

# Towards understanding cosmological correlators from boundary perspective

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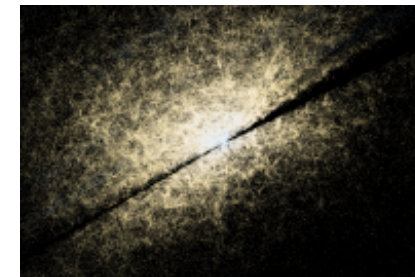
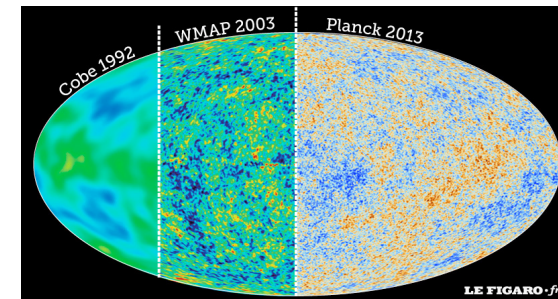
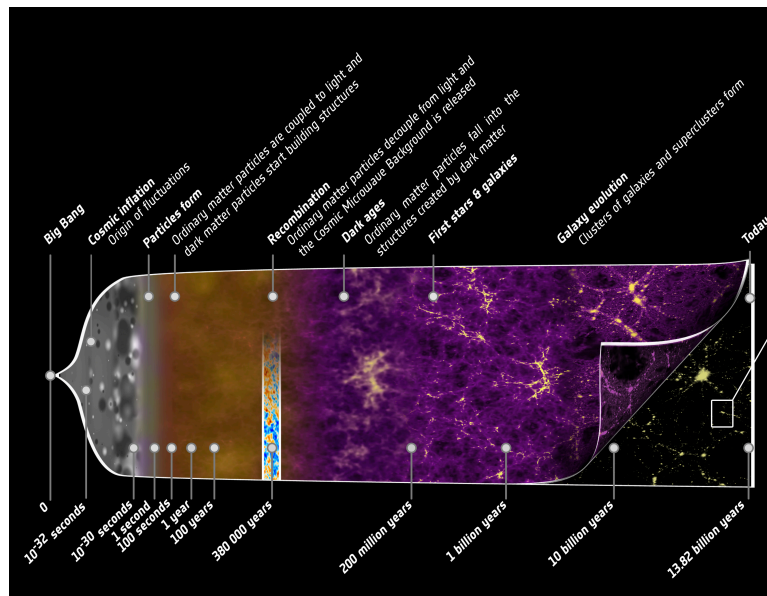
Yuko Urakawa (Bielefeld Univ./Nagoya Univ.)

# Acceleration just after “beginning” of Universe

## Cosmic inflation

Recall Yasusada's talk

$$H^2 = \left( \frac{1}{a(t)} \frac{da(t)}{dt} \right)^2 = \frac{8\pi G}{c^2} \rho = \frac{\Lambda c^2}{3} \longrightarrow a(t) = a_i e^{\sqrt{\frac{\Lambda}{3}} c(t-t_i)}$$



Observations

inflation (almost) established

# Cosmic inflation as natural laboratory

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$$H_{\text{inf}} < 2.7 \times 10^{-5} M_{\text{pl}} \sim 6.6 \times 10^{13} \text{ GeV} \quad \text{PLANCK18}$$

Natural laboratory to experiment with high energy physics (HEP)

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

## UV modification of gravity

- Radiative corrections

$$G_{\mu\nu} = \frac{8\pi G}{c^4} \langle T_{\mu\nu} \rangle \quad \langle T_{\mu\nu} \rangle = \frac{k_2}{2990\pi^2} \left( R_{\mu\rho} R^{\rho\nu} - \frac{2}{3} R R_{\mu\nu} + \frac{1}{4} g_{\mu\nu} R^2 - \frac{1}{2} g_{\mu\nu} R_{\rho\sigma} R^{\rho\sigma} \right) + \dots$$

Davies (77)  $\rightarrow$  Starobinsky (80)

- GR in 4 dim is not UV complete

e.g., Horava (09), string theory....

## BSM physics

- inflaton

- spectator fields

mass  $m \ll H$  or  $m > H$

spin  $s=0, 1, 2, \dots$

eg.  $s=1$ , axionic inflation



# de Sitter (dS) spacetime

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dS and EAdS are both embedded (d+1)-dim hyperboloid

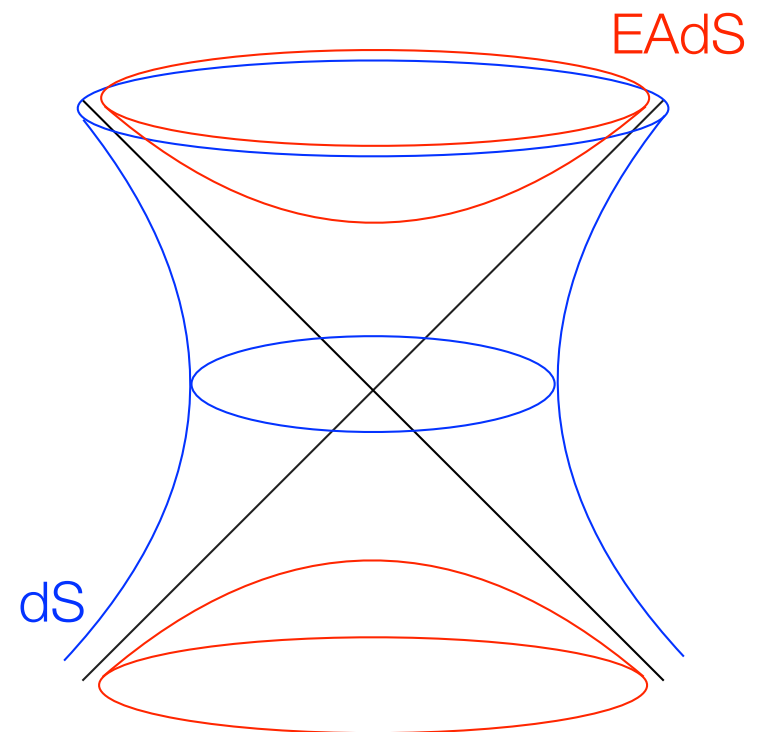
$$\epsilon_- X_-^2 + \epsilon_0 X_0^2 + \sum_{i=1}^d X_i^2 = \epsilon_- l^2$$

in (d+1, 1) flat space

$$ds^2 = \epsilon_- dX_-^2 + \epsilon_0 dX_0^2 + \sum_{i=1}^d dX_i^2$$

$$\left\{ \begin{array}{ll} \text{EAdS} & \epsilon_- = -1, \epsilon_0 = 1 \\ \text{dS} & \epsilon_- = 1, \epsilon_0 = -1 \end{array} \right.$$

