Flavour symmetries and SUSY soft-breaking at LHC

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with : L. Calibbi, J. Jones-Perez, A. Masiero, J. Park

Flavour physics: who ordered that??

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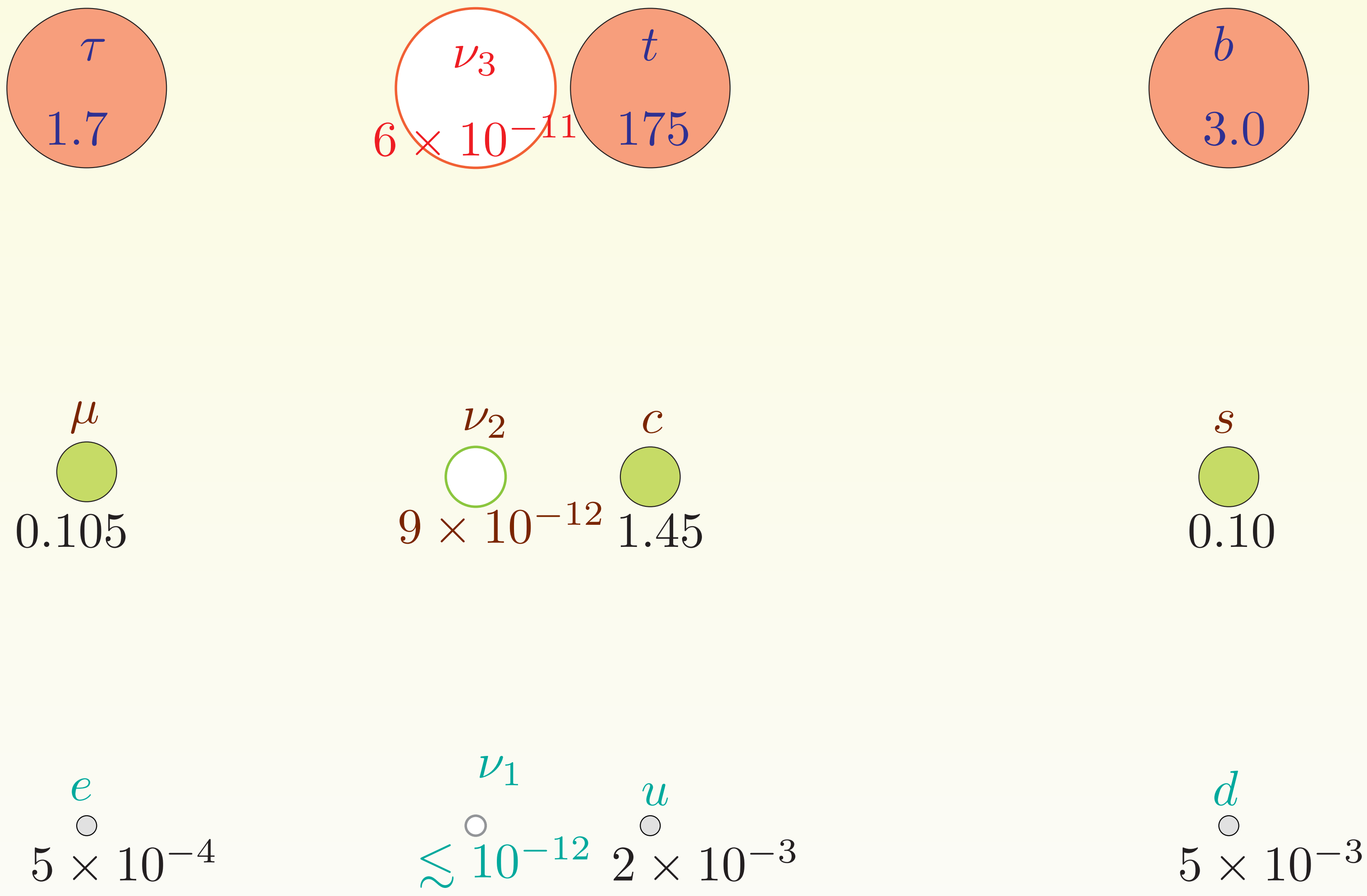

