

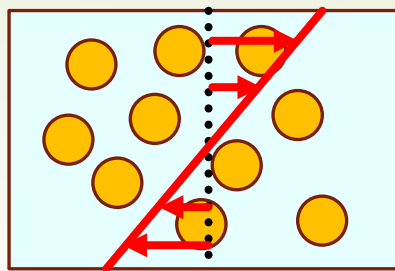
Discontinuous shear thickening of a moderately dense inertial suspension of hydrodynamically interacting frictionless soft particles



① Satoshi Takada¹ ② Kazuhiro Hara¹
③ Hisao Hayakawa²

¹Tokyo Univ. of Agri. & Tech. ²YITP, Kyoto Univ.

Introduction



- Viscosity $\eta(\phi) \equiv \sigma(\phi)/\dot{\gamma}$: characterizes noneq. transport

Dilute case (Einstein, 1906): $\eta_s(\phi)/\eta_0 = 1 + (5/2)\phi$ ($\phi \leq 0.03$)

Dense case (near jamming): $\eta_s(\phi)/\eta_0 = (1 - \phi/\phi_m)^{-2}$ (empirical)

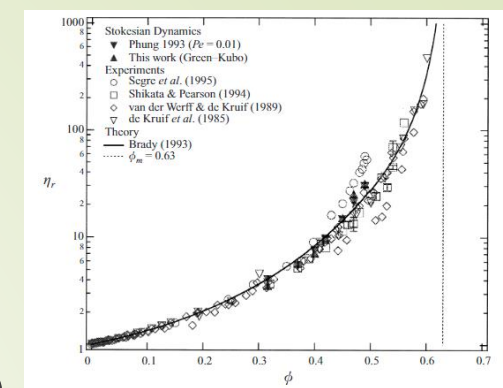
- Shear rate dependent viscosity (non-Newtonian fluid)

- Shear thickening (thinning):
Viscosity becomes large (small) as $\dot{\gamma}$ increases.

- Discontinuous shear thickening (DST)
can occur. (many mechanisms are proposed.)

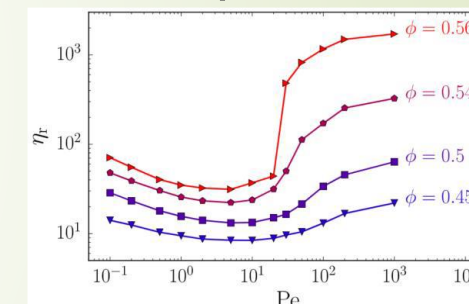


- If "inertia" is not ignored, system is called as
"inertial suspensions"
(a model of aerosols or colloid)



solvent viscosity: η_0

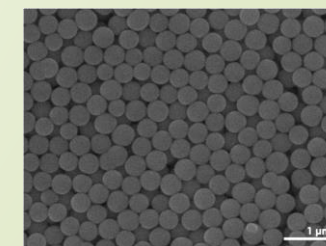
shear rate: $\dot{\gamma}$



"Inertial effect" is often **IGNORED**.
(Overdamped is assumed.)



COVID-19



Previous studies on inertial suspension

3

System: frictionless, Stokes' drag

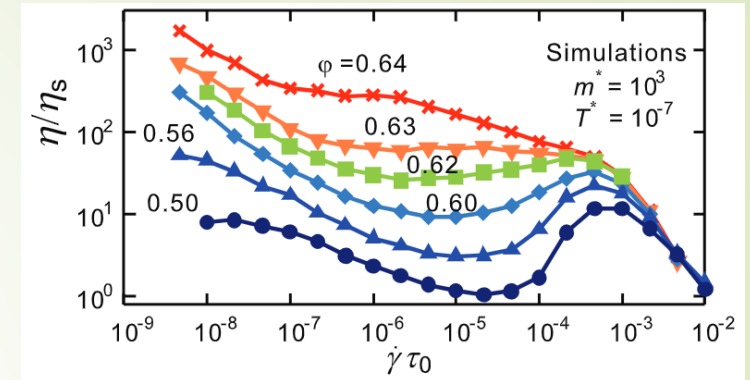
➤ Dense suspension (soft-core particles)

☞ Kawasaki, Ikeda, & Berthier, EPL (2014)

➤ Thinning → thickening → thinning for $\varphi \lesssim 0.60$

➤ No thickening for $\varphi \gtrsim 0.63$

✂ Only contact contribution is considered.



