

Holographic domain walls, multi boundary traversable wormholes and Replica wormholes

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Based on [arXiv:2011.12962](https://arxiv.org/abs/2011.12962) + work in progress

Introduction

Recently the page curve of evaporating black hole is reproduced from semiclassical gravity computation:

[Almheiri-Engerhardt-Marolf-Maxfield, 19]

[Penington, 19]

[Almheiri-Mahajan-Maldacena-Zhao, 19]

Spacetime wormholes play important roles

[Almheiri-Hartman-Maldacena-Shaghoulian-Tajdini, 19]

[Penington-Stanford-Shenker-Yang, 19]

One of important aspects:

Application of RT/HRT/EW formula for “holographic states”
in non-gravitational system, or entangled with gravitational system

[Harlow, 16] [Hayden-Penington, 18]

[AMMZ, 19]

Today we consider holographic state in free field (or Ising models),
based on bra-ket wormholes obtained from traversable wormholes.

Using JT gravity + matter system as domain walls between
holographic CFTs and free CFTs.

Traversable wormholes in JT gravity + matter system

Traversable wormhole solution:

2d Gravity

3d Gravity

QM

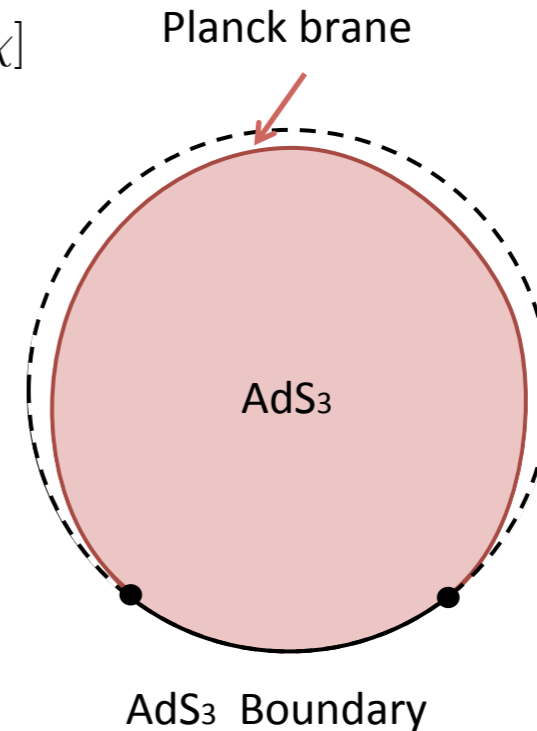
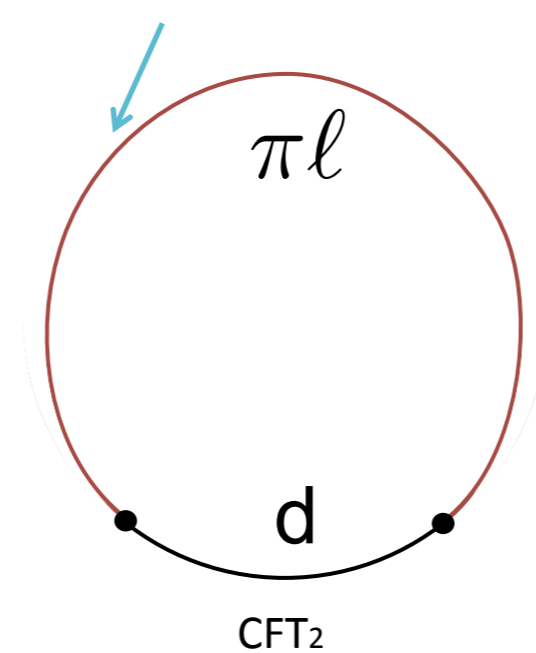
ansatz:

$$ds^2 = \frac{-dt^2 + d\sigma^2}{\ell^2 \sin^2 \frac{\sigma}{\ell}}$$

$$\phi(\sigma) = \frac{2\bar{\phi}_r}{\pi\ell} \left[\frac{\frac{\pi}{2} - \sigma}{\tan \frac{\sigma}{\ell}} + 1 \right]$$

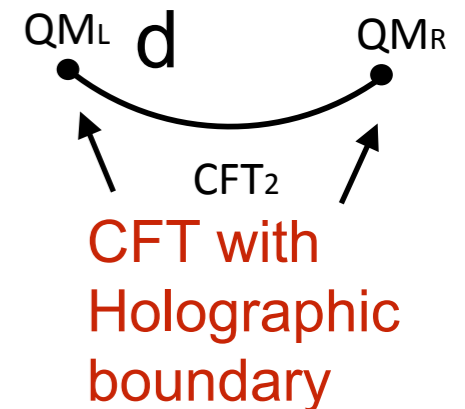
$$\langle T_{++}^{mat} \rangle = \frac{c}{48\pi\ell^2} - \frac{\pi^2 c}{12(\pi\ell + d)^2}$$

$$I_{\text{grav}}[g_{ij}^{(2)}, \phi] + I_{\text{CFT}}[g_{ij}^{(2)}, \chi]$$



Ground state in coupled system

$|G\rangle$



ℓ : “wormhole length”, dynamically determined by EOM

Alternatively: use variational method (approximate by TFD)

$$E(\ell) = 2 \times \underbrace{\frac{\pi\bar{\phi}_r}{4G_N} T_H^2}_{\text{BH mass}} + \underbrace{\frac{c}{24\ell}}_{\text{Weyl anomaly}} - \underbrace{\frac{c\pi}{6(\pi\ell + d)}}_{\text{Casimir energy}}, \quad \ell = \frac{1}{2\pi T_H}$$