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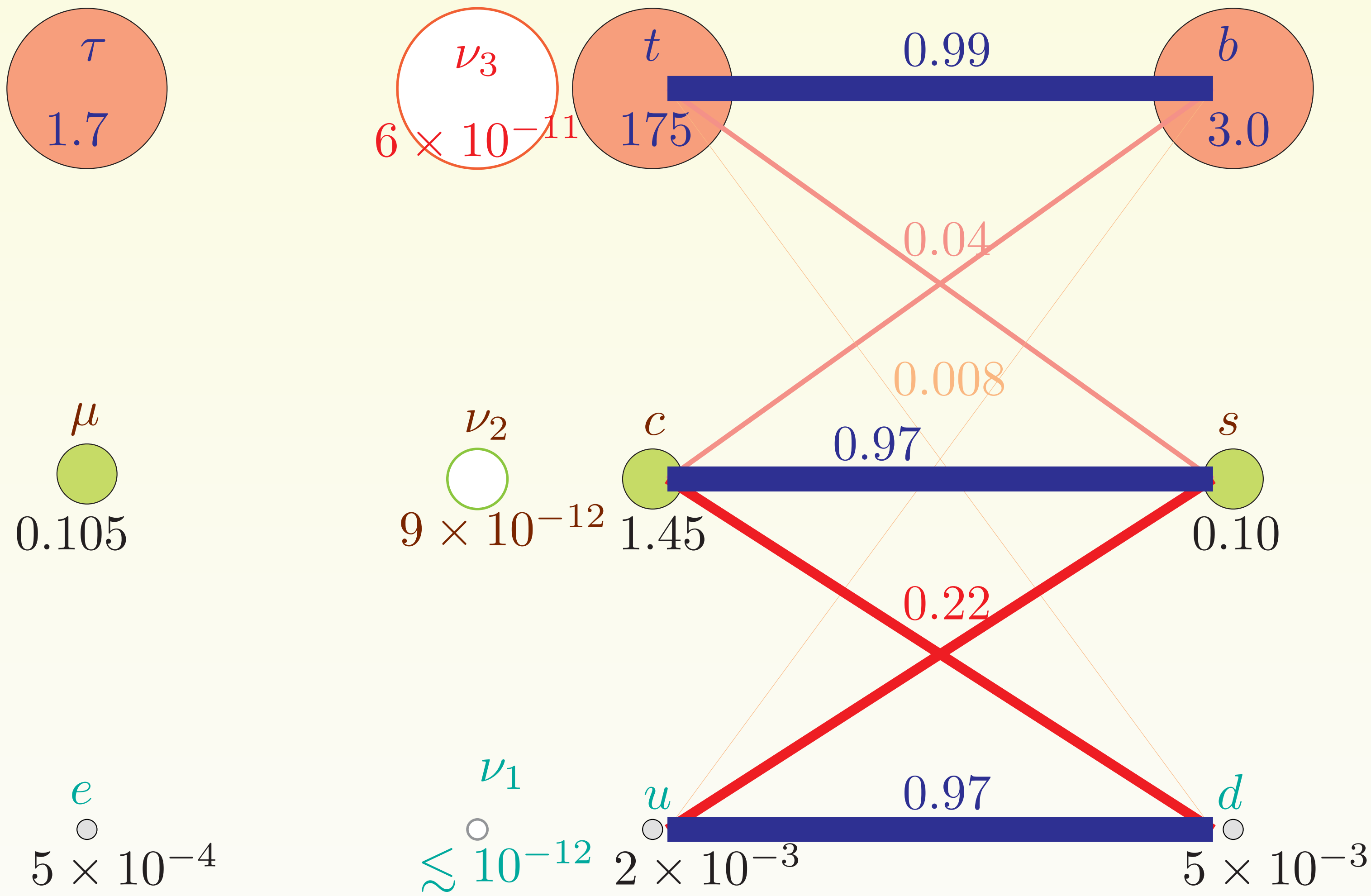