➤ Dilute to moderately dense inertial suspension (**hard-core**)

☞ Hayakawa & Takada, PTEP (2019),

Takada, Hayakawa, Santos, & Garzó, PRE (2020)

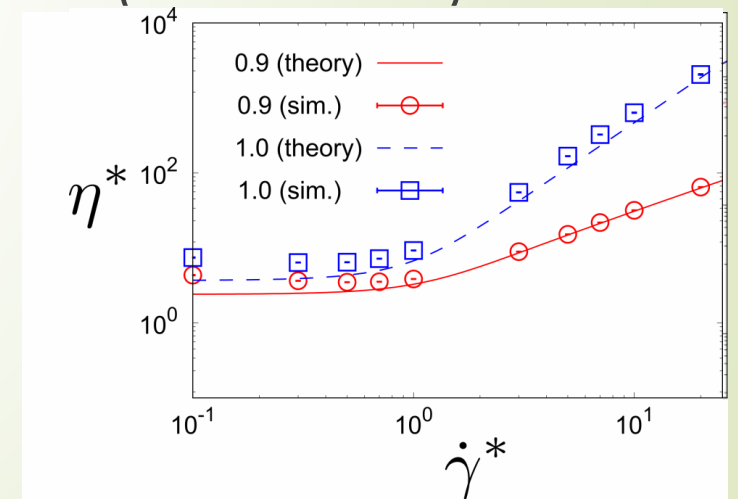
➤ **DST-like behavior** for dilute systems
(\cong ignited-quenched transition)

➤ Change to **CST-like behavior** at $\varphi \simeq 0.0176$

➤ Agreement for $\varphi \lesssim 0.5$

➤ Mpemba effect (**hotter** can become **colder**)

☞ Prof. Santos's talk)



Previous studies on inertial suspension

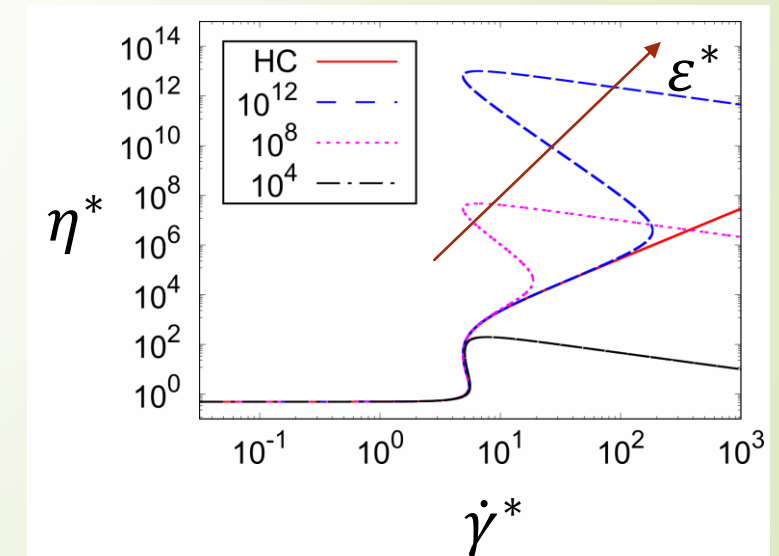
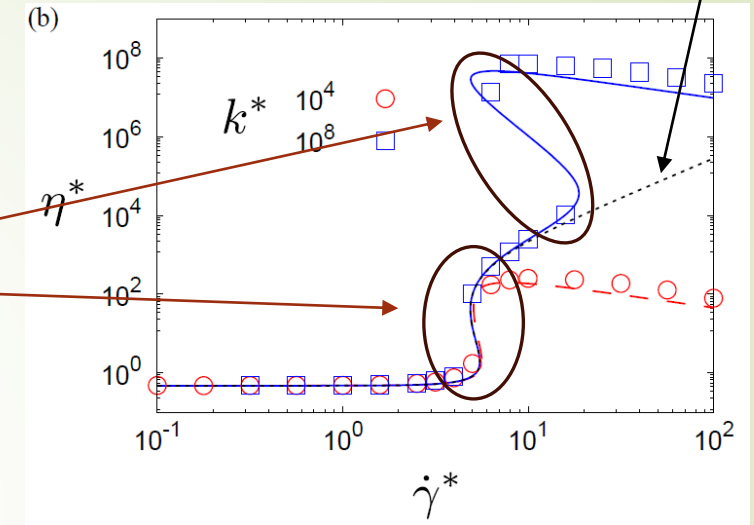
➤ Dilute inertial suspension (**soft-core**)

☞ Sugimoto & Takada, JPSJ (2020)

- **DST-like behaviors** can occur **TWICE!!**
- Kinetic theory is also developed.
(Good agreement with simulations.)