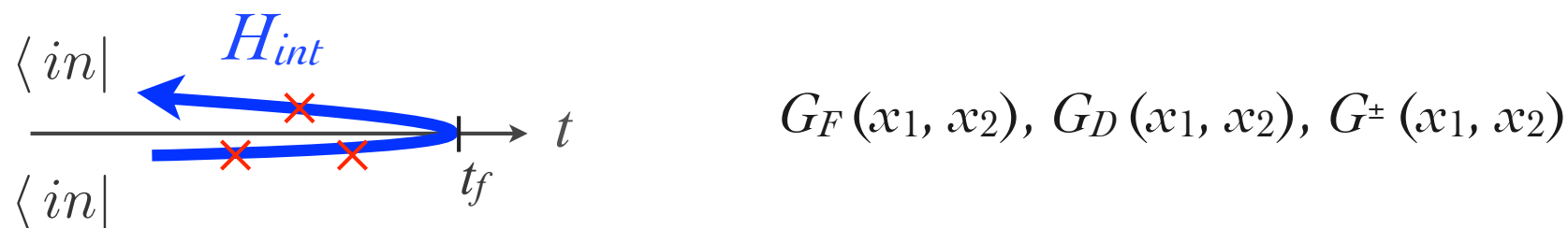
c.f. AdS  $\epsilon_- = \epsilon_0 = -1$



# In-in formalism and BD vacuum (Euclidean vac.)

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cosmological correlators are expectation values



## Euclidean vacuum (Bunch-Davies vac., Adiabatic vac.,.....)

All the  $n$ -point fns.  $\langle \phi(x_1) \dots \phi(x_{n-1}) \phi(x_n) \rangle$  should become regular in the limit  $t_i \rightarrow -\infty(1 \pm i\epsilon)$

free-level: dS invariant + Hadamard

B. Allen (85)

# 4D AdS and dS

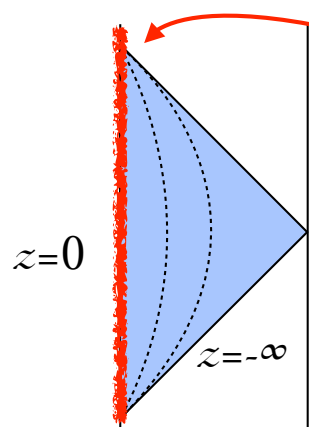
## Anti de Sitter (AdS)

Vacuum with  $\Lambda < 0$

in  $\mathbb{R}^{2,3}$   $(-, -, +, +, +)$   $SO(2,3)$

$$-X_0^2 - X_1^2 + \sum_{a=2,3,4} X_a^2 = -l^2$$

$$ds^2 = l_{\text{AdS}}^2 \left( \frac{-dt^2 + dx^2 + dy^2 + dz^2}{z^2} \right)$$



Boundary

.....  
z:const,  $\mathbb{R}^3$



$$l_{\text{AdS}} \rightarrow i l_{\text{dS}}$$

$$z \rightarrow i\eta$$

$$t \rightarrow -i w$$

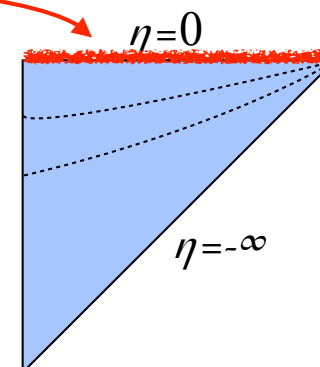
## de Sitter (dS)

Vacuum with  $\Lambda > 0$

in  $\mathbb{R}^{1,4}$   $(-, +, +, +, +)$   $SO(1,4)$

$$-X_0^2 + X_1^2 + \sum_{a=2,3,4} X_a^2 = l^2$$

$$ds^2 = l_{\text{dS}}^2 \left( \frac{-d\eta^2 + dx^2 + dy^2 + dw^2}{\eta^2} \right)$$



.....  
 $\eta$ :const,  $\mathbb{R}^3$

# dS/CFT conjecture

*Strominger(01), Witten(01), Maldacena(02), ...*

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dS and AdS are mathematically similar, yet physically different.

- Holographic direction is timelike, CFT lives on spatial boundary.
- Dual CFT is non-unitary.

e.g., $(\square - m^2)\phi = 0$	4dim AdS	$\Delta = \frac{3}{2} \pm \sqrt{\frac{9}{4} + m^2 l_{\text{AdS}}^2}$
	4dim dS	$\Delta = \frac{3}{2} \pm \sqrt{\frac{9}{4} - m^2 l_{\text{dS}}^2}$

Principal series in dS ~ Below BF found in AdS *Isono, Liu, Noumi(20)*

- Lack of concrete examples.

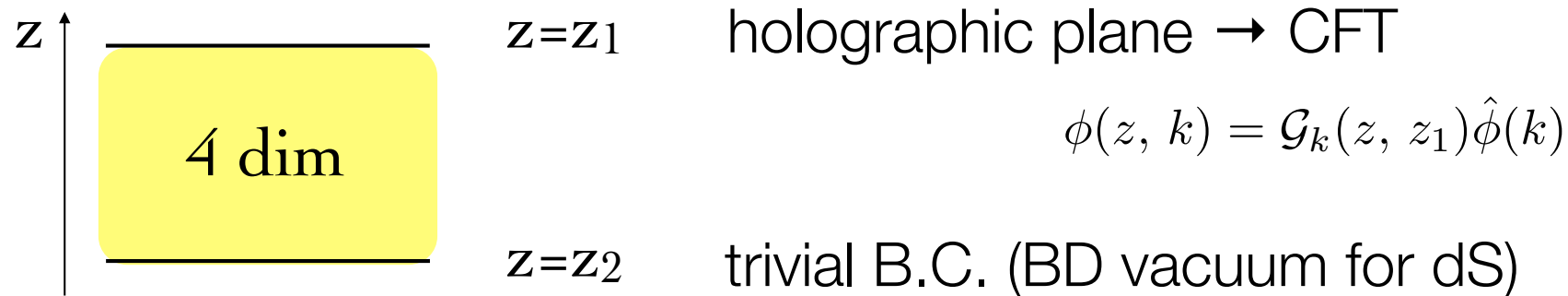
... yet, *Anninos, Hartman, Strominger(11)*

- Relation with area law?

(...yet Tetsuya's talk)



# Bulk to Boundary



using eom for 4-dim bulk theory  $\delta S \sim \mathcal{L} dz \Big|_{z=z_2}^{z=z_1}$

## - Massless free scalar field on EAdS

with  $Z \sim e^{-S_{\text{EAdS}}}$   $\langle \mathcal{O}(\mathbf{k}) \mathcal{O}(\mathbf{k}') \rangle \sim \delta(\mathbf{k} + \mathbf{k}') l_{\text{AdS}}^2 k^3$

## - Massless free scalar field on dS

with  $Z \sim e^{iS_{\text{dS}}}$   $\langle \mathcal{O}(\mathbf{k}) \mathcal{O}(\mathbf{k}') \rangle \sim \delta(\mathbf{k} + \mathbf{k}') l_{\text{dS}}^2 (-k^3)$

