

Semi-Analytic Calculation of Gravitational Wave Spectrum induced from Primordial Curvature Perturbations

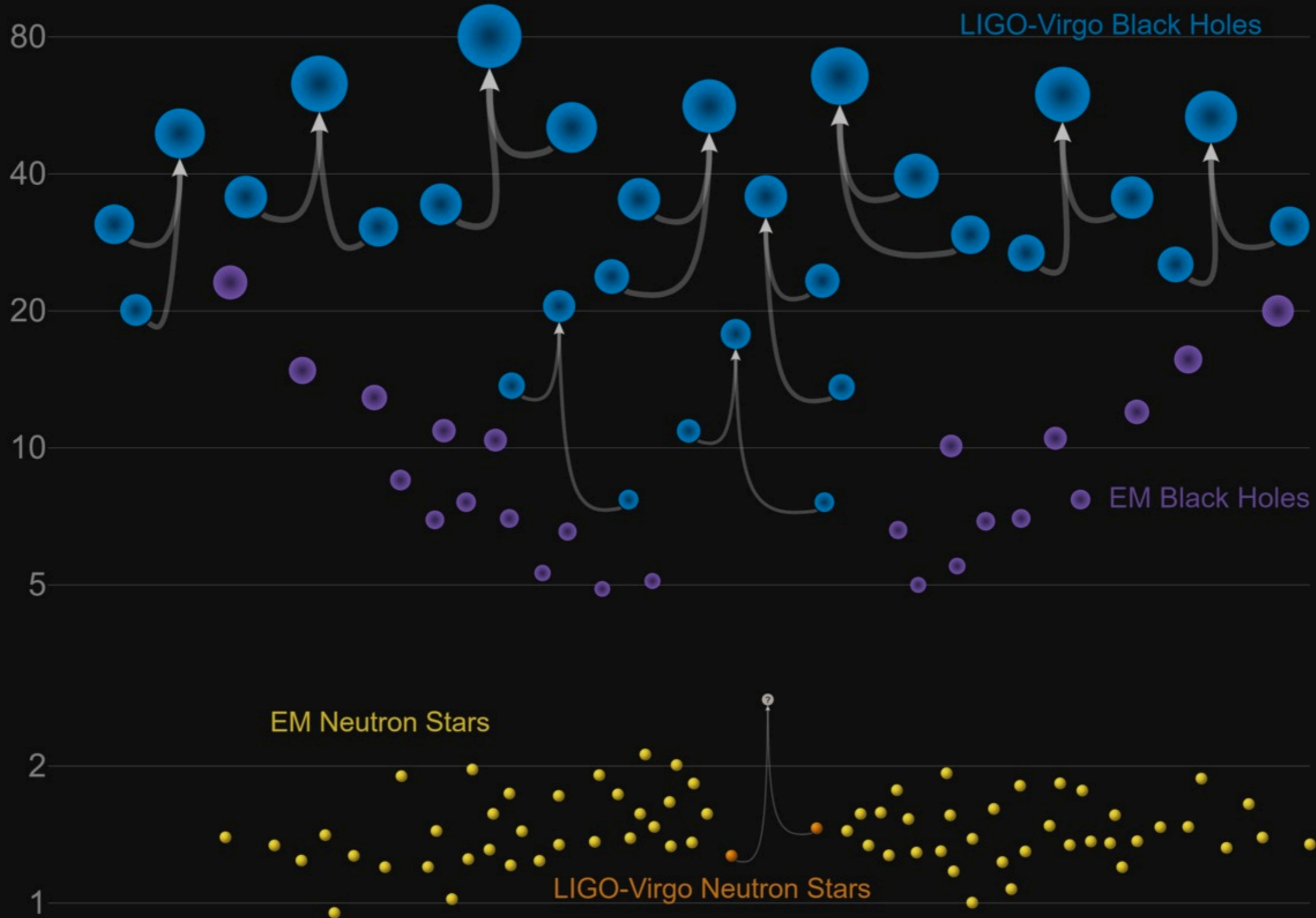
寺田 隆広

(KEK, JSPS fellow)

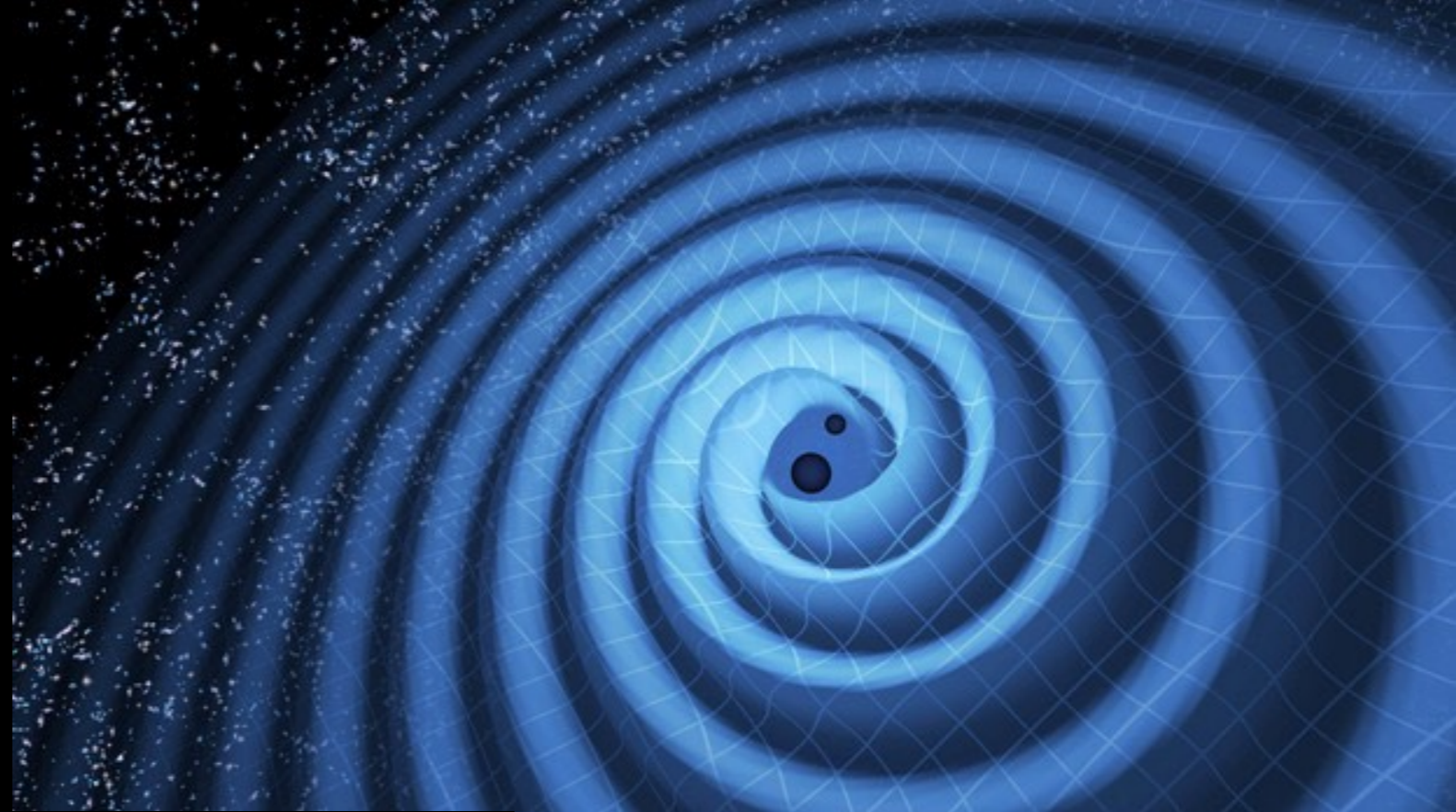
[Kohri, Terada, 1804.08577 [gr-qc], PRD 97 (2018) 123532]

Masses in the Stellar Graveyard

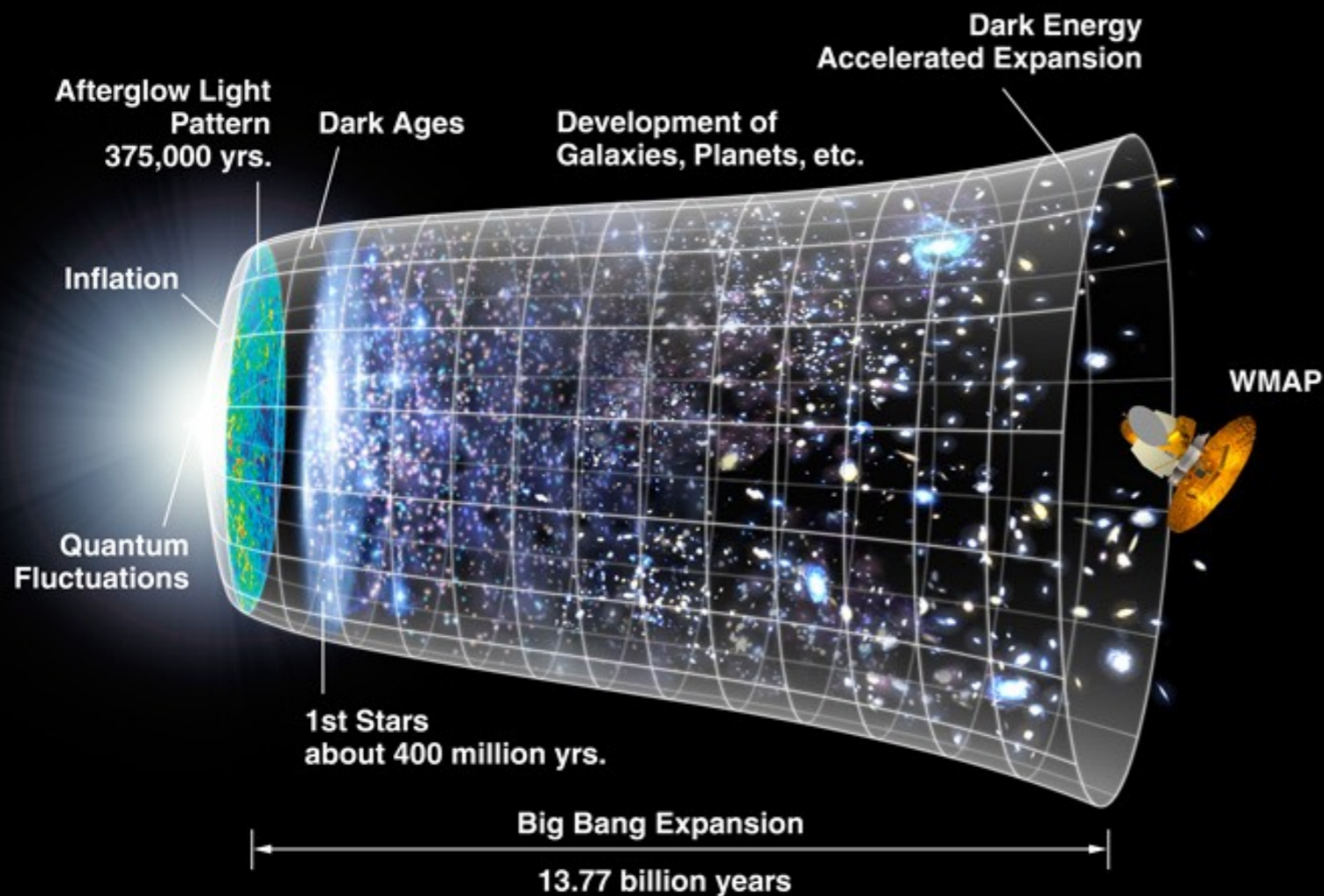
in Solar Masses



GW Astrophysics
↓
GW Cosmology



[Image credit: LIGO/T. Pyle]



We can probe
the very early Universe
or High Energy Physics.

[Image credit: NASA/WMAP Science Team]

Beginning of a Trend?

- Ananda, Clarkson, Wands, gr-qc/0612013
- Baumann, Steinhardt, Takahashi, Ichiki, hep-th/0703290
- Assadullahi, Wands, 0901.0989
- Inomata, Kawasaki, Mukaida, Tada, Yanagida, 1611.06130

- Espinosa, Racco, Riotto, 1804.07732
- Kohri, Terada, 1804.08577
- Cai, Pi, Sasaki, 1810.11000
- Bartolo, De Luca, Franciolini, Peloso, Riotto, 1810.12218
- Bartolo, De Luca, Franciolini, Peloso, Racco, Riotto, 1810.12224
- Byrnes, Cole, Patil, 1811.11158
- Inomata, Nakama, 1812.00674

- For other source of stochastic gravitational waves, see e.g. Kuroyanagi, Chiba, Takahashi, 1807.00786.

What is the secondary GW?

[Ananda, Clarkson, Wands, gr-qc/0612013]

[Baumann, Steinhardt, Takahashi, Ichiki, hep-th/0703290]

[Inomata, Kawasaki, Mukaida, Tada, Yanagida, 1611.06130]

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- Inflation produces
scalar and tensor (1st-order) perturbations

$$ds^2 = -a^2(1 + 2\Phi)d\eta^2 + a^2 \left((1 - 2\Phi)\delta_{ij} + \frac{1}{2}h_{ij} \right) dx^i dx^j$$

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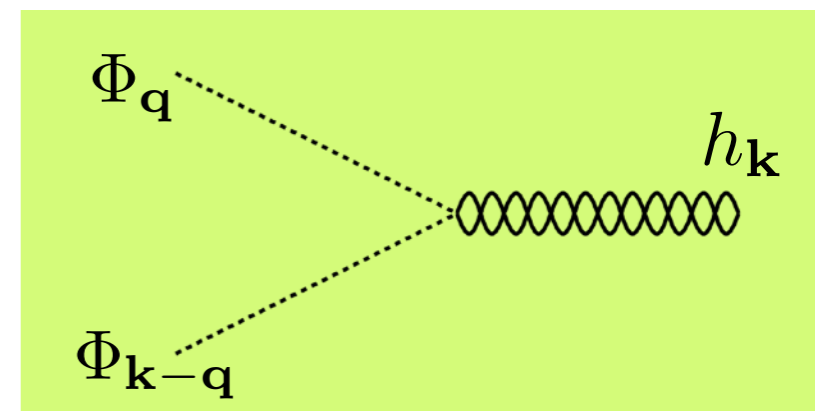
[Inomata, Kawasaki, Mukaida, Tada, Yanagida, 1611.06130]

- Inflation produces **scalar** and **tensor** (1st-order) perturbations

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- (2nd-order **tensor**) = (1st-order **scalar**)²

$$h''_{\mathbf{k}}(\eta) + 2\mathcal{H}h'_{\mathbf{k}}(\eta) + k^2 h_{\mathbf{k}}(\eta) = 4S_{\mathbf{k}}(\eta)$$

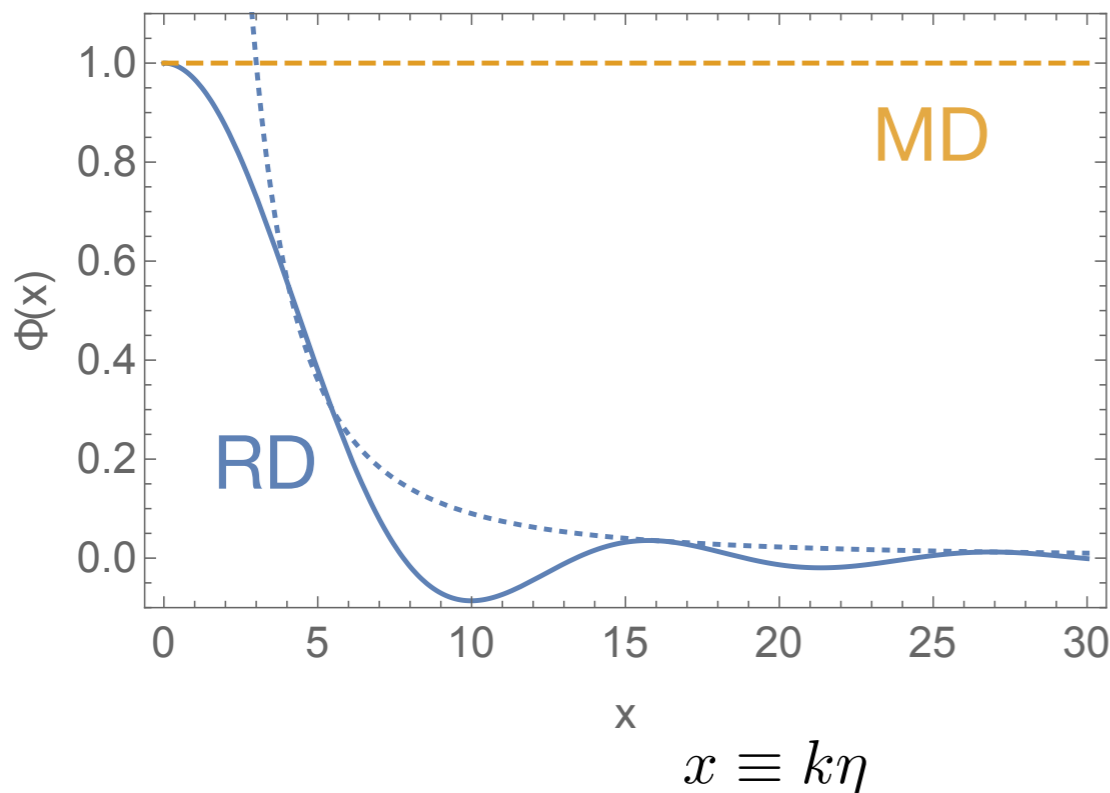
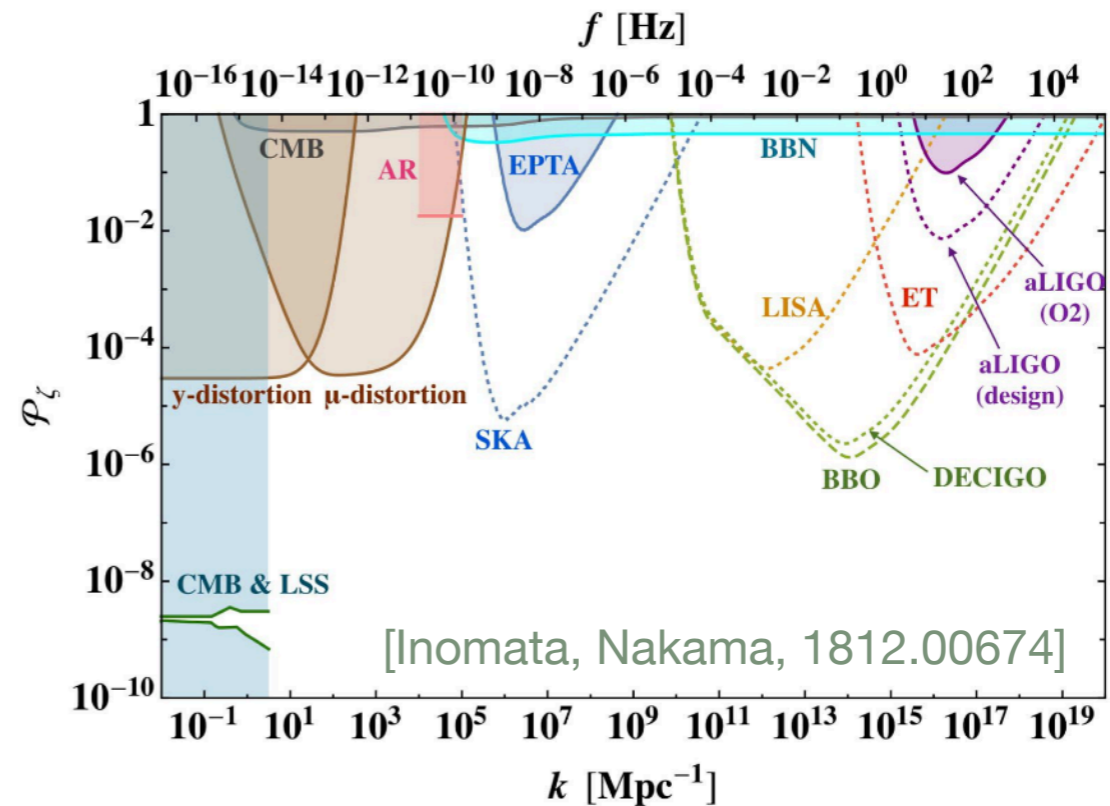


$$S_{\mathbf{k}} = \int \frac{d^3q}{(2\pi)^{3/2}} e_{ij}(\mathbf{k}) q_i q_j \left(2\Phi_{\mathbf{q}}\Phi_{\mathbf{k}-\mathbf{q}} + \frac{4}{3(1+w)} (\mathcal{H}^{-1}\Phi'_{\mathbf{q}} + \Phi_{\mathbf{q}}) (\mathcal{H}^{-1}\Phi'_{\mathbf{k}-\mathbf{q}} + \Phi_{\mathbf{k}-\mathbf{q}}) \right)$$

Enhanced scalar perturbations

Initial condition

Enhanced primordial curvature perturbations are allowed at small scales.

