

2 day mini-workshop:

Axion Cosmology

Yukawa Institute for Theoretical Physics, Kyoto University

May 14 - 15, 2019

**Exploring the string axiverse and parity violation in gravity
with gravitational waves**

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Collaborator : Jiro Soda

Kobe University

Menu

- 1. Traces of the string axiverse**
- 2. Axion dark matter**
- 3. Dynamical Chern-Simons gravity**
- 4. Parametric Resonance**
- 5. Estimations**
- 6. Conclusion**

Menu

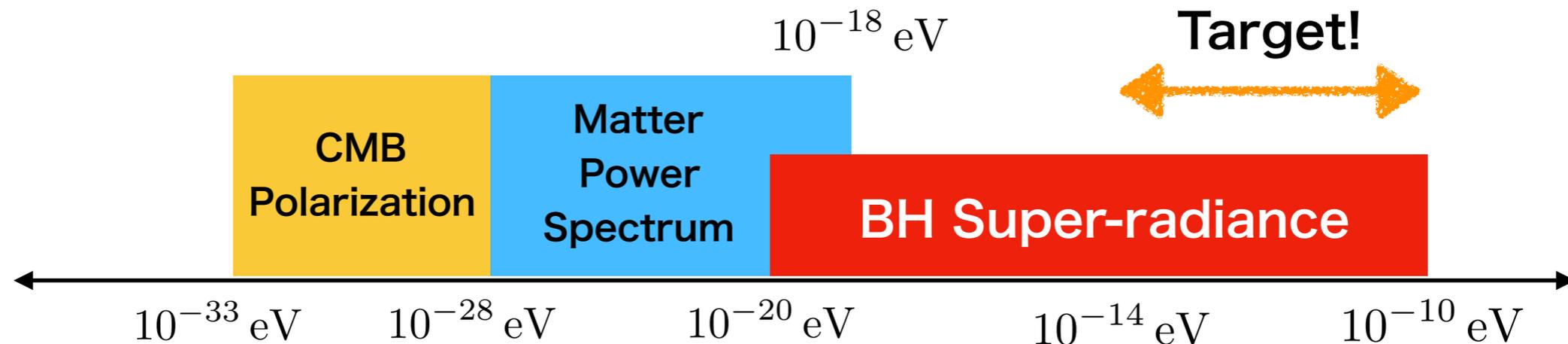
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Traces of String Axiverse

- String theory gives a lot of massive pseudo-scalar fields (Axion).

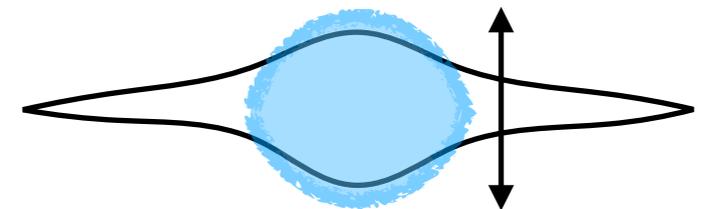
A. Arvanitaki, et al (2010),
P. Svrcek and E. Witten (2006)

- Their mass is $10^{-33} \sim 10^{-10}$ eV. Its range is the very wide.



- It is indicated that the axion can behave as the cold dark matter.

W. Hu, R. Barkana, and A. Gruzinov (2000)

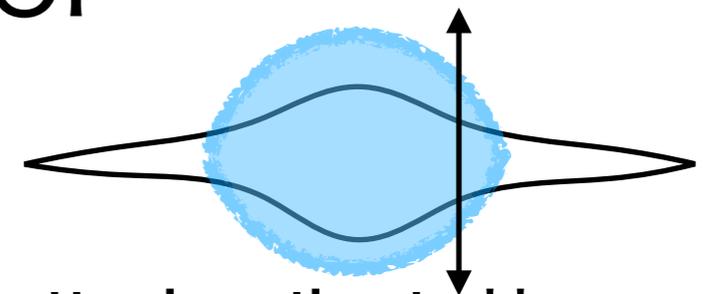


- The string axiverse gives the parity-violated interactions as follows, $\Phi \tilde{F}F$, $\Phi \tilde{R}R$. In this talk, we concentrate on the modified gravity.

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Axion Dark Matter



- The cold dark matter is the dust which is pressureless.
- The occupation number of the state of the axion dark matter is estimated by

$$\frac{N}{\Delta x^3 \Delta p^3} \sim \frac{n}{p^3} = \frac{\rho_{\text{DM}}}{m p^3} \sim 10^{43} \left(\frac{\rho_{\text{DM}}}{0.3 \text{ GeV/cm}^3} \right) \left(\frac{10^{-10} \text{ eV}}{m} \right)^4.$$

A. Khmelnitsky & V. Rubakov, 2014

- The oscillation of the axion as a classical field is given by $\Phi(t) \simeq \Phi_0 \cos(mt)$.
 $v \sim 10^{-3}, E \simeq m + \frac{1}{2}mv^2 \sim m$

→ Energy density : $T_{00} = \rho_{\text{DM}} = \frac{1}{2}\dot{\Phi}^2 + \frac{1}{2}m^2\Phi^2 \sim \frac{1}{2}m^2\Phi_0^2$

$$\Phi_0 \simeq 2.1 \times 10^7 \text{ eV} \left(\frac{10^{-10} \text{ eV}}{m} \right) \sqrt{\frac{\rho}{0.3 \text{ GeV/cm}^3}}$$

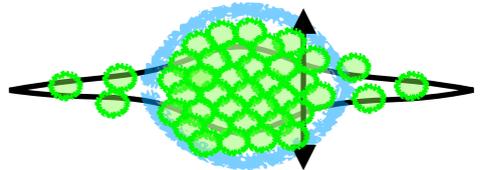
Pressure : $p_{\text{DM}} = \frac{1}{2}\dot{\Phi}^2 - \frac{1}{2}m^2\Phi^2 \sim -\frac{1}{2}m^2\Phi_0^2 \cos 2mt$

- The period of the oscillation of the axion can be estimated as $1/m$. Then, by the average under the long time scale $T \gg 1/m$, the pressure of the oscillating axion can be regarded as zero. Then, the oscillating axion behaves as the CDM.

