# DEMON DRIVEN BY GEOMETRICAL PHASE 

Based on H Hayakawa and R Yoshii, arXiv:2205.15193 (2022).
(H Hayakawa, VMM Paasonen, R Yoshii, arXiv:2112.12370 (2021).)

# Ryosuke Yoshii <br> (Sanyo-Onoda City University) 

Collaborator Hisao Hayakawa (YITP)

(Ville M. M. Paasonen)


## Introduction: Thouless pumping

Classical pump


Mean bias $=0 \Rightarrow$ mean current $=0$
Topological pump
Mean current can flow w/ bias voltages
Mean current!



Thouless, 1983

## Experimental implementation

## Quantum dot system

Quantum dot turnstile



Cold atom system
M. Switkes, et.al., Science, 283, 1905 (1999).

| MMaranadin <br> WWWWWalow <br> MMOMMOMOMOM <br> NWWWWWWM <br>  |
| :---: |
|  |  |
|  |  |



## Geometrical interpretation

Non-trivial curvature in parameter space


Gauss-Bonnet type argument
$\square$ Current ~ integration of the curvature inside trajectory Berry-Sinitsyn-Nemenman(BSN) curvature
N. A. Sinitsyn and I. Nemenman PRL 99, 220408 (2007).

## Connection with thermodynamics

Entropy production depends on the trajectory
T. Sagawa and H. Hayakawa PRE 84, 051110 (2011).




Vector potential in thermodynamics
K. Tomita and H. Tomita PTP 51, 6 (1974).

Related topics
-Brandner-Saito (PRL 2020) : single-bath -Hino-Hayakawa (PRR2021): 2-baths - Ito-Dechant (PRX2020)

## Purpose

Experimentally realizable system
Quantum dot system
Quantum dot + reservoirs


Fermion with the Coulomb interaction
Quantum effect?
Many body interaction?
Fermi statistics?
Geometrical effect
Work? Efficiency? Entropy production?

## Formalism: Master equation

Time evolution of the density matrix $\rho$

$$
\begin{array}{r}
\frac{d \rho}{d \theta}=\hat{\mathscr{K}} \rho \quad \hat{\mathscr{K}}: \text { linear operator acting on } \rho \\
\text { (CPTP property is satisfied) }
\end{array}
$$

$\theta$ : time $\quad \hat{\mathscr{K}}(\theta)$ : modulated in time

Recasting the density matrix into the vector form

$$
\begin{array}{cc}
\frac{d}{d \theta}|\hat{\rho}\rangle=\hat{\mathscr{K}}|\hat{\rho}\rangle & \left.\left|r_{i}(\theta)\right\rangle\right\rangle: \text { right eigenstate of } \hat{\mathscr{K}} \\
& \left\langle\left\langle\ell_{i}(\theta)\right|: \text { left eigenstate of } \hat{\mathscr{K}}\right. \\
\varepsilon_{i} \text { : eigenvalues of } \hat{\mathscr{K}}
\end{array}
$$

## Geometrical state

$$
\begin{aligned}
& \left.\left.|\rho(\theta)\rangle\rangle=\left|r_{0}(\theta)\right\rangle\right\rangle-\sum_{i \neq 0} \int_{0}^{\theta} d \phi e^{\int_{\phi}^{\theta} d \chi \epsilon^{-1} \varepsilon_{i}(\chi)} \mathscr{A}_{\mu} \frac{d \Lambda^{\mu}}{d \phi}\left|r_{i}(\theta)\right\rangle\right\rangle \\
& \left.\mathscr{A}_{\mu}=\left\langle\left.\left\langle\ell_{i}(\phi)\right| \frac{\partial}{d \Lambda^{\mu}} \right\rvert\, r_{0}(\phi)\right\rangle\right\rangle \quad \Lambda: \operatorname{par} \\
& \mathscr{F}_{\mu \nu}^{i} \equiv \frac{\partial \mathscr{A}_{\nu}^{i}}{\partial \Lambda_{\mu}}-\frac{\partial \mathscr{A}_{\mu}^{i}}{\partial \Lambda_{\nu}}
\end{aligned}
$$

## Geometrical state

$$
\begin{aligned}
|\rho(\theta)\rangle\rangle= & \left.\frac{\left.\left|r_{0}(\theta)\right\rangle\right\rangle}{\uparrow}-\sum_{i \neq 0} \int_{0}^{\theta} \frac{d \phi e^{\rho_{\phi}^{\theta} d x \epsilon^{-1} \varepsilon_{i}(x)}}{\uparrow} \mathscr{A}_{\mu} \frac{d \Lambda^{\mu}}{d \phi}\left|r_{i}(\theta)\right\rangle\right\rangle \\
\mathscr{A}_{\mu}= & \left.\left\langle\left.\left\langle\ell_{i}(\phi)\right| \frac{\partial}{d \Lambda^{\mu}} \right\rvert\, r_{0}(\phi)\right\rangle\right\rangle \\
\mathscr{F}_{\mu \nu}^{i} \equiv & \frac{\partial \mathscr{A}_{\nu}^{i}}{\partial \Lambda_{\mu}}-\frac{\partial \mathscr{A}_{\mu}^{i}}{\partial \Lambda_{\nu}} \\
& \text { } \quad \text { : parameters } \\
&
\end{aligned}
$$