minimize variational energy gives the same ℓ

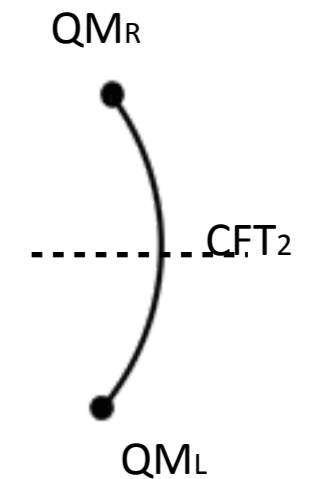
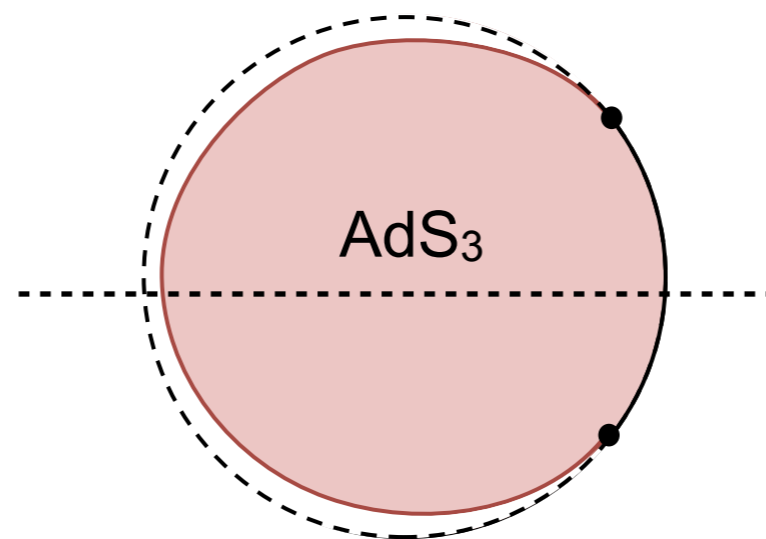
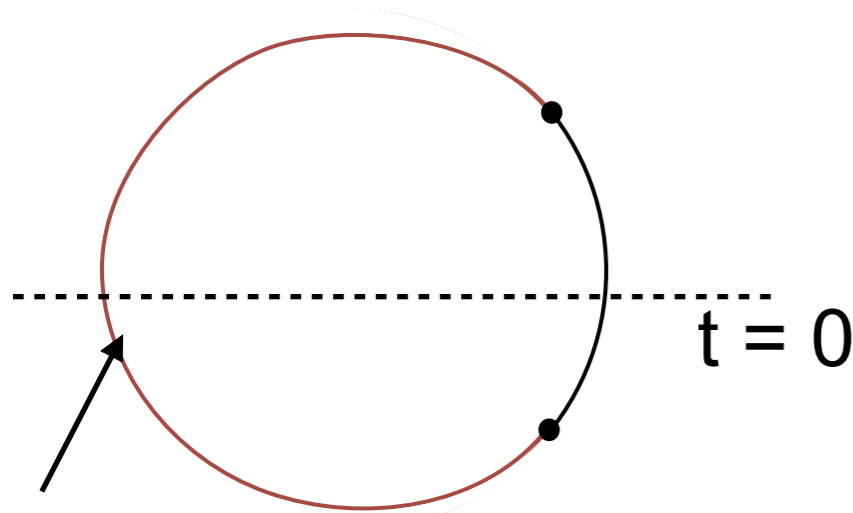
Bra-ket wormhole interpretation

- start from traversable wormhole, exchange (Euclidean) time and space

2d gravity

3d gravity

QM



wormhole connects bra and ket

[cf: Page, 86] [Chen-Gorbenko-Maldacena, 20]

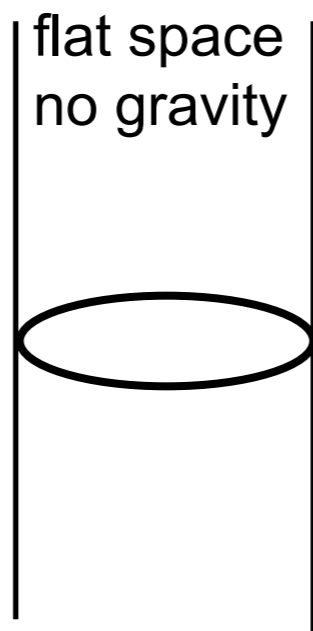
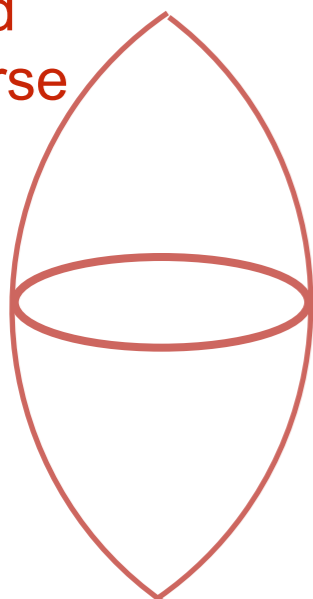
$$|\Psi\rangle \in \mathcal{H}_{CFT}$$

- wick rotation $\tau = it$ in 2d gravity description

FLRW
closed
universe

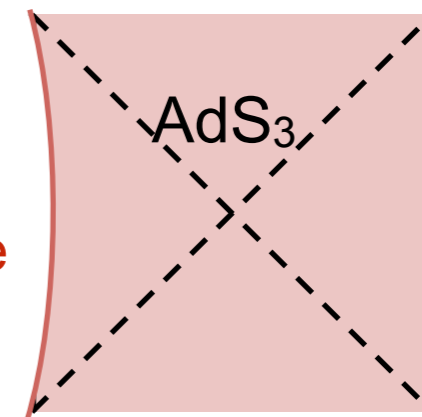
flat space
no gravity

entangled



3d gravity description

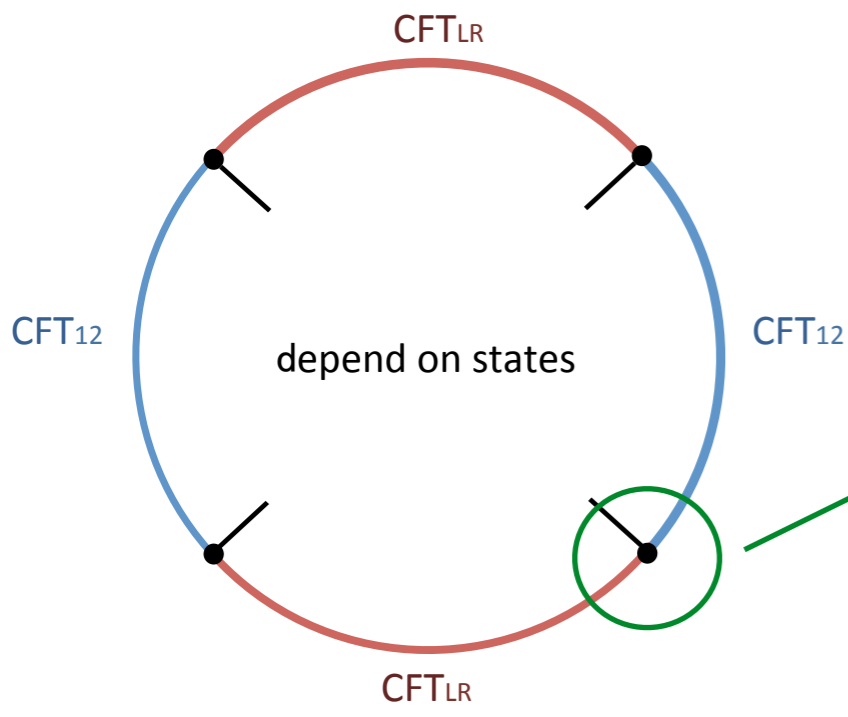
EOW
brane



[Gubser, 99]