Size of Axion Dark Matter

- In space, the matters are collected by the gravitational forces.
 - It appears the equilibrium between the characteristic pressure of matters and the gravitational force induced by matters themselves.
 - This equilibrium is expressed by the size of the halo of matters.

→ Jeans length of the axion dark matter is given by

$$r_J = 4.3 \times 10^{-3} \text{pc} \times \left(\frac{0.3 \text{ GeV/cm}^3}{\rho} \right)^{1/4} \left(\frac{10^{-10} \text{ eV}}{m} \right)^{1/2}$$

 A diagram illustrating a dark matter halo. It shows a central cluster of green particles, surrounded by a blue shaded region representing the halo. A vertical double-headed arrow indicates the size of the halo, and a horizontal double-headed arrow indicates the extent of the halo along the x-axis.

- Jeans length is the maximum size which the matters can exist stably in space.
 - In this talk, we assume the axion dark matters have the enough energy density and behave as the classical oscillating field.
 - In low energy scale, the potential of the axion can be approximated by

$$V \sim \frac{1}{2} m^2 \Phi^2.$$

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Dynamical Chern-Simon Gravity^{04/14}

- The action is given by

$$S = \int_{\mathcal{V}} dx^4 \sqrt{-g} \left[\kappa R + \frac{1}{4} \alpha \Phi \tilde{R} R - \frac{1}{2} (\nabla_\mu \Phi) (\nabla^\mu \Phi) - V(\Phi) \right]$$

$$\alpha = \sqrt{\frac{\kappa}{2}} \ell^2 \quad \tilde{R} R \equiv \frac{1}{2} \underbrace{\epsilon^{\gamma\delta\rho\sigma} R^\alpha_{\beta\rho\sigma} R^\beta_{\alpha\gamma\delta}}_{\equiv \tilde{R}^\alpha_{\beta\gamma\delta}}$$

- Dynamical or non-dynamical

→ This is switched by the existence of the dynamical term of the axion field.

$$i\epsilon_{ilm} n_l e_{mj}^{R/L}(\mathbf{n}) = \pm e_{ij}^{R/L}(\mathbf{n})$$

- Because the axion field is the pseudo scalar, this interaction term induces the circular polarized gravitational waves.

- The upper bound of the coupling constant ℓ is given as follows,

$$\ell \sim 10^8 \text{ km}$$

in non dynamical theory by Y. Ali-Haimoud & Y. Chen (2011).

in dynamical theory by Y. Nakamura et. al. (2018).

- The ghost modes exist in this theory. We assume this theory as the effective theory of gravity.

Equations of Motion of dCS Gravity

- The action of the dynamical Chern-Simons gravity is

$$S = \kappa \int_{\mathcal{V}} dx^4 \sqrt{-g} R + \frac{1}{4} \alpha \int_{\mathcal{V}} dx^4 \sqrt{-g} \Phi \tilde{R} R - \frac{1}{2} \int_{\mathcal{V}} dx^4 \sqrt{-g} [g^{\mu\nu} (\nabla_\mu \Phi)(\nabla_\nu \Phi) + 2V(\Phi)]$$

- The equations of motion are derived by

$$G_{\mu\nu} + \frac{\alpha}{\kappa} C_{\mu\nu} = \frac{1}{2\kappa} T_{\mu\nu} \quad \text{for the gravitational field and}$$

$$C^{\mu\nu} \equiv (\nabla_\alpha \Phi) \epsilon^{\alpha\beta\gamma(\mu} \nabla_\gamma R^{\nu)}_\beta + (\nabla_\alpha \nabla_\beta \Phi) \tilde{R}^{\beta(\mu\nu)\alpha}$$

$$\nabla_\mu \nabla^\mu \Phi - \frac{dV(\Phi)}{d\Phi} = -\frac{\alpha}{4} \tilde{R} R \quad \text{for the axion field.}$$

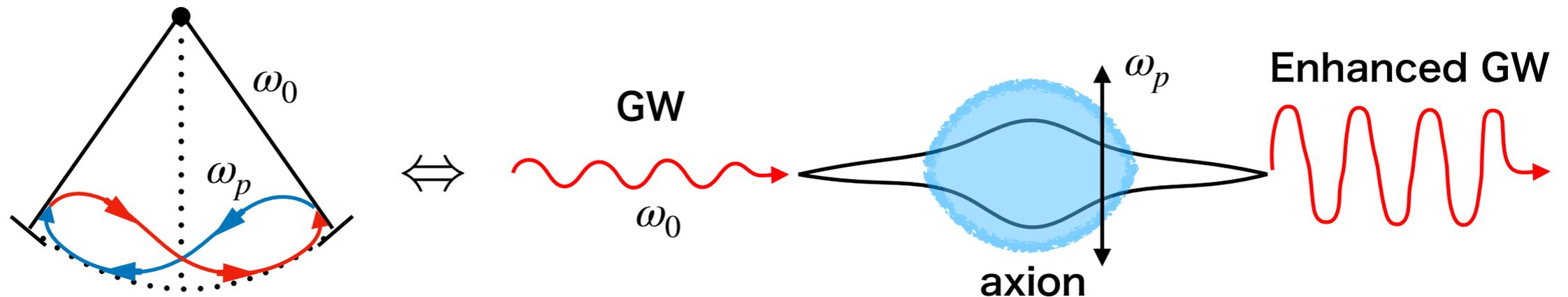
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Parametric Resonance

- The Swings

→ The parameters vary at roughly twice the natural frequency of the arms of the swing, the amplitude of it will grow.