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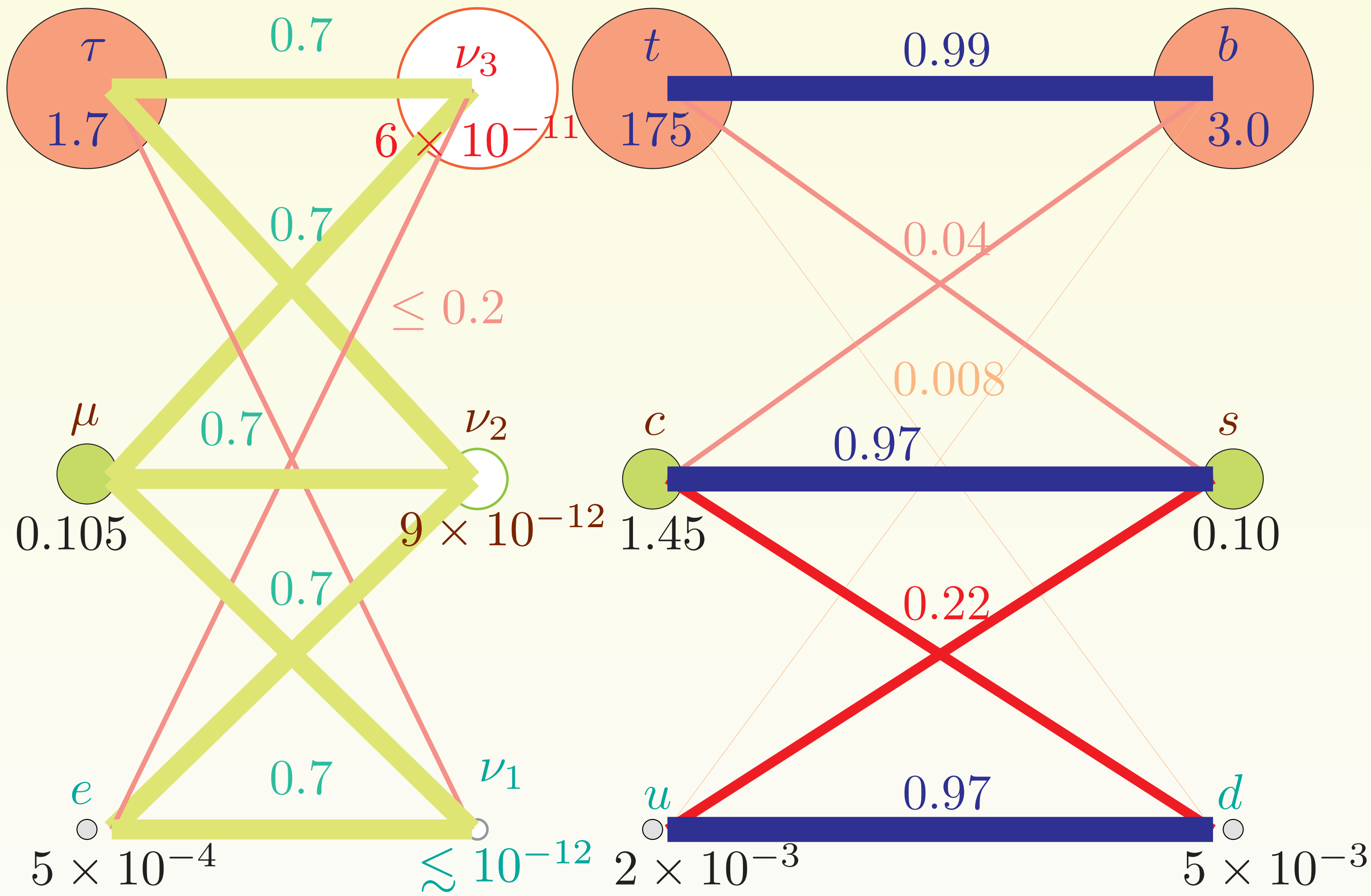

Standard Model

All flavour physics originate in Yukawa couplings:

$$\mathcal{L}_Y = H\, \bar{Q}_i\, Y^d_{ij}\, d_j + H^*\, \bar{Q}_i\, Y^u_{ij}\, u_j$$

In absence of Yukawas, $\mathcal{L}_{SM}$ invariant under global $(U(3))^5$

$\Longrightarrow$ quark masses and CKM mixings only observables in SM

Not enough information to determine the full Yukawa matrices

Supersymmetry

New flavour dependent interactions (sfermions/gauginos)

$\Longrightarrow$ new experimental information on flavour (urgently needed)!!

Flavour and CP problems

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SUSY Flavour and CP

Soft masses fixed by $m_{3/2}$. $O(m_{3/2})$ elements in soft matrices.

$\Longrightarrow$ **Severe FCNC problem !!!**

CP broken, we can expect all complex paramaters have $O(1)$ phases.

$\Longrightarrow$ **Too large EDMs !!!**

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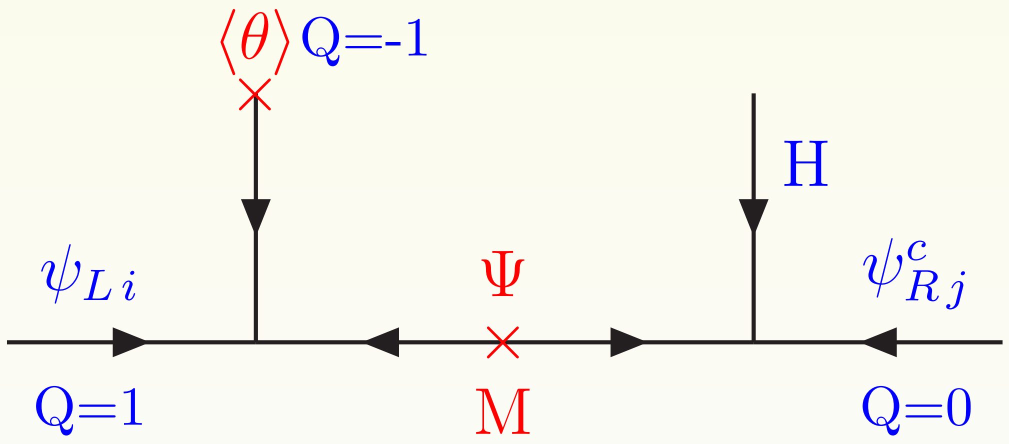
SM Flavour and CP

Fermion masses fixed by M_W. If $O(1)$ elements in Yukawa matrices

$\Rightarrow$ **Impossible reproduce masses, mixings !!!**

Flavour symmetries in SUSY

- Very different elements in Yukawa matrices: $y_t \simeq 1$, $y_u \simeq 10^{-5}$
- Expect couplings in a "fundamental" theory $\mathcal{O}(1)$
- Small couplings generated at higher order or function of small vevs.
- Froggatt-Nielsen mechanism and flavour symmetry to understand small Yukawa elements. Example: $U(1)_{fl}$



$$\Rightarrow Y_{ij} = \left(\frac{\langle\theta\rangle}{M}\right) \ll 1$$

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$\Downarrow$

We can relate the structure in Yukawa matrices to the nonuniversality in Soft Breaking masses !!!

