

The equation of state after Inflation



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MAX-PLANCK-GESELLSCHAFT

Outline

1.) The equation of state after inflation KL and M. Amin (2016,17,19)

2.) Oscillons and baryogenesis

KL and M. Amin (2014)

3.) Gauge fields, inflation and reheating кL and M. Amin (2016) **к**L, A. Maleknejad and E. Komatsu (2018) L. Mirzagoli, A. Maleknejad and **к**L (2019)

4.) Cosmology of fine-tuning (in particle physics)

M. Amin, J. Fan, KL and M. Reece (2019)

The equation of state after inflation

KL and M. Amin, arXiv:1902.06736 KL and M. Amin, PRD 97 023533 (2017) KL and M. Amin, PRL 119 061301 (2016)

 $w = \frac{\text{pressure}}{\text{energy density}}$

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \Big]$$

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assumption: self-couplings dominate over others

at sufficiently late times:

$$w = \begin{cases} 0 & \text{if } n = 1\\ 1/3 & \text{if } n > 1 \end{cases}$$

(even without couplings to other fields!)



Inflaton dynamics



oscillatory phase

Inflaton (homogeneous) dynamics ϕ $V(\phi)$ inflation $\phi \times a^{3/(n+1)}$ 0.0 inflation ends: 50 50 oscillatory phase x 100 $\bar{\phi}(t) \propto a(t)^{-3/(n+1)}$

Inflaton (homogeneous) dynamics ϕ $V(\phi)$ inflation $w \equiv \frac{p}{\rho} = \frac{\phi^2/2 - (\nabla\phi)^2/6 - V(\phi)}{\dot{\phi}^2/2 + (\nabla\phi)^2/2 + V(\phi)}$ inflation ends: oscillatory phase $\bar{\phi}(t) \propto a(t)^{-3/(n+1)}$





- parametric resonance of $\delta \phi(t, \mathbf{x})$
- $\bar{\phi}$ fragments

KL and M. Amin (2016) KL and M. Amin (2017)

Equation of state



(e-folds after inflation) ΔN

$$\Delta N \equiv \int_{a_{\rm end}}^{a} d\ln a$$

Equation of state

 $-M \ll m_{\rm pl}$ (efficient resonance)

 $-M \sim m_{\rm pl}$ (inefficient resonance) \mathfrak{M} $2.5 \quad 3.1$ 10.71213.23.6 $\left(\right)$ ()(equation of state) 1.01.0 n-1 $w_{\rm hom}$ n + 1 $w_{\rm rad} =$ 0.50.50.0 0.0 -0.5-0.5n = 1n > 1-1.0-1.02 1.51 2 3 1 ()

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 $m^2 \equiv V'(\bar{\phi}_{\rm osc})/\bar{\phi}_{\rm osc}$

Non-perturbative decay (parametric self-resonance)



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inflaton field wavenumber



 $\delta\phi_k \propto \exp(\pm\mu_k t)$

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inflaton field wavenumber

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inflaton field wavenumber

Non-perturbative decay (parametric self-resonance)



inflaton field wavenumber

 $V(\bar{\phi}) \propto |\bar{\phi}|^{2n}$ $\bar{\phi} \sim M \sim m_{\rm pl}$

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n=3Power spectrum: inflaton vev amplitude 2.0 $\Re(\mu_k)$ $\left< \delta \phi(x)^2 \right> \equiv \int \mathcal{P}_{\delta \phi} d \ln k$ रु mconst 0.036 10^{-1} 1.5 ${\cal P}_{\delta\phi}/ar{\phi}^2_{
m osc}$ 10^{-3} $\phi_{\rm osc}/M$ 10^{-5} 1.0 10^{-7} 0.018 10^{-9} 0.5 10^{-11} 0.000 10^{-13} 0.00.51.01.52.0 10^{-3} 10^{-2} 0.0 10^{-1} k/m_0 $\kappa = k/am$ inflaton field wavenumber comoving wavenumber