[cf: Cooper-Rozali-Swingle-Raamsdonk-Waddell-Wakeham, 18]

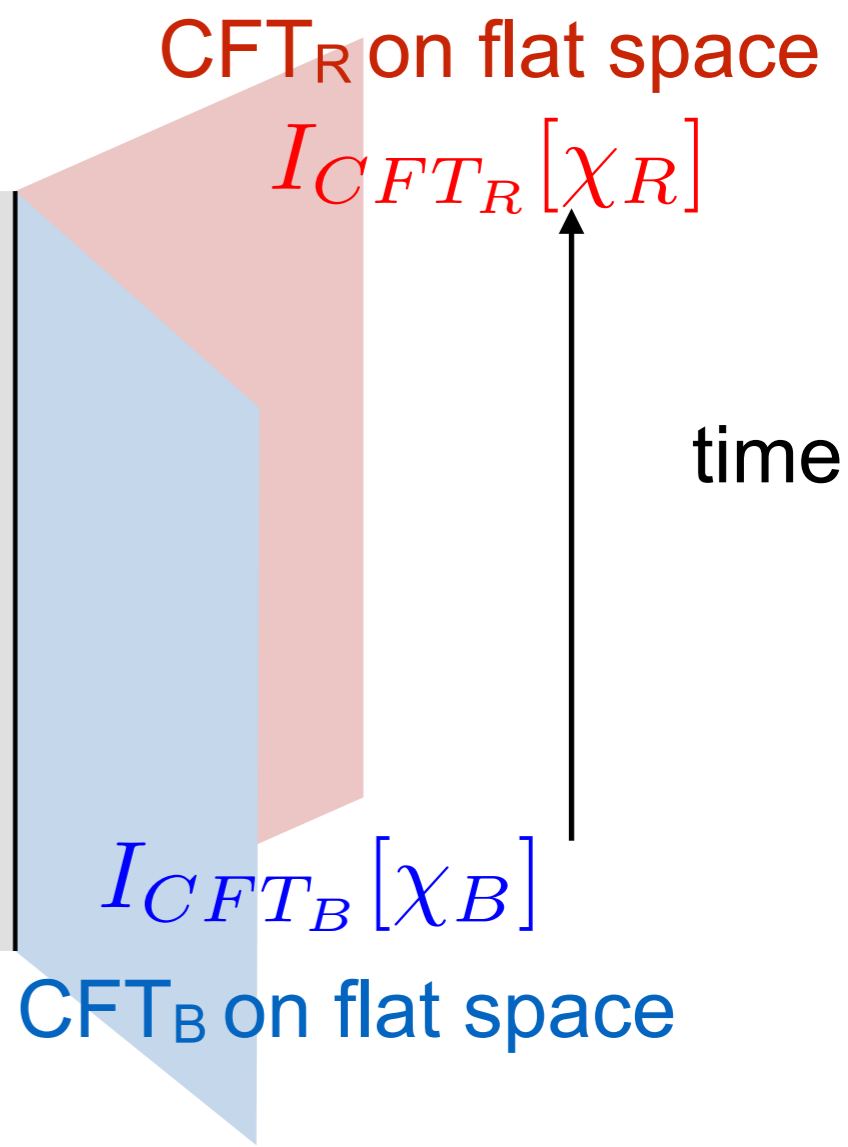
Generalization to Holographic domain wall



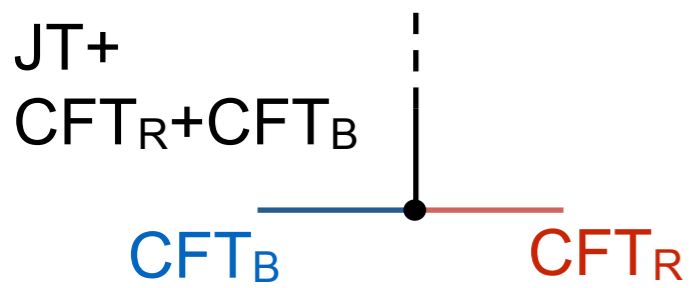
time slice

JT gravity
+CFT_B
+CFT_R

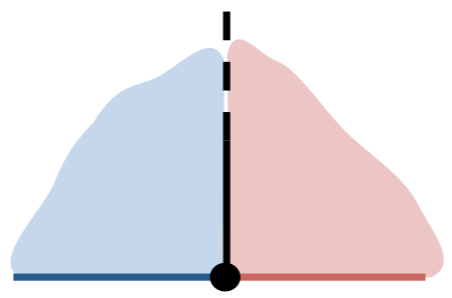
$$I_{\text{grav}}[g_{ij}^{(2)}, \phi] + I_{\text{CFT}_R}[g_{ij}^{(2)}, \chi_R] + I_{\text{CFT}_B}[g_{ij}^{(2)}, \chi_B]$$