- The parametric resonance is induced by Mathieu equation.

$$\ddot{x}(t) + \omega_0^2 (1 - f_0 \sin(\omega_p t)) x(t) = 0$$

- The properties of the Mathieu equation are noted.

→ The (first) resonance frequency is given by

$$\omega_p \simeq 2\omega_0$$

$$h \propto e^{\Gamma \eta}$$

→ The growth rate of the amplitude is given by

$$\Gamma_{\max} = \frac{1}{4} |f_0| \omega_0$$

Gravitational Wave in dCS Gravity^{07/14}

- We choose the ansatz of the spacetime below,

$$ds^2 \simeq a(\eta)^2 (-d\eta^2 + \delta_{ij} dx^i dx^j + h_{ij} dx^i dx^j)$$

- For the interest in the circular polarization, we take the circular polarization basis which are defined by

$$h_{ij}^{\text{TT}} = \sum_A h_A e_{ij}^A$$

$$e_{ij}^{\text{R}}(\mathbf{n}) = \frac{1}{\sqrt{2}} (e_{ij}^+(\mathbf{n}) + ie_{ij}^\times(\mathbf{n})) \quad h_{\text{R}} = \frac{1}{\sqrt{2}} (h_+ - ih_\times)$$

$$e_{ij}^{\text{L}}(\mathbf{n}) = \frac{1}{\sqrt{2}} (e_{ij}^+(\mathbf{n}) - ie_{ij}^\times(\mathbf{n})) \quad h_{\text{L}} = \frac{1}{\sqrt{2}} (h_+ + ih_\times)$$

Then, the equations of the gravitational wave is derived by

$$\left(a^2 - \epsilon_A \Phi' \frac{\alpha}{\kappa} k \right) h''_A + \left(2aa' - \epsilon_A \Phi'' \frac{\alpha}{\kappa} k \right) h'_A + \left(a^2 - \epsilon_A \Phi' \frac{\alpha}{\kappa} k \right) k^2 h_A = 0$$

- We can neglect the effect of the expansion of space in the time scale when the gravitational waves pass Galaxy. Then, the scale factor can be estimated as

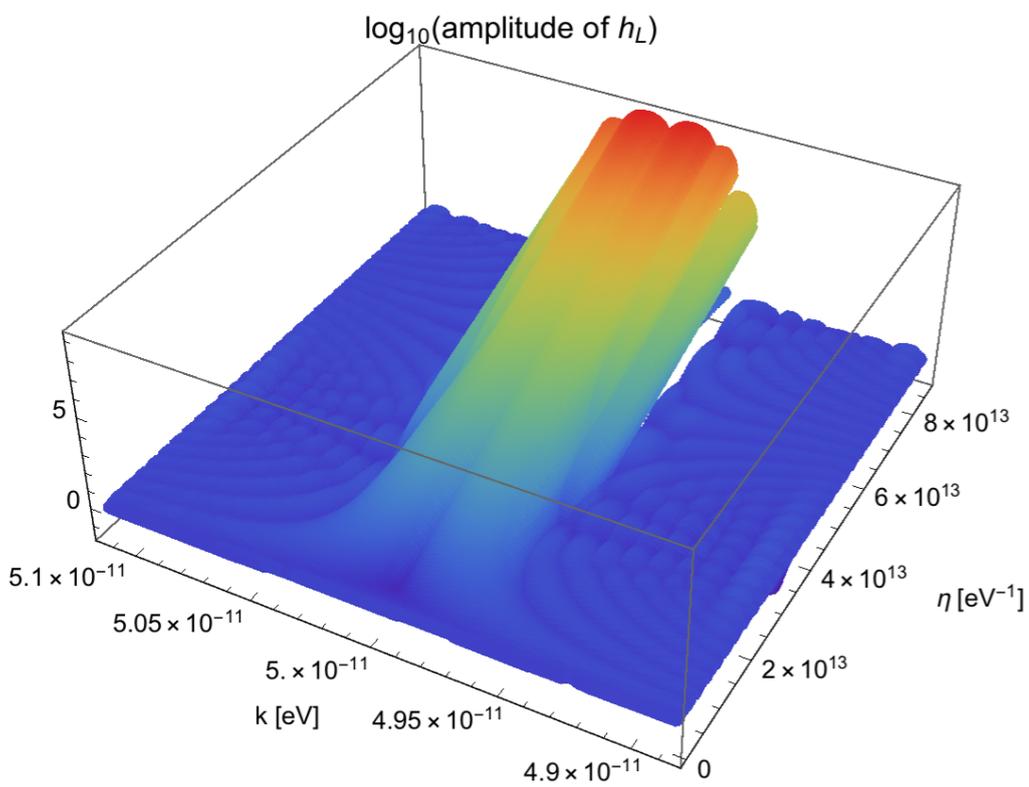
$$a(\eta) \simeq 1. \text{ Then, we can get the solution of the axion field, } \Phi \simeq \Phi_0 \cos(m\eta).$$