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• $\bar{\phi}$ fragments quickly

 $M \sim m_{\rm pl}$

- slow production of $\delta \phi(t,\mathbf{x})$

• $\bar{\phi}$ fragments gradually

$$\Delta N_{\rm fr} \approx \frac{n+1}{3} \ln \left(10^3 \frac{M}{m_{\rm pl}} \right)$$

at sufficiently late times:

 ϕ virialized + turbulent $\rightarrow w = \frac{1}{3}$

Spectral index: $n_{\rm s}$

Tensor-to-scalar ratio: r

Spectral index:
$$n_{
m s}=n_{
m s}(M,n,N_{*})$$

Tensor-to-scalar ratio: $r=r(M,n,N_{*})$

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$$50 \le N_* \le 60$$

Planck Collaboration (2015)

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$$N_* = 66.89 - \frac{1}{12} \ln g_{\rm th} + \frac{1}{4} \ln \frac{V_*^4}{m_{\rm pl}^4 \rho_{\rm end}} - \ln \frac{k_*}{a_0 H_0} + \frac{3\bar{w}_{\rm int} - 1}{4} \Delta N_{\rm rad}$$

$$\Delta N_{\rm rad} \equiv \int_{a_{\rm end}}^{a_{\rm rad}} d\ln a$$

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reheating

$$\Delta N_{\rm rad} \equiv \int_{a_{\rm end}}^{a_{\rm rad}} d\ln a$$

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$$\Delta N_{\rm fr} = \begin{cases} 1 & \text{if} & M \ll m_{\rm pl} \\ \frac{n+1}{3} \ln \left(10\frac{\kappa}{\Delta\kappa}\frac{M}{m_{\rm pl}}\right) & \text{if} & M \sim m_{\rm pl} \end{cases}$$

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$$w_{\rm int}(\Delta N) = \begin{cases} \frac{n-1}{n+1} & \text{if} & 0 < \Delta N < \Delta N_{\rm rad} \\ \frac{1}{3} & \text{if} & \Delta N > \Delta N_{\rm rad} \end{cases}$$

$$\Delta N_{\rm rad} \equiv \int_{a_{\rm end}}^{a_{\rm rad}} d\ln a$$

$$V(\phi) \propto \tanh^{2n} \left(\frac{|\phi|}{M}\right)$$
 α -attractors (T-models)



Kallosh and Linde (2013) Carrasco, Kallosh and Linde (2015)









spectral index

Planck Collaboration (2015)



tensor-to-scalar ratio

$$V(\phi) \propto \left| 1 - e^{-\phi/M} \right|^{2n}$$
 α -attractors (E-models)











Planck Collaboration (2015)



$$V(\phi) \propto \left[1 + \left|\frac{\phi}{M}\right|^{2n}\right]^{\frac{q}{2n}} - 1$$
 Monodromy



Silverstein and Westphal (2008) McAllister, et al. (2014)

$$V(\phi) \propto \left[1 + \left|\frac{\phi}{M}\right|^{2n}\right]^{\frac{1}{4n}} - 1$$
 Monodromy
 $q = 1/2$



Silverstein and Westphal (2008) McAllister, et al. (2014)










$$V(\phi) \propto \left[1 + \left|\frac{\phi}{M}\right|^{2n}\right]^{\frac{1}{2n}} - 1$$
 Monodromy
 $q = 1$



Silverstein and Westphal (2008) McAllister, et al. (2014)











Equation of state

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KL and M. Amin (2016)

Non-perturbative decay (parametric self-resonance)

n = 1



inflaton field wavenumber

See also Amin, Hertzberg + (2011), Fukunaga, Kitajima and Urakawa (2019) $m^2 \equiv V'(ar{\phi}_{
m osc})/ar{\phi}_{
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Matter domination? n = 1 KLa

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n = 1





• $\overline{\phi}$ forms oscillons (stable) See also Amin, Hertzberg et al (2011), Urakawa's talk, Kitajima's talk, Torrenti's talk matter-like eos: w = 0

couplings to other fields?