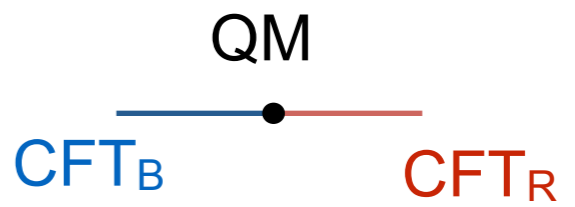
2d gravity



3d gravity



QM

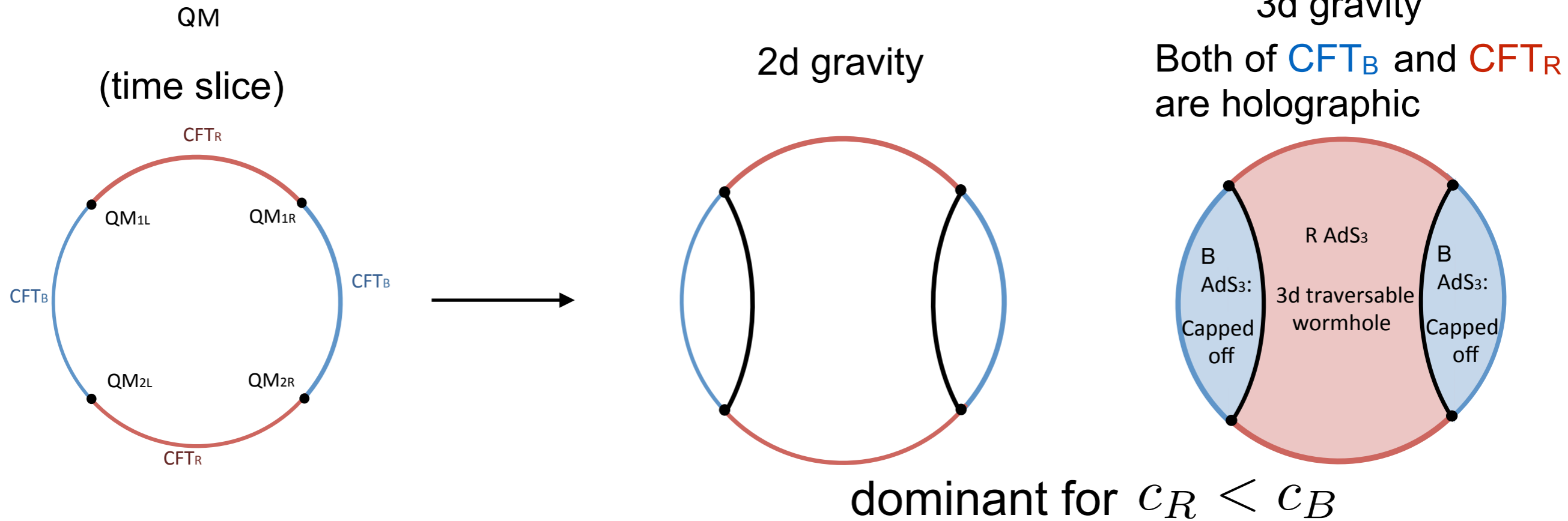


- two CFTs are joined by domain walls with holographic duals

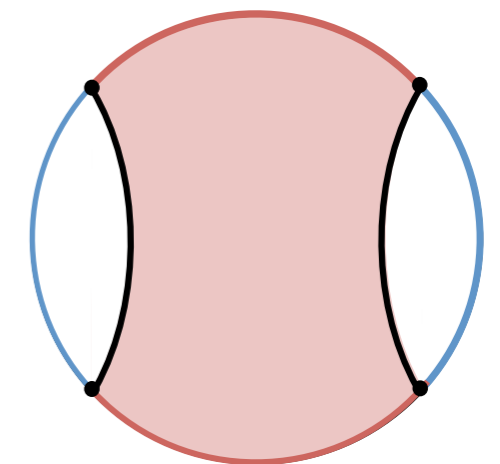
[cf: Chen-Myers-Neuenfeld-Reyes-Sandor ,20]
[Ooguri-Takayanagi, 20]

2 traversable Wormhole solution with Holographic matters:

- start from four domain walls



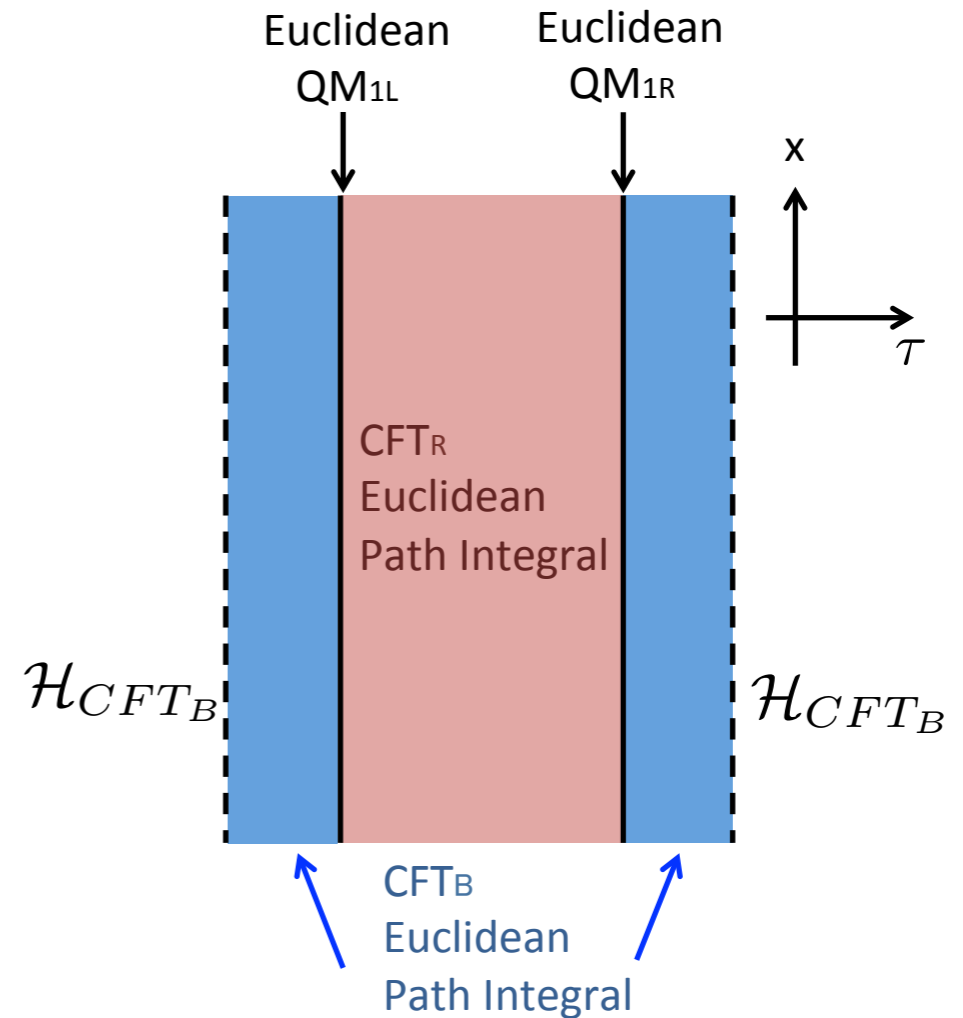
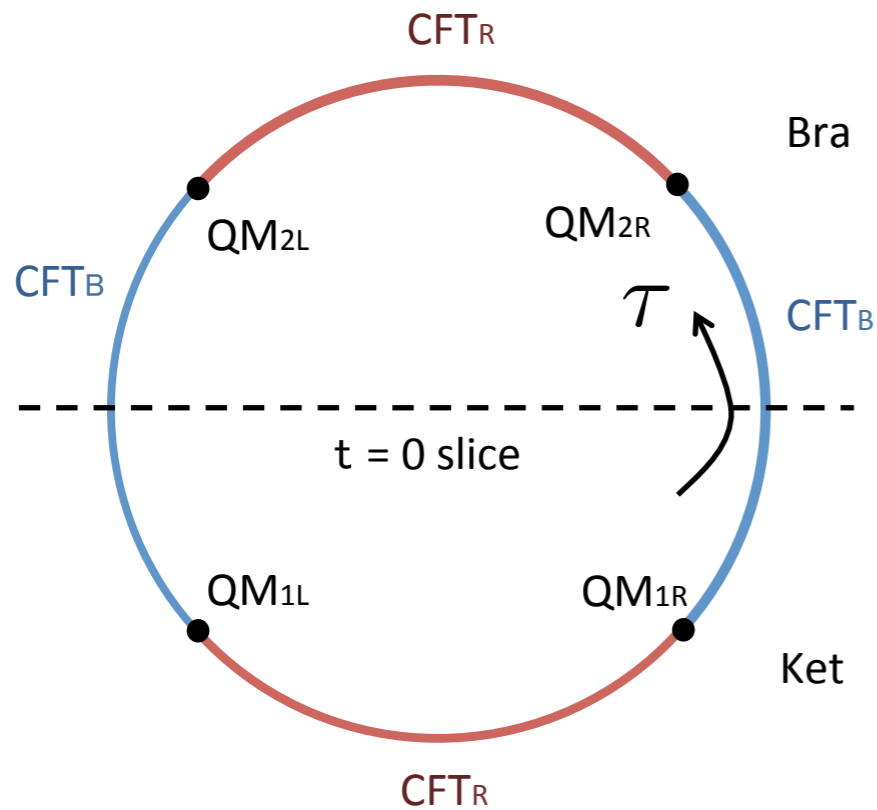
JT Gravity + Holographic CFT + other CFT
is similar to Randall-Sundram II
(4d gravity + 5d gravity (= 4d holographic CFT)
+ Standard model)