Yukawa textures

- Masses and mixings in terms of a few fundamental parameters.
- Small mixing due to smallness of offdiagonal vs diagonal entries.
- Approximate texture zeros (GST) $\Rightarrow$ relate masses and mixings

Phenomenological fits:

$$Y_d \propto \begin{pmatrix} \leq \bar{\varepsilon}^5 & a\,\bar{\varepsilon}^3 & b\,\bar{\varepsilon}^3 \\ a\,\bar{\varepsilon}^3 & \bar{\varepsilon}^2 & c\,\bar{\varepsilon}^2 \\ \leq \bar{\varepsilon} & \leq 1 & 1 \end{pmatrix}, \qquad Y_u \propto \begin{pmatrix} \leq \varepsilon^4 & \varepsilon^3 & \mathcal{O}(\varepsilon^3) \\ \leq \varepsilon^3 & \varepsilon^2 & \mathcal{O}(\varepsilon^2) \\ \leq \varepsilon & \leq 1 & 1 \end{pmatrix}$$

with $\varepsilon \simeq 0.05$ and $\bar{\varepsilon} \simeq 0.15$

Symmetric texture

- Non-Abelian flavour symmetries.

$$Y^d = \begin{pmatrix} 0 & 1.5\ \varepsilon^3 & 0.4\ \varepsilon^3 \\ 1.5\ \varepsilon^3 & \varepsilon^2 & 1.3\ \varepsilon^2 \\ 0.4\ \varepsilon^3 & 1.3\ \varepsilon^2 & 1 \end{pmatrix} y_b$$

- Universal sfermion masses in in unbroken limit:

$$\mathcal{L}_{m^2} = m_0^2 \Phi^\dagger \Phi = m_0^2 \left(\phi_1\ \phi_2\ \phi_3\right)^* \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}$$

- After symmetry breaking:

$$M^2_{\tilde{D}_R} \simeq \begin{pmatrix} 1+\bar{\varepsilon}^3 & \bar{\varepsilon}^3 & 0 \\ \bar{\varepsilon}^3 & 1+\bar{\varepsilon}^2 & \bar{\varepsilon}^2 \\ 0 & \bar{\varepsilon}^2 & 1+\bar{\varepsilon} \end{pmatrix} m_0^2$$

Asymmetric texture

- Abelian flavour symmetries.

$$Y^d = \begin{pmatrix} \varepsilon^4 & \varepsilon^3 & \varepsilon^3 \\ \varepsilon^3 & \varepsilon^2 & \varepsilon^2 \\ \varepsilon & 1 & 1 \end{pmatrix} y_b$$

- In principle nonuniversal masses in unbroken symmetry:

$$\mathcal{L}_{m^2} = m_1^2\ \phi_1^*\phi_1 + m_2^2\ \phi_2^*\phi_2 + m_3^2\ \phi_3^*\phi_3$$

- After symmetry breaking:

$$M^2_{\tilde{D}_R} \simeq \begin{pmatrix} 1 & \bar{\varepsilon} & \bar{\varepsilon} \\ \bar{\varepsilon} & c & b \\ \bar{\varepsilon} & b & a \end{pmatrix} m_0^2$$

$SU(3)$ Flavour model

- $Q, L \sim \mathbf{3}$ and $d^c, u^c, e^c \sim \mathbf{3}$; flavon fields: $\theta_3, \theta_{23} \sim \overline{\mathbf{3}}$, $\overline{\theta}_3, \overline{\theta}_{23} \sim \mathbf{3}$
- Family Symmetry breaking: $SU(3) \xrightarrow{\langle\theta_3\rangle} SU(2) \xrightarrow{\langle\theta_{23}\rangle} \emptyset$

$$\theta_3, \overline{\theta}_3 = \begin{pmatrix} 0 \\ 0 \\ a_3 \end{pmatrix}, \quad \theta_{23}, \overline{\theta}_{23} = \begin{pmatrix} 0 \\ b \\ b \end{pmatrix} \text{with} \ \left(\frac{a_3}{M}\right) \sim \mathcal{O}(1), \left(\frac{b}{M_u}\right) \simeq \left(\frac{b}{M_d}\right)^2 = \varepsilon \sim 0.05.$$

- Yukawa superpotential: $W_Y = H\psi_i\psi_j^c\left[\theta_3^i\theta_3^j + \theta_{23}^i\theta_{23}^j\left(\theta_3\overline{\theta_3}\right) + \epsilon^{ikl}\overline{\theta}_{23,k}\overline{\theta}_{3,l}\theta_{23}^j\left(\theta_{23}\overline{\theta_3}\right)\right]$

$$Y^f = \begin{pmatrix} 0 & a\ \varepsilon^3 & b\ \varepsilon^3 \\ a\ \varepsilon^3 & \varepsilon^2 & c\ \varepsilon^2 \\ b\ \varepsilon^3 & c\ \varepsilon^2 & 1 \end{pmatrix} \frac{|a_3|^2}{M^2},$$