Origin of two DSTs:

- 1st DST: occurs at the same $\dot{\gamma}^*$ as hard-core (same mechanism)
⇒ disappears as φ : ↗
- 2nd DST: occurs depending on the softness of particles
⇒ **Softness (hardness) induced DST**



5

Does the second DST in dilute systems survive for dense systems?

Both numerical to theoretical

- Numerical:
Langevin simulation
- Theoretical (for ①):
(Kinetic theory of inertial suspension)

		Only Stokes' drag		Hydro.
		Hard-core	Soft-core	
Theory & Sim.	Dilute	Hayakawa & Takada, PTEP (2019)	Sugimoto & Takada, JPSJ (2020)	
	Mod. dense	Hayakawa, Takada, & Garzó, PRE (2017), Takada, Hayakawa, Santos, & Garzó, PRE (2020)		
Sim.	Dense		Kawasaki, Ikeda, & Berthier, EPL (2014)	Mari <i>et al.</i> , PNAS (2015)

Model and setup

- System = Particle + Solvent
- Particle:
monodisperse (mass m , diameter d)

$$\begin{aligned}\frac{d\mathbf{r}_i}{dt} &= \frac{\mathbf{p}_i}{m} + \dot{\gamma} y_i \hat{\mathbf{e}}_x \\ \frac{d\mathbf{p}_i}{dt} &= \sum_{j \neq i} \mathbf{F}_{ij}^{(\text{el})} + \mathbf{F}_i^{\text{H}} + \boldsymbol{\xi}_i(t)\end{aligned}$$

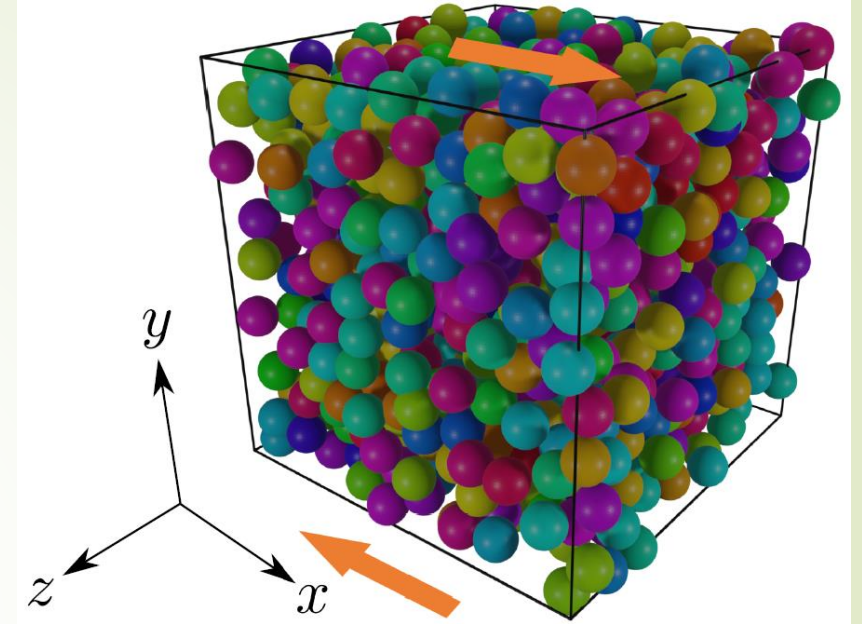
Shear (shear rate: $\dot{\gamma}$)

Hydrodynamic interaction
from the solvent: **two models**

Interparticle interaction
= harmonic potential

$$\mathbf{F}_{ij}^{(\text{el})} = -\frac{\partial U(r_{ij})}{\partial \mathbf{r}_i}, U(r_{ij}) = \frac{\varepsilon}{2} \left(1 - \frac{r_{ij}}{d}\right)^2 \Theta\left(1 - \frac{r_{ij}}{d}\right)$$

Noise term
(satisfies fluctuation-dissipation theorem)



- Solvent (fluid): viscosity η_0 , temperature T_{env} are kept constant.

