



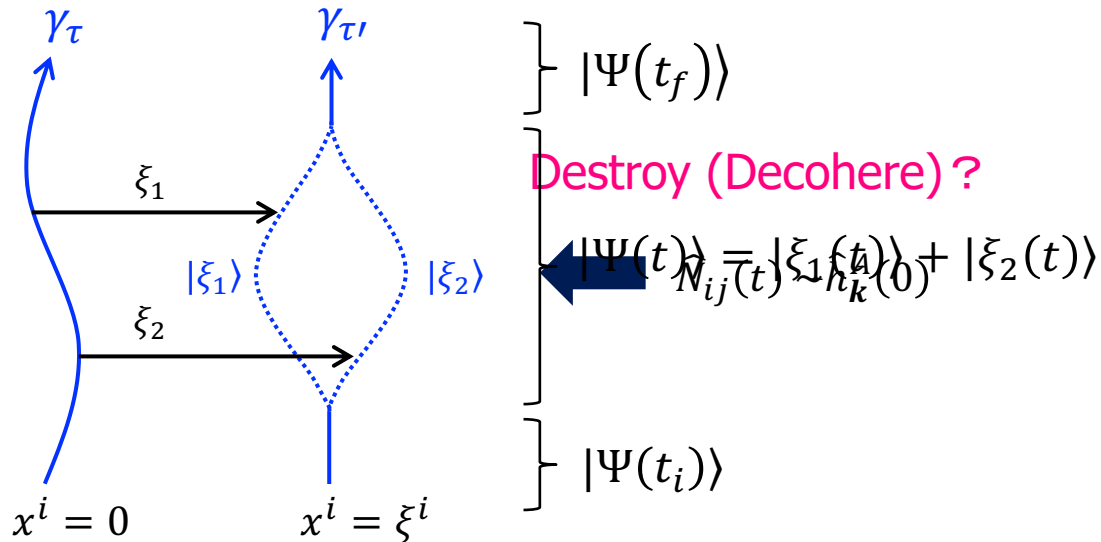
The strain: $h_{\text{eff}} \sim 10^{-24} \text{ Hz}^{-\frac{1}{2}}$ at $f \sim 0.1 \text{ Hz}$



Can we detect the noise of gravitons
in a laboratory?

Effect of noise on a superposition state in a lab

It's possible to create a superposition state of two spatially separated locations



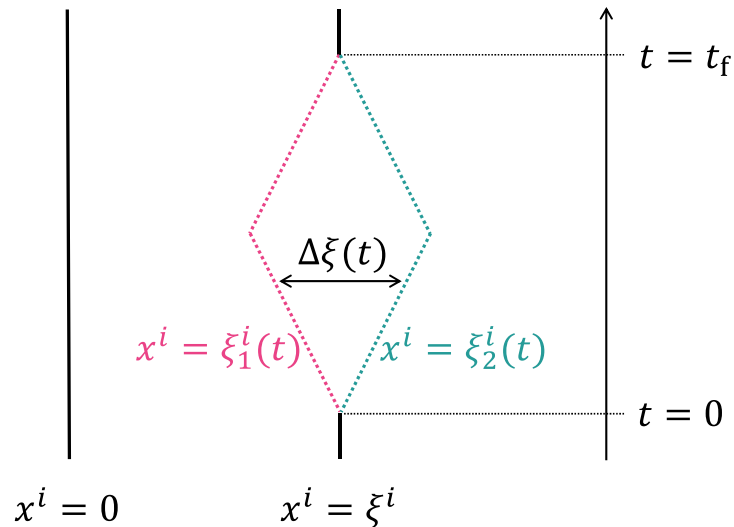
The decoherence factor

$$\Gamma \approx \frac{m^2}{8} \int_0^{\tau_{\text{dec}}} dt \Delta(\xi^i \xi^j)(t) \int_0^{\tau_{\text{dec}}} dt' \Delta(\xi^k \xi^\ell)(t') \langle \{ \hat{N}_{ij}(t), \hat{N}_{k\ell}(t') \} \rangle \geq 1$$

The loss of quantum coherence between $|\xi_1(t)\rangle$ and $|\xi_2(t)\rangle$ occurs

By using this decoherence factor, we can predict the decoherence time due to gravitons

A configuration of the superposition state



$$\Delta\xi(t) = \begin{cases} 2vt & \text{for } 0 < t \leq t_f/2 \\ 2v(t_f - t) & \text{for } t_f/2 < t \leq t_f \end{cases}$$

The decoherence factor $\Gamma(t_f) \approx \frac{2m^2 v^2}{5\pi^2 M_{\text{pl}}^2} G(\Omega_m t_f) \quad G(\Omega_m t_f) \sim \mathcal{O}(1)$

$mv \ll M_{\text{pl}} \approx 4 \times 10^{-6} \text{ g} \quad \Rightarrow \quad \Gamma(t_f) \ll 1 : \text{Decoherence does not occur}$

$mv \gg M_{\text{pl}} \approx 4 \times 10^{-6} \text{ g} \quad \Rightarrow \quad \Omega_m \tau_{\text{dec}} \cong 10 \times \left(\frac{M_{\text{pl}}}{mv}\right)^{\frac{1}{2}} \text{ s}$