- Soft mass coupling $\Phi^\dagger \Phi$ invariant $\Rightarrow$ common soft mass for the triplet
- Universality guaranteed in the exact symmetry limit.
- After breaking new entries proportional to (complex) flavon vevs. Minimal terms:

$$M^2_{ij\,\mathrm{min}} = m_0^2 \Big(\delta^{ij} + \frac{1}{M^2}[\theta_3^{i\dagger}\theta_3^j + \overline{\theta}_3^{i\dagger}\overline{\theta}_3^j + \theta_{23}^{i\dagger}\theta_{23}^j + \overline{\theta}_{23}^{i\dagger}\overline{\theta}_{23}^j] + \frac{1}{M^4}[(\epsilon^{ikl}\overline{\theta}_{3,k}\overline{\theta}_{23,l})^\dagger(\epsilon^{jmn}\overline{\theta}_{3,m}\overline{\theta}_{23,n}) + (\epsilon_{ikl}\theta_3^k\theta_{23}^l)^\dagger(\epsilon_{jmn}\theta_3^m\theta_{23}^n)] + \dots\Big)$$

- However, Kähler is real function not holomorfic and new combinations (if) allowed by discrete symmetries possible without spoiling the superpotential:

$$M^2_{ij} = M^2_{ij\,\mathrm{min}} + m_0^2\,\frac{1}{M^2}[\theta_3^i\overline{\theta}_{23,j} + h.c.]$$

- Soft matrices now change in the (2,3) sector, from the minimal:

$$M^2_{\tilde{D}_R}{}^{\rm SCKM} \simeq 6\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1+\bar{\varepsilon}^3 & \bar{\varepsilon}^3 & \bar{\varepsilon}^3 \\ \bar{\varepsilon}^3 & 1+\bar{\varepsilon}^2 & \bar{\varepsilon}^2 \\ \bar{\varepsilon}^3 & \bar{\varepsilon}^2 & 1+a \end{pmatrix} m_0^2$$

$(a \sim O(1))$ to the following texture (with O(1) phases):

$$M^2_{\tilde{D}_R}{}^{\rm SCKM} \simeq 6\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1+\bar{\varepsilon}^3 & \bar{\varepsilon}^3 & \bar{\varepsilon}^3 \\ \bar{\varepsilon}^3 & 1+\bar{\varepsilon}^2 & \bar{\varepsilon} \\ \bar{\varepsilon}^3 & \bar{\varepsilon} & 1+a \end{pmatrix} m_0^2$$

$$\Downarrow$$

Possible interesting phenomenology in the B_S sector !!!

- Soft mass coupling $\Phi^\dagger\Phi$ invariant $\Longrightarrow$ common soft mass for the triplet
- Universality guaranteed in the exact symmetry limit.
- After symmetry breaking offdiagonal entries proportional to (complex) flavon vevs

$$M_{ij}^2 = m_0^2 \Big(\delta^{ij} + \frac{1}{M^2}[\theta_3^{i\dagger}\theta_3^j + \overline{\theta}_3^{i\dagger}\overline{\theta}_3^j + \theta_{23}^{i\dagger}\theta_{23}^j + \overline{\theta}_{23}^{i\dagger}\overline{\theta}_{23}^j] + \frac{1}{M^4}[(\epsilon^{ikl}\overline{\theta}_{3,k}\overline{\theta}_{23,l})^\dagger(\epsilon^{jmn}\overline{\theta}_{3,m}\overline{\theta}_{23,n}) + (\epsilon_{ikl}\theta_3^k\theta_{23}^l)^\dagger(\epsilon_{jmn}\theta_3^m\theta_{23}^n)] + \ldots \Big)$$

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(with $\bar\varepsilon \simeq 0.15$, $\varepsilon \simeq 0.05$)

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$$M^2_{\tilde{D}_R}{}^{\mathrm{SCKM}} \simeq 6\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1 & 0.003 & 0.003 \\ 0.003 & 1 & 0.02 \\ 0.003 & 0.02 & 1 \end{pmatrix} m_0^2$$

(with $\bar{\varepsilon} \simeq 0.15, \varepsilon \simeq 0.05$)