We consider two cases for hydrodynamic interaction:

7

1. Scalar resistance model: $F_i^H = -\zeta \mathbf{p}_i$

Only **Stokes'** drag ($\zeta = 3\pi d\eta_0/m$)

Theoretical treatment is available.

☞ Enskog kinetic equation for the inertial suspension

2. Stokes' + lubrication model: $F_i^H = -\sum_j \overleftrightarrow{\zeta}_{ij} \mathbf{p}_j$

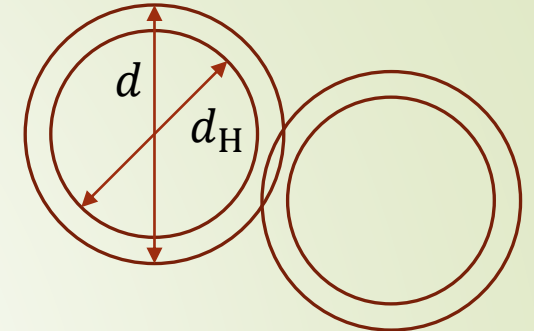
$\overleftrightarrow{\zeta}_{ij}$: non-diagonal, diverge at $r \rightarrow d$

“**Roughness parameter**” (dimple) is introduced.

$\delta \equiv \frac{d-d_H}{d_H} \sim 1 \sim 10\%$: Magnitude of **dimple**

(☞ Mari et al., J. Rheol. **58**, 1693 (2014);

Pradipto & Hayakawa, Soft Matter **16**, 945 (2020), etc.)



d : collision diameter

d_H : lubrication diameter

$$\zeta_{ij,\alpha\beta} = \begin{cases} \frac{3\pi d\eta_0}{m} \delta_{\alpha\beta} + \sum_{k \neq i} \frac{1}{m} A_{ik,\alpha\beta}^{(1,1)} \Theta(r_c - r_{ik}) & (i = j) \\ -\frac{1}{m} A_{ij,\alpha\beta}^{(1,1)} \Theta(r_c - r_{ij}) & (i \neq j) \end{cases}$$

$A_{ij,\alpha\beta}^{(1,1)}$: function of $\hat{\mathbf{k}} \equiv \mathbf{r}_{ij}/|\mathbf{r}_{ij}|$

(☞ Kim & Karrila, “Microhydrodynamics”)

$r_c \equiv d_H + \lambda$: cutoff length ($\lambda = 0.25d$)

(Dimensionless) control parameters:

① Packing fraction: φ ② Shear rate: $\dot{\gamma}^* \equiv \dot{\gamma}/\zeta$

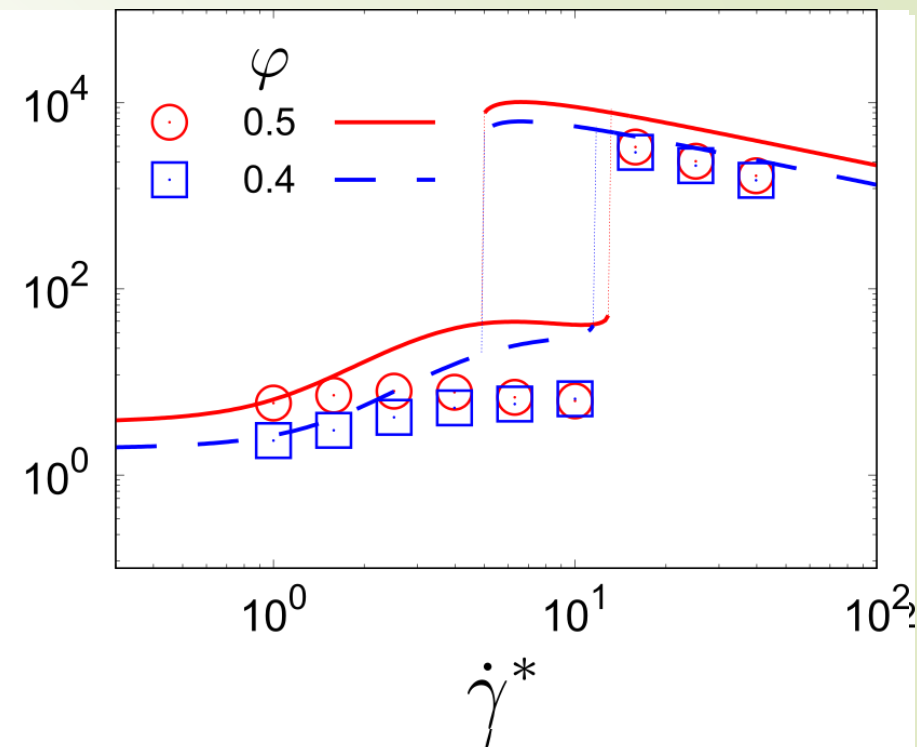
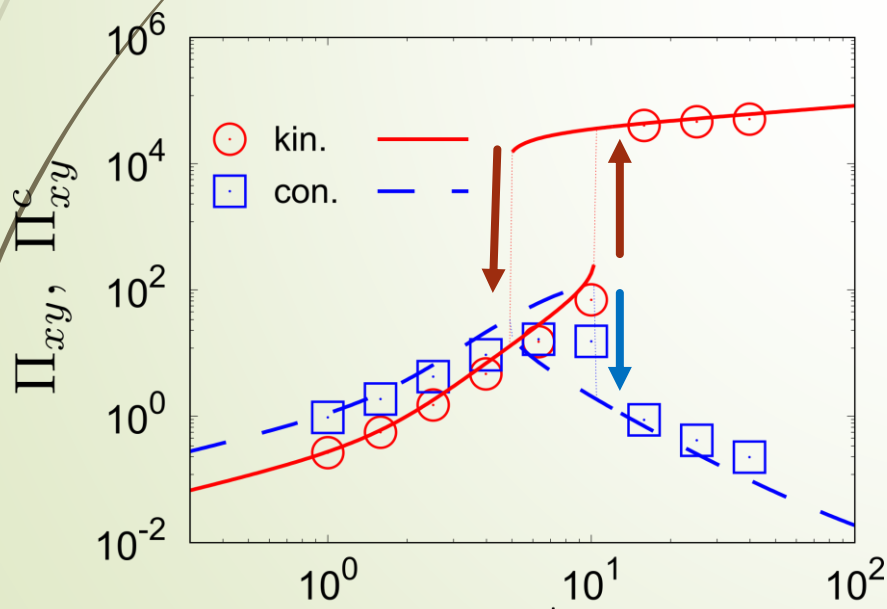
③ Particle hardness: $\varepsilon^* \equiv \frac{\varepsilon}{m\sigma^2\zeta^2}$ ④ Envir. temp.: $\xi_{\text{env}} \equiv \sqrt{\frac{T_{\text{env}}}{m}} \frac{1}{\zeta\sigma}$

⑤ Magnitude of dimple: δ (only for 2nd case)

Results ①: Scalar model

Parameters:
 $\varphi = 0.10, 0.20, 0.30$
 $\varepsilon^* = 10^4, \xi_{\text{env}} = 1.0$

- **(2nd) DST-like behavior survives** even for finite density φ .
 (\Leftrightarrow CST-like for hard-core system)
- **Shear thinning** in high shear regime $\dot{\gamma}^*$
- (Kinetic theory reproduces the sim. results.)