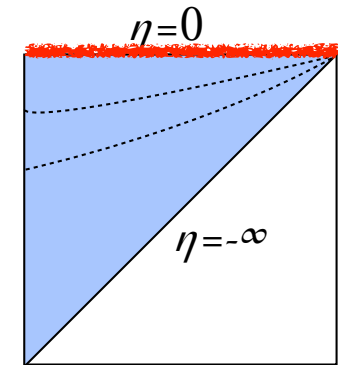
$$l_{\text{dS}} = i l_{\text{AdS}}$$

# Boundary to Bulk (dS)

- Wave function on future boundary

*Maldacena (02)*

$$\Psi[\hat{\phi}] \sim Z[\hat{\phi}] = \langle e^{-\hat{\phi} \cdot \mathcal{O}} \rangle$$



Bulk correlators

$$\langle \hat{\phi}(\mathbf{k}_1) \cdots \hat{\phi}(\mathbf{k}_n) \rangle = \int D\hat{\phi} \hat{\phi}(\mathbf{k}_1) \cdots \hat{\phi}(\mathbf{k}_n) |\Psi[\hat{\phi}]|^2$$

e.g.  $\langle \hat{\phi}(\mathbf{k}) \hat{\phi}(-\mathbf{k}) \rangle' = -\frac{1}{2\text{Re}[\langle \mathcal{O}(\mathbf{k}) \mathcal{O}(-\mathbf{k}) \rangle']}$

Harrison-Zeldovich spectrum (← 0th approx. of CMB spectrum)

$$\langle \mathcal{O}(\mathbf{k}) \mathcal{O}(-\mathbf{k}) \rangle' \propto k^{6-\Delta} \propto k^3 \quad \longrightarrow \quad \langle \hat{\phi}(\mathbf{k}) \hat{\phi}(-\mathbf{k}) \rangle' \propto 1/k^3$$

$\uparrow \Delta_+ = \frac{3}{2} + \sqrt{\frac{9}{4} - m^2 l_{\text{dS}}^2} \sim 3$

c.f. Inflation w/deformed CFT

*McFadden, Skenderis (09, 10, 11, ...) Bzowski et al. (12), Schalm et al. (12) Garriga & Y.U. (13, 16), Garriga, Skenderis, Y.U. (14), ...*

# Cosmological bootstrap

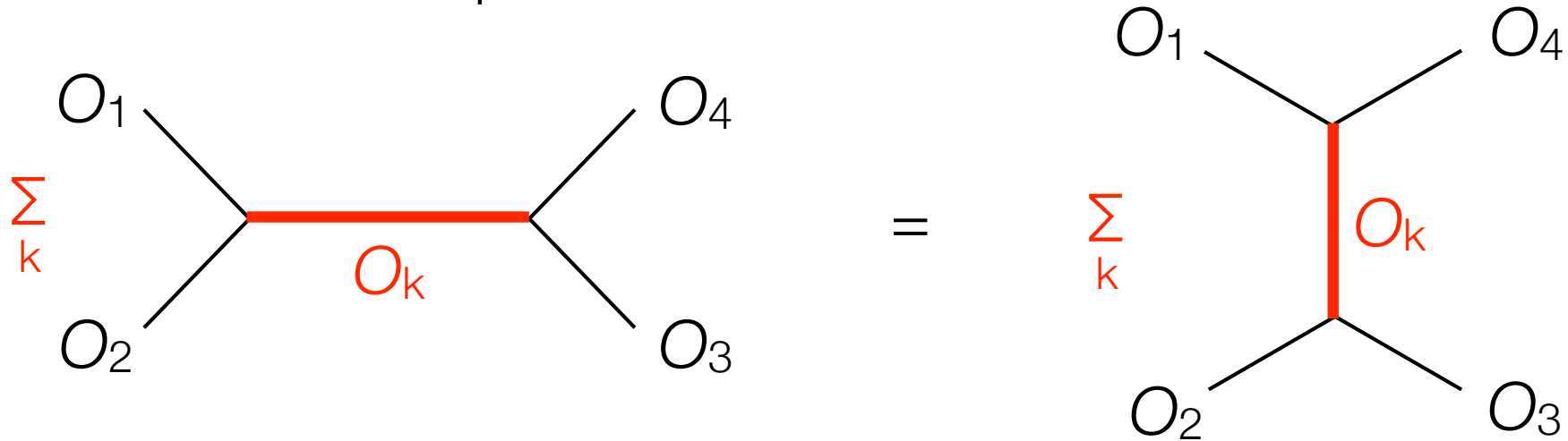
*Arkani-Hamed, Baumann, Lee, & Pimental JHEP 04 (20) 105*

*Baumann, Pueyo, Joyce, Lee, & Pimental JHEP 12 (20) 204*

*Baumann, Pueyo, Joyce, Lee, & Pimental 2005.04234*

# Cosmological bootstrap

## Conformal bootstrap



## (Goal of) Cosmological bootstrap

*Arkani-Hamed, Baumann, Lee, & Pimental (18)*

- Imposing de Sitter symmetry.
- Imposing correct UV energy singularities.

(+ Choosing Bunch-Davies vacuum, ..... )

→ Possible (4pt) exchange diagrams

# Scattering amplitude

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- Getting around redundancies in Lagrangian.
- Bootstrapping scattering amplitude from physical principles.

## Massless spinning particles

Compute  $A(1^{h_1} \dots n^{h_n}) = e_{\mu_1}^{h_1} \dots e_{\mu_n}^{h_n} A^{\mu_1 \dots \mu_n},$

using spin helicity variables  $[ij] = \tilde{\lambda}_{i\dot{\alpha}} \tilde{\lambda}_{j\dot{\beta}} \epsilon^{\dot{\alpha}\dot{\beta}}, \langle ij \rangle = \lambda_{i\alpha} \lambda_{j\beta} \epsilon^{\alpha\beta}$

## (MHV) Gluon scattering amplitude

$$A(\dots i^- \dots j^- \dots) = \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$



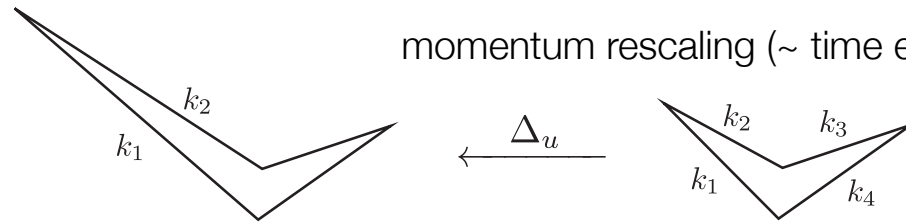
# Cosmological colliders

Arkani-Hamed & Maldacena (15)