At M_W in the SCKM basis:

$$M^2_{\tilde{D}_L} \simeq 6\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1+\varepsilon^3 & \varepsilon^2\bar{\varepsilon} & \varepsilon^2\bar{\varepsilon} + c_{\rm run}\ \bar{\varepsilon}^3 \\ \varepsilon^2\bar{\varepsilon} & 1+\varepsilon^2 & \varepsilon^2 + c_{\rm run}\ \bar{\varepsilon}^2 \\ \varepsilon^2\bar{\varepsilon} + c_{\rm run}\ \bar{\varepsilon}^3 & \varepsilon^2 + c_{\rm run}\ \bar{\varepsilon}^2 & 1+\bar{\varepsilon} \end{pmatrix} m_0^2$$

$$M^2_{\tilde{E}_R} \simeq 0.15\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1+\bar{\varepsilon}^3 & \frac{\bar{\varepsilon}^3}{3} & \bar{\varepsilon}^3 \\ \frac{\bar{\varepsilon}^3}{3} & 1+\bar{\varepsilon}^2 & \bar{\varepsilon}^2 \\ \bar{\varepsilon}^3 & \bar{\varepsilon}^2 & 1+\bar{\varepsilon} \end{pmatrix} m_0^2$$

$$M^2_{\tilde{E}_L} \simeq 0.5\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1+\varepsilon^3 & \frac{\varepsilon^2\bar{\varepsilon}}{3} & \varepsilon^2\bar{\varepsilon} + c_{\rm run}\ \bar{\varepsilon}^3 \\ \frac{\varepsilon^2\bar{\varepsilon}}{3} & 1+\varepsilon^2 & \varepsilon^2 + 3\ c_{\rm run}\ \bar{\varepsilon}^2 \\ \varepsilon^2\bar{\varepsilon} + c_{\rm run}\ \bar{\varepsilon}^3 & \varepsilon^2 + 3\ c_{\rm run}\ \bar{\varepsilon}^2 & 1+\varepsilon \end{pmatrix} m_0^2$$

At M_W in the SCKM basis:

$$M^2_{\tilde{D}_L} \simeq 6\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1 & 4\times 10^{-4} & 7\times 10^{-4} \\ 4\times 10^{-4} & 1 & 4\times 10^{-3} \\ 7\times 10^{-4} & 4\times 10^{-3} & 1 \end{pmatrix} m_0^2$$

$$M^2_{\tilde{E}_R} \simeq 0.15\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1 & 0.001 & 0.003 \\ 0.001 & 1 & 0.02 \\ 0.003 & 0.02 & 1 \end{pmatrix} m_0^2$$

$$M^2_{\tilde{E}_L} \simeq 0.5\, M^2_{1/2}\, \mathbb{1} + \begin{pmatrix} 1 & 1\times 10^{-4} & 7\times 10^{-4} \\ 1\times 10^{-4} & 1 & 1\times 10^{-2} \\ 7\times 10^{-4} & 1\times 10^{-2} & 1 \end{pmatrix} m_0^2$$

FCNC constraints

- Large offdiagonal entries in sfermion mass matrices generally overproduce FCNC and CP Violation transitions

$\Longrightarrow$ SUSY flavour problem

- Strong phenomenological bounds on Mass Insertions

$$\left(\delta^f_A\right)_{ij} = \frac{(m^2_{\tilde{f}_A})_{ij}}{m^2_{\tilde{f}}}$$

- Very stringent bounds on $d \to s$ transitions from ΔM_k and ε_k:

$$\mathrm{Re}\{\left(\delta^d_R\right)_{12}\} \leq 4 \times 10^{-2}, \qquad \mathrm{Im}\{\left(\delta^d_R\right)_{12}\} \leq 3.2 \times 10^{-3}$$

- Less stringent bounds from $b \to d$ and $b \to s$ transitions

$$\mathrm{Re}\{\left(\delta^d_R\right)_{13}\},\ \mathrm{Im}\{\left(\delta^d_R\right)_{13}\} \leq 0.1$$

($\Rightarrow$ Simple abelian models not allowed by ΔM_k and ε_k)

Spontaneous CP violation

- CP spontaneously broken in the flavour sector by complex flavon vevs.

$$\langle\theta_3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} a_3^u & 0 \\ 0 & a_3^d\ e^{i\chi} \end{pmatrix}, \quad \langle\bar{\theta}_3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} a_3^u\ e^{i\alpha_u} & 0 \\ 0 & a_3^d\ e^{i\alpha_d} \end{pmatrix},$$

$$\langle\theta_{23}\rangle = \begin{pmatrix} 0 \\ b_{23} \\ b_{23}\ e^{i\beta_3} \end{pmatrix}, \quad \langle\bar{\theta}_{23}\rangle = \begin{pmatrix} 0 \\ b_{23}\ e^{i\beta_2'} \\ b_{23}\ e^{i(\beta_2'-\beta_3)} \end{pmatrix}.$$

- Model dependent: Explicit example from Ross, Velasco-Sevilla and O.V.

$$M^2_{\tilde{E}_R}(M^2_{\tilde{E}_L}) \propto \begin{pmatrix} 1+\bar{\varepsilon}^3 & \frac{\bar{\varepsilon}^3}{3} & \bar{\varepsilon}^3 e^{-i(\beta_3-\chi)} \\ \frac{\bar{\varepsilon}^3}{3} & 1+\bar{\varepsilon}^2 & \bar{\varepsilon}^2 e^{-i(\beta_3-\chi)} \\ \bar{\varepsilon}^3 e^{-i(\beta_3-\chi)} & \bar{\varepsilon}^2 e^{-i(\beta_3-\chi)} & 1+\bar{\varepsilon} \end{pmatrix} m_0^2$$

Solution to the SUSY CP problem

- $SU(3)_{fl}$ and CP spontaneously broken in Yukawas at $M \ll M_{\rm Planck}$
- At $M_{\rm Planck}$, Kähler (soft masses) real and universal and Guidice-Masiero μ term real.
- After $SU(3)$ breaking, Yukawa matrices and offdiagonal elements in soft masses contain $\mathcal{O}(1)$ CP violating phases (δ_{CKM})
- Trilinear couplings, Y^A, same (leading order) structure as Y.