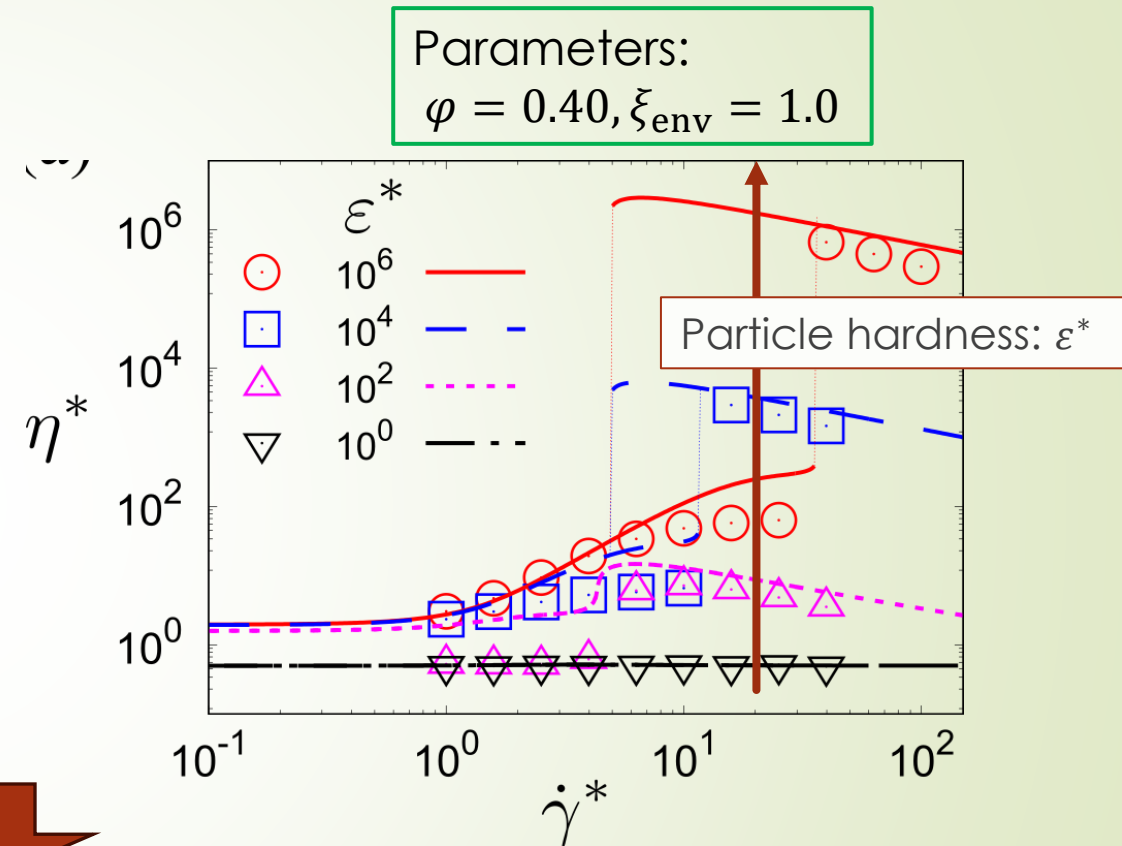
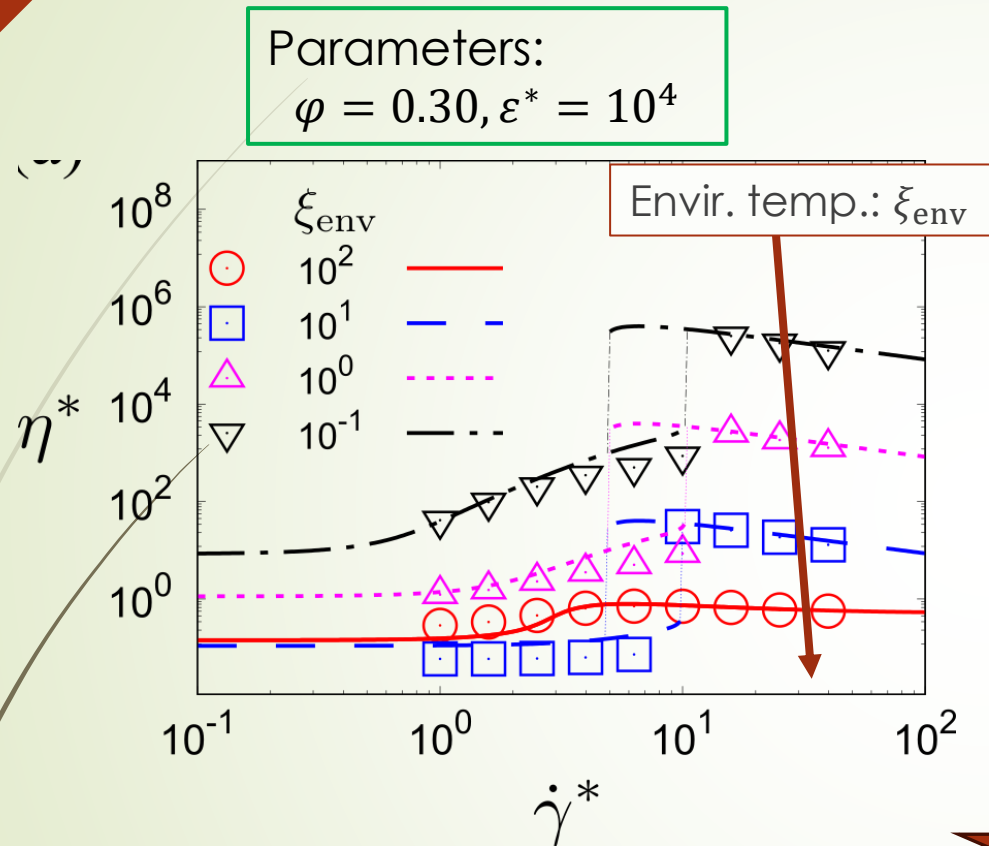