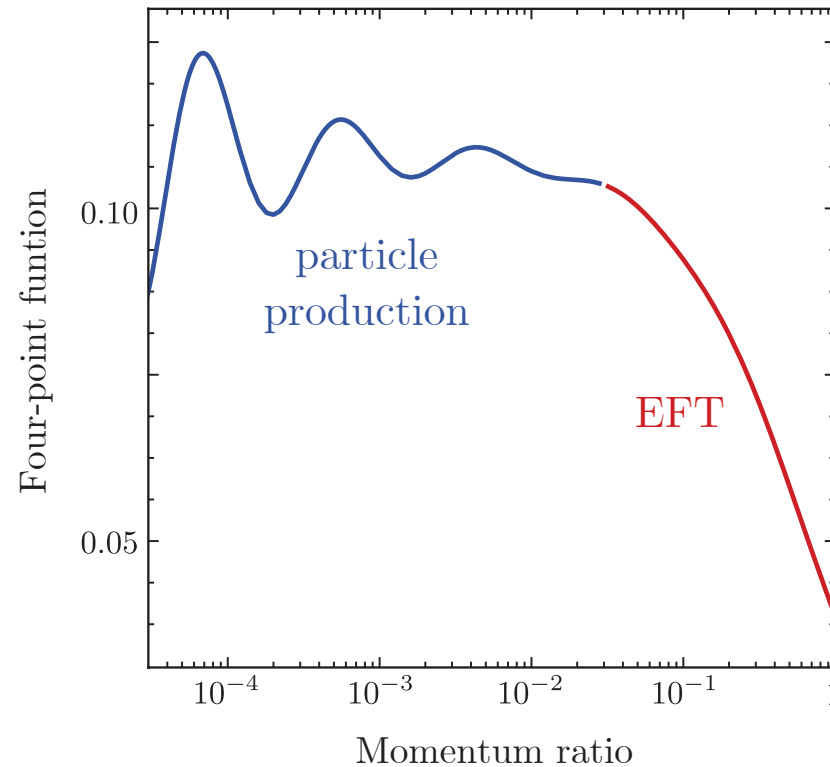
Arkani-Hamed et al. (18)

$$u \equiv \frac{s}{k_1 + k_2}, \quad v \equiv \frac{s}{k_3 + k_4}$$

$$s \equiv |\mathbf{k}_1 + \mathbf{k}_2|$$



$$\Delta_u \equiv u^2(1-u^2)\partial_u^2 - 2u^3\partial_u$$



ext. conformal scalar

Excitation of massive scalar *Chen & Wang (09), Noumi, Yamaguchi, Yokoyama (12).....,*

# Mellin space approach

Sleight JHEP 01 (20) 090

Sleight & Taronna JHEP 02 (20) 098

Sleight & Taronna arXiv:2007.09993



# Mellin transformation

Various advantages to analyze CFT (on  $\mathbf{R}^d$ ).

Mack(08)

encoding unitarity, causality, ... Recall Toshifumi's talk

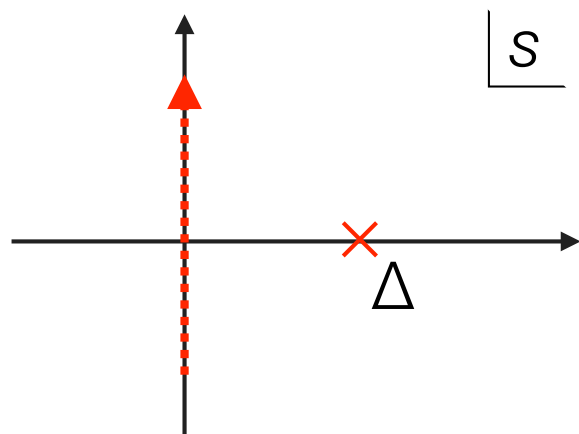
c.f. Fourier trans.    2pt fn.     $1/(x^2)^\Delta \longleftrightarrow (p^2)^{\Delta-d/2}$     branch cut at origin

## Mellin transformation

$$f(x) = \int_{-i\infty}^{i\infty} ds x^{-s} \tilde{f}(s)$$

ICTP lecture by Gopakumar

Picking out different scaling behaviors for power-law decomposition.



for  $\tilde{f}(s) \sim \frac{1}{s - \Delta} \quad (\Delta > 0)$

$x > 1 \quad f(x) \sim 1/x^\Delta$

# Understanding dS from AdS

Sleight(19)

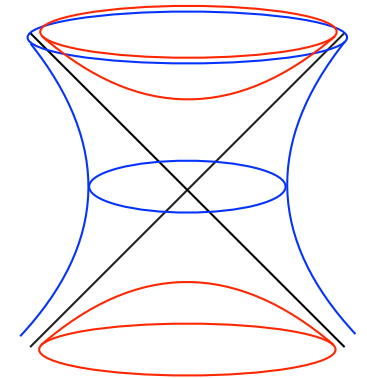
Sleight & Taronna(19)

Sleight & Taronna(20)

- AdS
- 1) Well defined notion of Unitarity
  - 2) Non-perturbative understanding of singularities

Using AdS as a guide to understand dS.

- Bulk-boundary/Bulk-Bulk propagator in Mellin-Barnes representation for dS.



dS ↔ EAdS

analytic continuation

$$z = -e^{\pm i \frac{\pi}{2} \eta}$$

# Mellin-Barnes (MB) representation

Sleight(19)

Sleight & Taronna(19)

Harmonic function in AdS  $(\nabla_{\text{AdS}}^2 - m^2) \Omega_\nu(x_1, x_2) = 0$

$$\Omega_\nu(x_1, x_2) = \int_{\partial\text{AdS}} \text{Diagram} \xrightarrow{\text{Fourier trans.}} \Omega_{\nu, \vec{k}}(z_1; z_2) = \frac{\nu^2}{\pi} K_{\frac{d}{2}+i\nu}(z_1, \vec{k}) K_{\frac{d}{2}-i\nu}(z_2, -\vec{k})$$

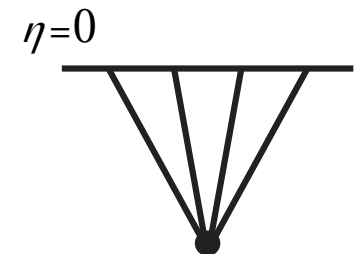
K: bulk-boundary propagator

Expressing K in MB rep, Performing analytic continuation  $z = -e^{\pm i\frac{\pi}{2}} \eta$

## MB rep. of Wightman fn. for BD vac. in dS

$$G_{\mathbf{k}}(\eta_1, \eta_2) \propto (\eta_1 \eta_2)^{d/2} \prod_{i=1}^2 \int_{-i\infty}^{i\infty} \frac{du_i}{2\pi i} (\dots) \prod_{j=1}^2 \underbrace{\Gamma(u_j + i\frac{\nu}{2}) \Gamma(u_j - i\frac{\nu}{2})}_{\text{poles}} \left(-\frac{k\eta_j}{2}\right)^{-2u_j}$$

- bulk-boundary(bulk) propagator by sending  $\eta_1 \rightarrow 0$
- contact diagram by integrating bulk-boundary propagator



Sleight(19)

Sleight & Taronna(19)

Sleight & Taronna(20)

# Understanding dS from AdS

AdS 1) Well defined notion of Unitarity

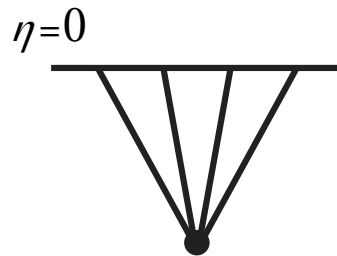
2) Non-perturbative understanding of singularities

## Using AdS as a guide to understand dS.

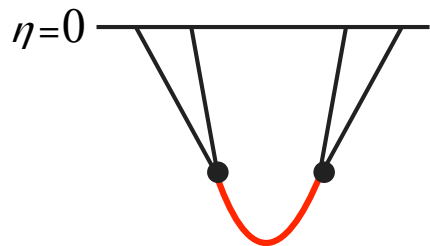
\* external scalar field is general. Recall Arkani-Hamed et al. (18)

• Propagators in Mellin-Barnes rep. for dS.