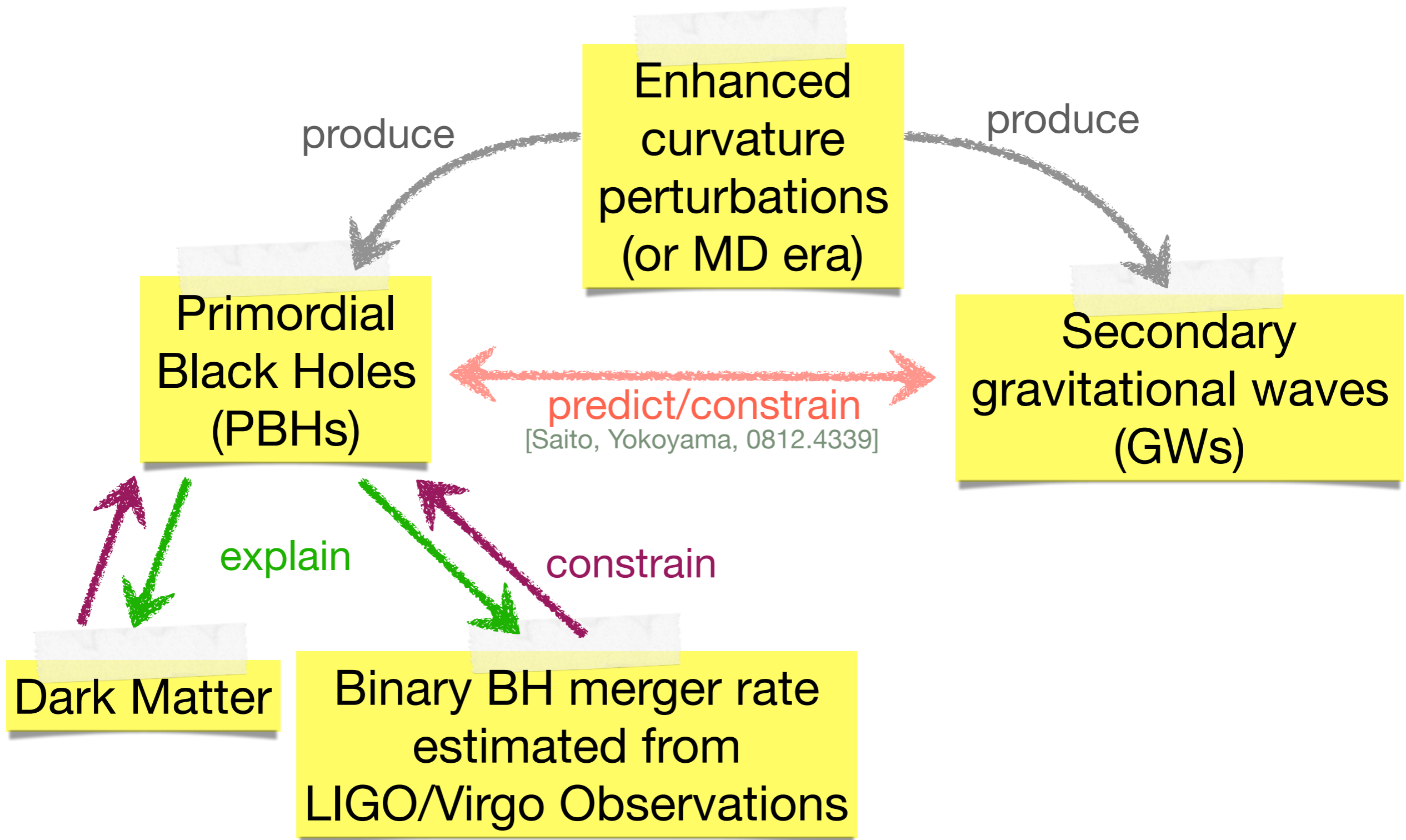


Dynamics

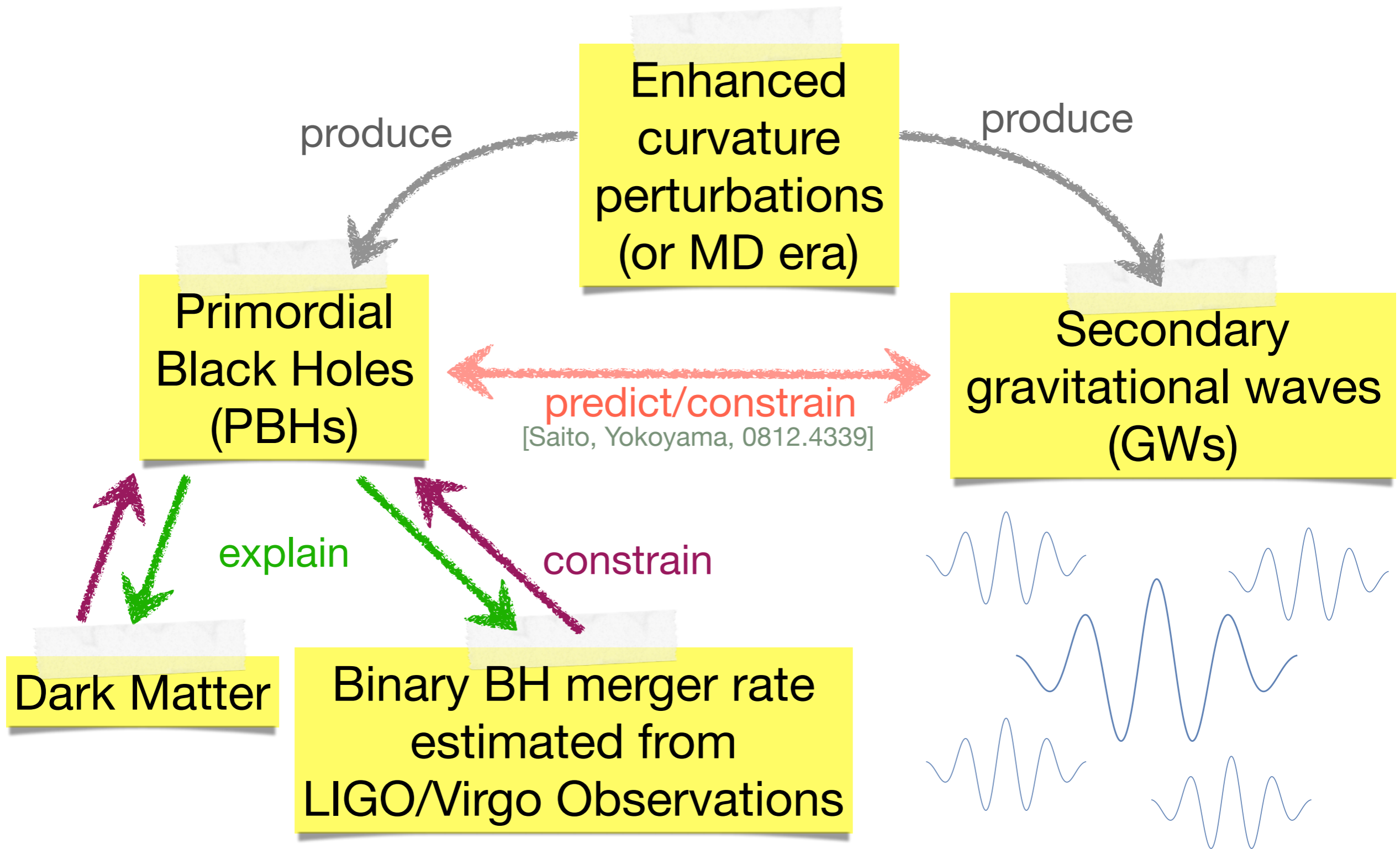
In a radiation-dominated (RD) era, the source decays.

In a matter-dominated (MD) era, the source does not decay.

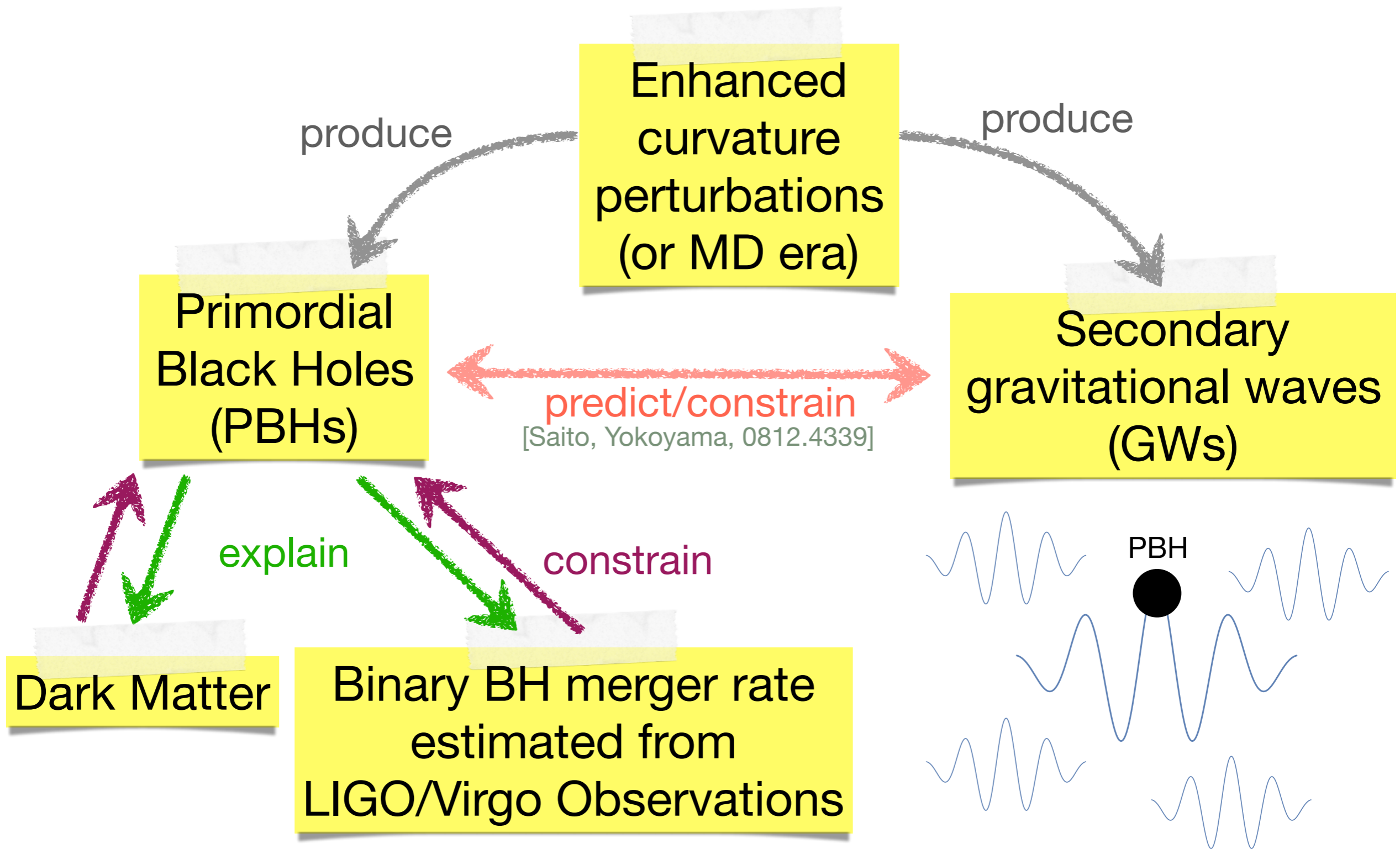
Relation to PBH scenario



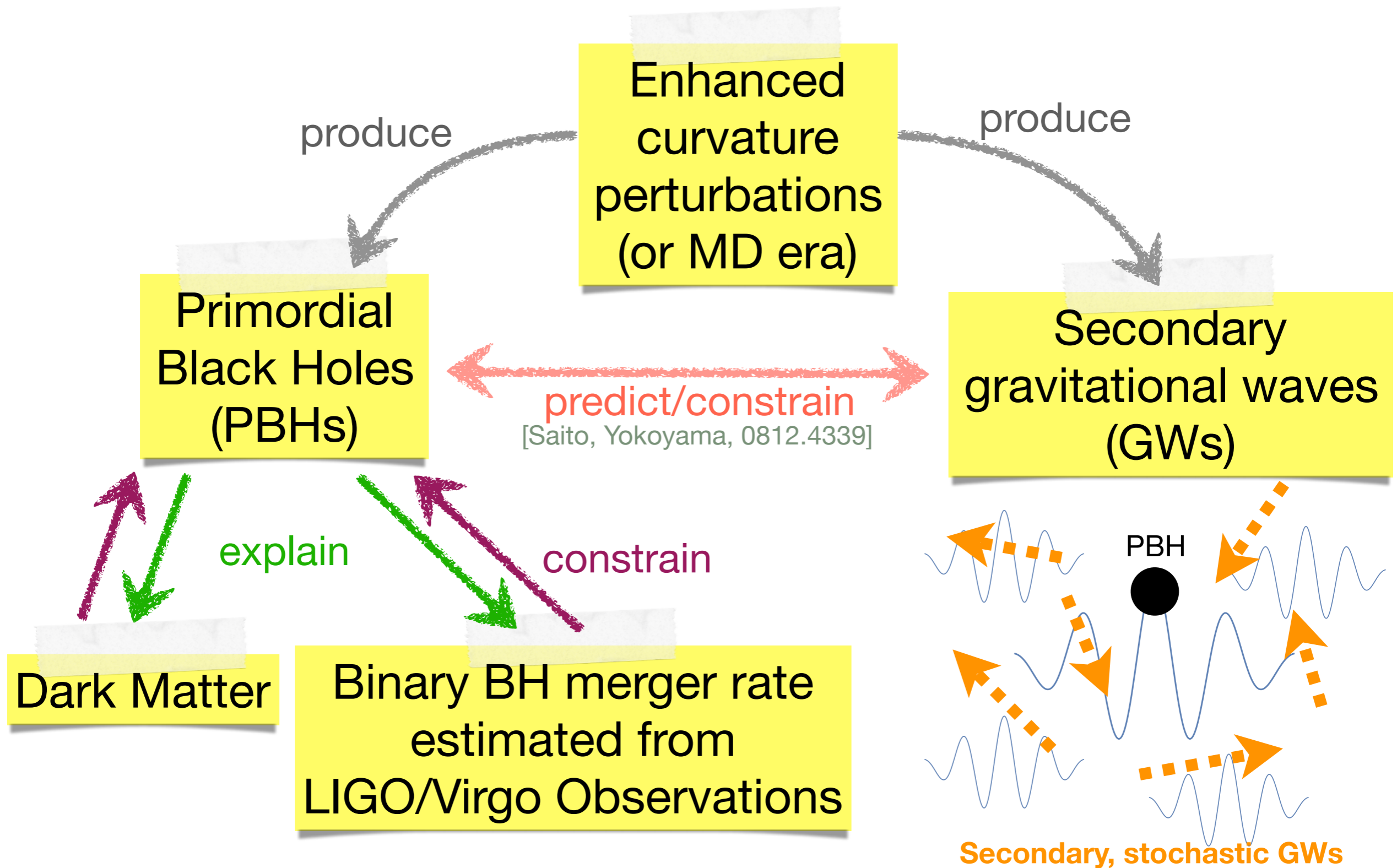
Relation to PBH scenario

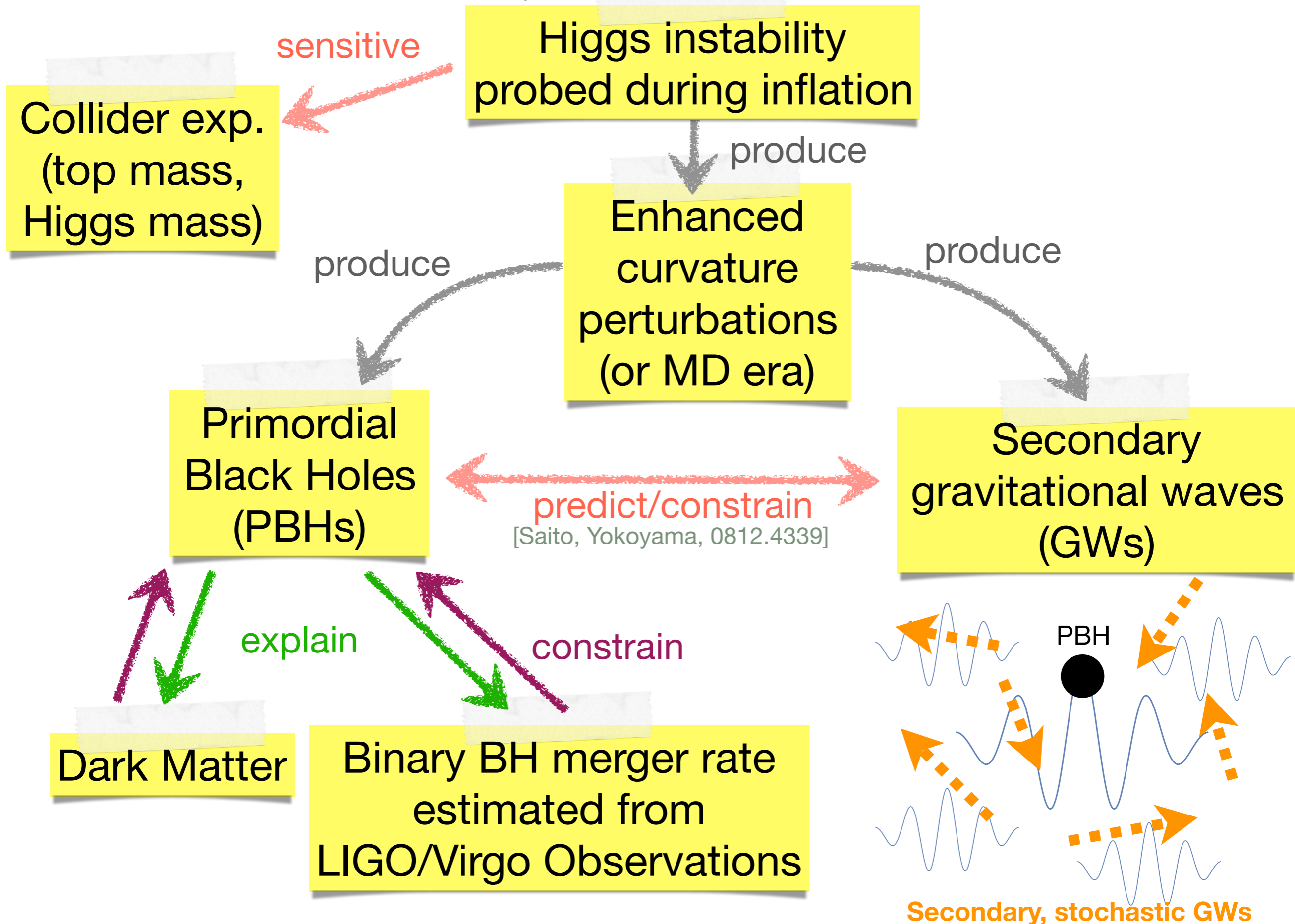


Relation to PBH scenario



Relation to PBH scenario





Motivation (What we do)