Only CFT_R is
holographic

Wick rotation:

QM description



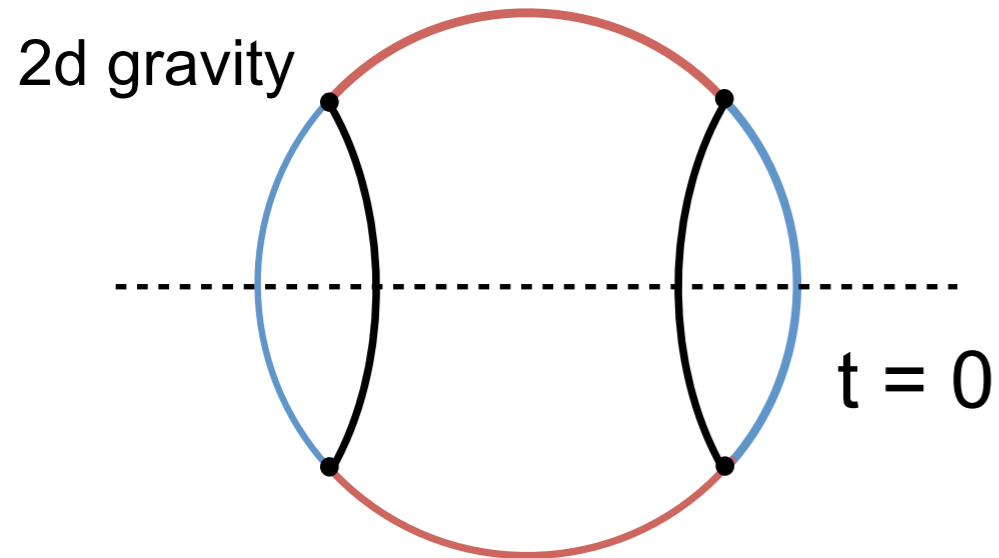
The path integral prepare a state

$$|\Psi\rangle \in \mathcal{H}_{CFT_B} \otimes \mathcal{H}_{CFT_B}$$

- We can also add bulk particles by inserting some operators in **CFT_R**
- Similar to problem in quantum quenches in QFT [cf: Calabrese-Cardy, 05]

