

Effective viscosity and elasticity of dynamically jammed region and their role in the hopping motion on dense suspensions

arXiv:2205.13822

14pH121-13 Pradipto* and Hisao Hayakawa

Yukawa Institute for Theoretical Physics Kyoto University

JPS Autumn Meeting 2022 Tokyo Institute of Technology 2022/09/14

*Tagawa Lab, Tokyo University of Agriculture and Technology (Starting from October)

Introduction Dense suspensions under impact

Non-Brownian Suspensions

Mixture of macroscopic ($1\mu m - 500\mu m$) undissolved particles in a Newtonian solvent



Impact-induced hardening



source: https://www.youtube.com/watch?v=hP88C-_LgnE&list=PLVjiIPzFTOLpxiwdwrFuPYSIBpr6jFm32

Able to run



Brown, et. al., Rep. Prog. Phys 77, 046602 (2014).

Liquid body armor



https://www.youtube.com/watch?v=L5Ts9IYZIDk

Introduction Dense suspensions under impact

Waitukaitis and Jaeger, Nature 487, 205 (2012)



Han et al, Nat. Comms. 2016



3

Dynamically Jammed Region (DJR)

- Localized and transient solid-like region beneath the impactor Boundary \rightarrow peak of
- strain rate

Fracture

Roche et al, PRL **110**, 148304 (2013)



Elastic rebound



Egawa and Katsuragi, Phys. Fluids **31**, 053304 (2019)

Previous works: Recovered rebound from percolating force chains

Pradipto and H. Hayakawa, Phys. Fluids **33**, 093110 (2021) Pradipto and H. Hayakawa, Phys. Rev. Fluids **6**, 033301 (2021)



- Calculate number of percolating force chains n(t)
- Incorporate elastic response from n(t)
- Do elastic response exists when no force chains percolate?
- What about actual running motion?

LBM-DEM Simulation

Contacting suspended particles form networks \rightarrow force chains



Obtaining fields inside the hardening suspensions





Particles velocity (discrete) Force chains (contact between particles)

Coarse graining J. Zhang, R. P. Behringer, I. Goldhirsch, Prog. Theor. Phys. Supp 184, 16 (2010)

Tensorial basis adapted to local flow conditions Giusteri and Seto, J. Rheol. 62, 713 (2018)

Strain rate tensor

$$\overleftrightarrow{D} = \frac{1}{2} (\nabla \mathbf{u} + \nabla \mathbf{u}^T)$$

Scalar strain rate

$$\dot{\varepsilon} = \sqrt{2\overleftrightarrow{D}:\overleftrightarrow{D}}$$

Strain tensor (from particle overlap)

$$\overleftrightarrow{L} = \frac{1}{2} (\nabla \delta_n + \nabla \delta_n^T)$$

Scalar strain

$$\varepsilon = \sqrt{2\overleftrightarrow{L}} : \overleftrightarrow{L}$$

Dynamically jammed region → region with finite rigidity



Time-dependent effective viscosity and elasticity from the DJR

$$\eta_{\text{eff}} = \eta_0 + \frac{1}{V_{\text{djr}}} \int_d \eta dV \qquad G_{\text{eff}} = \frac{1}{V_{\text{djr}}} \int_d G dV$$

DJR model (data-assisted)

$$m_I \ddot{z}_I = -m_I g - 3\pi \eta_{\text{eff}} z \dot{z}_I + C \frac{A_{DJR}}{H_{DJR}} G_{\text{eff}} z$$



 $z = z_I - a_I$

Force chains vs DJR model

Floating + force chains model



Count number of percolating force chains n(t)

4.0 20 3.5 15 $G_{ m eff}(t)/k_n$ Finite rigidity n(t)even when no percolating 1.5 5 force chains 1.0 0.5 0.00 0.05 0.10 0.15 0.20 0.25 0.05 0.10 0.15 0.20 0.25 t/t_g t/t_g

DJR model



7

$$G_{\rm eff} = \frac{1}{V_{\rm djr}} \int_d G dV$$

Force chains vs DJR model



Both model work well

DJR model works better

Mild elastic effect captured

Hopping robot on top of dense suspensions

Spring-mass model for one legged hopping robot Raibert and Tello, IEEE Expert 1, 89 (1986) (Raibert's hopper)



Source: <u>https://youtu.be/XFXj81mvInc</u>





Motion of spring-mass-body system



Foot undergoes multiple rebounds ($u_p > 0$) and also hops ($z_p > 0$)



- Dynamically jammed region (DJR) is formed inside the suspensions after the impact → Region with finite rigidity inside the suspensions
- Elastic response exists even without percolating force chains → We proposed model that contains effective viscosity and elasticity of the DJR
- Simulation of foot-spring-body system to mimic hopping
 →observed multiple rebounds and hops