The secondary GW power spectrum is obtained by **multiple integral** of an **oscillatory function**.

$$\mathcal{P}_h \sim \int dk \int dk' \left(\int dt f(k, k', t) \right)^2 \mathcal{P}_\zeta(k) \mathcal{P}_\zeta(k')$$

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describing time evolution primordial quantity

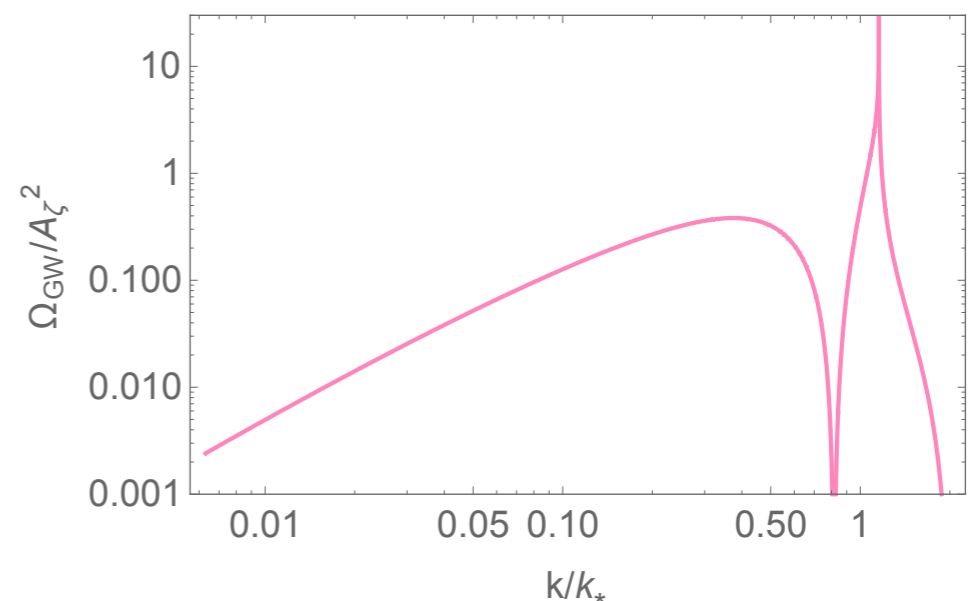
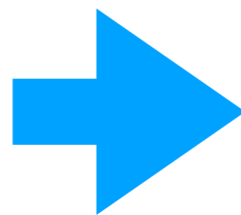
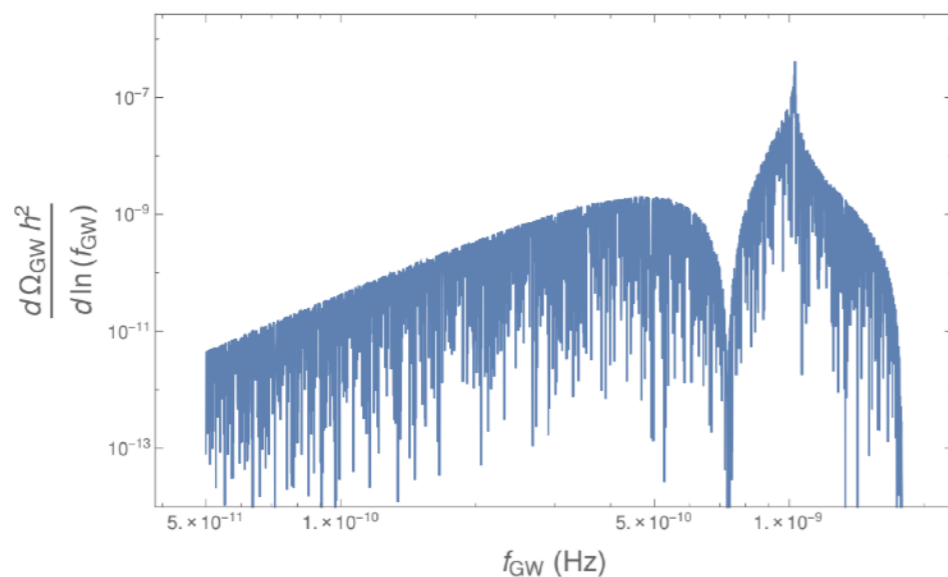
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[Orlofsky, Pierce, Wells, 1612.05279]



Details (1/3)

[Ananda, Clarkson, Wands, gr-qc/0612013] [Baumann, Steinhardt, Takahashi, Ichiki, hep-th/0703290]
[Inomata, Kawasaki, Mukaida, Tada, Yanagida, 1611.06130]

$$\mathcal{P}_h(\eta, k) = 4 \int_0^\infty dv \int_{|1-v|}^{1+v} du \left(\frac{4v^2 - (1 + v^2 - u^2)^2}{4vu} \right)^2 I^2(v, u, x) \mathcal{P}_\zeta(kv) \mathcal{P}_\zeta(ku)$$

where $x \equiv k\eta$ $u = |\mathbf{k} - \tilde{\mathbf{k}}|/k$ $v = \tilde{k}/k$

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where $x \equiv k\eta$ $u = |\mathbf{k} - \tilde{\mathbf{k}}|/k$ $v = \tilde{k}/k$

We want to analytically calculate this function.

$$\text{Definitions: } I(v, u, x) = \int_0^x d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} k G_k(\eta, \bar{\eta}) f(v, u, \bar{x})$$

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$$f(v, u, \bar{x}) = \frac{6(w+1)}{3w+5} \Phi(v\bar{x})\Phi(u\bar{x}) + \frac{6(1+3w)(w+1)}{(3w+5)^2} (\bar{x}\partial_{\bar{\eta}}\Phi(v\bar{x})\Phi(u\bar{x}) + \bar{x}\partial_{\bar{\eta}}\Phi(u\bar{x})\Phi(v\bar{x})) \\ + \frac{3(1+3w)^2(1+w)}{(3w+5)^2} \bar{x}^2 \partial_{\bar{\eta}}\Phi(v\bar{x})\partial_{\bar{\eta}}\Phi(u\bar{x})$$

We have used $\mathcal{H} = aH = 2/((1+3w)\eta)$ $\bar{x} \equiv k\bar{\eta}$ $w \equiv P/\rho$

Details (2/3)

For definiteness, consider $w = 1/3$ (RD era).

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$$\Phi(x) = \frac{9}{x^2} \left(\frac{\sin(x/\sqrt{3})}{x/\sqrt{3}} - \cos(x/\sqrt{3}) \right)$$

$$f_{\text{RD}}(v, u, x) = \frac{12}{u^3 v^3 x^6} \left(18uvx^2 \cos \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} + (54 - 6(u^2 + v^2)x^2 + u^2 v^2 x^4) \sin \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} \right. \\ \left. + 2\sqrt{3}ux(v^2 x^2 - 9) \cos \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} + 2\sqrt{3}vx(u^2 x^2 - 9) \sin \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} \right)$$

Details (3/3)

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$$\begin{aligned}
 I_{\text{RD}}(v, u, x) = & \frac{3}{4u^3v^3x} \left(-\frac{4}{x^3} \left(uv(u^2 + v^2 - 3)x^3 \sin x - 6uvx^2 \cos \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} \right. \right. \\
 & + 6\sqrt{3}ux \cos \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} + 6\sqrt{3}vx \sin \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} - 3(6 + (u^2 + v^2 - 3)x^2) \sin \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} \left. \right) \\
 & + (u^2 + v^2 - 3)^2 \left(\sin x \left(\text{Ci} \left(\left(1 - \frac{v-u}{\sqrt{3}} \right) x \right) + \text{Ci} \left(\left(1 + \frac{v-u}{\sqrt{3}} \right) x \right) \right. \right. \\
 & - \text{Ci} \left(\left| 1 - \frac{v+u}{\sqrt{3}} \right| x \right) - \text{Ci} \left(\left(1 + \frac{v+u}{\sqrt{3}} \right) x \right) + \log \left| \frac{3 - (u+v)^2}{3 - (u-v)^2} \right| \left. \right) \\
 & + \cos x \left(-\text{Si} \left(\left(1 - \frac{v-u}{\sqrt{3}} \right) x \right) - \text{Si} \left(\left(1 + \frac{v-u}{\sqrt{3}} \right) x \right) \right. \\
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$$\text{Si}(x) = \int_0^x d\bar{x} \frac{\sin \bar{x}}{\bar{x}}$$

$$\text{Ci}(x) = - \int_x^\infty d\bar{x} \frac{\cos \bar{x}}{\bar{x}}$$

See also [Espinosa, Racco, Riotto, 1804.07732]

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Oscillation average in the late-time limit

$$\begin{aligned}
 \overline{I_{\text{RD}}^2}(v, u, x \rightarrow \infty) = & \frac{1}{2} \left(\frac{3(u^2 + v^2 - 3)}{4u^3v^3x} \right)^2 \left(\left(-4uv + (u^2 + v^2 - 3) \log \left| \frac{3 - (u+v)^2}{3 - (u-v)^2} \right| \right)^2 \right. \\
 & \left. + \pi^2 (u^2 + v^2 - 3)^2 \Theta(v + u - \sqrt{3}) \right)
 \end{aligned}$$

See also [Espinosa, Racco, Riotto, 1804.07732]

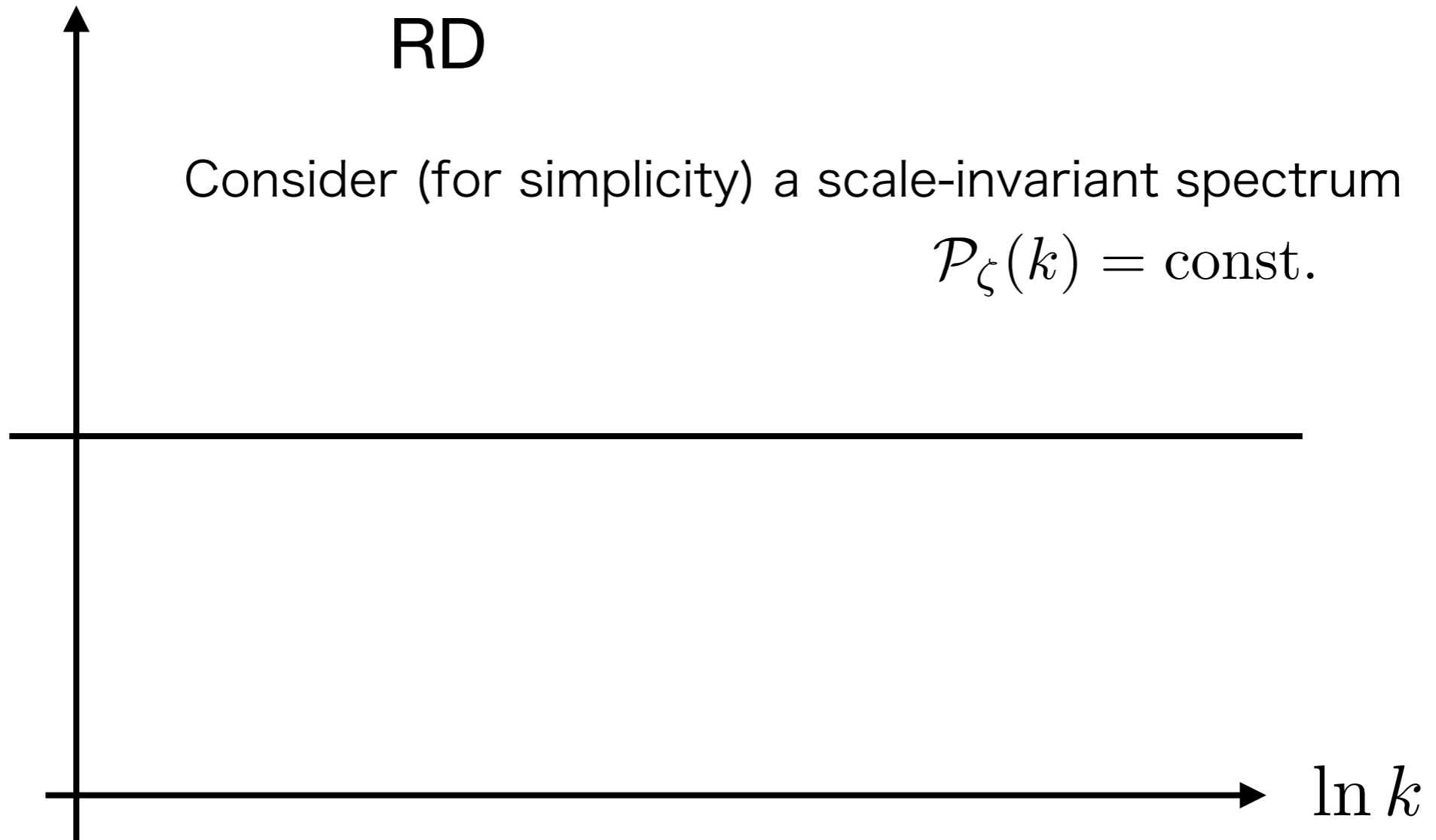
General cases

$\ln \Omega_{\text{GW}}(k)$

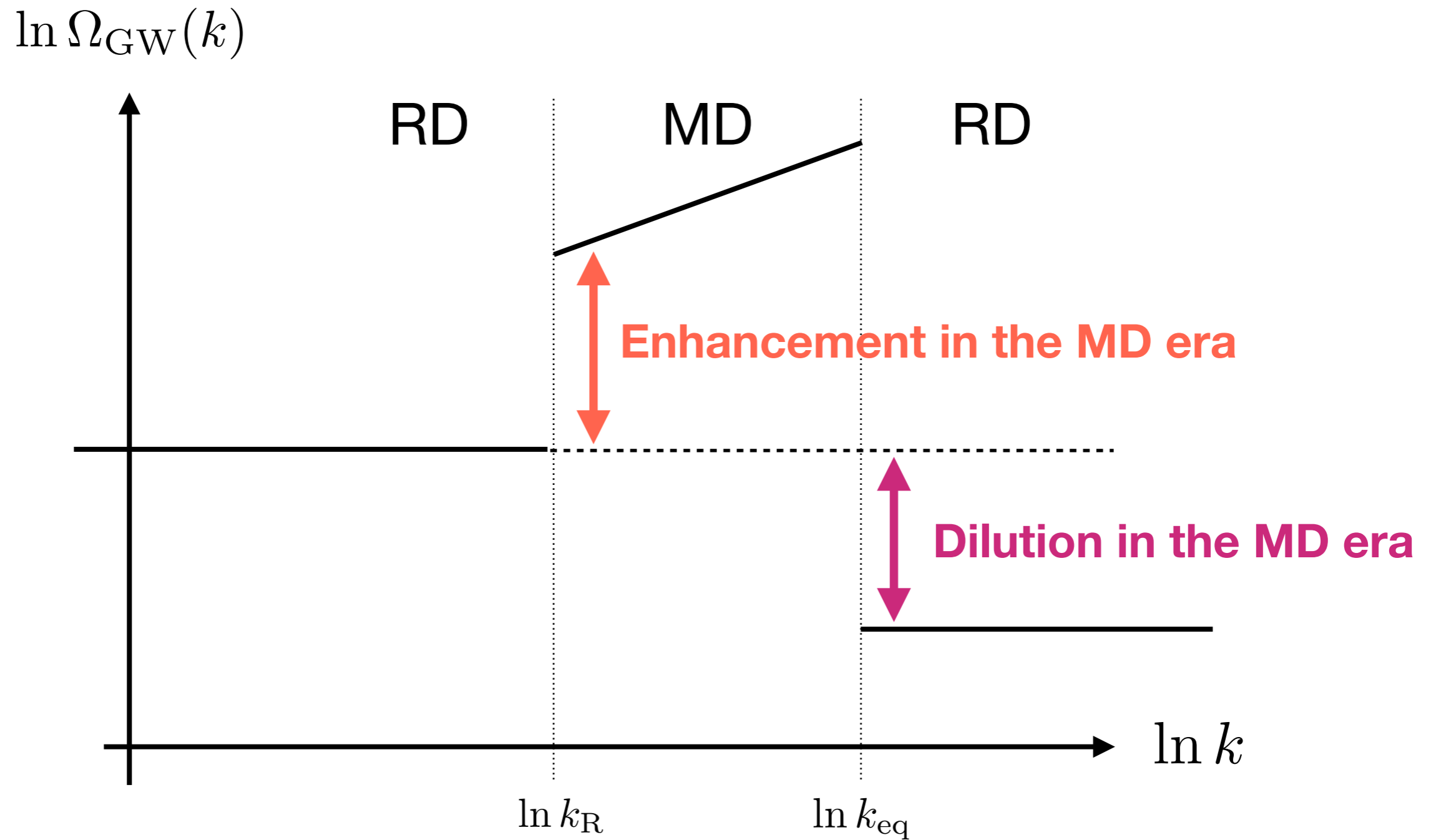
RD

Consider (for simplicity) a scale-invariant spectrum

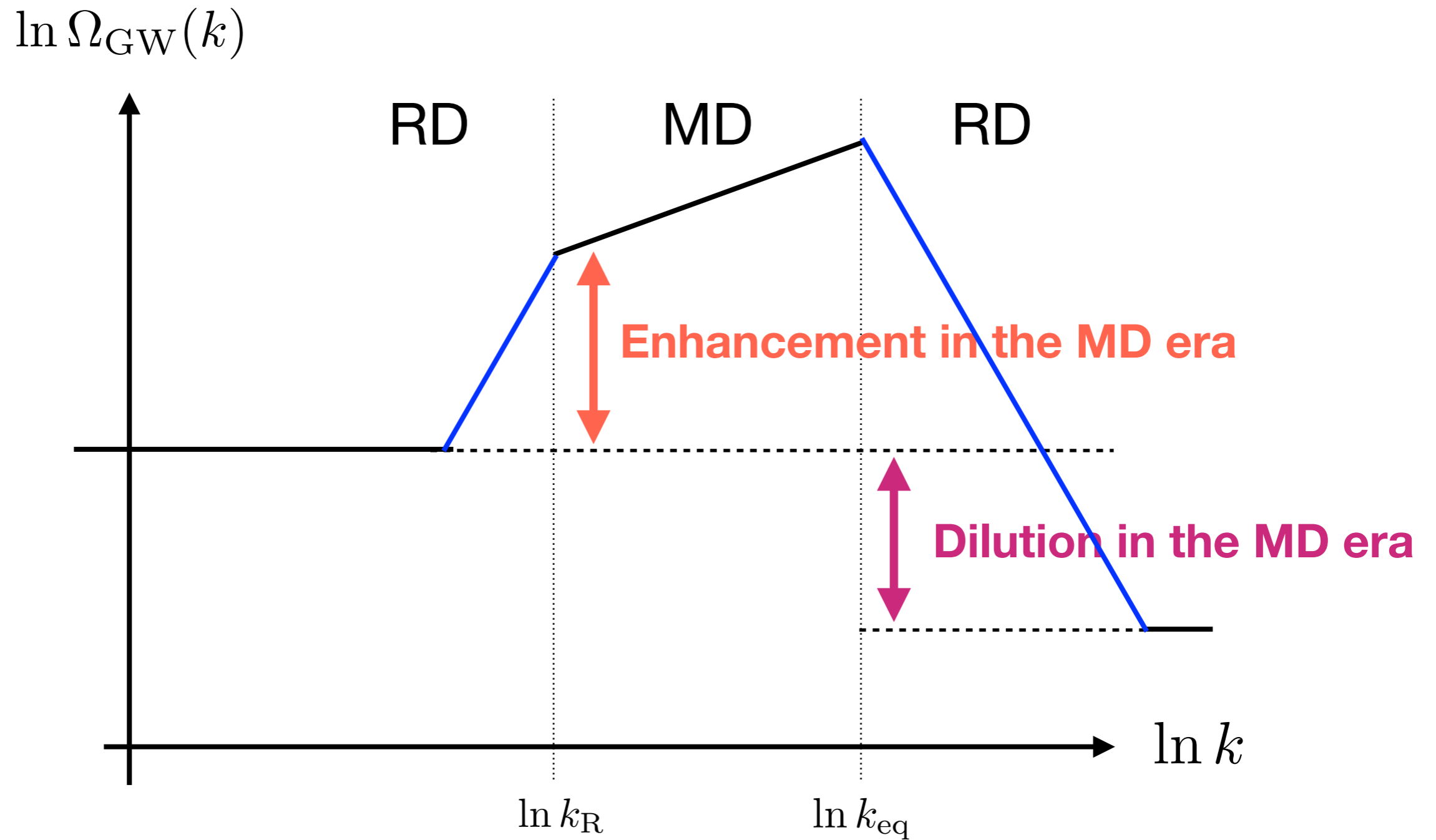
$$\mathcal{P}_\zeta(k) = \text{const.}$$