- Finally, the equations of the gravitational wave is derived as

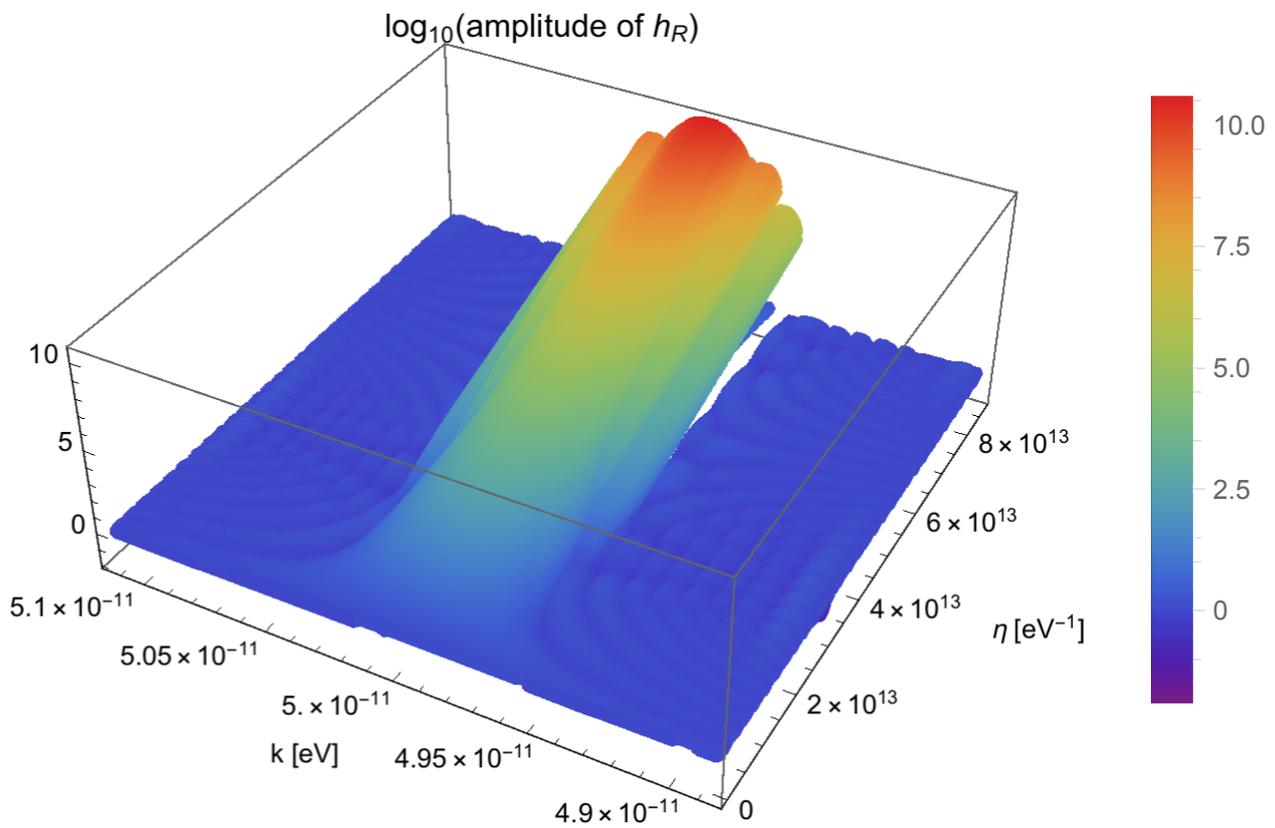
$$h''_A + \frac{\epsilon_A \delta \cos(m\eta)}{1 + \epsilon_A \frac{k}{m} \delta \sin(m\eta)} k h'_A + k^2 h_A = 0 \quad \epsilon_A \equiv \begin{cases} 1 & : A = \text{R} \\ -1 & : A = \text{L} \end{cases} \quad \delta \equiv \frac{\alpha}{\kappa} m^2 \Phi_0$$

Growth of Amplitudes of GWs

- These Plots shows the growth of the amplitude in the circular polarization basis.
 We used the values below $\ell = 10^8 \text{ km}$, $m = 10^{-10} \text{ eV}$, $\rho = 0.3 \times 10^6 \text{ GeV/cm}^3$
 $\delta \simeq 0.02$



h_L



h_R

→ There is the difference between the growth of the amplitude in R- and L-circular polarization basis. In these conditions, the amplitude of h_R becomes 10^4 times bigger than the amplitude of h_L . So, the fully circular polarized wave appears.

Strength of Circular Polarization

- This plot shows the time evolution of the strength of the circular polarization.

We used the same values before slide.