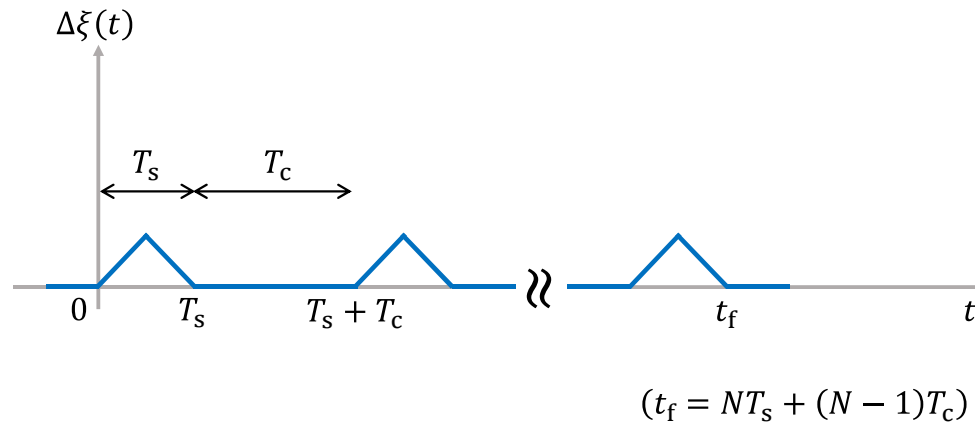
$(c = 1)$
Need to exceed the speed of light...

Currently, the spatial superposition state is only created up to $m = 10^{-20} \text{ g}$ in a lab

Repeated configuration of the superposition state

Kanno, Soda & Tokuda (2020)

Previous configuraton repeats n times

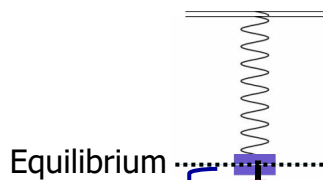


The decoherence factor $\Gamma(t_f) \approx n \times \frac{2m^2 v^2}{5\pi^2 M_{\text{pl}}^2} G(\Omega_m t_f) \quad G(\Omega_m t_f) \sim \mathcal{O}(1)$

If we take sufficiently large n , decoherence occurs even for $mv \ll M_{\text{pl}} \approx 4 \times 10^{-6} \text{ g}$

If the decoherence is observed and the decoherence time in a lab agrees with the theoretical prediction, it turns out to be a discovery of gravitons

Effect of noise on an entangled state in a lab

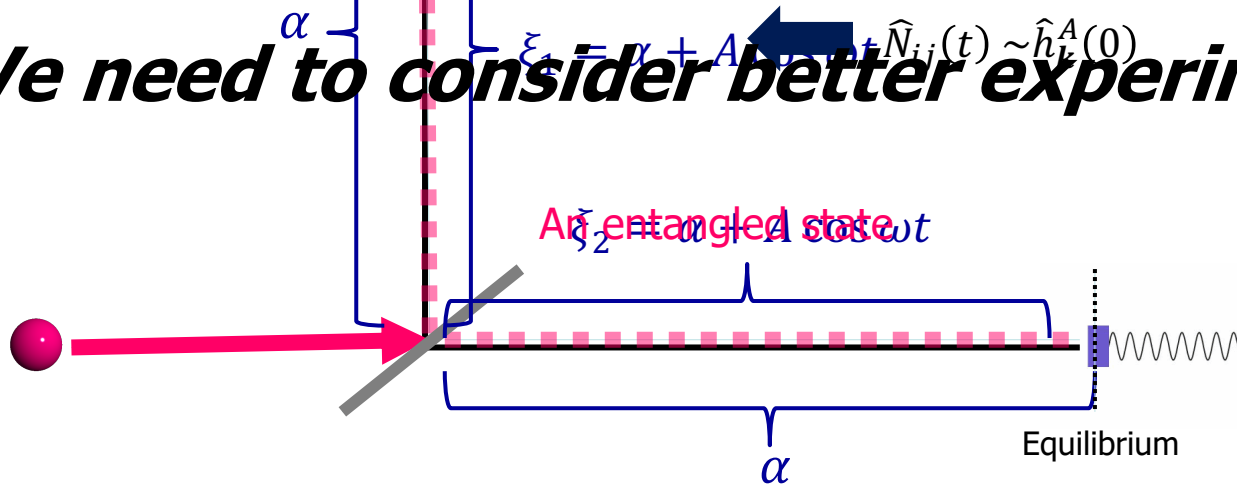


$|1\rangle$: Oscillation $|0\rangle$: Equilibrium

The system is entangled

$$|\psi\rangle = \frac{1}{\sqrt{2}} |1\rangle_{\xi_1} |0\rangle_{\xi_2} + \frac{1}{\sqrt{2}} |0\rangle_{\xi_1} |1\rangle_{\xi_2}$$

We need to consider better experimental designs



An entangled state

The decoherence factor

$$\Gamma(t_f) \approx \left(\frac{m}{M_{\text{pl}}}\right)^2 \left(\frac{A}{\alpha}\right)^2 N \left(\frac{f_c}{f}\right)^4 \sim 1$$

Decoherence time ≈ 1 y **Preliminary**

$m \sim 40$ kg
 $A \sim 10^{-17}$ cm
 $\alpha \sim 40$ km
 $f \sim 10^4$ Hz
 $N \sim 10^{10}$

Summary and future work

We estimated the quantum noise of gravitons on the mirrors of gravitational wave interferometers

We found that the quantum noise of gravitons might be detectable by the future DECIGO

By using the decoherence time due to gravitons in a lab, we might be able to discover gravitons indirectly in the future