General cases

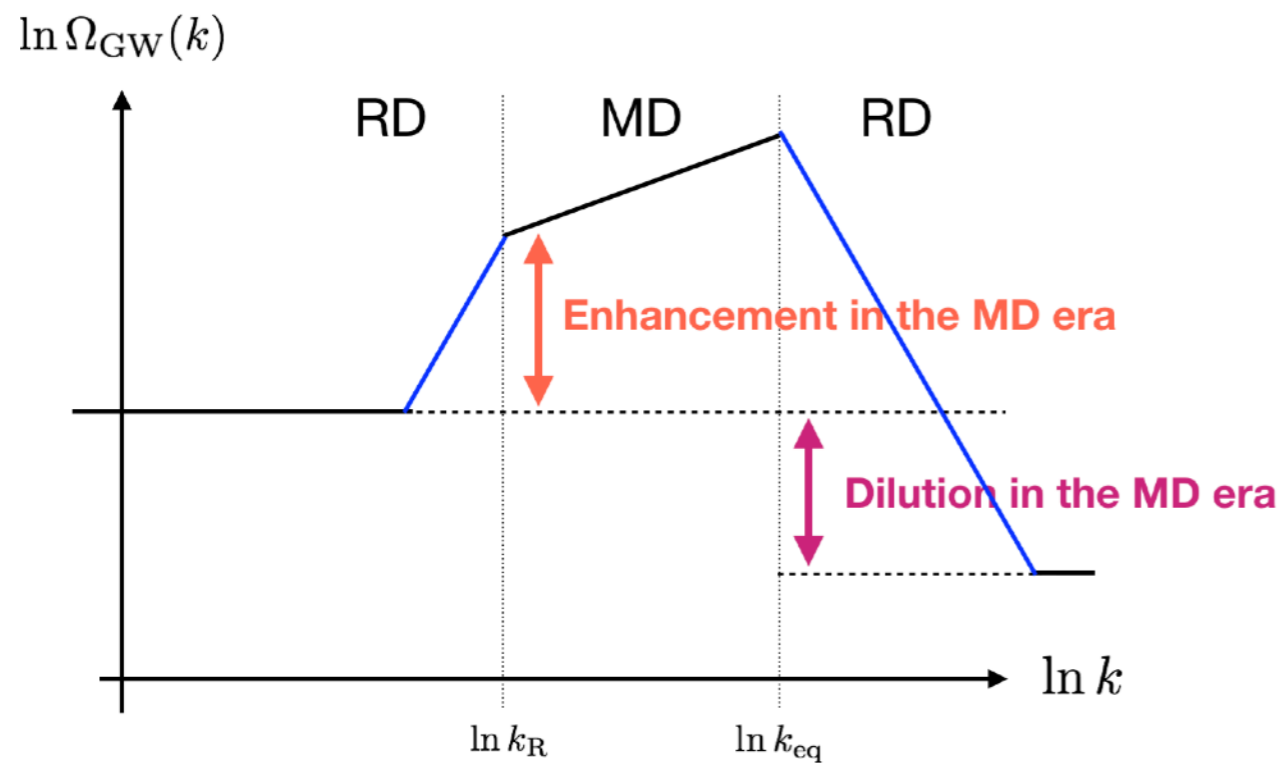


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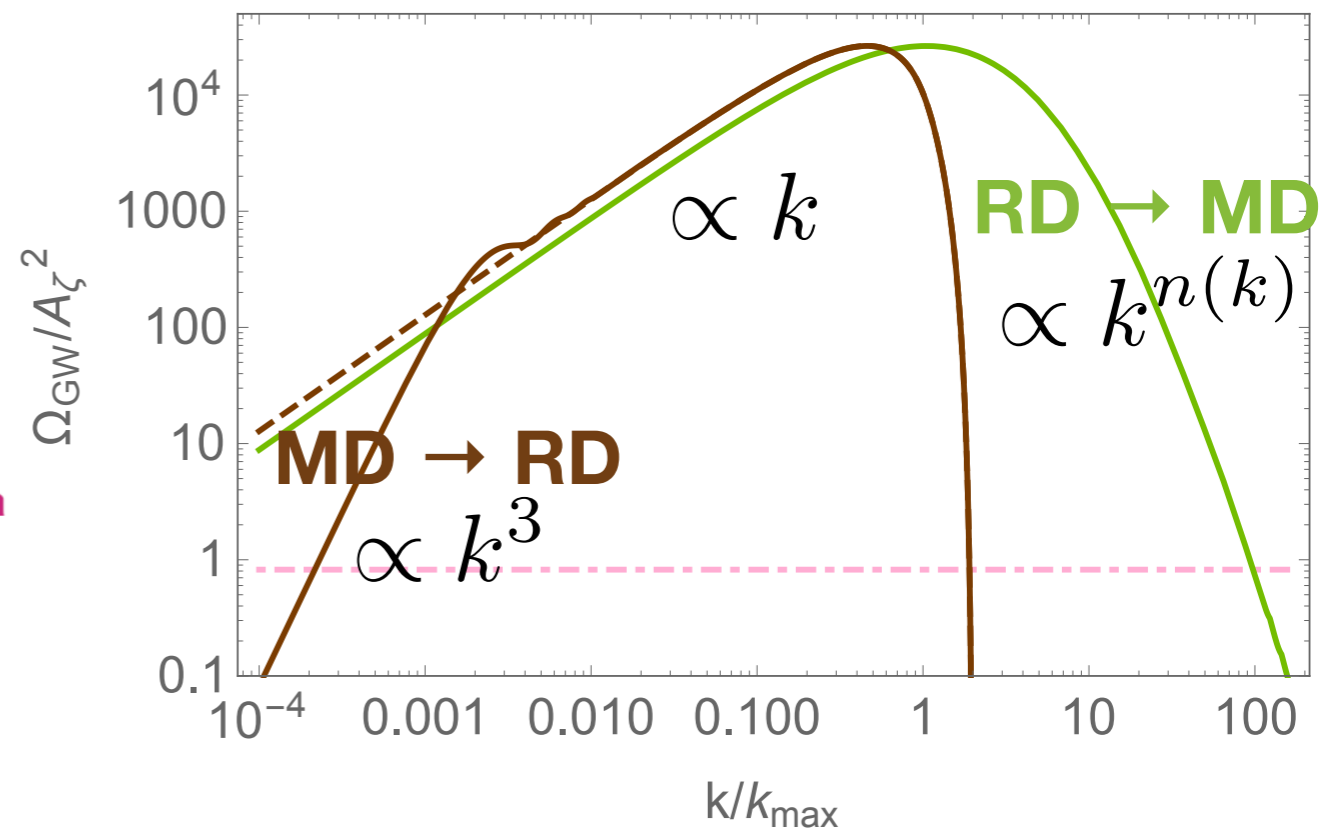


General cases

Schematic



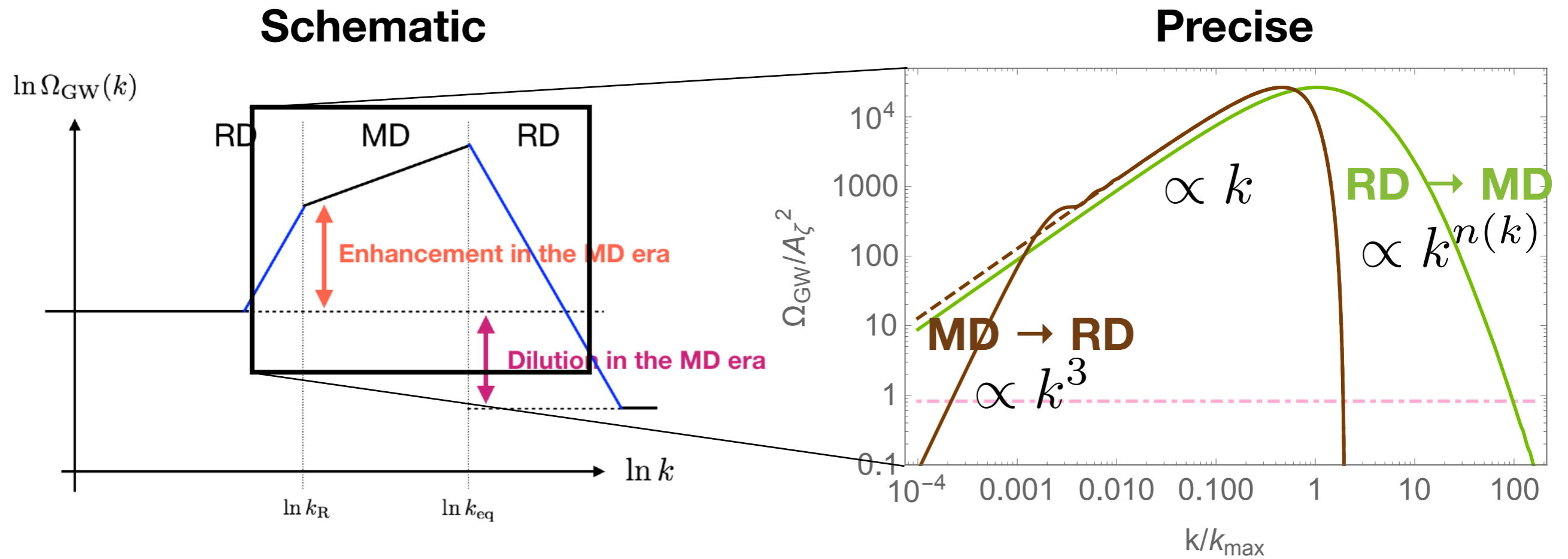
Precise



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Scale-invariant primordial spectrum is assumed as a simple example.

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Summary

- **Secondary GW** is induced by primordial curvature perturbations. (important e.g. in PBH scenarios)
- We **analytically calculated** the **universal part** of the power spectrum of the secondary GW.
- Applicable to GWs induced in **RD**, **MD**, and **more general cases**.
- **Please use our results to write your papers!**

Appendix:

More quantitative explanations

- Basic things
- Radiation-dominated (RD) era
- Matter-dominated (MD) era
- More general cases

Observable & Basic definitions

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Energy fraction

$$\Omega_{\text{GW}}(\eta, k) = \frac{1}{\rho_{\text{tot}}(\eta)} \frac{d\rho_{\text{GW}}(\eta)}{d \ln k} = \frac{1}{24} \left(\frac{k}{a(\eta)H(\eta)} \right)^2 \overline{\mathcal{P}_h(\eta, k)}$$

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Power spectrum

$$\langle h_{\mathbf{k}}^\lambda(\eta) h_{\mathbf{k}'}^{\lambda'}(\eta) \rangle = \delta_{\lambda\lambda'} \delta^3(\mathbf{k} + \mathbf{k}') \frac{2\pi^2}{k^3} \mathcal{P}_h(\eta, k)$$

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Gravitational field: Green's function method

$$a(\eta)h_{\mathbf{k}}(\eta) = 4 \int^{\eta} d\bar{\eta} G_{\mathbf{k}}(\eta, \bar{\eta}) a(\bar{\eta}) S_{\mathbf{k}}(\bar{\eta})$$

$$G_{\mathbf{k}}''(\eta, \bar{\eta}) + \left(k^2 - \frac{a''(\eta)}{a(\eta)} \right) G_{\mathbf{k}}(\eta, \bar{\eta}) = \delta(\eta - \bar{\eta})$$

Power spectrum & “universal part”

[Ananda, Clarkson, Wands, gr-qc/0612013] [Baumann, Steinhardt, Takahashi, Ichiki, hep-th/0703290]
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 [Inomata, Kawasaki, Mukaida, Tada, Yanagida, 1611.06130]

$$\mathcal{P}_h(\eta, k) = 4 \int_0^\infty dv \int_{|1-v|}^{1+v} du \left(\frac{4v^2 - (1 + v^2 - u^2)^2}{4vu} \right)^2 I^2(v, u, x) \mathcal{P}_\zeta(kv) \mathcal{P}_\zeta(ku)$$

where $x \equiv k\eta$ $u = |\mathbf{k} - \tilde{\mathbf{k}}|/k$ $v = \tilde{k}/k$

We want to analytically calculate this function.

$$\text{Definitions: } I(v, u, x) = \int_0^x d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} k G_k(\eta, \bar{\eta}) f(v, u, \bar{x})$$

$$f(v, u, \bar{x}) = \frac{6(w+1)}{3w+5} \Phi(v\bar{x})\Phi(u\bar{x}) + \frac{6(1+3w)(w+1)}{(3w+5)^2} (\bar{x}\partial_{\bar{\eta}}\Phi(v\bar{x})\Phi(u\bar{x}) + \bar{x}\partial_{\bar{\eta}}\Phi(u\bar{x})\Phi(v\bar{x})) \\ + \frac{3(1+3w)^2(1+w)}{(3w+5)^2} \bar{x}^2 \partial_{\bar{\eta}}\Phi(v\bar{x})\partial_{\bar{\eta}}\Phi(u\bar{x})$$

We have used $\mathcal{H} = aH = 2/((1+3w)\eta)$ $\bar{x} \equiv k\bar{\eta}$ $w \equiv P/\rho$

Analytic Calculation

For definiteness, consider $w = 1/3$ (RD era).

$$I(v, u, x) = \int_0^x d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} k G_k(\eta, \bar{\eta}) f(v, u, \bar{x})$$

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$$kG_k(\eta, \bar{\eta}) = \sin(x - \bar{x})$$

Analytic Calculation

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$$a(\bar{\eta})/a(\eta) = \bar{x}/x$$

$$kG_k(\eta, \bar{\eta}) = \sin(x - \bar{x})$$

$$\Phi(x) = \frac{9}{x^2} \left(\frac{\sin(x/\sqrt{3})}{x/\sqrt{3}} - \cos(x/\sqrt{3}) \right)$$

$$f_{\text{RD}}(v, u, x) = \frac{12}{u^3 v^3 x^6} \left(18uvx^2 \cos \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} + (54 - 6(u^2 + v^2)x^2 + u^2 v^2 x^4) \sin \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} \right. \\ \left. + 2\sqrt{3}ux(v^2 x^2 - 9) \cos \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} + 2\sqrt{3}vx(u^2 x^2 - 9) \sin \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} \right)$$

Analytic Calculation

For definiteness, consider $w = 1/3$ (RD era).

$$\begin{aligned}
 I_{\text{RD}}(v, u, x) = & \frac{3}{4u^3v^3x} \left(-\frac{4}{x^3} \left(uv(u^2 + v^2 - 3)x^3 \sin x - 6uvx^2 \cos \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} \right. \right. \\
 & + 6\sqrt{3}ux \cos \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} + 6\sqrt{3}vx \sin \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} - 3(6 + (u^2 + v^2 - 3)x^2) \sin \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} \left. \right) \\
 & + (u^2 + v^2 - 3)^2 \left(\sin x \left(\text{Ci} \left(\left(1 - \frac{v-u}{\sqrt{3}} \right) x \right) + \text{Ci} \left(\left(1 + \frac{v-u}{\sqrt{3}} \right) x \right) \right. \right. \\
 & - \text{Ci} \left(\left| 1 - \frac{v+u}{\sqrt{3}} \right| x \right) - \text{Ci} \left(\left(1 + \frac{v+u}{\sqrt{3}} \right) x \right) + \log \left| \frac{3 - (u+v)^2}{3 - (u-v)^2} \right| \left. \right) \\
 & + \cos x \left(-\text{Si} \left(\left(1 - \frac{v-u}{\sqrt{3}} \right) x \right) - \text{Si} \left(\left(1 + \frac{v-u}{\sqrt{3}} \right) x \right) \right. \\
 & \left. \left. + \text{Si} \left(\left(1 - \frac{v+u}{\sqrt{3}} \right) x \right) + \text{Si} \left(\left(1 + \frac{v+u}{\sqrt{3}} \right) x \right) \right) \right)
 \end{aligned}$$

$$\text{Si}(x) = \int_0^x d\bar{x} \frac{\sin \bar{x}}{\bar{x}}$$

$$\text{Ci}(x) = - \int_x^\infty d\bar{x} \frac{\cos \bar{x}}{\bar{x}}$$

Analytic Calculation

For definiteness, consider $w = 1/3$ (RD era).

$$\begin{aligned}
 I_{\text{RD}}(v, u, x) = & \frac{3}{4u^3v^3x} \left(-\frac{4}{x^3} \left(uv(u^2 + v^2 - 3)x^3 \sin x - 6uvx^2 \cos \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} \right. \right. \\
 & + 6\sqrt{3}ux \cos \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} + 6\sqrt{3}vx \sin \frac{ux}{\sqrt{3}} \cos \frac{vx}{\sqrt{3}} - 3(6 + (u^2 + v^2 - 3)x^2) \sin \frac{ux}{\sqrt{3}} \sin \frac{vx}{\sqrt{3}} \left. \right) \\
 & + (u^2 + v^2 - 3)^2 \left(\sin x \left(\text{Ci} \left(\left(1 - \frac{v-u}{\sqrt{3}} \right) x \right) + \text{Ci} \left(\left(1 + \frac{v-u}{\sqrt{3}} \right) x \right) \right. \right. \\
 & - \text{Ci} \left(\left| 1 - \frac{v+u}{\sqrt{3}} \right| x \right) - \text{Ci} \left(\left(1 + \frac{v+u}{\sqrt{3}} \right) x \right) + \log \left| \frac{3 - (u+v)^2}{3 - (u-v)^2} \right| \left. \right) \\
 & + \cos x \left(-\text{Si} \left(\left(1 - \frac{v-u}{\sqrt{3}} \right) x \right) - \text{Si} \left(\left(1 + \frac{v-u}{\sqrt{3}} \right) x \right) \right. \\
 & \left. \left. + \text{Si} \left(\left(1 - \frac{v+u}{\sqrt{3}} \right) x \right) + \text{Si} \left(\left(1 + \frac{v+u}{\sqrt{3}} \right) x \right) \right) \right)
 \end{aligned}$$

$$\text{Si}(x) = \int_0^x d\bar{x} \frac{\sin \bar{x}}{\bar{x}}$$

$$\text{Ci}(x) = - \int_x^\infty d\bar{x} \frac{\cos \bar{x}}{\bar{x}}$$

Oscillation average in the late-time limit

$$\begin{aligned}
 \overline{I_{\text{RD}}^2}(v, u, x \rightarrow \infty) = & \frac{1}{2} \left(\frac{3(u^2 + v^2 - 3)}{4u^3v^3x} \right)^2 \left(\left(-4uv + (u^2 + v^2 - 3) \log \left| \frac{3 - (u+v)^2}{3 - (u-v)^2} \right| \right)^2 \right. \\
 & \left. + \pi^2 (u^2 + v^2 - 3)^2 \Theta(v + u - \sqrt{3}) \right)
 \end{aligned}$$

