

Running on liquid: Impact-induced hardening in dense suspensions

Y25.00002

Pradipto and H. Hayakawa, Phys. Rev. Fluids 6, 033301 (2021).

Pradipto and H. Hayakawa, Phys. Fluids 33, 093110 (2021).

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Y25 CONSTRAINT-BASED RHEOLOGY OF DENSE SUSPENSIONS AND GRANULAR MATERIALS II

Dense suspensions under impact





Fracture Roche et al, PRL 2013

Source: Itai Cohen Group on YouTube https://youtu.be/hP88C- LgnE

Impact-induced hardening

- Occurs on hard particles suspended in Newtonian solvent
- Cannot be observed in fluids or granular particles alone
- A person can run on top of dense suspensions

Physical explanations remain elusive

- Inherently far-from-equilibrium
- Highly dissipative —> transient
- Can shear thickening (CST and DST) explain such strong solidification?

Dense suspensions under impact have stronger response than the one observed in shear thickening suspensions

Waitukaitis and Jaeger, Nature 487, 205 (2012) Han et al, Nat. Comms. 2016





Dynamically Jammed Region

Localized and transient solid-like region beneath the impactor

Elastic rebound Egawa and Katsuragi, Phys. Fluids **31**, 053304 (2019)



Ingredients of our simulation



Rebound depends on impact velocity and frictional interactions between particles



Impact velocity



We can run on top suspension but we'll sink if we walk

Frictional interaction increases the contact duration between particles that leads to a stronger hardening

Friction coefficients



Inside the hardening suspension



Force chains

- Formed by contacting suspended particles
- Transmitting force from the impactor to the bottom boundary
- Sustain the impactor \rightarrow rebound

Coarse-grained fields J. Zhang, R. P. Behringer, I. Goldhirsch, Prog. Theor. Phys. Supp 184, 16 (2010)

1.0e+02

- 75

50

25

0.0e+00



High normal stress instead of shear stress c.f DST

Force on the impactor reach its peak before the rebound





Rebound depends on the depth of the container

The origin of F_{max} is the **floating** jammed region

The origin of rebound is the transmission of force from the impactor to the bottom plate and vice versa

Phenomenology Floating + force chains model

$$m_{I} \frac{d^{2} z_{I}}{dt^{2}} = -m_{I} \tilde{g} + 3\pi \eta_{\text{eff}} \dot{z}_{I} |z| + n(t) k_{n} z_{I}$$

From floating jammed region

Drag on the impactor is proportional to its depth before completely immersed



From force chains Number of percolating

$$n(t) \longrightarrow$$
 force chains from the impactor to the bottom boundary



From analytical solution of floating model in $u_0 \gg 1/\eta_{\rm eff}$ one can obtain

$$F_{\rm max} \propto u_0^{\frac{3}{2}}$$
 and $t_{\rm max} \propto u_0^{-\frac{1}{2}}$

Model vs simulations

Power-law relationships between F_{max} , t_{max} , and u_0



- **Crossover** from low u_0 to high u_0 regime
- High u_0 regime: $F_{\rm max} \propto u_0^{\alpha}$ and $t_{\rm max} \propto u_0^{\beta}$
 - Also observed in experiments

Brassard, et. al, JFM 923, A38 (2021)

Independent of system size



• Elastic term from percolating force chains is necessary to recover rebound



Conclusions

The strength of the impact-induced hardening depends on volume fraction, impact velocity, and the frictional contact between particles

High normal stress (not shear stress)is observed in the dynamically jammed region

Our phenomenology shows that the power-law relationship between $F_{\rm max}$, $t_{\rm max}$, and u_0 arises solely from the viscous process

Elastic rebound is originated from elastic force due to the percolating force chains

Future prospects

Dynamics and effective elasticity of the jammed region



Simulations of actual running model e. g Raibert hopper on dense suspensions