After the DST occurs,
 Kinetic contribution \gg Contact contribution

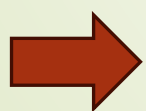
$$\left(P_{xy}^k = \frac{1}{V} \sum_i m v_{i,x} v_{i,y} \gg P_{xy}^c = \frac{1}{V} \sum_i \sum_{j \neq i} x_{ij} f_{ij,y} \right)$$

\Rightarrow **Inertia plays a crucial role for DST!!**
 (= cannot be ignored)

Viscosity vs. shear rate for various sets of parameters



We can evaluate the parameters' space where the DST occurs.



Question:

Does this DST survive when the lubrication forces exist?

Results ②: Stokes' + lubrication model

10

- Answer is Yes.

Scaled viscosity $\tilde{\eta} \equiv \eta/\eta_1$ against the Peclet number $Pe \equiv \frac{3\pi\eta_0 d^3}{4T_{\text{env}}} \dot{\gamma}$ shows DST-like behavior.

$$\eta \equiv P_{xy}/\dot{\gamma}, \eta_1 = \eta_0(1 + 2.5\varphi + 4\varphi^2 + 42\varphi^3):$$

η_1 : Empirical expression of the apparent viscosity in the low shear limit

Parameters:

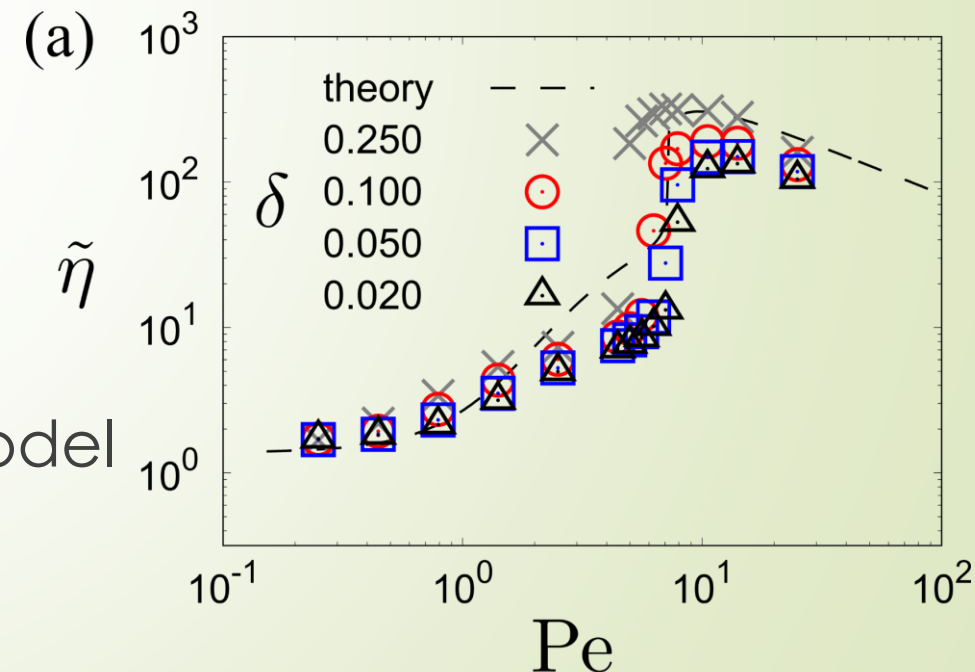
$$\varphi = 0.30$$

$$\varepsilon^* = 10^4, \xi_{\text{env}} = 1.0$$

- Even for small δ (small dimple), **DST** occurs at $Pe \simeq 10$.