Analytically

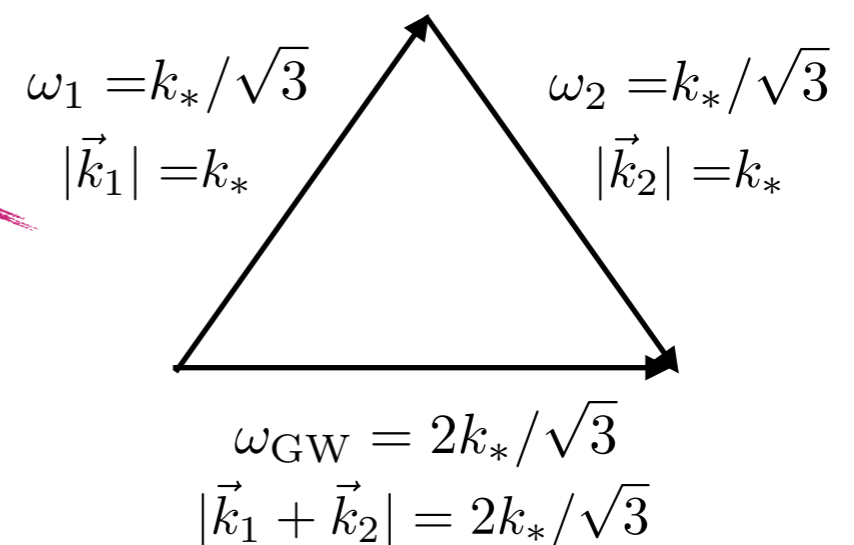
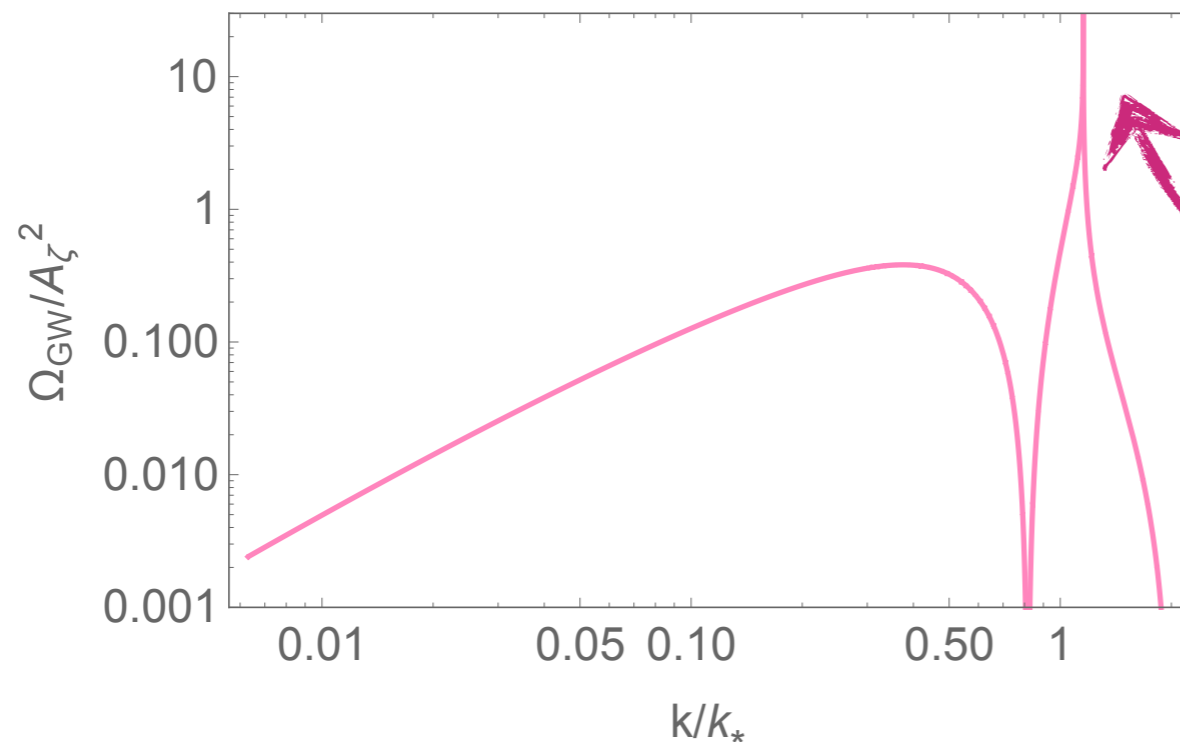
Calculable GW spectrum

Example 1 in RD: Monochromatic case

$$\mathcal{P}_\zeta(k) = A_\zeta \delta(\log k/k_*)$$

$$\Omega_{\text{GW}}(\eta, k) = \frac{3A_\zeta^2}{64} \left(\frac{4 - \tilde{k}^2}{4} \right)^2 \tilde{k}^2 (3\tilde{k}^2 - 2)^2 \times \left(\pi^2 (3\tilde{k}^2 - 2)^2 \Theta(2\sqrt{3} - 3\tilde{k}) + \left(4 + (3\tilde{k}^2 - 2) \log \left| 1 - \frac{4}{3\tilde{k}^2} \right| \right)^2 \right) \Theta(2 - \tilde{k})$$

$$\tilde{k} \equiv k/k_*$$



Analytically

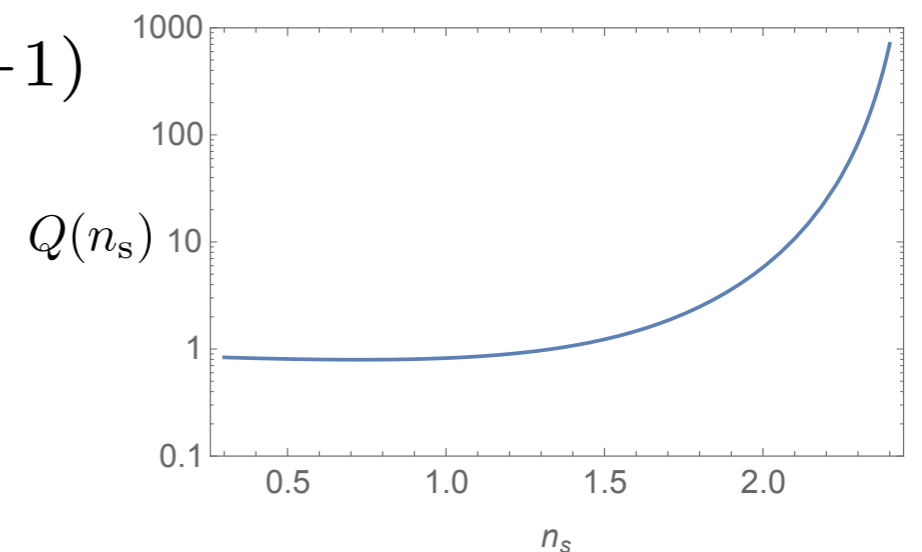
Calculable GW spectrum

Example 2 in RD: Scale-invariant case $\mathcal{P}_\zeta(k) = A_\zeta$

$$\Omega_{\text{GW}}(\eta, k) \simeq 0.8222 A_\zeta^2$$

Example 3 in RD: Power-law case $\mathcal{P}_\zeta = A_\zeta \left(\frac{k}{k_*} \right)^{n_s - 1}$

$$\Omega_{\text{GW}}(\eta, k) = Q(n_s) A_\zeta^2 \left(\frac{k}{k_*} \right)^{2(n_s - 1)}$$



Enhancement in MD era

[Assadullahi, Wands, 0901.0989] [Baumann, Steinhardt, Takahashi, Ichiki, hep-th/0703290]

Next, consider $w = 0$ (MD era).


$$I(v, u, x) = \int_0^x d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} k G_k(\eta, \bar{\eta}) f(v, u, \bar{x})$$

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$$kG_{\mathbf{k}}(\eta, \bar{\eta}) = \frac{1}{x\bar{x}} ((1 + x\bar{x}) \sin(x - \bar{x}) - (x - \bar{x}) \cos(x - \bar{x}))$$

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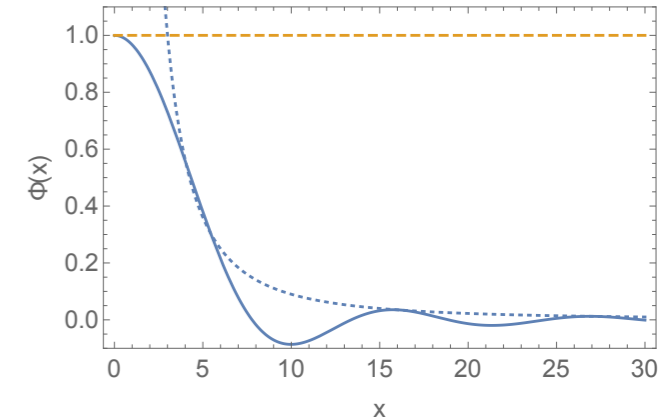
$$I(v, u, x) = \int_0^x d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} kG_k(\eta, \bar{\eta}) f(v, u, \bar{x})$$

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$$\Phi(x) = 1$$

$$f_{\text{MD}}(v, u, x) = \frac{6}{5}$$



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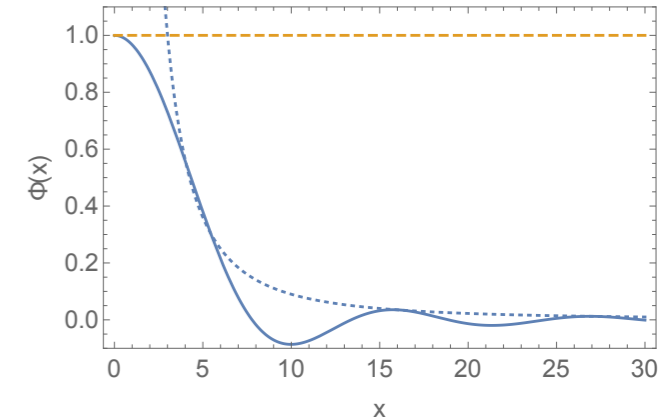
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Non-decaying source!



Enhancement in MD era

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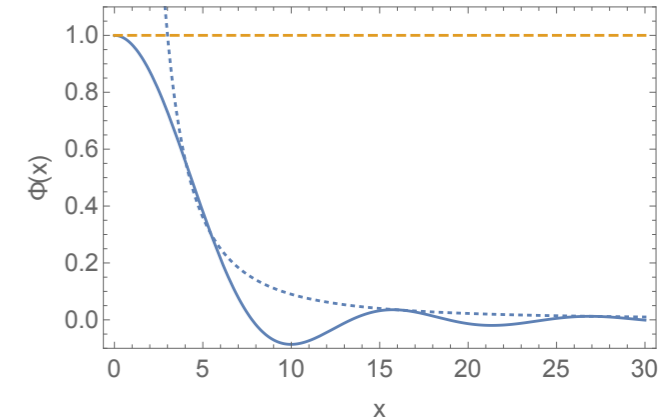
$$\Phi(x) = 1$$

$$f_{\text{MD}}(v, u, x) = \frac{6}{5}$$

Non-decaying source!

$$\overline{I_{\text{MD}}^2(v, u, x \rightarrow \infty)} = \frac{18}{25}$$

← Not diluted!



Enhancement in MD era

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Next, consider $w = 0$ (MD era).

$$I(v, u, x) = \int_0^x d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} kG_k(\eta, \bar{\eta}) f(v, u, \bar{x})$$

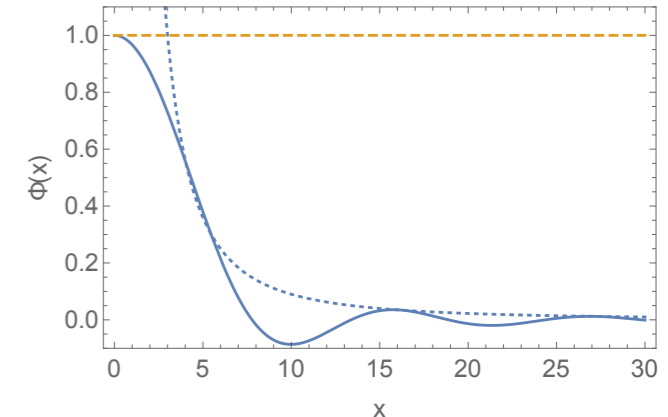
$$a(\bar{\eta})/a(\eta) = \bar{x}^2/x^2$$

$$kG_{\mathbf{k}}(\eta, \bar{\eta}) = \frac{1}{x\bar{x}} ((1 + x\bar{x}) \sin(x - \bar{x}) - (x - \bar{x}) \cos(x - \bar{x}))$$

$$\Phi(x) = 1$$

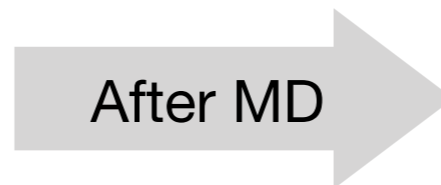
$$f_{\text{MD}}(v, u, x) = \frac{6}{5}$$

Non-decaying source!



$$\overline{I_{\text{MD}}^2(v, u, x \rightarrow \infty)} = \frac{18}{25} \quad \leftarrow \text{Not diluted!}$$

Constant metric distortion



Propagating GW

Calculable GW spectrum

Example 1 in MD: Monochromatic case $\mathcal{P}_\zeta(k) = A_\zeta \delta(\log k/k_*)$

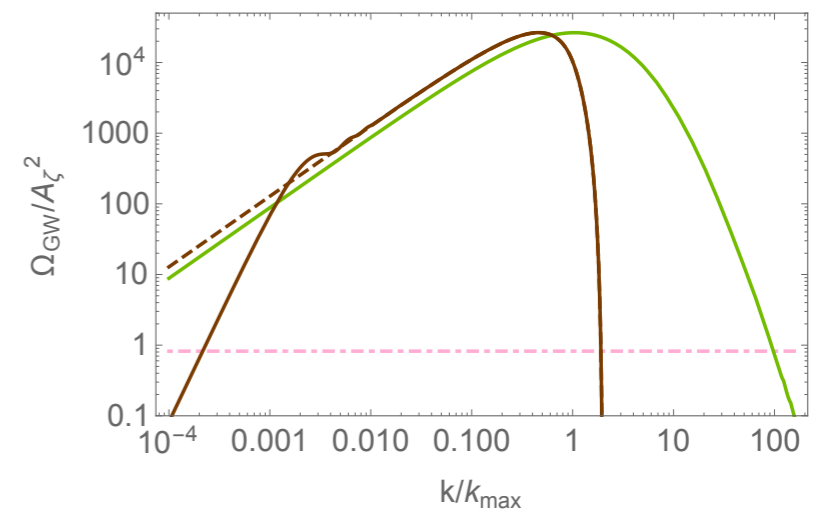
$$\Omega_{\text{GW}} = \frac{3}{25} \left(\frac{k_*}{aH} \right)^2 \left(1 - \left(\frac{k}{2k_*} \right)^2 \right)^2 A_\zeta^2 \Theta(2k_* - k)$$

Example 2 in MD: Scale-invariant case with a cutoff

$$\mathcal{P}_\zeta(k) = A_\zeta \Theta(k_{\text{max}} - k)$$

$$\Omega_{\text{GW}} = \frac{A_\zeta^2}{14000} \left(\frac{k}{aH} \right)^2 \times \begin{cases} \left(1792\tilde{k}^{-1} - 2520 + 768\tilde{k} + 105\tilde{k}^2 \right) & (0 < k \leq k_{\text{max}}) \\ \left(1 - 2\tilde{k}^{-1} \right)^4 \left(105\tilde{k}^2 + 72\tilde{k} + 16 - 32\tilde{k}^{-1} - 16\tilde{k}^{-2} \right) & (k_{\text{max}} < k \leq 2k_{\text{max}}) \end{cases}$$

$$\tilde{k} \equiv k/k_*$$

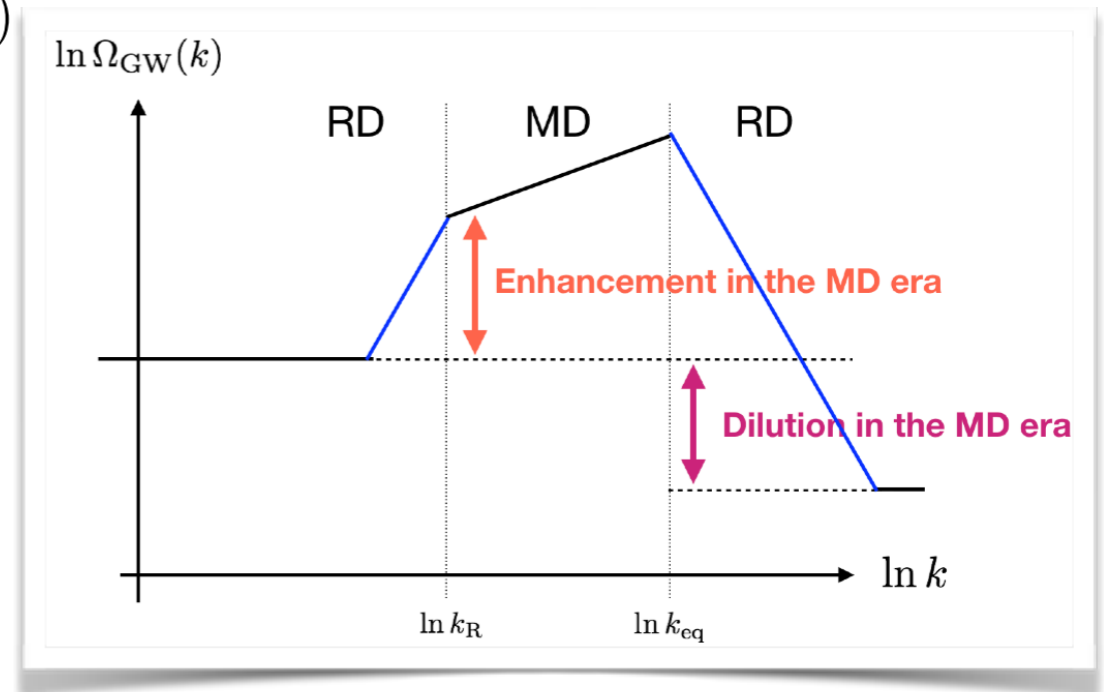


General cases

The MD → RD transition

$$I(v, u, x) = \int_0^{x_R} d\bar{x} \left(\frac{x_R}{x}\right) \left(\frac{\bar{x}}{x_R}\right)^2 k G_k^{\text{MD} \rightarrow \text{RD}}(\eta, \bar{\eta}) f_{\text{MD}}(v, u, \bar{x})$$

$$+ \int_{x_R}^x d\bar{x} \left(\frac{\bar{x}}{x}\right) k G_k^{\text{RD}}(\eta, \bar{\eta}) f_{\text{MD} \rightarrow \text{RD}}(v, u, \bar{x})$$



The RD → MD transition

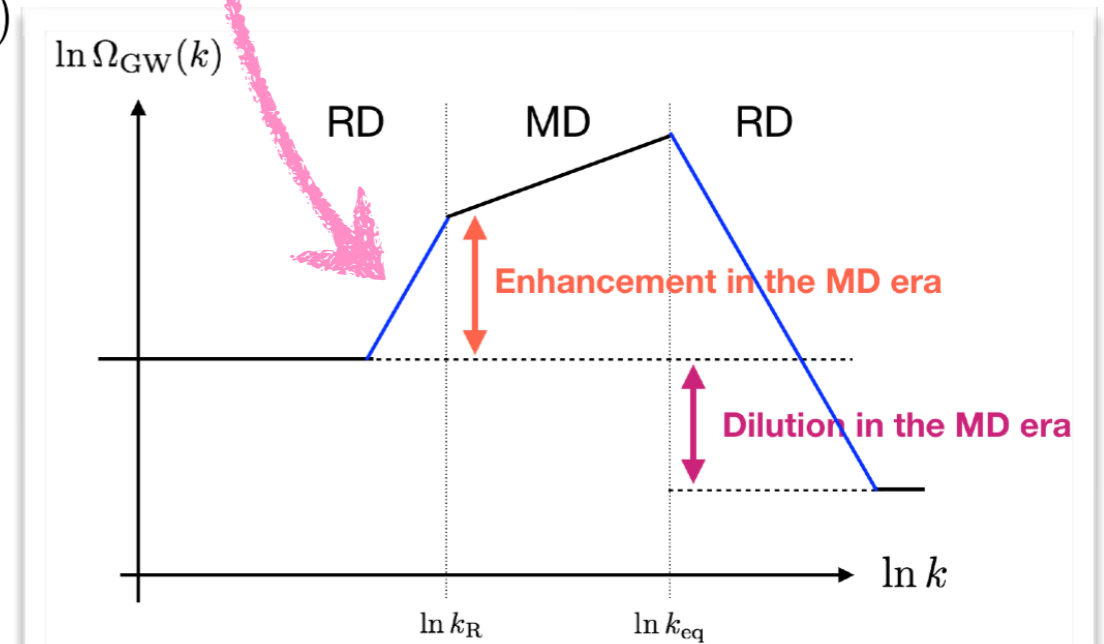
$$I(v, u, x) = \int_0^{x_{\text{eq}}} d\bar{x} \left(\frac{x_{\text{eq}}}{x}\right)^2 \left(\frac{\bar{x}}{x_{\text{eq}}}\right) k G_k^{\text{RD} \rightarrow \text{MD}}(\eta, \bar{\eta}) f_{\text{RD}}(v, u, \bar{x})$$

$$+ \int_{x_{\text{eq}}}^x d\bar{x} \left(\frac{\bar{x}}{x}\right)^2 k G_k^{\text{MD}}(\eta, \bar{\eta}) f_{\text{RD} \rightarrow \text{MD}}(v, u, \bar{x})$$