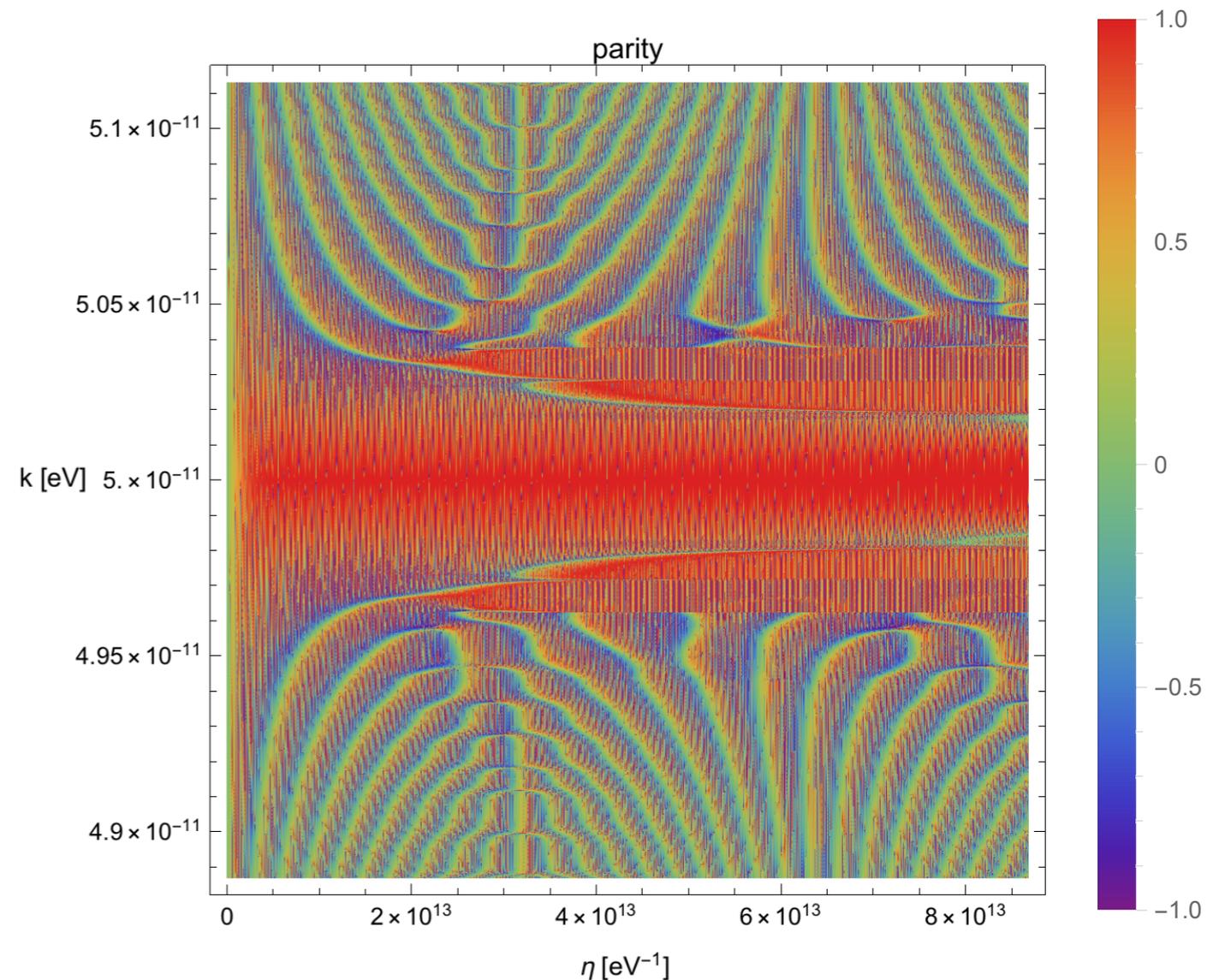
→ The horizontal axis shows the time η .

→ The vertical axis shows the frequency k of the GWs.

→ The plotted function is defined by

$$\text{parity}(\eta) \equiv \frac{|h_R|^2 - |h_L|^2}{|h_R|^2 + |h_L|^2}$$

parity $\sim \pm 1$: R or L



- The color shows the strength of the circular polarization.
 - At the resonance frequency, the GW is completely circular polarized.

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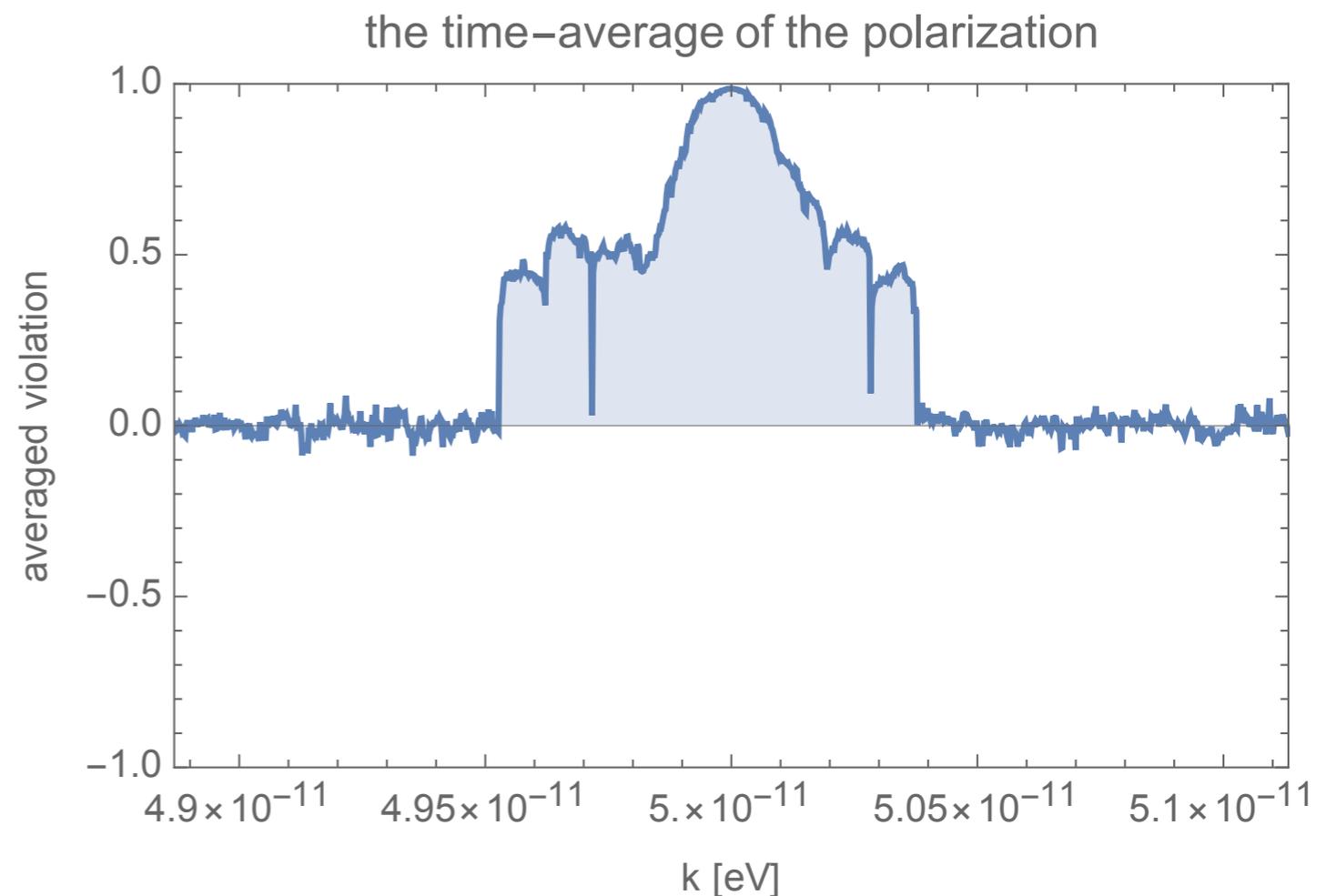
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Proceedings, 4th Workshop on Cosmology and the Quantum Vacuum: Segovia, Spain, September 4-8, 2017