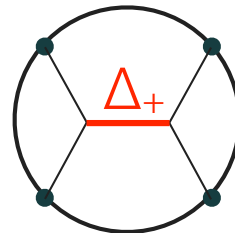
• Contact interaction



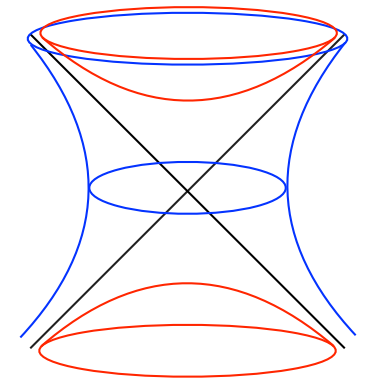
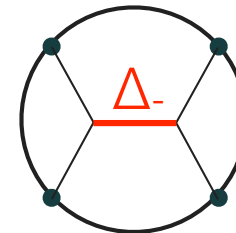
• Relation between exchange diagrams for dS and AdS



= (.....)

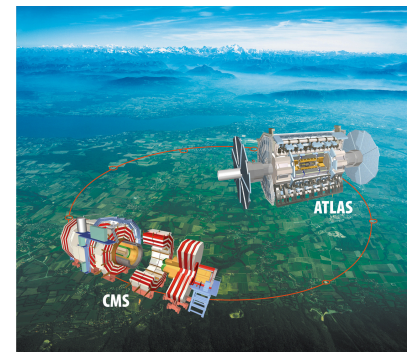
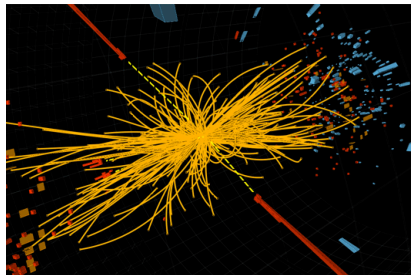


+ (.....)



dS ↔ EAdS  
analytic continuation

# Detectability of spinning particle



# 3pt fn. from spinning particle



$$k_3/k_1, k_3/k_2 \ll 1$$

Particle creation

$$\langle \delta\Phi\delta\Phi\delta\Phi \rangle \propto (\mathbf{k}_1 \cdot \mathbf{k}_3)^s$$

Arkani-Hamed & Maldacena (15)

# Minimum extension from dS

$$\sim \lambda^2 \dot{\phi}_{\text{bkg}} e^{-M/T_{\text{H}}} \left(\frac{k_1}{k_3}\right)^{\Delta} (\hat{\mathbf{k}}_1 \cdot \hat{\mathbf{k}}_3)^s \mathcal{P}(k_1) \mathcal{P}(k_3)$$

*Arkani-Hamed & Maldacena (15)*

$$k_3/k_1, k_3/k_2 \ll 1$$

$$\Delta = \frac{3}{2} \pm i \sqrt{\left(\frac{M}{H}\right)^2 - \left(s - \frac{1}{2}\right)^2} \quad \mathcal{P}(k) : \text{power spectrum}$$

1) Boltzmann suppression

Higuchi bound  $M > O(H)$

2) Dilution of massive fields

dS *Bordin et al. (18)*      global rotation *Kehagias & Riotto (17)*

3) Weak coupling

Large coupling tends to make 1) more severe.

*Wang & Xianyu (19, 20)*

# 3pt fn. from spinning particle



$$k_3/k_1, k_3/k_2 \ll 1$$

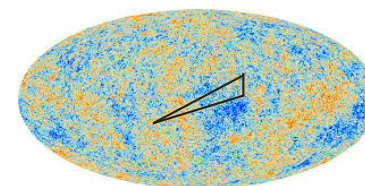
Particle creation

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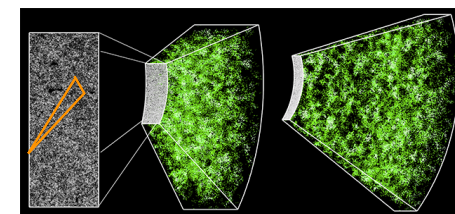
## 1) Cosmic microwave background $\langle \delta T \delta T \delta T \rangle$

PLANCK18, Bartolo et al. (17), Franciolini et al. (18),  
Bordin & Cabass (19)



## 2) Galaxy number count $\langle \delta n \delta n \delta n \rangle$

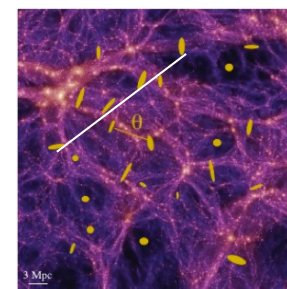
Moradinezhad et al (18<sup>1</sup>, 18<sup>2</sup>)



## 3) Galaxy shape correlation $\langle (\text{shape}) (\text{shape}) \rangle$

Schmidt et al (15, 16), Kogai, Matsubara, Nishizawa, Y.U. (18)

Kogai, Akitsu, Schmidt, Y.U. (20)

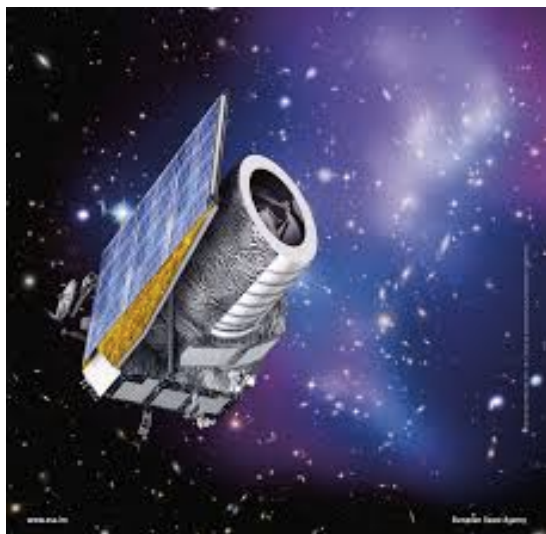




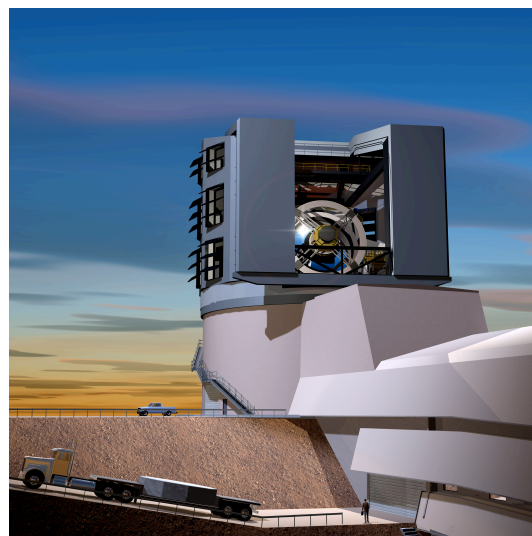
# Large Scale Structure surveys in the next decade