Exponential decay

## Geometrical state

$$
\begin{aligned}
& \left.|\rho(\theta)\rangle\rangle=\frac{\left.\left|r_{0}(\theta)\right\rangle\right\rangle}{\uparrow}-\sum_{i \neq 0} \int_{0}^{\theta} \frac{d \phi e^{\int_{\phi}^{\theta} d \chi \epsilon^{-1} \varepsilon_{i}(\chi)}}{\uparrow} \mathscr{A}_{\mu} \frac{d \Lambda^{\mu}}{d \phi}\left|r_{i}(\theta)\right\rangle\right\rangle \\
& \text { ^: parameters } \\
& \text { Steady state } \\
& \text { ^: parameters }
\end{aligned}
$$

Only $\theta-\epsilon \leq \phi \leq \theta$ contributes
Quasi local in time
"quasi" steady state


## Application to Impurity Anderson Model

$$
\hat{H}^{\mathrm{tot}}=\hat{H}+\hat{H}^{r}+\hat{H}^{\mathrm{int}}
$$

$\hat{H}^{\prime}$

$\hat{H}^{\prime \prime}$

$$
\hat{H}=\sum_{\sigma} \epsilon_{0} \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma}+U \hat{n}_{\uparrow} \hat{n}_{\downarrow} \quad\left(\hat{n}_{\uparrow / \downarrow}=\hat{d}_{\uparrow / \downarrow}^{\dagger} \hat{d}_{\uparrow \downarrow \downarrow}\right)
$$

$\hat{d}_{\sigma}^{\dagger}\left(\hat{d}_{\sigma}\right)$ : creation and annihilation operator in dot (spin $\sigma$ )

$$
\hat{H}^{r}=\sum_{\alpha, k, \sigma} \epsilon_{k} \hat{a}_{\alpha, k, \sigma}^{\dagger} \hat{a}_{\alpha, k, \sigma}
$$

$\hat{a}_{\alpha, k, \sigma}^{\dagger}\left(\hat{a}_{\alpha, k, \sigma}\right)$ : creation and annihilation operator in leads
(spin $\sigma$, wave number $k, \alpha=$ left or right)

$$
\begin{aligned}
\hat{H}^{\text {int }}= & \sum_{\alpha, k, \sigma} V_{\alpha} \hat{d}_{\sigma}^{\dagger} \hat{a}_{\alpha, k, \sigma}+\text { h.c. }, \\
& \left(V_{L}=V_{R} \text { for simplicity }\right)
\end{aligned}
$$

## Modulating parameters


$\delta$ : phase deference


- Temperature is difficult to be controlled
- Parameter in quantum dot must be tuned to extract work


## Modulating parameters


$\delta$ : phase deference


- Temperature is difficult to be controlled
- Parameter in quantum dot must be tuned to subtract work



## Absorb and release of heat \& work

$Q_{A}=\oint \frac{\mathscr{P}_{\mu}+\left|\mathscr{P}_{\mu}\right|}{2} d \Lambda^{\mu}, \quad Q_{R}=\oint \frac{\mathscr{P}_{\mu}-\left|\mathscr{P}_{\mu}\right|}{2} d \Lambda^{\mu}, \quad W=\oint \mathscr{P}_{\mu} d \Lambda^{\mu}, \quad \mathscr{P}_{\mu}=\operatorname{Tr}\left(\rho \frac{\partial H}{\partial \Lambda^{\mu}}\right)$



Negative work
Efficiency: $\eta \equiv \frac{|W|}{Q_{A}}=\frac{\left|Q_{A}+Q_{R}\right|}{Q_{A}}$



## Time evolution of relative entropy

Relative entropy: $S^{\mathrm{HS}}\left(\rho(\theta) \| \rho^{\mathrm{SS}}(\theta)\right)=\operatorname{Tr} \rho(\theta)\left[\ln \rho(\theta)-\ln \rho^{\mathrm{SS}}(\theta)\right]$
Initial decay toward quasi steady state
$\theta$ : time



Exp. Decay

Initial relaxation $\rightarrow$ negative $\Delta S$


$$
\begin{aligned}
& \Delta S=-S^{\mathrm{HS}}\left(\rho(2 \pi) \| \rho^{\mathrm{SS}}(2 \pi)\right) \\
& +S^{\mathrm{HS}}\left(\rho^{\mathrm{SS}}(0) \| \rho^{\mathrm{SS}}(0)\right) \\
& 2 \pi: \text { one-cycle }
\end{aligned}
$$

## Interpretation of the results

Negative work
Extracted from the geometrical phase

## "Geometrical demon"



Cf. Maxwell demon utilizes information

Relative entropy decrease in the CPTP process


No modulation


With modulation

State cannot relax to a simple "steady state" under a modulation => Non-adiabatic effect.

## Remaining issues (on going)

First and Second laws of thermodynamics?


The cost is not considered
(Interaction inside the system is modulated)


Joule heating is generated
(Second law would be maintained)

## SUMMARY

- Quasi steady state is realized under parameter modulation $|\rho(\theta)\rangle\rangle$ : quasi local in time
- Work can be extracted by utilizing the geometrical phase


## "Geometrical Demon"

- Relative entropy can be decreased by geometrical effect Initial relaxation process
- Extracted Work exponentially decays

One needs to wait to initialize the state

- For details, please see
H. Hayakawa and R. Yoshii, arXiv:2205.15193 (2022).
H. Hayakawa, VMM Paasonen, R. Yoshii, arXiv:2112.12370 (2021).