⇔ DST occurs at $Pe \simeq 20$
for frictional Brownian suspension
☞ Mari *et al.*, PNAS **112**, 15326 (2015)

- For larger δ , tends to the previous model



Disc.: Estimation of quantities for model ②

11

➤ Aerosol

$$d \sim 10^{-6} \text{ m}, \rho \sim 1 \text{ g/cm}^3, E \sim 10 \text{ GPa} \Rightarrow m \sim 10^{-14} \text{ kg}$$

$$\text{Viscosity of air: } \eta_0 \sim 10^{-5} \text{ Pa} \cdot \text{s}$$

$$\text{DST takes place at } \dot{\gamma}_c \sim 10^4 \text{ 1/s}$$

$$\Rightarrow \text{shear speed } 10^2 \text{ m/s if } L = 1 \text{ cm}$$

➤ Colloid

$$d \sim 2 \times 10^{-6} \text{ m}, \rho \sim 1 \text{ g/cm}^3, E \sim 1 \text{ GPa}$$

$$\Rightarrow m \sim 10^{-15} \text{ kg}$$

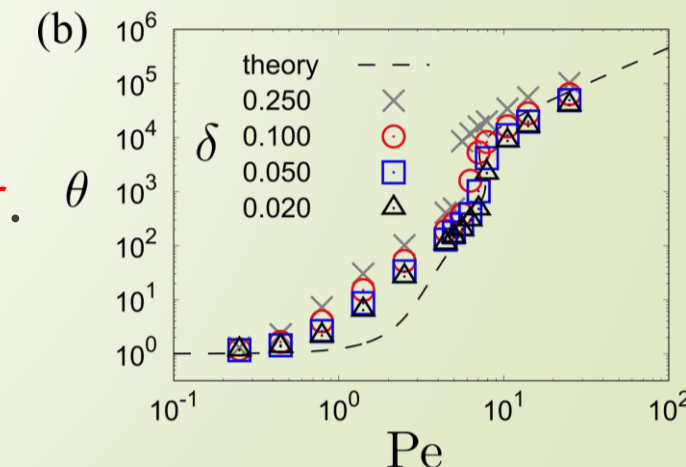
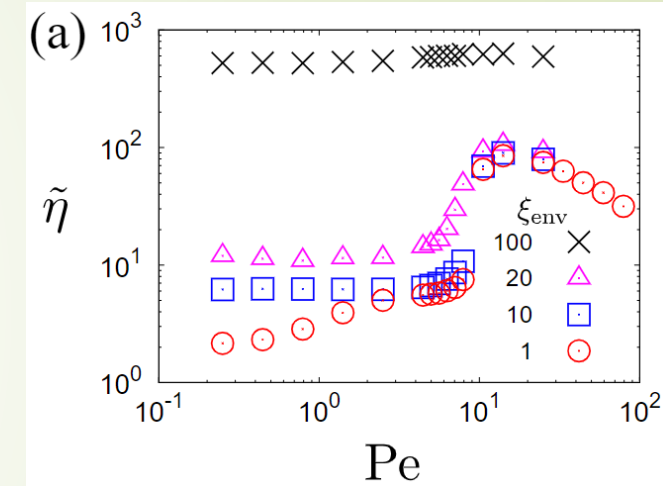
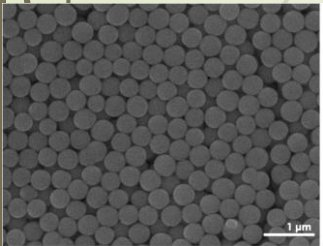
$$\text{Viscosity of water: } \eta_0 \sim 10^{-3} \text{ Pa} \cdot \text{s}$$

$$\text{Room temperature (300K): } \xi_{\text{env}} \sim 2 \times 10^{-4}$$

➤ Kinetic temperature becomes 10^2 times larger.

Is it possible to achieve this?

\Rightarrow This will open for discussion.



Summary

- Softness induced DST-like behaviors can exist for **frictionless** system.
- This DST survives even for finite densities.
- Scalar model
DST-like behaviors
Good agreement between sim. and theory.
- Lubrication model
DST-like behaviors survive
even for small roughness parameter

Future work

- Long range interaction
(inclusion of Lotne-Prager tensor)
- Verifiability in experiments