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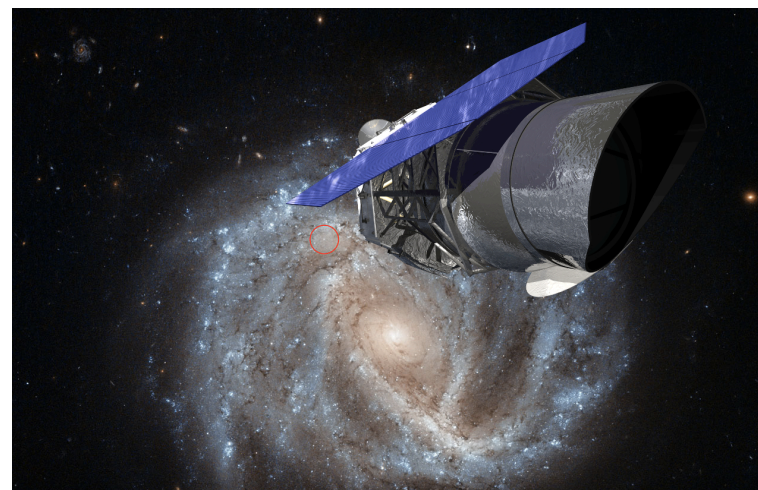
## Wide and Deep survey missions



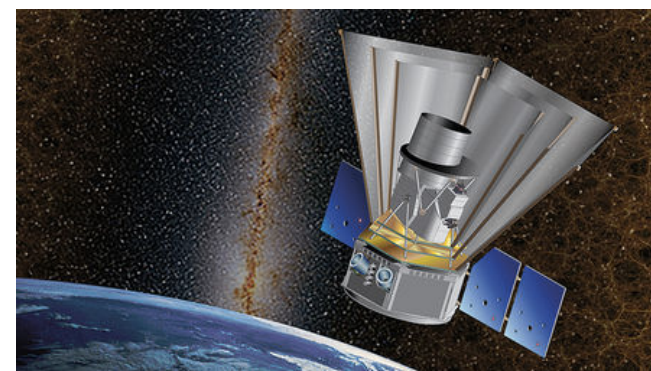
Euclid launch 2020



LSST science operation mid2020s



Roman space telescope launch mid2020s



SPHREx launch 2023

# Galaxy imaging survey as spin sensitive detector

Galaxy shape traces tidal force  $\Phi[\partial\partial\Phi]^{\text{TL}}$ ,  $\Phi[\partial\partial\partial\Phi]^{\text{TL}}$ ,  $\Phi[\partial\partial\partial\partial\Phi]^{\text{TL}}$ , ...

Angular dep. PNG of spin-s particle

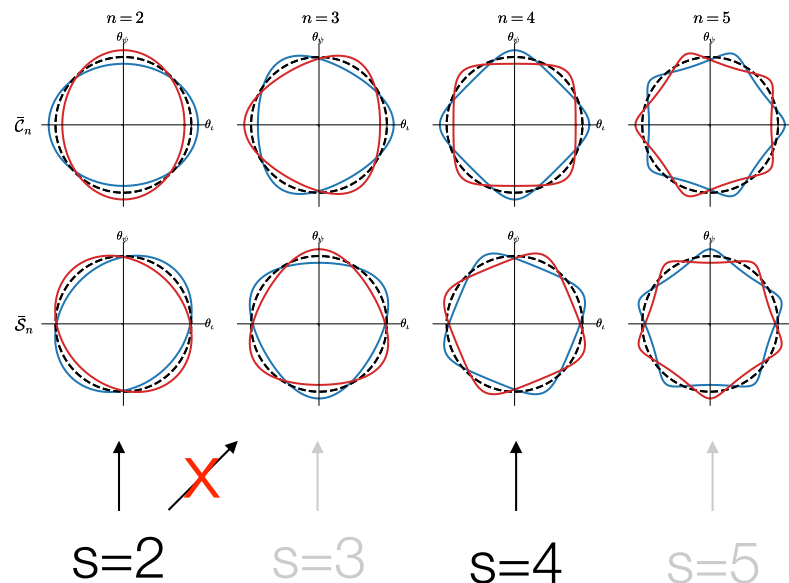
$$B_{\Phi}(\mathbf{k}_S, \mathbf{k}_L) \propto \mathcal{P}_s(\hat{\mathbf{k}}_L \cdot \hat{\mathbf{k}}_S)$$

Arkani-Hamed & Maldacena (15)

Kogai, Akitsu, Schmidt, & Y.U. (20)



decomposition



Imprints of PNG from spin s

LSST like Forecast

s=4 “massless”  $O(10)$ , massive  $O(10^5)$  needs deeper survey

What about w/o de Sitter symmetry?



~~w/dS~~, studies of dS are still useful??

# g $\delta$ N formalism

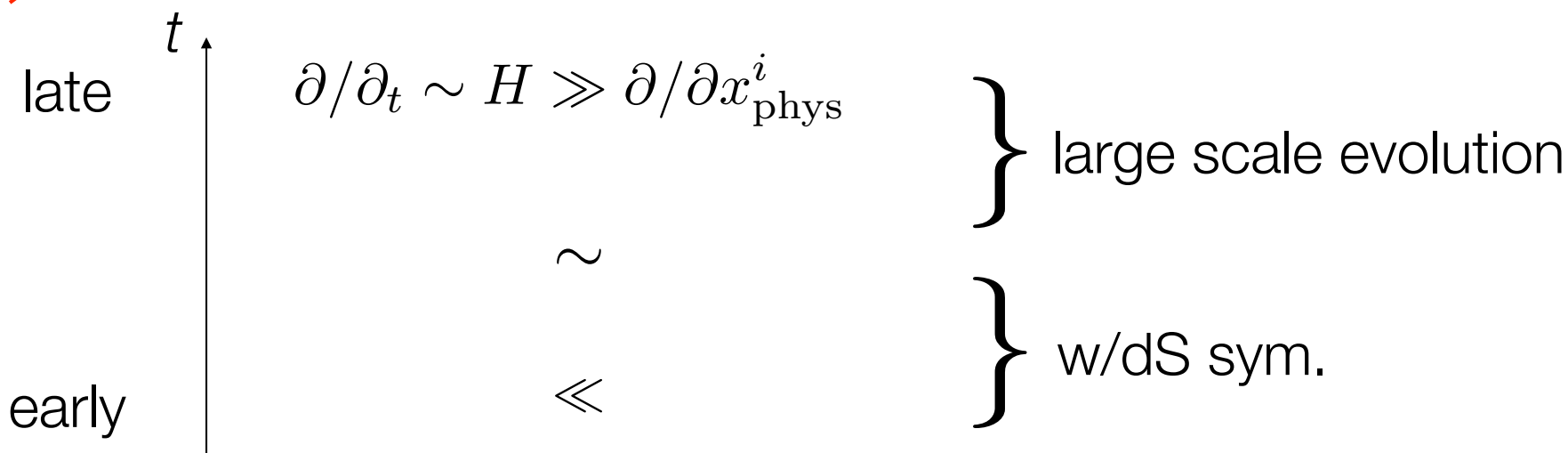
T.Tanaka @ Y.U. 2101.05707

w/Takahiro Tanaka (Kyoto, YITP)

# Large scale evolution

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~~dS~~, global rotation, etc..., often become significant at large scales.