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Tools for Estimations

- From the general theory of the parametric resonance, we can give the formulae for the estimations of the resonance frequency and the rate of the amplification in the length.

- Resonance

→ Resonance frequency : $k_r = \frac{m}{2} \Rightarrow f_r = \frac{k_r}{2\pi} \simeq 1.2 \times 10^4 \text{ Hz} \left(\frac{m}{10^{-10} \text{ eV}} \right)$

→ Width of the resonance : $\frac{m}{2} - \frac{m}{8}\delta \lesssim k_r \lesssim \frac{m}{2} + \frac{m}{8}\delta \quad \delta \equiv \frac{\alpha}{\kappa} m^2 \Phi_0 \quad \alpha = \sqrt{\frac{\kappa}{2}} \ell^2$

- Amplification of the amplitude of GW

We can convert Γ into the length $R_{\times 10}$ in which the amplitude of the GWs becomes 10 times bigger.

→ Typical length : $R_{\times 10} \simeq 5.2 \times 10^{-8} \text{ pc} \times \left(\frac{10^{-10} \text{ eV}}{m} \right)^2 \left(\frac{10^8 \text{ km}}{\ell} \right)^2 \sqrt{\frac{0.3 \text{ GeV/cm}^3}{\rho}}$

Simple Estimation

- The length in which the GWs becomes 10 times bigger is estimated as

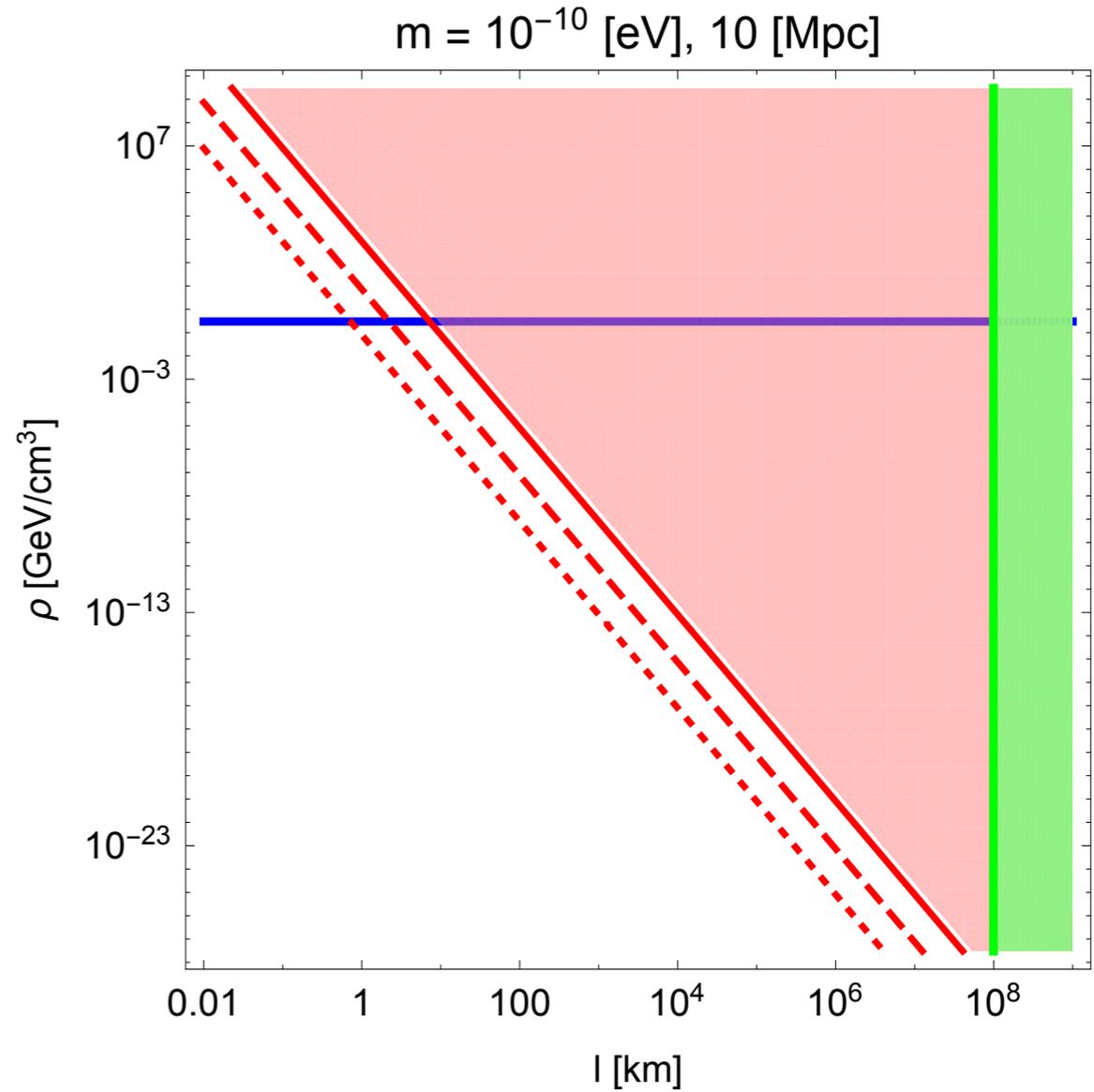
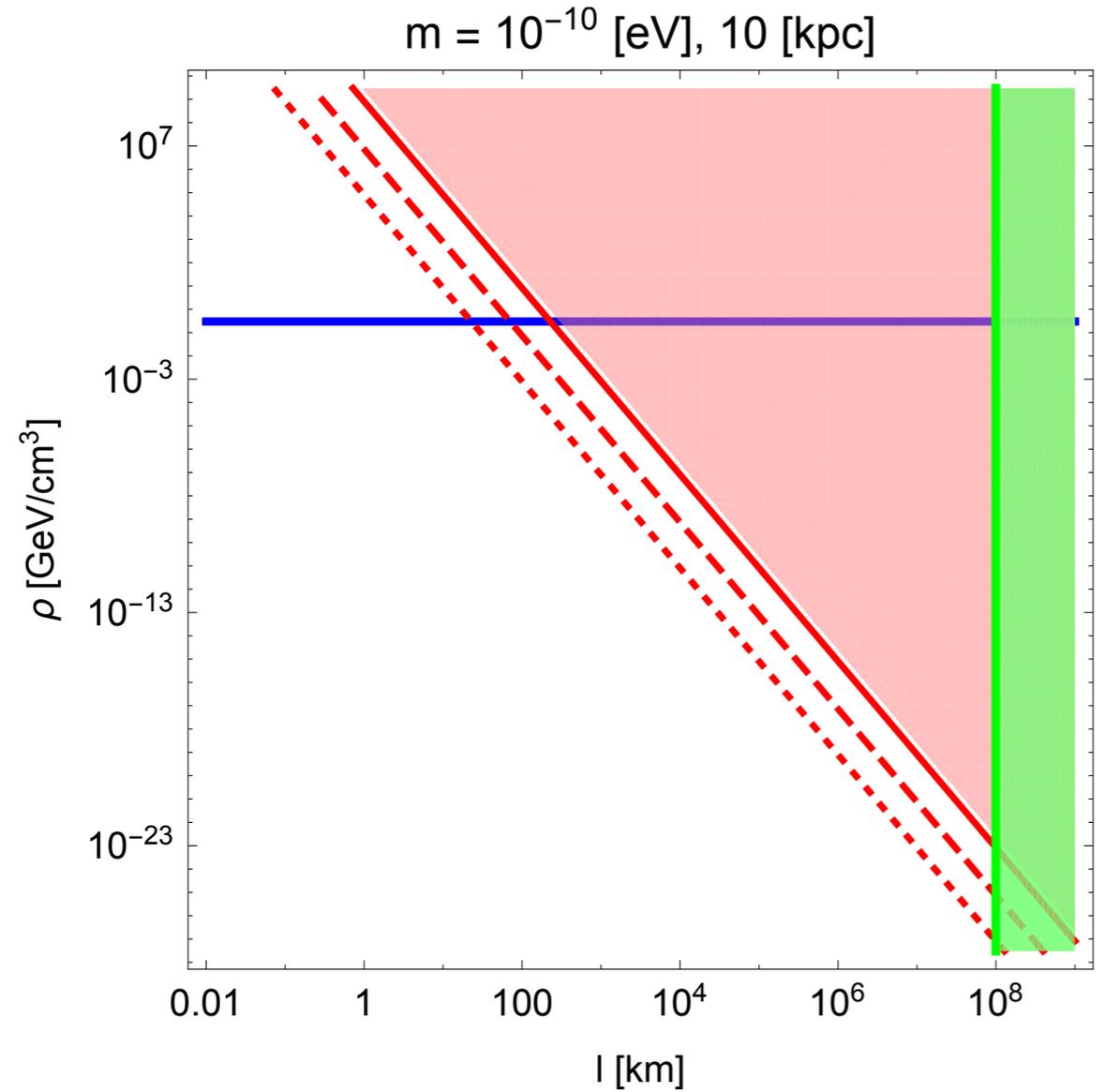
$$R_{\times 10} \simeq 5.0 \times 10^{-8} \text{ pc} \\ \times \left(\frac{10^{-10} \text{ eV}}{m} \right)^2 \left(\frac{10^8 \text{ km}}{\ell} \right)^2 \sqrt{\frac{0.3 \text{ GeV/cm}^3}{\rho}}$$

- If the GWs which have the resonance frequency through in Galaxy for 10 kpc, the amplitude of the GWs become $10^{10^{12}}$ times bigger.
- This effect will be detected in the difference between the waveform from **the interferometers** and from **the numerical template** of the GWs in pure GR.
- If we found no difference in the waveform, we can give the constraint to the coupling constants or the energy density of the axion dark matter.

→ If you believe $\ell \sim 10^8 \text{ km}$, then ρ will be constraint by $\rho \leq 10^{-26} \text{ GeV/cm}^3$.