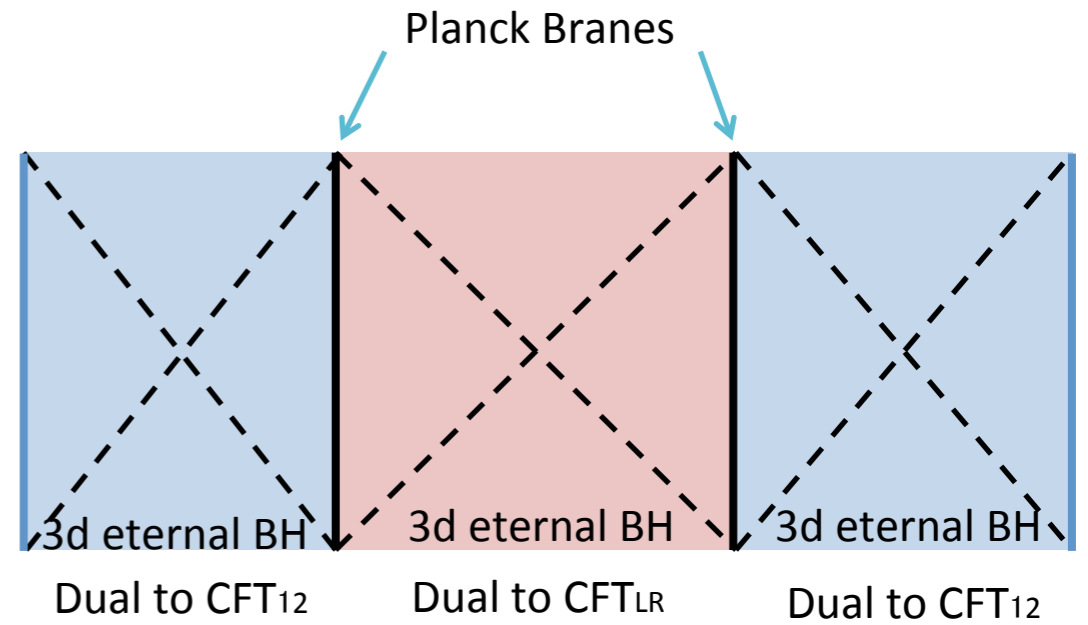
Lorentzian time continuation

bra-ket wormhole state

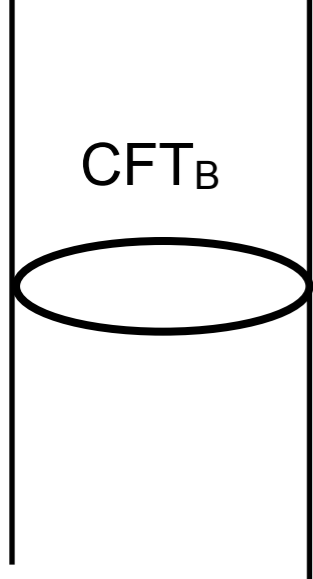


dominant saddle for $c_R < c_B$

3d Gravity



flat space
no gravity

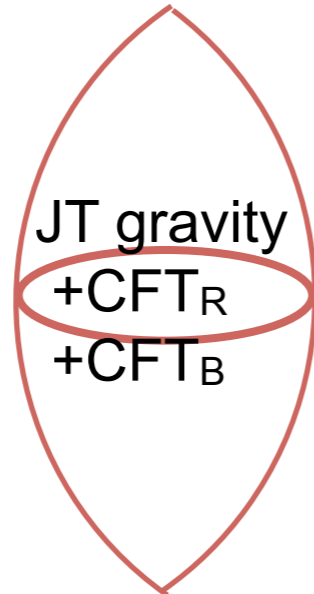


entangled

CFT_B



closed
universe

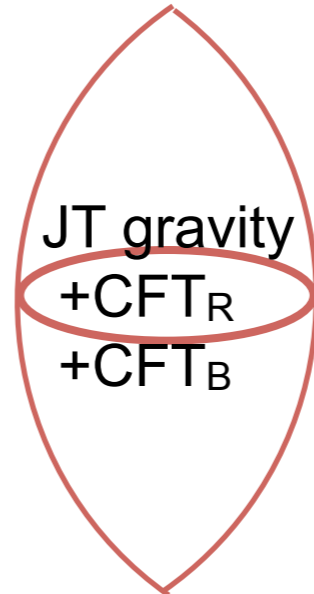


entangled

CFT_R



closed
universe

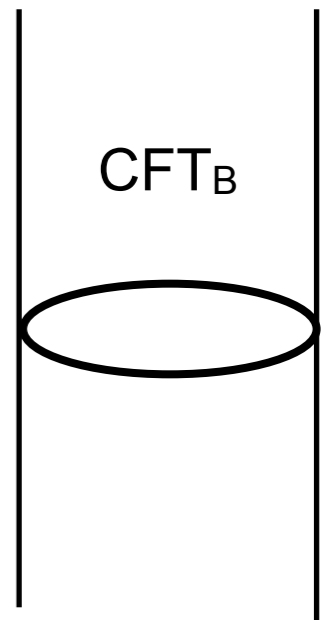


entangled

CFT_B

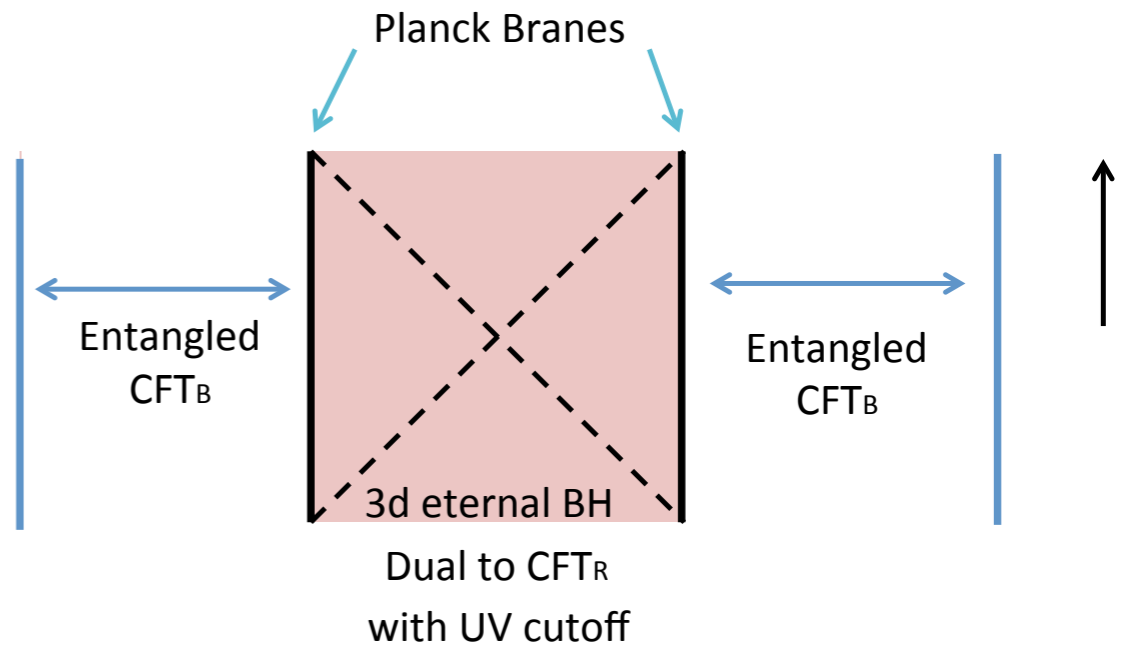
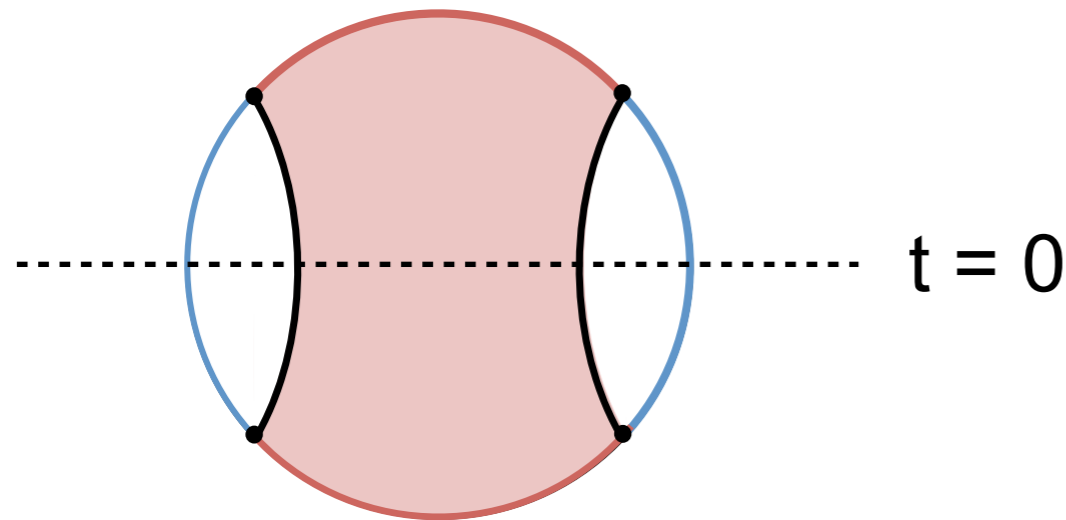


flat space
no gravity



Embed holographic states to free field Hilbert spaces:

“partially doubly holographic”



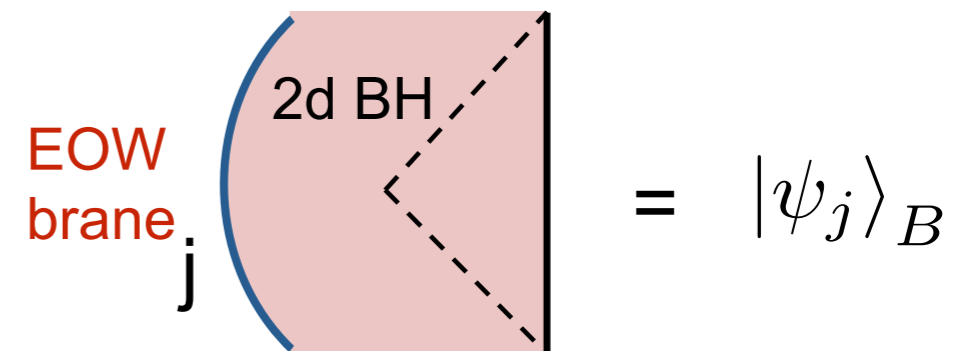
[cf: “Holo-ween” Simidzija-Raamsdonk, 20]

states in $\mathcal{H}_{CFT_B} \otimes \mathcal{H}_{CFT_B}$

- Take CFT_B to be bunch of free fields (or Ising CFTs)
 → realize holographic states in free fields Hilbert sp.

cf: JT gravity with auxiliary system

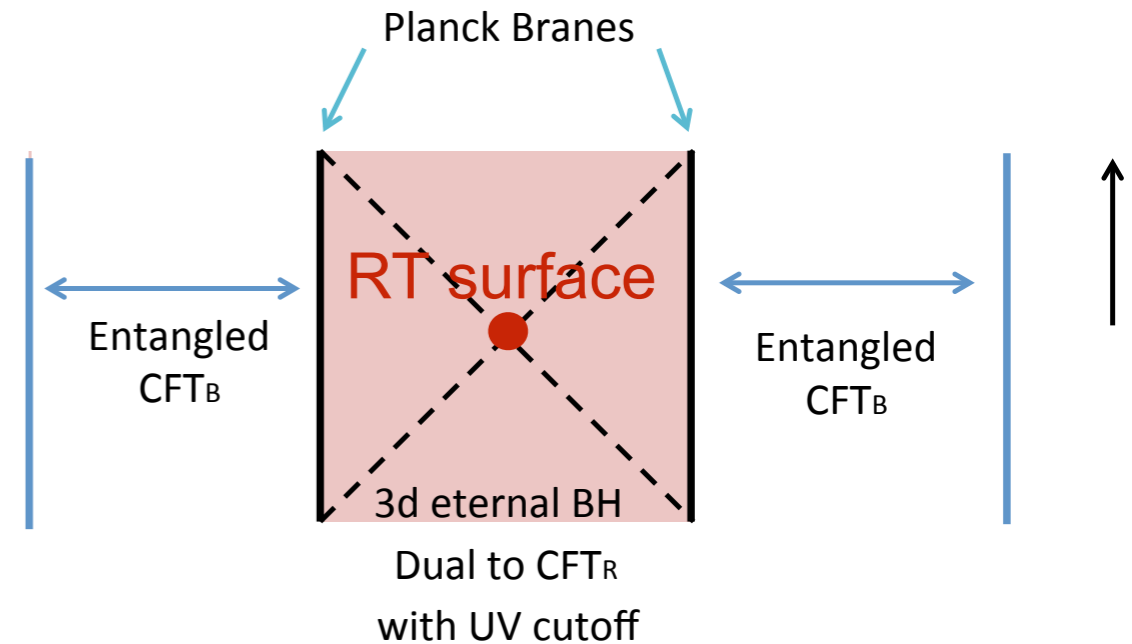
[Penington-Stanford-Shenker-Yang 19]



$$\rho_R \rightarrow |\rho_R\rangle = \sum \langle \psi_i | \psi_j \rangle_B |i\rangle_R \otimes |j\rangle_R \in \mathcal{H}_{aux} \otimes \mathcal{H}_{aux}$$

Entanglement entropy and Replica wormholes:

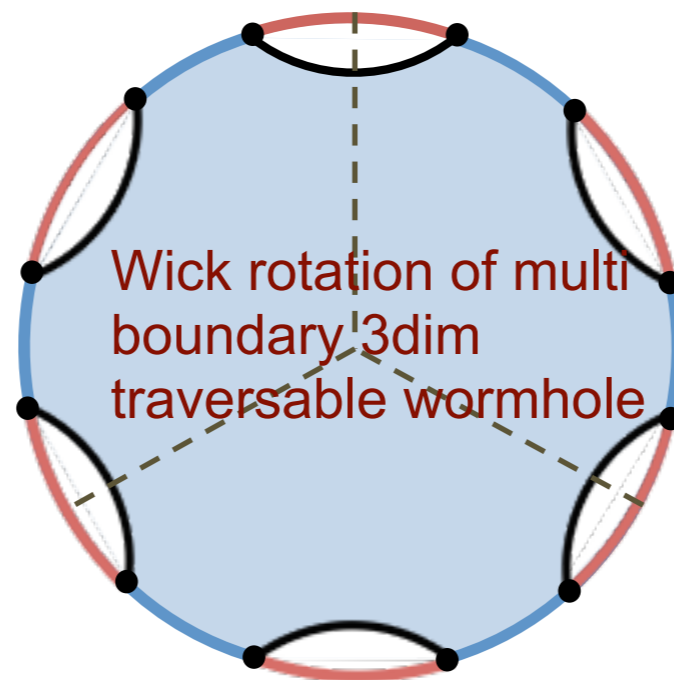
Entanglement entropy between two side is calculated using RT formula for states in non-holographic CFTs