General cases

The MD → RD transition

$$I(v, u, x) = \int_0^{x_R} d\bar{x} \left(\frac{x_R}{x} \right) \left(\frac{\bar{x}}{x_R} \right)^2 k G_k^{\text{MD} \rightarrow \text{RD}}(\eta, \bar{\eta}) f_{\text{MD}}(v, u, \bar{x}) + \int_{x_R}^x d\bar{x} \left(\frac{\bar{x}}{x} \right) k G_k^{\text{RD}}(\eta, \bar{\eta}) f_{\text{MD} \rightarrow \text{RD}}(v, u, \bar{x})$$



The RD → MD transition

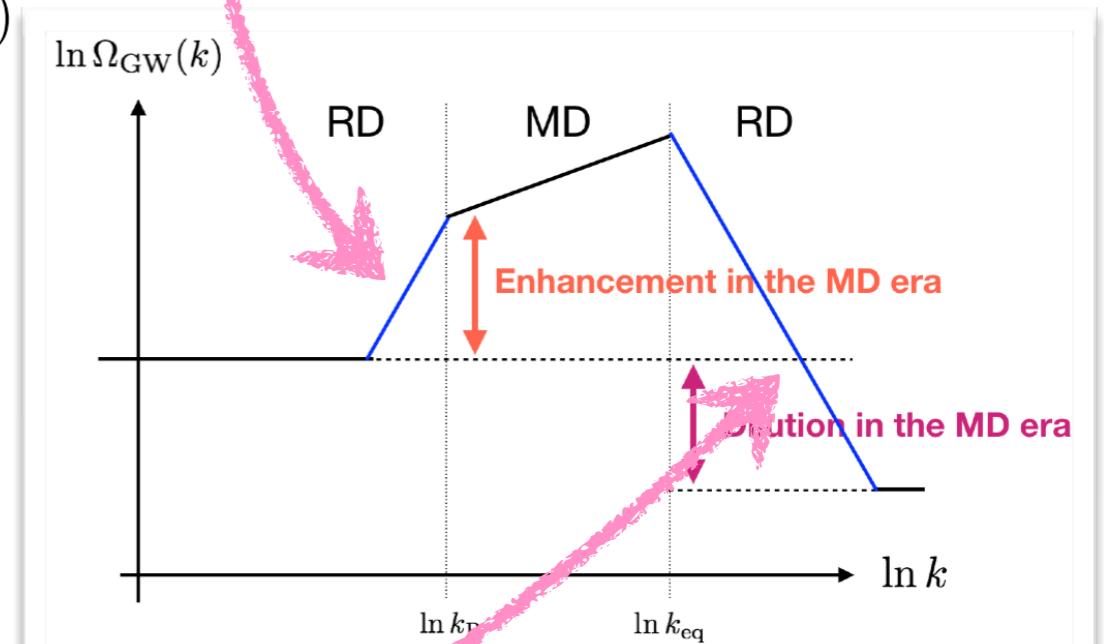
$$I(v, u, x) = \int_0^{x_{\text{eq}}} d\bar{x} \left(\frac{x_{\text{eq}}}{x} \right)^2 \left(\frac{\bar{x}}{x_{\text{eq}}} \right) k G_k^{\text{RD} \rightarrow \text{MD}}(\eta, \bar{\eta}) f_{\text{RD}}(v, u, \bar{x}) + \int_{x_{\text{eq}}}^x d\bar{x} \left(\frac{\bar{x}}{x} \right)^2 k G_k^{\text{MD}}(\eta, \bar{\eta}) f_{\text{RD} \rightarrow \text{MD}}(v, u, \bar{x})$$

General cases

The MD → RD transition

$$I(v, u, x) = \int_0^{x_R} d\bar{x} \left(\frac{x_R}{x}\right) \left(\frac{\bar{x}}{x_R}\right)^2 k G_k^{\text{MD} \rightarrow \text{RD}}(\eta, \bar{\eta}) f_{\text{MD}}(v, u, \bar{x})$$

$$+ \int_{x_R}^x d\bar{x} \left(\frac{\bar{x}}{x}\right) k G_k^{\text{RD}}(\eta, \bar{\eta}) f_{\text{MD} \rightarrow \text{RD}}(v, u, \bar{x})$$



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General cases

The MD → RD transition

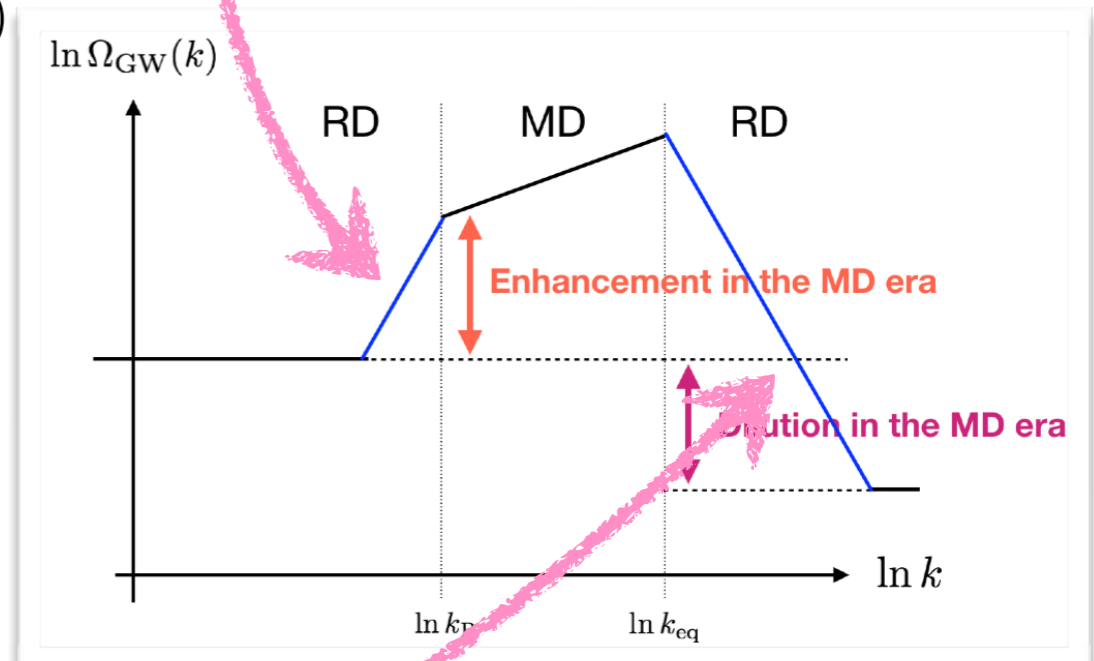
$$I(v, u, x) = \int_{x_R}^{x_{eq}} d\bar{x} \left(\frac{x_{eq}}{\bar{x}} \right)^2 \left(\frac{\bar{x}}{x_R} \right) k G_k^{RD \rightarrow MD}(\eta, \bar{\eta}) f_{MD \rightarrow RD}(v, u, \bar{x})$$

$$\frac{3}{5x x_R^3} (3(2x_R^2 - 1) \cos x - 6x_R \sin x + 2x_R^4 \cos(x - x_R) + 4x_R^3 \sin(x - x_R) + 3 \cos(x - 2x_R))$$

The RD → MD transition

$$I(v, u, x) = \int_0^{x_{eq}} d\bar{x} \left(\frac{x_{eq}}{\bar{x}} \right)^2 \left(\frac{\bar{x}}{x_{eq}} \right) k G_k^{RD \rightarrow MD}(\eta, \bar{\eta}) f_{RD}(v, u, \bar{x})$$

$$+ \int_{x_{eq}}^x d\bar{x} \left(\frac{\bar{x}}{x} \right)^2 k G_k^{MD}(\eta, \bar{\eta}) f_{RD \rightarrow MD}(v, u, \bar{x})$$



General cases

The MD → RD transition

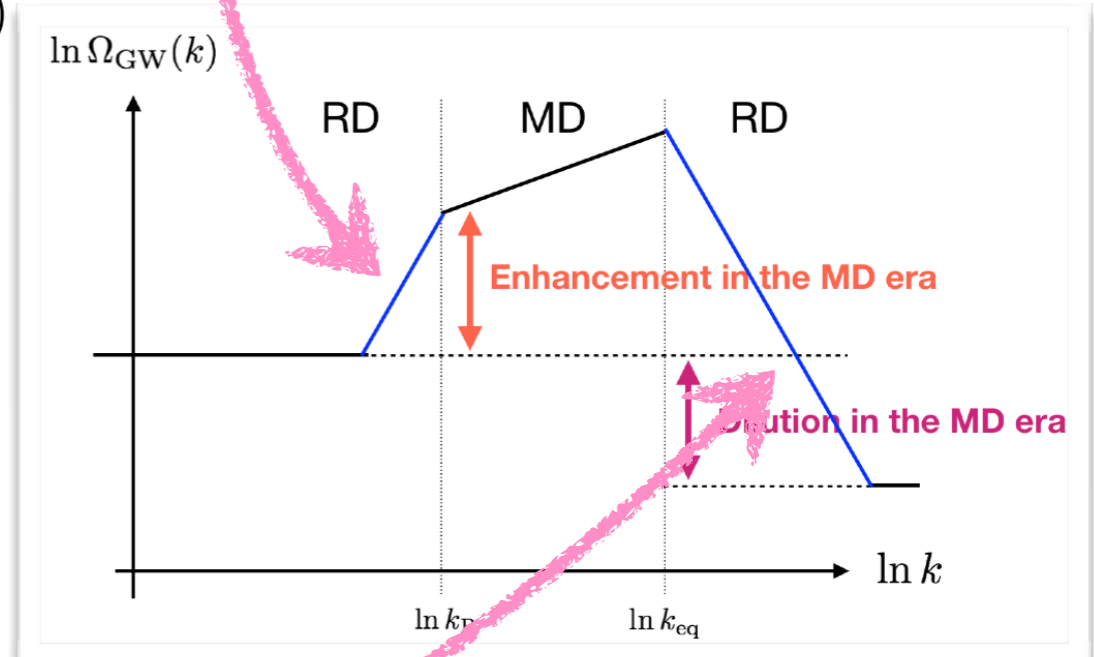
$$I(v, u, x) = \int_{x_R}^x d\bar{x} \left(\frac{x}{x_R} \right)^{3\Omega_{GW}^{RD}(\eta, \bar{\eta})} f_{MD \rightarrow RD}(v, u, \bar{x})$$

$$\frac{3}{5x x_R^3} \left(3(2x_R^2 - 1) \cos x - 6x_R \sin x + 2x_R^4 \cos(x - x_R) + 4x_R^3 \sin(x - x_R) + 3 \cos(x - 2x_R) \right)$$

The RD → MD transition

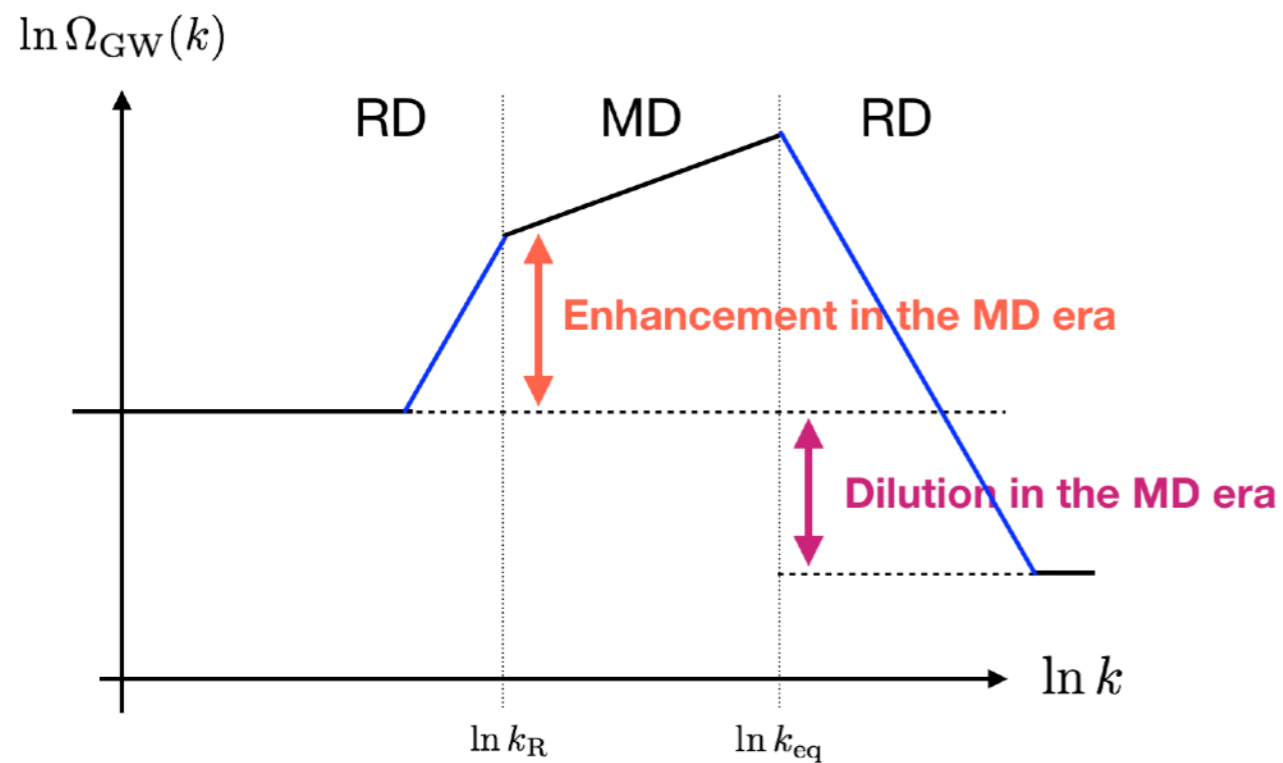
$$\frac{(x^3 - 3(x - x_{eq}) - x x_{eq}^2) \cos(x - x_{eq}) - (3 + 3x x_{eq} - x_{eq}^2) \sin(x - x_{eq})}{x^3} \times \frac{6 \ln(c_1 u x_{eq}) \ln(c_1 v x_{eq})}{5 (c_2 u x_{eq})^2 (c_2 v x_{eq})^2}$$

$$+ \int_{x_{eq}}^x d\bar{x} \left(\frac{x}{x_{eq}} \right)^{3\Omega_{GW}^{RD}(\eta, \bar{\eta})} f_{RD \rightarrow MD}(v, u, \bar{x})$$

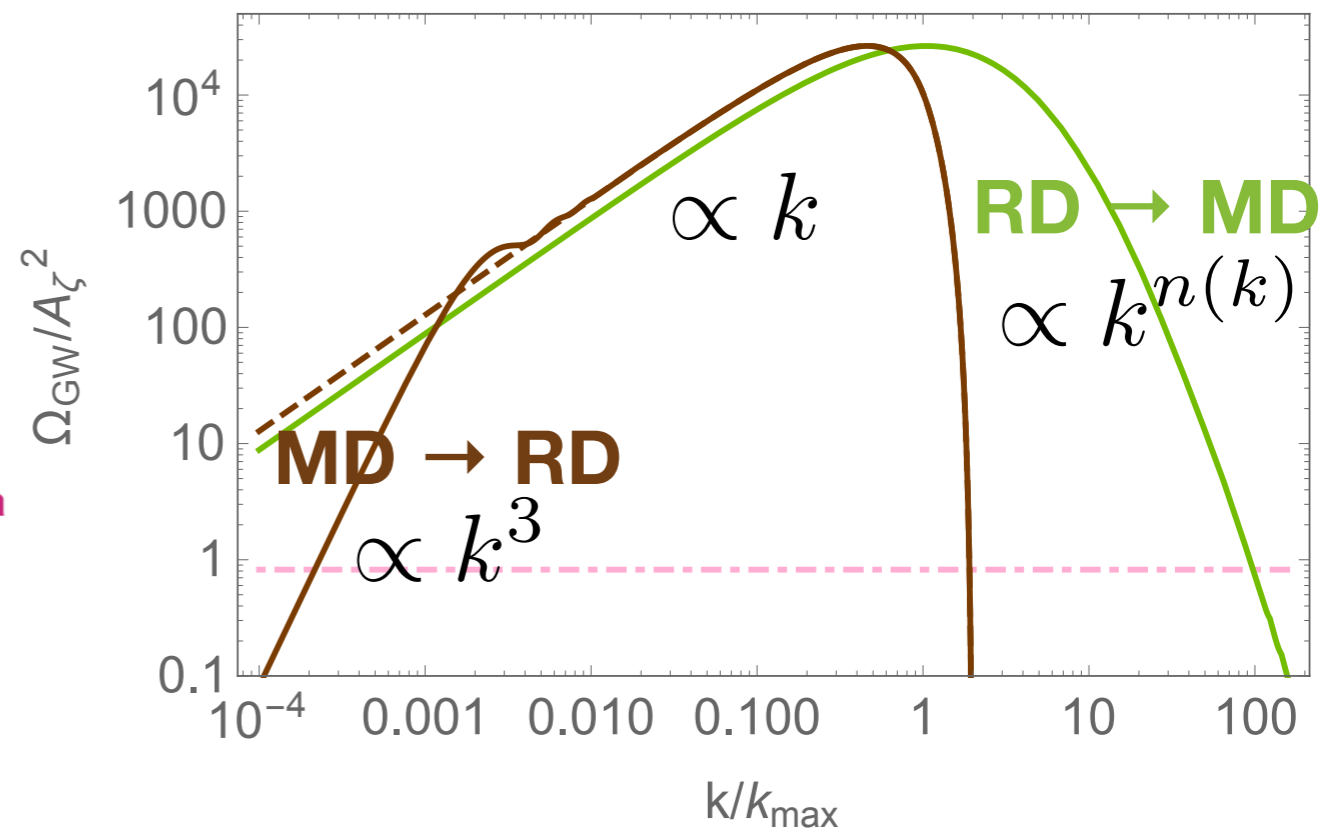


General cases

Schematic



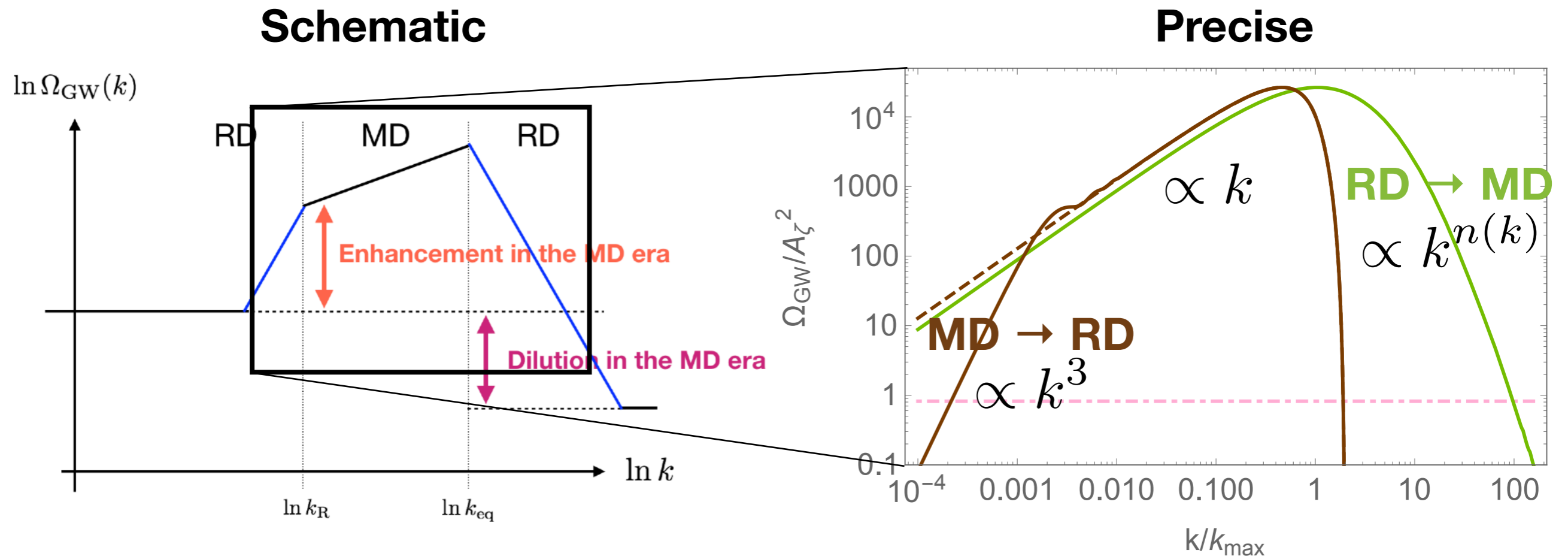
Precise



$$-6 \lesssim n(k) \lesssim -4$$

Scale-invariant primordial spectrum is assumed as a simple example.