See also Hertzberg (2010), Adshead et al (2015)

 $M \sim m_{\rm pl}$

- $\delta \phi(t, \mathbf{x})$ production shut off
- $\bar{\phi}_{\rm osc}(t) = {\rm pressureless \ dust}$



observationally favored $V(\phi)$



simple result

$$w = \begin{cases} 0 & \text{if } n = 1\\ 1/3 & \text{if } n > 1 \end{cases}$$

(even without couplings to other fields!)

Summary

(at sufficiently late times)

KL and M. Amin, PRD 97 023533 (2017) KL and M. Amin, PRL 119 061301 (2016)



Reduction in theoretical uncertainty of inflationary models

• stochastic GWs from fragmentation

KL and M. Amin (2019)

• stochastic GWs from fragmentation

KL and M. Amin (2019)

Oscillons



• stochastic GWs from fragmentation Transients

KL and M. Amin (2019)

 10^{-9} 10^{-10} $\Omega_{\rm GW,0} h_{100}^2$ 10^{-11} 10^{-12} 10^{-13} 10^{-14} 10^{8} 10^{9} 10^{10} f_0/Hz



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stochastic GWs from fragmentation

KL and M. Amin (2019)

• dark energy

P. Agrawal, L. Randall, et al (2019)

- stochastic GWs from fragmentation KL and M. Amin (2019)
- dark energy

P. Agrawal, L. Randall, et al (2019)

• dark matter

- stochastic GWs from fragmentation KL and M. Amin (2019)
- dark energy

P. Agrawal, L. Randall, et al (2019)

- dark matter
- matter-antimatter asymmetry...

KL and M. Amin (2014)
Oscillons and matter-antimatter asymmetry

KL and M. Amin, PRD 90, 083528 (2014)



The model

A variation of the Affleck-Dine Mechanism (1985) Hertzberg & Karouby (2013, 2014)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} + |\partial\phi|^2 - V(\phi, \phi^*) \right]$$

$$V(\phi,\phi^*) = V_{\rm s}(|\phi|) + V_{\rm b}(\phi,\phi^*)$$

respects U(1) symmetry

breaks U(1) symmetry

responsible for inflation

responsible for generating inflaton/antiinflaton asymmetry

$$V_{\rm b}(\phi,\phi^*) = \frac{c_n}{n}(\phi^n + \phi^{*n})$$

"small" symmetry breaking

- technically natural
- small during inflation
- small long after inflation

Inflaton asymmetry – baryon asymmetry

$$\Delta N_{\phi} = N_{\phi} - N_{\bar{\phi}} = i \int d^3x a^3 (\phi^* \dot{\phi} - \dot{\phi}^* \phi)$$

inflaton number (not conserved!)

• generated at end of inflation

 $\phi \rightarrow b$

decay

 $N_b - N_{\bar{b}} = b_\phi (N_\phi - N_{\bar{\phi}})$

baryon number

Inflaton dynamics $|\phi|$ $(|\phi|)$ inflation I I .

inflation ends: oscillatory phase



Inflaton (actual) dynamics



inflation ends: oscillatory phase

- parametric resonance of $\delta\phi$
- ϕ fragments
- forms oscillons



Inflaton asymmetry



Asymmetry-fragmentation

Non-linear structures DO AFFECT inflaton/anti-inflaton ASYMMETRY!



Where is the inflaton/anti-inflaton ASYMMETRY?



More than 70% of ASYMMETRY locked in oscillons!





 $c_3 \ll 1, M \ll m_{\rm Pl}$

Inflaton to baryons (qualitative)



caveats: uncertainty here!! particle physics details, inhomogeneous decay...