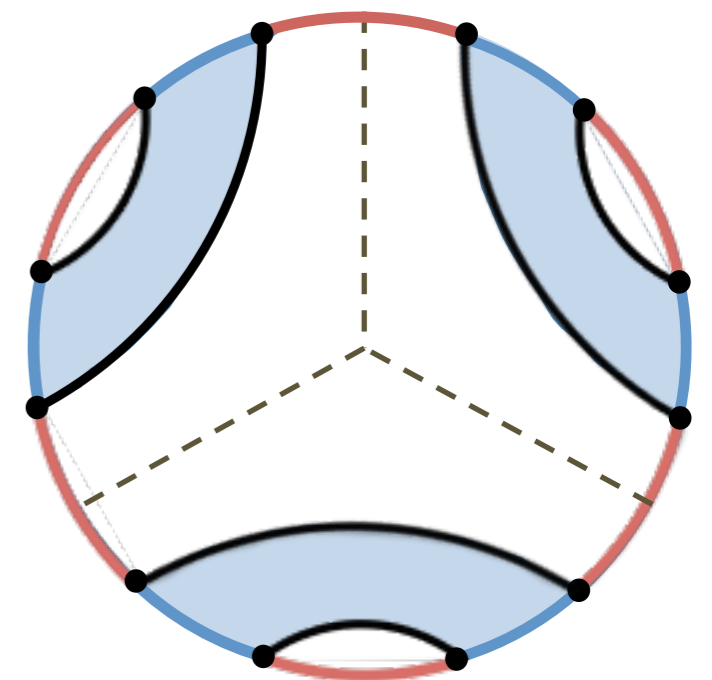
$n=3$ Renyi Entropy

Replica wormhole justify this calculation

CFTs with $4n$ domain walls



Replica wormhole



No wormhole

Lewkowycz-Maldacena derivation and holographic Renyi entropy:

[cf:Lewkowycz-Maldacena, 13]

Partition function: $Z_n \approx e^{-E_n L}$ for large L

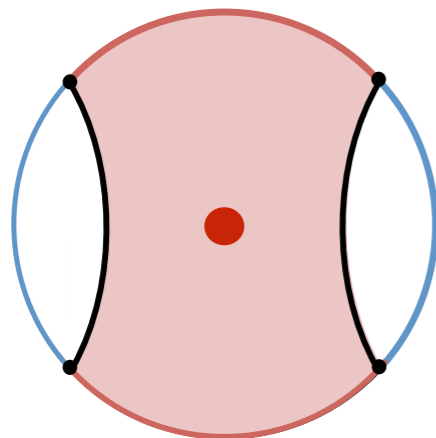
→ Calculation reduces to that of energy in traversable wormholes !

$$E_n = E(\ell_n, n) \quad \left. \frac{\partial E(\ell, n)}{\partial \ell} \right|_{\ell=\ell_n} = 0$$

$$E(\ell, n) = n \left(\frac{\bar{\phi}_r}{4\pi G_N \ell^2} + \frac{c_R + c_B}{12\ell} - \frac{c_R \pi}{3(\pi\ell + d_R)} - \frac{c_B \pi}{12n^2(\pi\ell + d_B)} \right)$$

$$\left(\ell_n = \frac{6\bar{\phi}_r}{\pi G_N} \frac{1}{3c_R - (1 - \frac{1}{n^2})c_B} \quad \text{for } d_R, d_B \ll \ell \right)$$

$$E(\ell_n, n) - nE(\ell_1, 1) = \underbrace{[E(\ell_n, n) - E(\ell_1, n)]}_{O((n-1)^2)} + \underbrace{[E(\ell_1, n) - nE(\ell_1, 1)]}_{\frac{c_B}{24} \left(n - \frac{1}{n}\right) \frac{\pi}{\pi\ell_1 + d_B}}$$



= dimension of twist operator !

Lewkowycz-Maldacena derivation and holographic Renyi entropy:

Energy par a Replica: $\hat{E}(\ell, n) = \frac{E(\ell, n)}{n}$

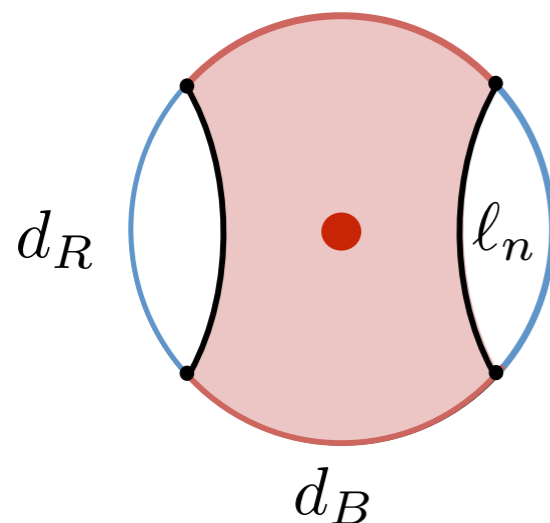
$$n^2 \frac{d}{dn} \hat{E}(\ell(n), n) = n^2 \frac{\partial \hat{E}}{\partial \ell} \frac{d\ell_n}{dn} + n^2 \frac{\partial \hat{E}}{\partial n}$$

$= 0$ $= \frac{c_B}{24} \frac{1}{\pi \ell_n + d_B} \frac{\partial}{\partial n} \left(1 - \frac{1}{n^2}\right)$:derivative of twist op dim

by saddle eq.

$$= \frac{c_B}{12} \frac{1}{n(\pi \ell_n + d_B)}$$

→ Similar result w/ holographic Renyi Entropy [Dong, 16] holds.



Analysis of other saddles are still work in progress...

Conclusion/Future works

- We discussed relations between traversable WH and bra-ket WH.
- Using JT gravity + matters as holographic defects on CFTs. They can be used to embed holographic 2d CFT state into Free fields Hilbert spaces.
- Even in free CFTs, entanglement entropy is calculated by RT formula for holographic states
- These are justified by Replica wormholes. Replica wormholes in this situation are wick rotation of traversable wormholes.
- It is interesting to consider other entanglement measures in this setup
- Tensor networks/ de Sitter version etc...