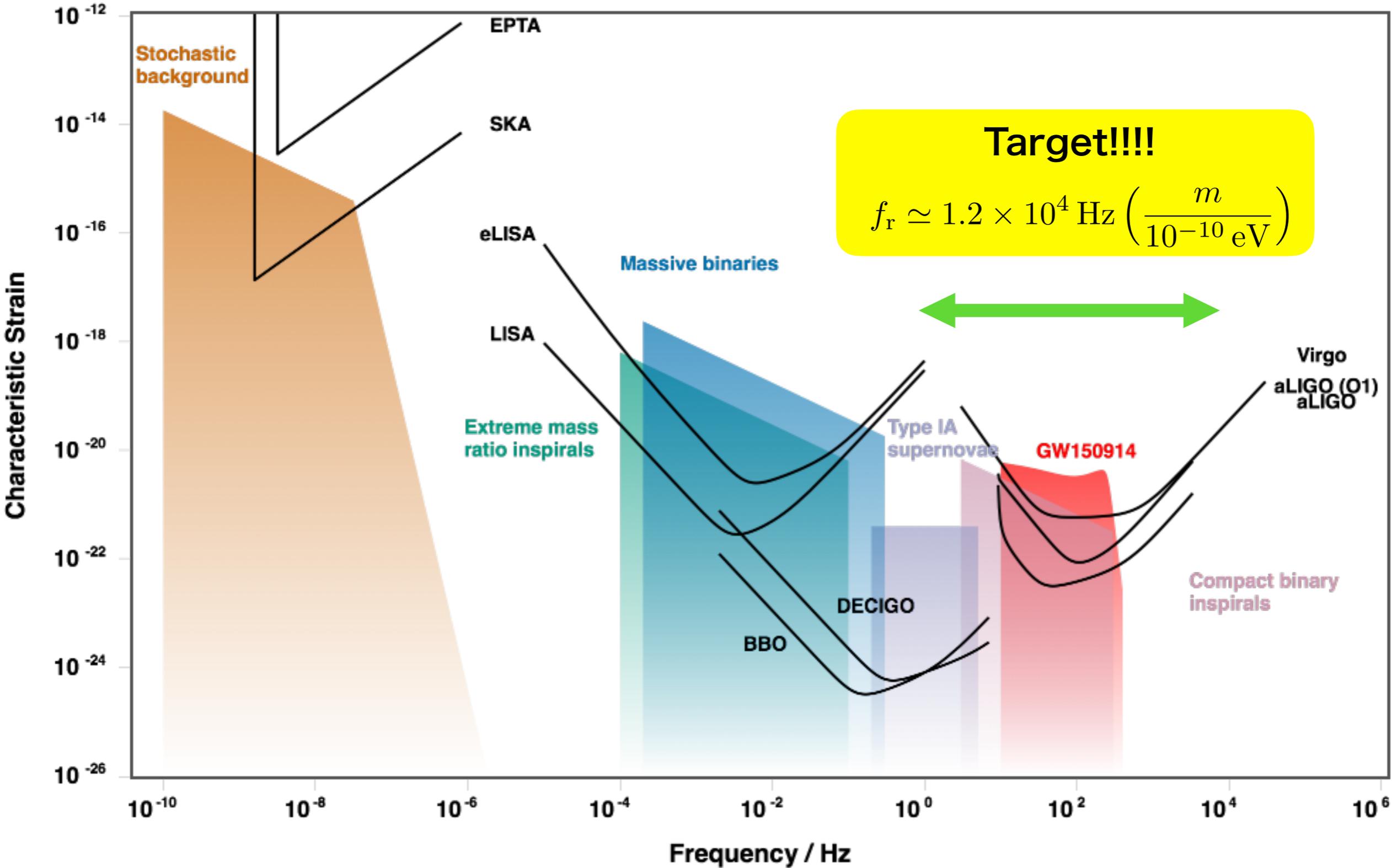
→ If you believe $\rho \sim 0.3 \text{ GeV/cm}^3$, then ℓ will be constraint by $\ell \leq 10 \text{ km}$.

Maps of Constraints



- The green region is excluded by the observation.**
- The blue line shows the local dark matter density 0.3 GeV/cm^3 .**
- The red line represents the 10 times enhancement of GWs.**
- The red dashed line represents the 0.1 times enhancement.**
- The red dotted line represents the 0.01 times enhancement.**

Sensitivity Curves and Axion Dark Matter



These Sources are from "<http://rhcole.com/apps/GWplotter/>".

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Conclusion

- String axiverse generates the axions which have the light mass and the dCS coupling between the axion field and the gravitational field.
- The coherently oscillating axion can behave as the cold dark matter well, so, through the above coupling, the monochromatic and parity violated gravitational waves are generally induced.
- This effect may use to detect the counterpart of the GR.
→ They might give the new constraint to the abundance of the axion dark matter or the CS-coupling.
- This analysis can connect the problems about the CDM and the modified gravity.

Related Researches

- **Articles**

“Electromagnetic memory effect induced by axion dark matter”

DY and Jiro Soda, Phys. Rev. D **96**, 064005 (2017)

“Exploring the string axiverse and parity violation in gravity with gravitational waves”,

DY and Jiro Soda, International Journal of Modern Physics D Vol. 27, No. 9 (2018) 1850096

“Electromagnetic waves propagating in the string axiverse”,

DY and Jiro Soda, Progress of Theoretical and Experimental Physics, Volume 2018, Issue 4, 1 April 2018, 041E01

- **Proceedings**

“Exploring String Axions with Gravitational Waves”,

Jiro Soda and DY, Galaxies 5 (2017) no.4, 96