Inflaton to baryons (qualitative)



sample numbers: $A_{\phi} \sim 10^{-4}, T_{\rm reh} \sim 10^{7} \,{\rm GeV}, \, m \sim 10^{14} \,{\rm GeV}$

caveats: uncertainty here!! particle physics details, inhomogeneous decay...



KL and M. Amin, PRD 90, 083528 (2014)



very different dynamics from homogeneous case!

matter-antimatter asymmetry $\eta \approx 6 \times 10^{-10}$

KL and M. Amin, JCAP 1606 032 (2016)

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - (D_{\mu}\phi)^{\dagger} D^{\mu}\phi - V(|\phi|) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \Big]$$

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wavenumber

inflaton amplitude

wavenumber

KL and M. Amin, JCAP 1606 032 (2016), see also Yamaguchi's and Torrenti's talks

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - (D_{\mu}\phi)^{\dagger} D^{\mu}\phi - V(|\phi|) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \Big]$$



(in progress)



with Eiichiro Komatsu and Mustafa Amin

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

Adshead and Wyman (2012)

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Non-trivial vevs:

Adshead and Wyman (2012)

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Non-trivial vevs:

 $\overline{\phi}(t)$

Adshead and Wyman (2012)

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Non-trivial vevs:

$$\overline{\phi}(t) \\ \overline{A_i^b}(t) = a(t)Q(t)\delta_i^b$$

See also Maleknejad (2011)

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \Big]$$

Non-trivial vevs:
$$\overline{\phi}(t)$$
$$\overline{A^b_i}(t) = a(t)Q(t)\delta^b_i$$
 See also Maleknejad (2012)
Drive inflation

$$\begin{split} S &= \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \Big] \\ & \text{Non-trivial vevs:} \\ & \overline{\phi}(t) \\ & \overline{A^b_i}(t) = a(t) Q(t) \delta^b_i \\ & \text{See also Maleknejad (2013)} \\ & \text{Drive inflation} \\ & (+\text{extensions}) \\ & \text{Adshead et al (2016, 2017)} \end{split}$$

$$\begin{split} S &= \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \Big] \\ & \text{Non-trivial vevs:} \\ & \overline{\phi}(t) \\ & \overline{A^b_i}(t) = a(t) Q(t) \delta^b_i \\ & \text{See also Maleknejad (2011)} \\ & \text{Drive inflation} \\ & \text{(+extensions)} \\ & \text{Adshead et al (2016, 2017)} \end{split}$$

Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

Non-trivial vevs:

 $\frac{\phi(t)}{\overline{A_i^b}(t) = a(t)Q(t)\delta_i^b}$

See also Maleknejad (2011)

Drive inflation (+extensions) Adshead et al (2016, 2017) Spectator sector Dimastrogiovanni et al (2016) Maleknejad (2016, 2018) Adshead et al (2017) Agrawal, Komatsu et al (2017, 2018) Soda and Urakawa (2017) Kitajima et al (2018)

Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \Big]$$

Non-trivial vevs:

$$\frac{\overline{\phi}(t)}{\overline{A_i^b}(t) = a(t)Q(t)\delta_i^b}$$

Linear coupling between Gauge Fields and GWs!

Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

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Linear coupling between Gauge Fields and GWs! δA_i^b

Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

Non-trivial vevs:

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Linear coupling between Gauge Fields and GWs! δA_i^b and

Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

Non-trivial vevs:

$$\overline{\phi}(t) \\ \overline{A_i^b}(t) = a(t)Q(t)\delta_i^b$$

Linear coupling between Gauge Fields and GWs! δA_i^b and h_{ij}

Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \Big]$$

Non-trivial vevs:

$$\frac{\overline{\phi}(t)}{\overline{A_i^b}(t) = a(t)Q(t)\delta_i^b}$$

Linear coupling between Gauge Fields and GWs!



Figure from Dimastrogiovanni et al. arXiv:1608.04216

Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

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Linear coupling between Gauge Fields and GWs!