## Large scale evolution for scalar field system

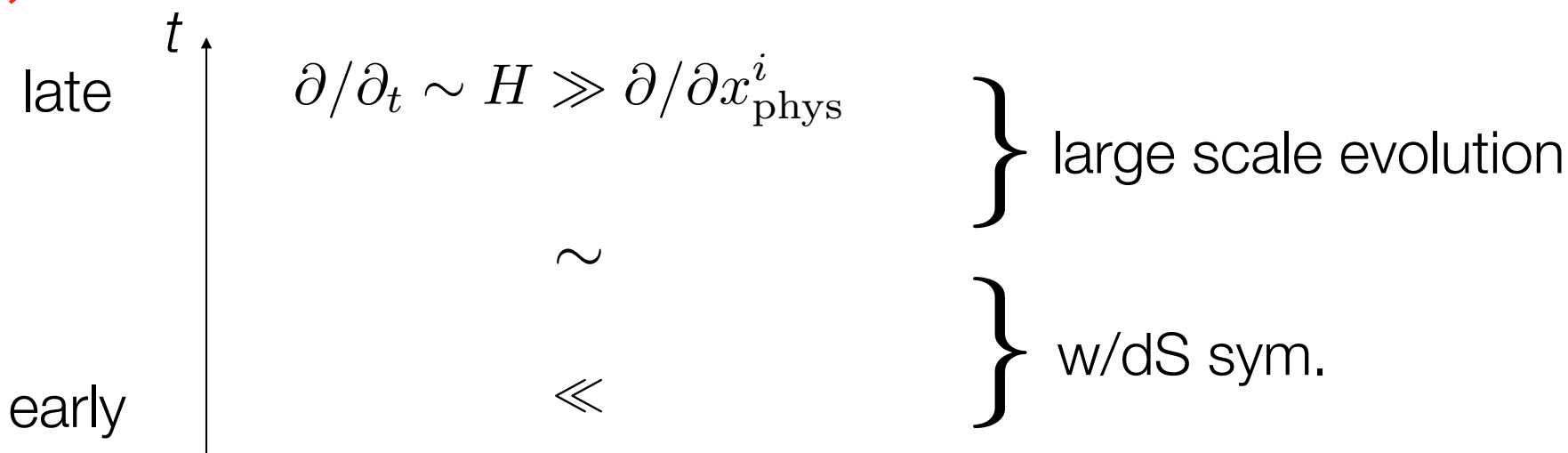
Gradient expansion (GE) → delta N formalism

Salopek & Bond (90)  
Shibata & Sasaki (90),  
Deruelle & Langlois (94),...

Starobinsky (82, 85),  
Sasaki & Stewart (95),  
Sasaki & Tanaka (98),  
Lyth, Malik, & Sasaki (04),...

# Large scale evolution

~~dS~~, global ~~rotation~~, etc..., often become significant at large scales.



Large scale evolution ~~for scalar field system~~

Gradient expansion (GE)  $\rightarrow$   ~~$\delta N$  formalism~~

g(eneralized) $\delta N$   
formalism

Salopek & Bond (90)  
Shibata & Sasaki (90),  
Deruelle & Langlois (94),...

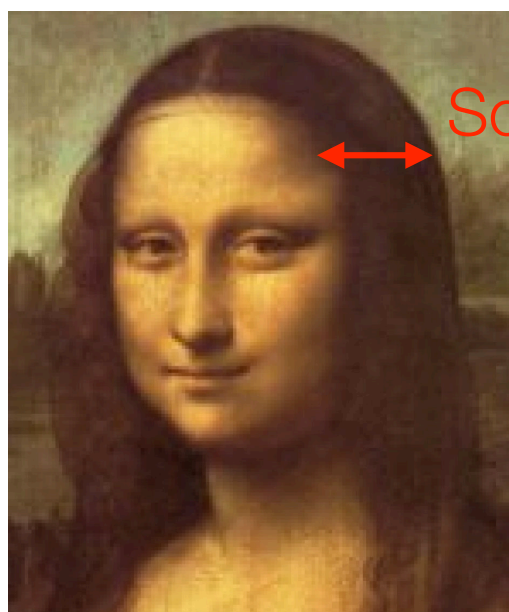
Starobinsky (82, 85),  
Sasaki & Stewart (95),  
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Lyth, Malik, & Sasaki (04),...

Tanaka & Y.U. (21)

# Separate Universe approach ~ Mosaic art

Basic assumption in gradient expansion

size of mono color stone  
= smoothing scale  $\lambda_s$



Scale of interest  
 $\lambda (\gg \lambda_s)$



**Fine-grained** view

**Coarse-grained** view

Let her get aged.



(☆) Her history in detailed view = Her history in coarse-grained view

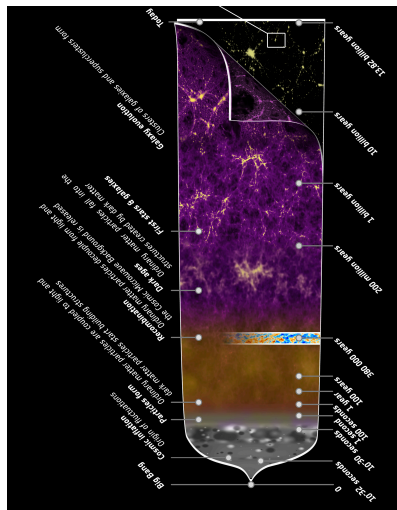
# Separate universe

Salopek & Bond (90)

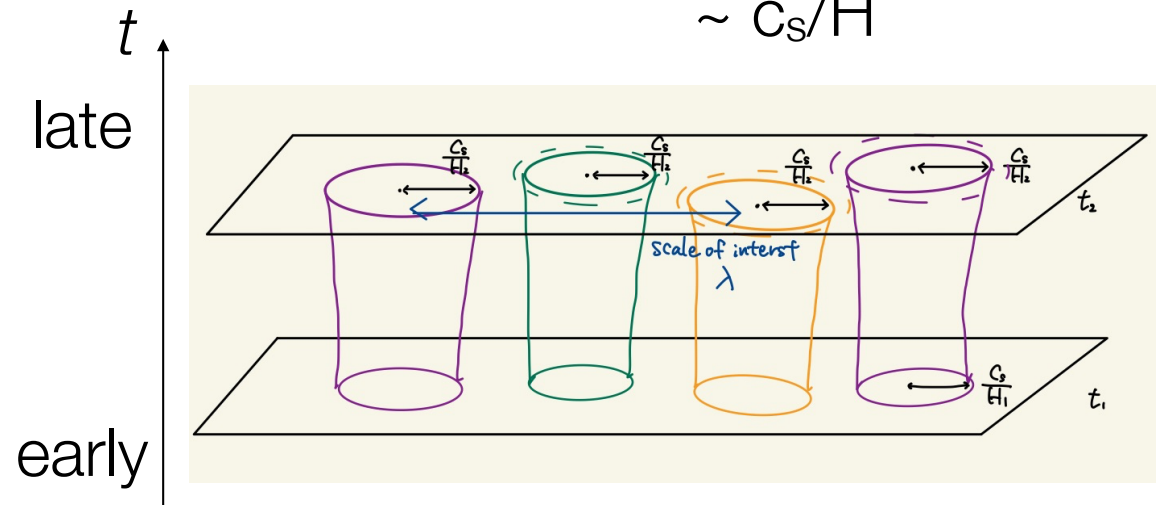
(☆) Evolution of inhomogeneous Universe  
 = Evolution of glued numerous homogeneous universes