General cases



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APPLICATION:
PPRIMORDIAL
BBLACK **H**HOLE
SCENARIOS

[Kohri, Terada, 1802.06785]

Current PBH abundance

PBH fraction in CDM

$$f_{\text{PBH}}(M) = \frac{1}{\rho_{\text{CDM}}} \frac{d\rho_{\text{PBH}}}{d \ln M}$$
$$= \left(\frac{g_*(T)}{g_*(T_{\text{eq}})} \frac{g_{*,s}(T_{\text{eq}})}{g_{*,s}(T)} \frac{T}{T_{\text{eq}}} \gamma \beta(\sigma(k(M))) \right) \Big|_{T=\text{Min}[T_M, T_R]} \frac{\Omega_m}{\Omega_{\text{CDM}}}$$

Current PBH abundance

PBH fraction in CDM

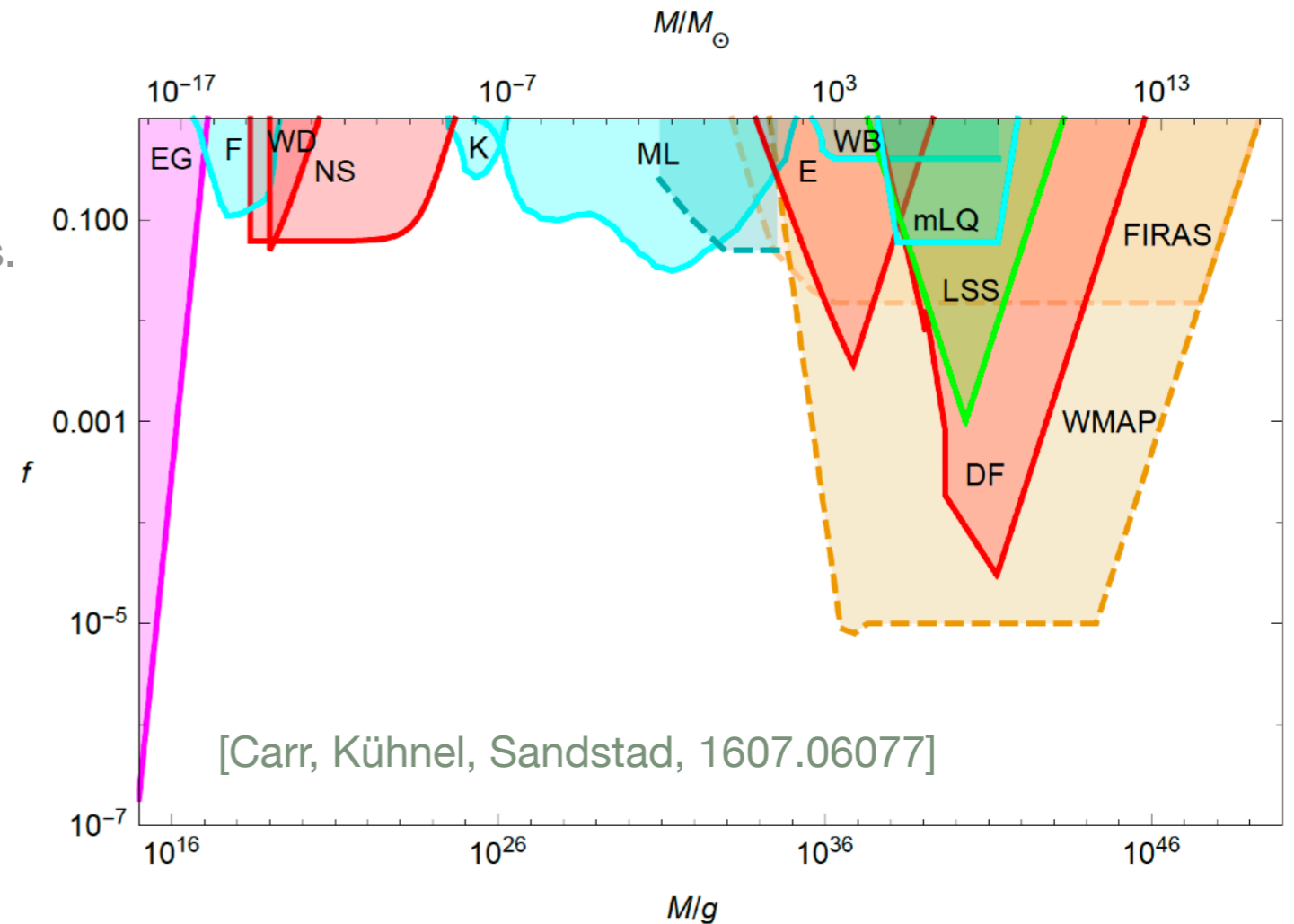
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PBH formation probability

Current PBH abundance

Observational constraints

Note: These are not the latest constraints.



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PBH formation probability

PBH formation probability

Coarse-grained perturbations

[Young, Byrnes, Sasaki, 1405.7023]
see also [Ando, Inomata, Kawasaki, 1802.06393]

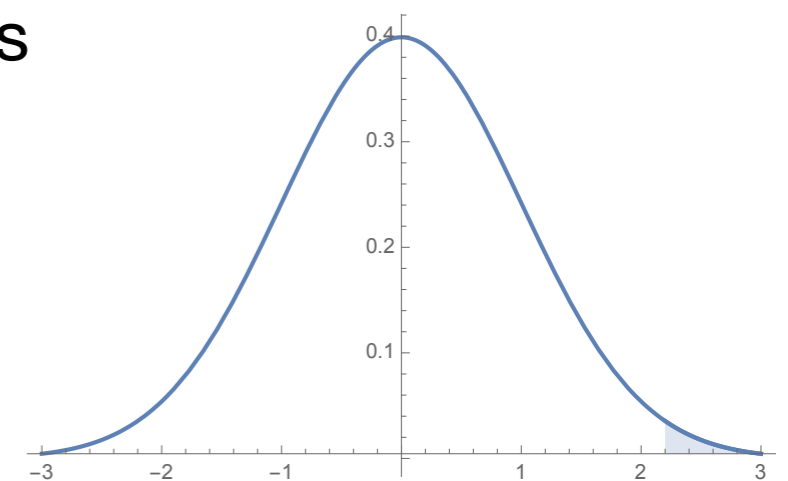
$$\sigma^2(k) = \int_{-\infty}^{\infty} d \ln q w^2 \left(\frac{q}{k}\right) \frac{4(1+w_{\text{eos}})^2}{(5+3w_{\text{eos}})^2} \left(\frac{q}{k}\right)^4 T^2(q/k) P_{\zeta}(q)$$
$$\sim \frac{2(1+w_{\text{eos}})^2}{(5+3w_{\text{eos}})^2} P_{\zeta}(k)$$

Formation probability in RD

Pressure prevents the formation → rare process

$$\beta(\sigma) = \int_{\delta_c}^{\infty} \frac{1}{\sqrt{2\pi}\sigma} \exp\left(-\frac{\delta^2}{2\sigma^2}\right) d\delta$$

[Press, Schechter, 1974]



(exaggerated figure for illustration)

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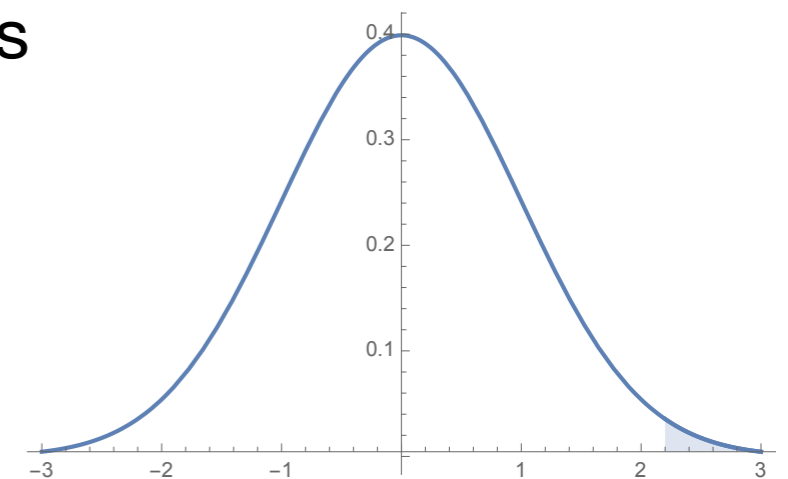
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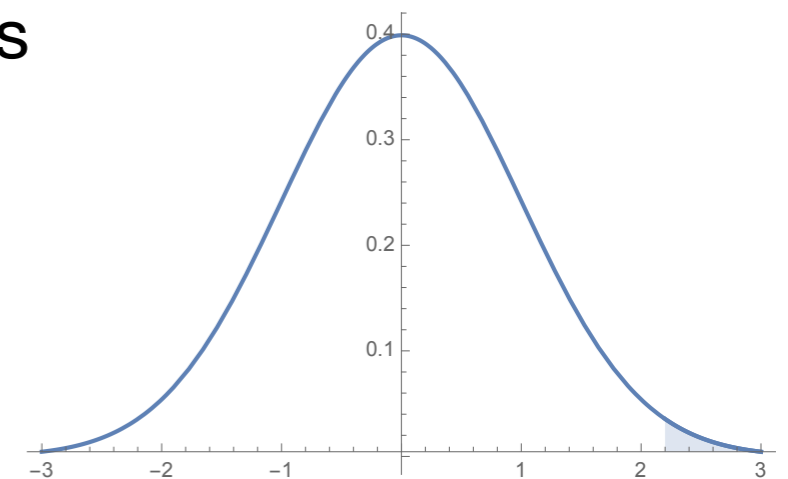
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Primordial curvature perturbations

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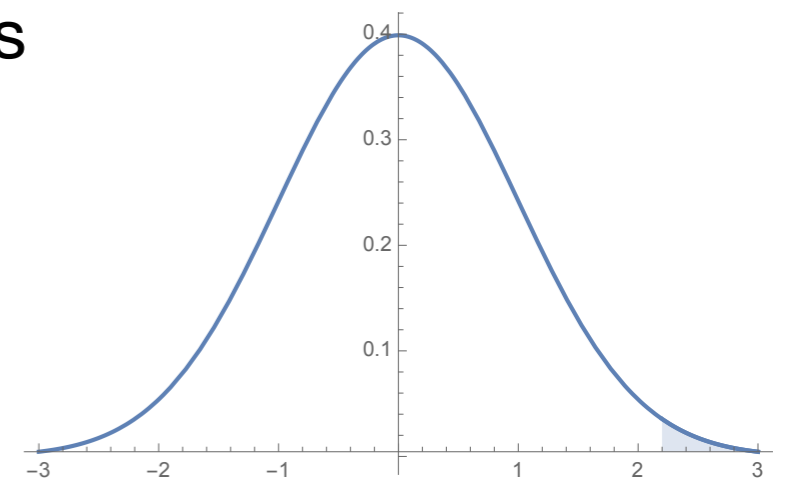
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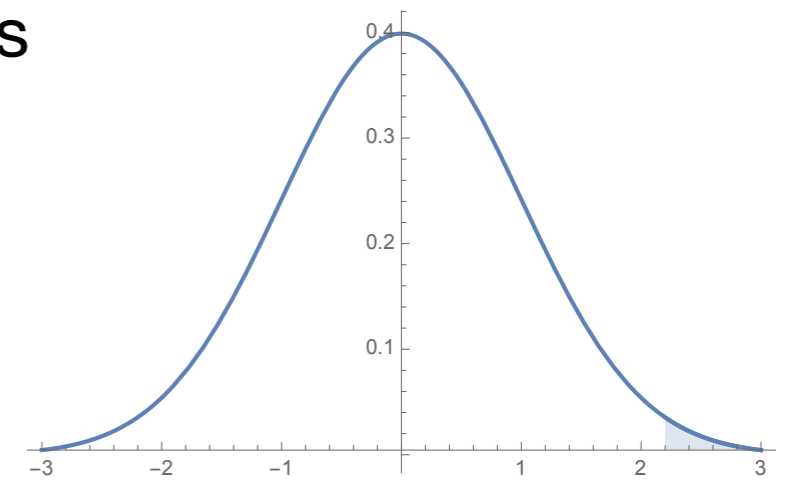
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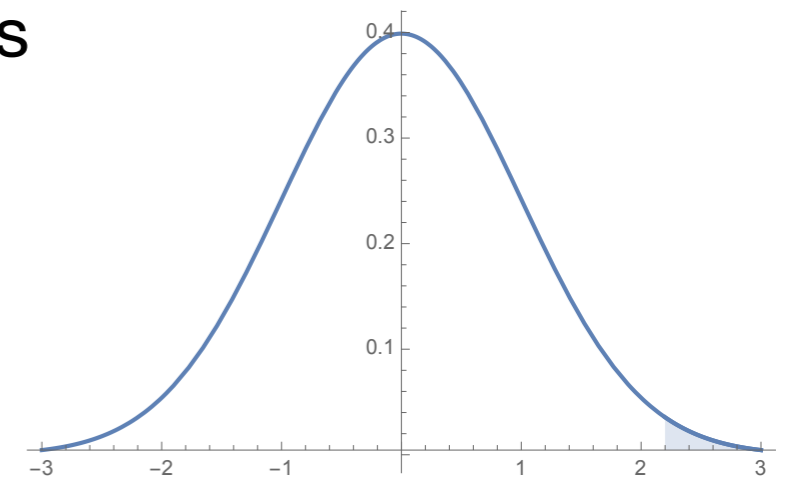
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Formation probability in MD

[Khlopov, Polnarev, 1980, 1985]

$$\frac{\delta\rho}{\rho} \propto a(t)$$

Perturbation grows in the MD era → low threshold

Anisotropy & angular momentum affects the formation rate.

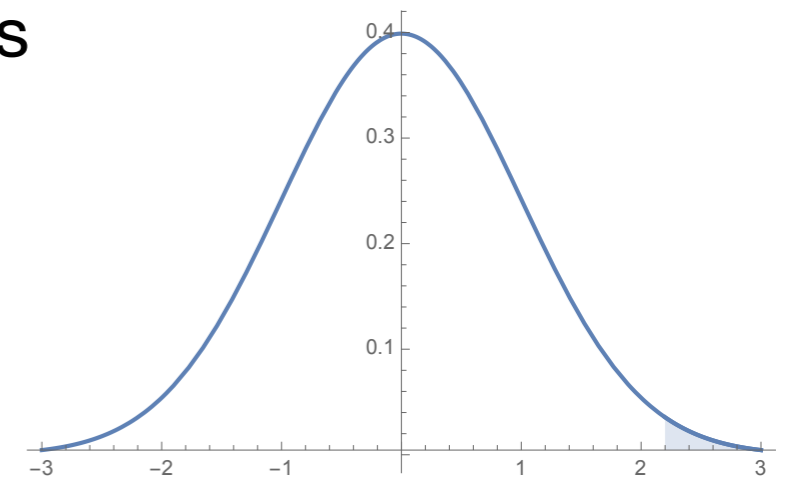
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[Harada, Yoo, Kohri, Nakao, Jhingan, 1609.01588]

[Harada, Yoo, Kohri, Nakao, 1707.03595]

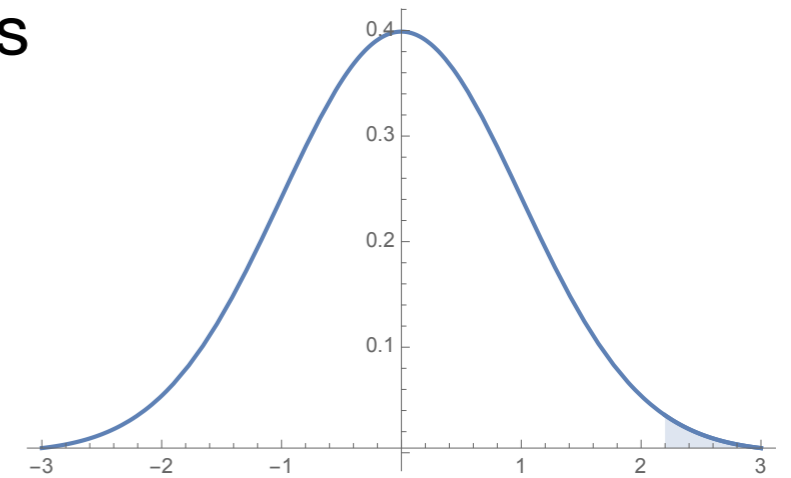
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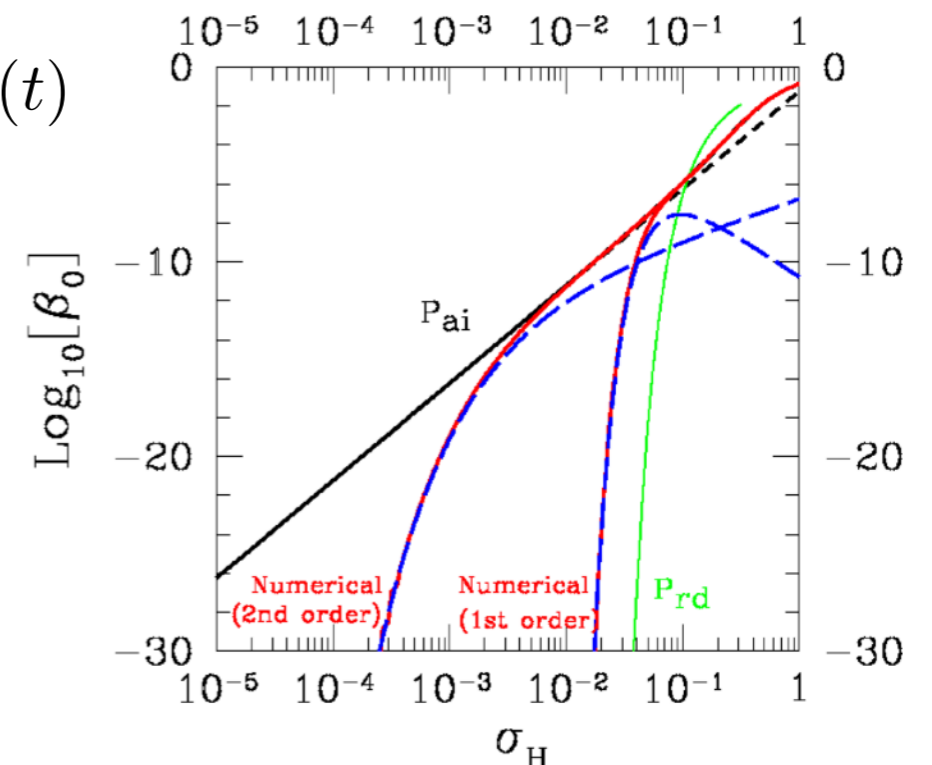
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Inflation with running spectrum

$$\mathcal{P}_{\mathcal{R}}(k) = A_s \left(\frac{k}{k_*} \right)^{n_s - 1 + \frac{\alpha_s}{2} \ln \frac{k}{k_*} + \frac{\beta_s}{6} \left(\ln \frac{k}{k_*} \right)^2 + \dots}$$