Adshead and Wyman (2012)

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Adshead and Wyman (2012)

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^a_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

Non-trivial vevs:

$$\overline{A_i^b}(t) = a(t)Q(t)\delta_i^b$$

Adshead and Wyman (2012)

Backreaction?

 $\overline{A_i^b}(t) = a(t)Q(t)\delta_i^b$
Adshead and Wyman (2012)

$$S = \int d^{4}x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - V(\phi) - \frac{1}{4} F^{a}_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^{a}_{\mu\nu} \tilde{F}^{a\mu\nu} \Big]$$

Non-trivial vevs:
$$\overline{\phi}(t)$$
$$Schwinger effect?
$$\overline{A^{b}_{i}(t)} = a(t)Q(t)\delta^{b}_{i}$$
Backreaction?$$

KL, Maleknejad and Komatsu (2018)

$$S_1 = \int d^4x \sqrt{-g} \Big[- (D_\mu \varphi)^{\dagger} D^\mu \varphi - V(|\varphi|) \Big]$$

Adshead and Wyman (2012)

 $\varphi =$

 $arphi^{ extsf{1}} \ arphi^{2}$

$$S_1 = \int d^4x \sqrt{-g} \Big[- (D_\mu \varphi)^{\dagger} D^\mu \varphi - V(|\varphi|) \Big]$$

Adshead and Wyman (2012)

KL, Maleknejad and Komatsu (2018)

 $\varphi = \left(\begin{array}{c} \varphi^1 \\ \varphi^2 \end{array} \right)$

$$S_1 = \int d^4x \sqrt{-g} \Big[-(D_\mu \varphi)^{\dagger} D^\mu \varphi - V(|\varphi|) \Big]$$

Mirzagoli, Maleknejad and KL (2019)

$$S_2 = \int d^4x \sqrt{-g} \left[i\overline{\psi}\gamma^{\mu}D_{\mu}\psi - m\overline{\psi}\psi - \lambda \frac{\partial_{\mu}\phi}{f}J^{5\mu} \right]$$

Adshead and Wyman (2012)

KL, Maleknejad and Komatsu (2018)

 $\varphi = \left(\begin{array}{c} \varphi^1 \\ \varphi^2 \end{array} \right)$

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Adshead and Wyman (2012)

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See also Domcke et al (2018), Adshead et al (2015, 2018, 2019)

Adshead and Wyman (2012)

$$S = \int d^{4}x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - V(\phi) - \frac{1}{4} F^{a}_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^{a}_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

Non-trivial vevs:

$$\begin{array}{c} \hline \phi(t) \\ A^{b}_{i}(t) = a(t)Q(t)\delta^{b}_{i} \end{array}$$
Schwinger effect?
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Mirzagoli, Maleknejad and KL (2019)

$$\psi = \begin{pmatrix} \psi^{1} \\ \psi^{2} \end{pmatrix}$$

Gauge field coupling strength

Adshead and Wyman (2012)

$$S = \int d^{4}x \sqrt{-g} \left[\frac{1}{16\pi G} \mathcal{R} - \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - V(\phi) - \frac{1}{4} F^{a}_{\mu\nu} F^{a\mu\nu} - \frac{\phi}{4f} F^{a}_{\mu\nu} \tilde{F}^{a\mu\nu} \right]$$

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Mirzagoli, Maleknejad and KL (2019)

$$\psi = \begin{pmatrix} \psi^{1} \\ \psi^{2} \end{pmatrix}$$

Cosmology of a Fine-Tuned Higgs

Mustafa Amin, JiJi Fan, KL and Matthew Reece, PRD 99 035008 (2019)



Conclusions

Reheating at the end of inflation:

• very rich dynamics

Future plans:

- more realistic models (e.g. gauge fields)
- observational signatures

 expansion history, relics, GWs





Strings

$$S = \int d^4x \sqrt{-g} \Big[\frac{1}{16\pi G} \mathcal{R} - (D_{\mu}\phi)^{\dagger} D^{\mu}\phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \Big]$$