Scale of interest  $\lambda \gg$  Smoothing scale  $\lambda_s >$  Size of casual patch

$$\sim c_s/H$$



=  
 (☆)



**Fine-grained** view

Solving PDEs

**Coarse-grained** view

Solving ODEs

(Inhomogeneity: Different ICs)





# gδN formalism

Tanaka & Y.U. (21)

a local theory w/gravity, spatial Diff

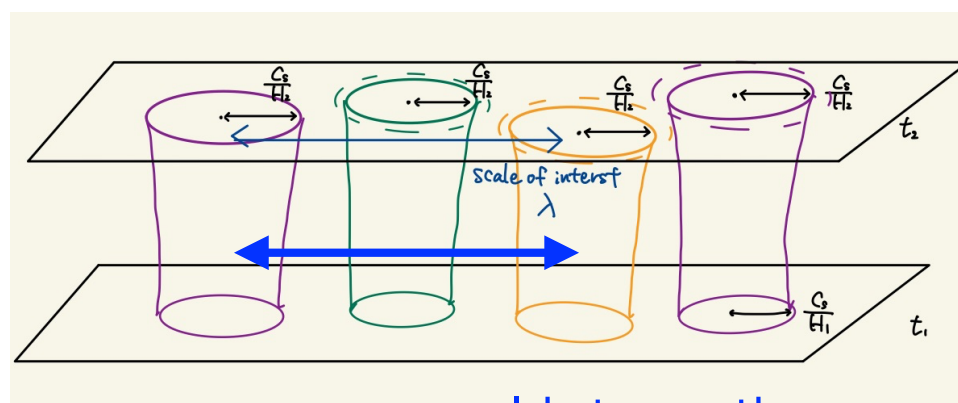
fields in  $\mathcal{L}$

→  
solving constraints (C)

physical fields  $\{\phi_{phys}\}$

non-local

e.g., MC  $\mathcal{H}_i \equiv \frac{\partial(N\sqrt{g}\mathcal{L})}{\partial N^i} = 0$        $\partial_i(\dots) + (\dots)\partial_i(\dots) = 0$



acausal interaction



# gδN formalism

Tanaka & Y.U. (21)

a local theory w/gravity, spatial Diff

fields in  $\mathcal{L}$   $\longrightarrow$   $\{\varphi_{phys}, \varphi_{auxi}\}$   $\longrightarrow$  physical fields  $\{\varphi_{phys}\}$

solving local C

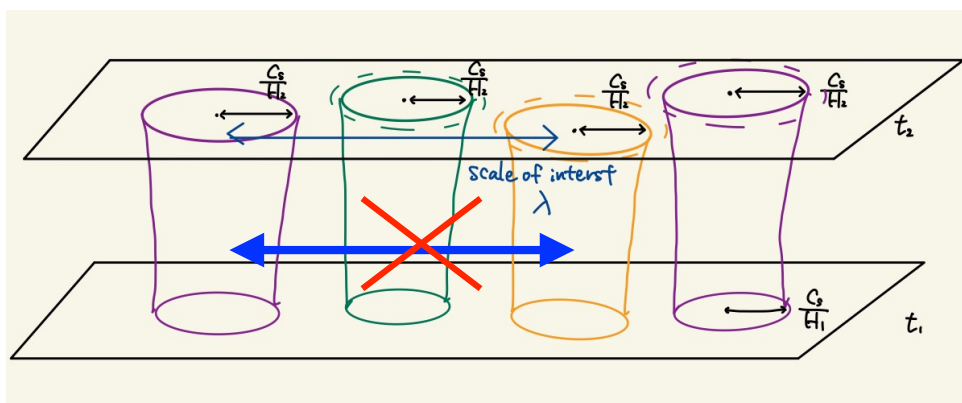
solving nonlocal C

non-local

t

late

early



(☆) ✓

IC by  $\{\varphi_{phys}, \varphi_{auxi}\}$   
 at  $\partial/\partial t \sim H \sim \partial/\partial x^i_{phys}$

density pert.  $\zeta(t_{final}) = \zeta_{hom}(t_{final}; \{\varphi_{phys}, \varphi_{auxi}\})$

incl. spinning particles

GW  $\gamma_{ij}(t_{final}) = \gamma_{ij\,hom}(t_{final}; \{\varphi_{phys}, \varphi_{auxi}\})$

\* Initial correlators should satisfy nonlocal C.

# Infrared universality

N.B. non-local terms from MC decay with  $1/V_{\text{phys}}$  for scalar field system

Sugiyama, Komatsu, & Futamase (12)

Garriga, Y.U., & Vernizzi (16)



Noether charge for spatial Diff. invariance. Tanaka & Y.U. (21)

if there is no other non-local  $C$ , approximate locality for  $\{\phi_{\text{phys}}\}$

Similar structure for cosmological correlators yet w/nontrivial cond.

Tanaka & Y.U. (17, in prep)

\* Condition for  $\exists$ WAM

Weinberg (02)

Strominger+ (13,14, ...)

Recall Hotta-san's talk

MEMORY EFFECT

Christodoulou (1991)

FOURIER TRANSFORM VACUUM TRANSITION

WARD IDENTITY

SOFT THEOREM

ASYMPTOTIC SYMMETRY

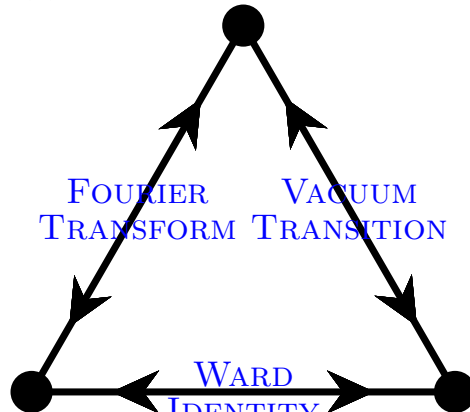
Bondi, Metzner, Sachs (1962)

Cancellation of IR divergence

Faddeev-Kulish (1970)

Bloch & Nordsieck (1937)

Weinberg (1965)



# Summary: Comprehensive studies of inflation