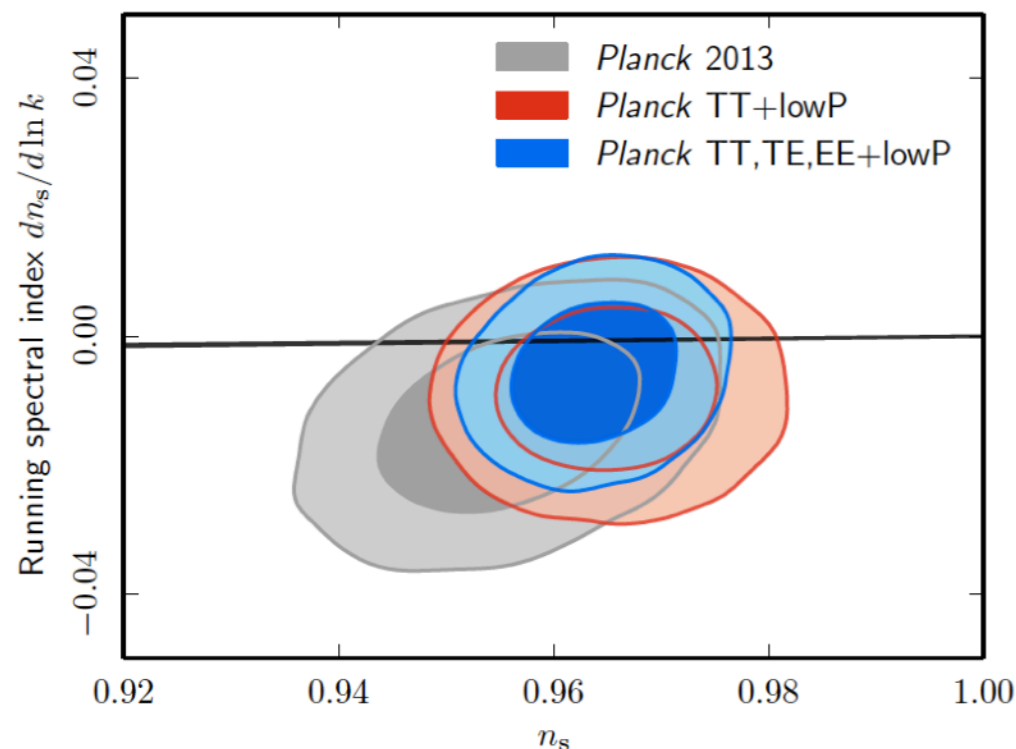


Fig. 4. Marginalized joint 68 % and 95 % CL for $(n_s, dn_s/d \ln k)$ using *Planck* TT+lowP and *Planck* TT,TE,EE+lowP. For comparison, the thin black stripe shows the prediction for single-field monomial chaotic inflationary models with $50 < N_* < 60$.

Planck TT+lowP (TT,TE,EE+lowP)

$$n_s = 0.9569 \pm 0.0077 \quad (0.9586 \pm 0.0056),$$

$$dn_s/d \ln k = 0.011^{+0.014}_{-0.013} \quad (0.009 \pm 0.010),$$

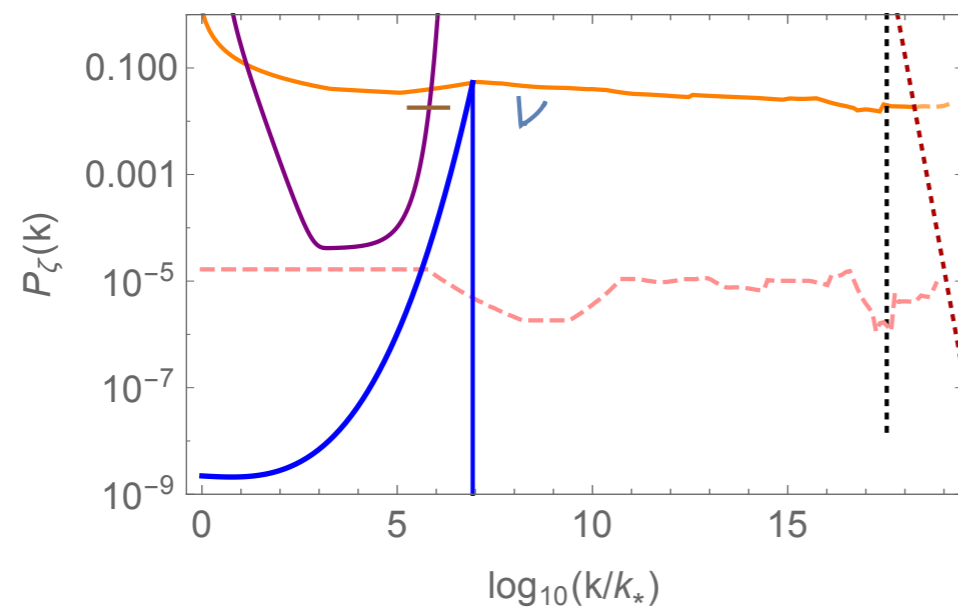
$$d^2 n_s/d \ln k^2 = 0.029^{+0.015}_{-0.016} \quad (0.025 \pm 0.013),$$

at 68% CL at the pivot scale $k^*=0.05/\text{Mpc}$.

Interplay between PBH & GW

PBH-for-LIGO scenario

Curvature perturbations



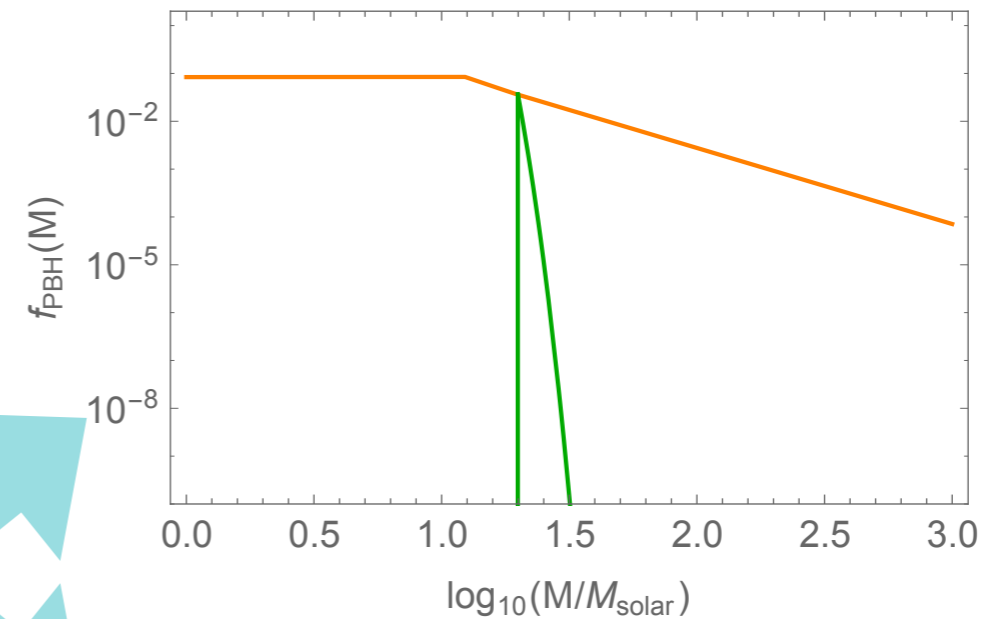
$$n_s = 0.96$$

$$\alpha_s = 0$$

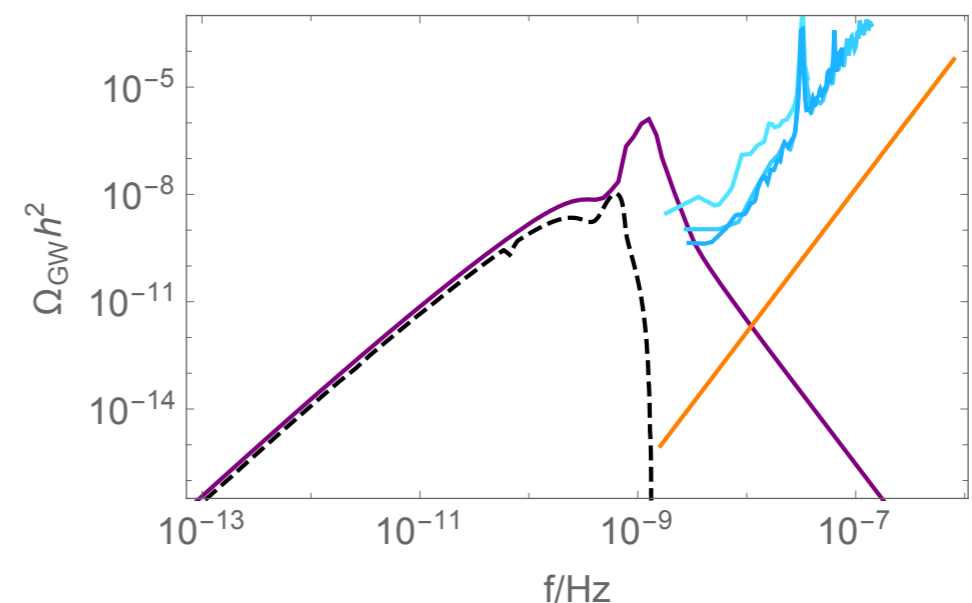
$$\beta_s = 0.026$$

$$T_R = 10^9 \text{ GeV}$$

PBH mass spectrum



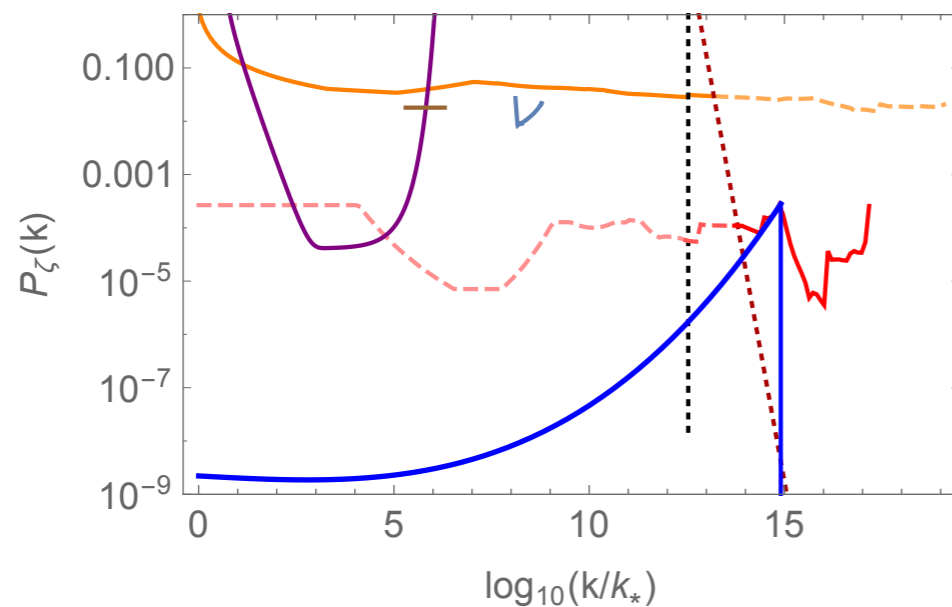
Secondary GW spectrum



Interplay between PBH & GW

PBH-DM scenario

Curvature perturbations



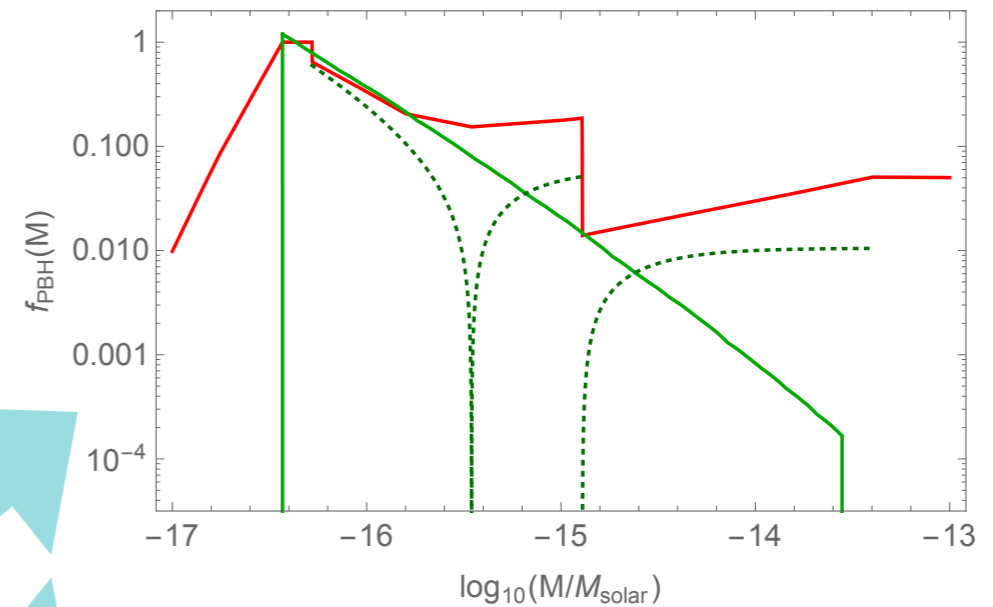
$$n_s = 0.96$$

$$\alpha_s = 0$$

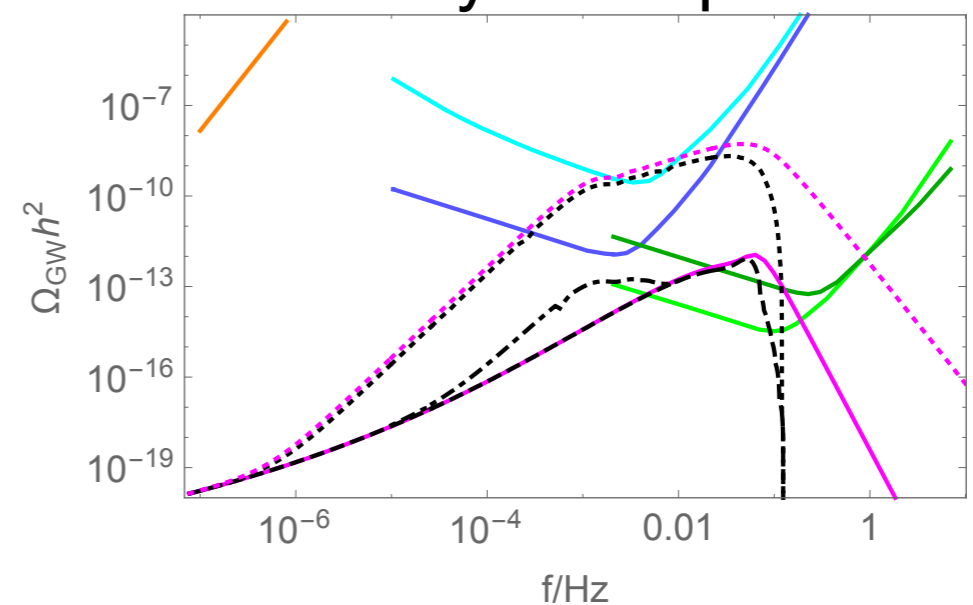
$$\beta_s = 0.0019485$$

$$T_R = 10^4 \text{ GeV}$$

PBH mass spectrum

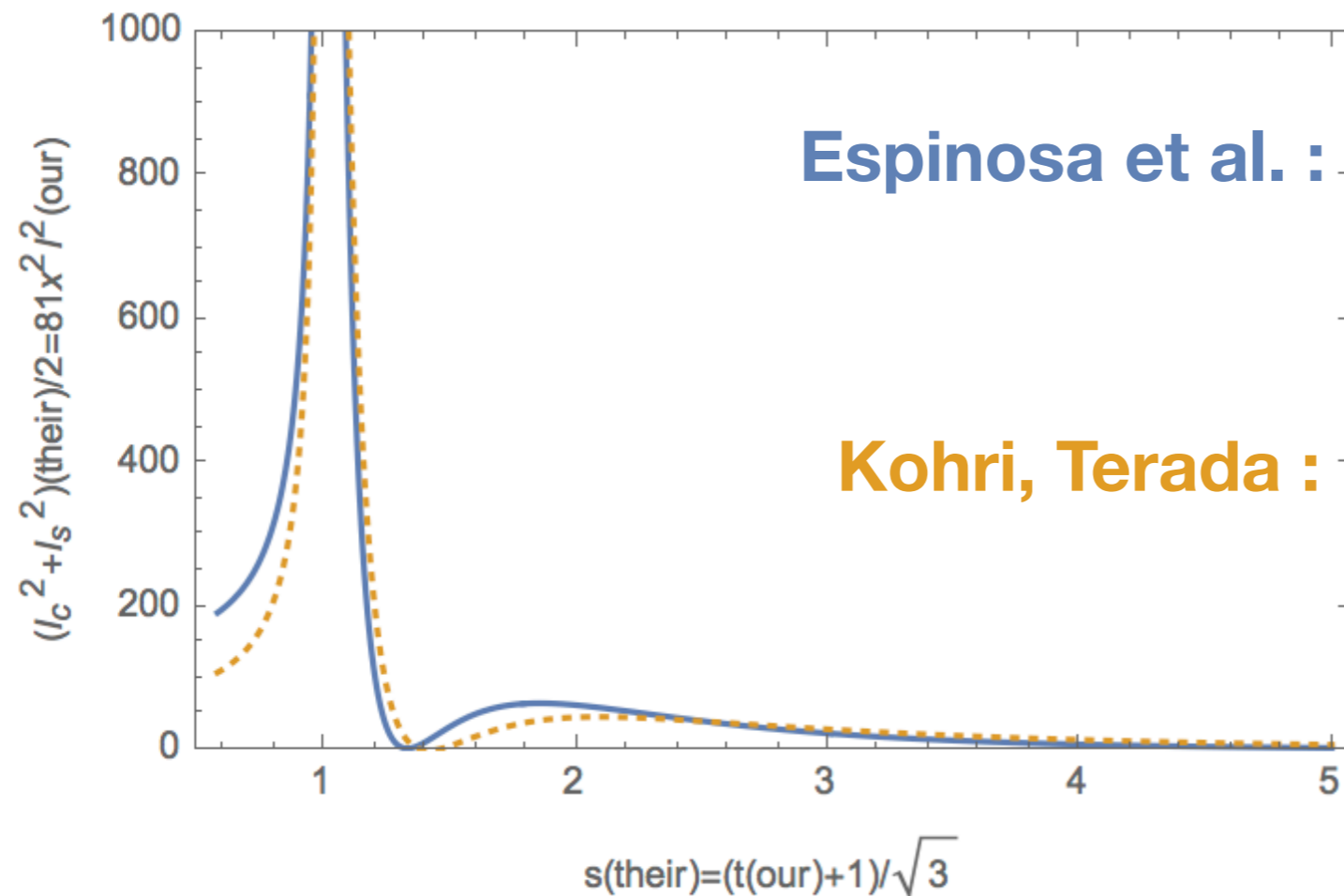


Secondary GW spectrum



Comparison with Espinosa et al.

$$I(v, u, x) = \int_{x_{\min}}^x d\bar{x} \frac{a(\bar{\eta})}{a(\eta)} k G_k(\eta, \bar{\eta}) f(v, u, \bar{x})$$



Espinosa et al. :

$$x_{\min} = k\eta_{\min} = 1$$

GW production right after inflation

Kohri, Terada :

$$x_{\min} = k\eta_{\min} = 0$$

inflation → RD/MD → GW production

[Espinosa, Racco, Riotto, 1804.07732]

[Kohri, Terada, 1804.08577]