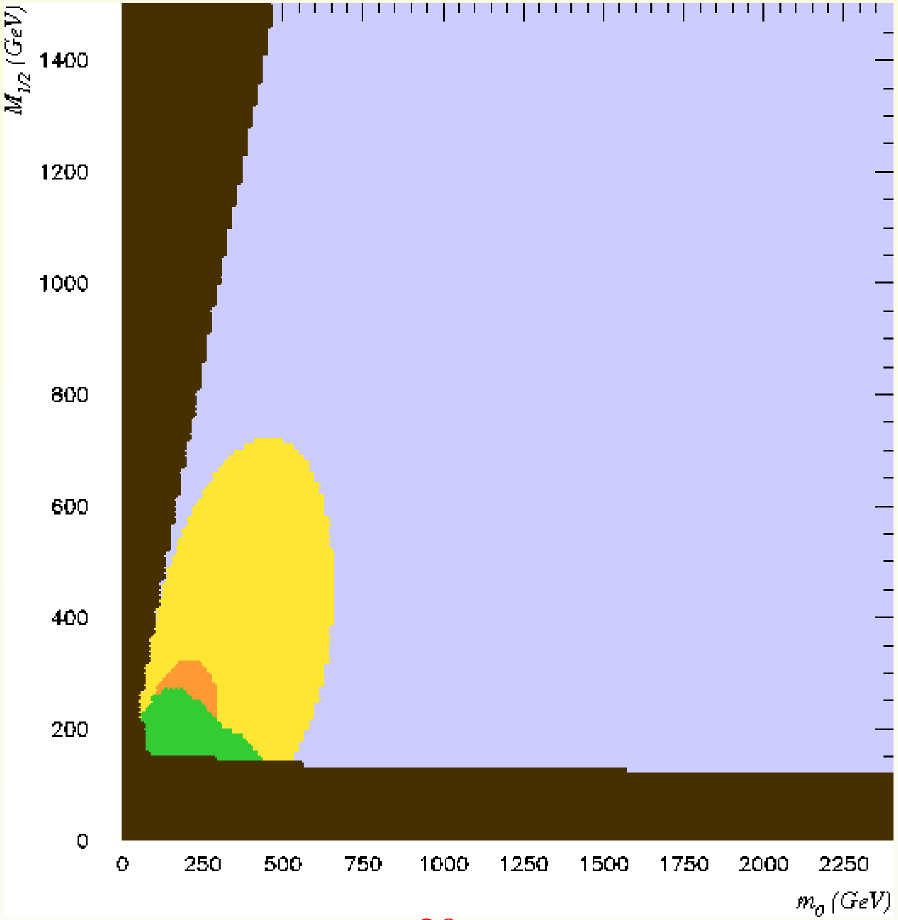
$\Longrightarrow$ Diagonal elements in Y^A are real at leading order in SCKM basis.

Still contributions to EDMs from offdiagonal elements in sfermion masses:

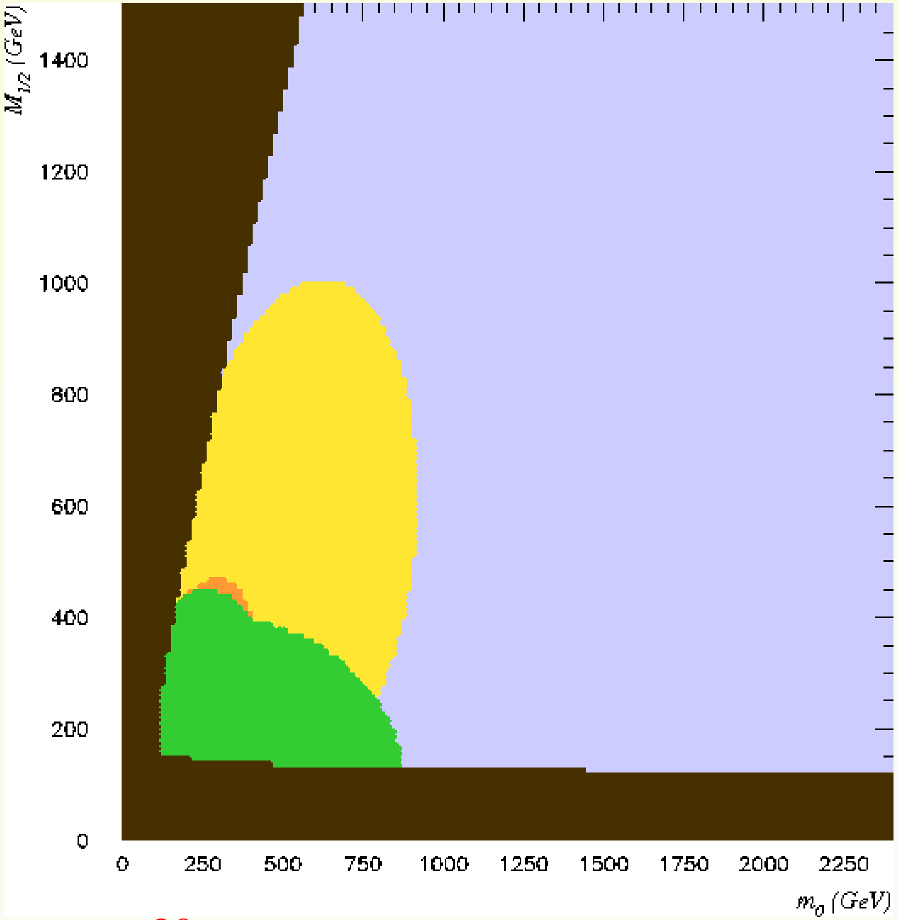
$$d_e \propto (\delta^e_{LL})_{1i}(\delta^e_{LR})_{i1}\, f_1 + (\delta^e_{LR})_{1i}(\delta^e_{RR})_{i1}\, f_2 + (\delta^e_{LL})_{1i}(\delta^e_{LR})_{ij}(\delta^e_{RR})_{j1} f_3$$

lepton EDMs

RVV model 1 $\tan\beta = 10$, $A_0 = 0$



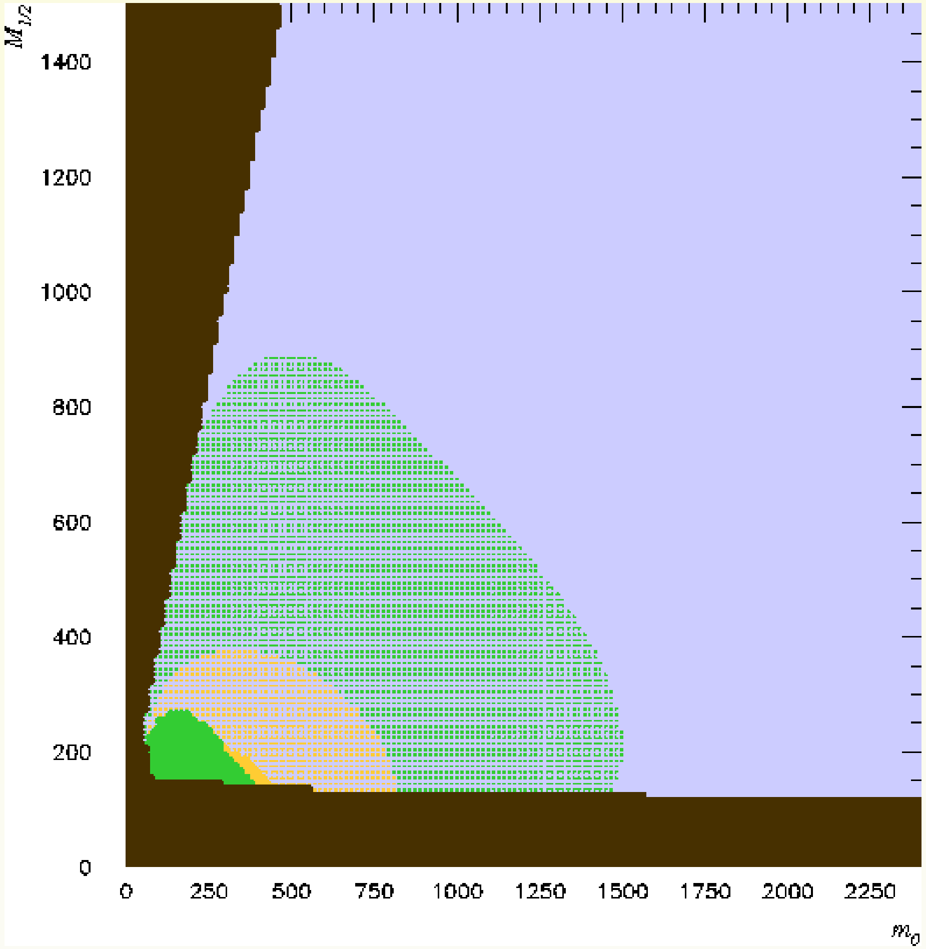
RVV model 1 $\tan\beta = 30$, $A_0 = 0$



Yellow: $d_e = 10^{-29}$ e cm, Orange: $d_e = 5 \times 10^{-29}$ e cm, green: LFV constraints.

Lepton Flavour Violation

$\tan\beta = 10,\ A_0 = 0$



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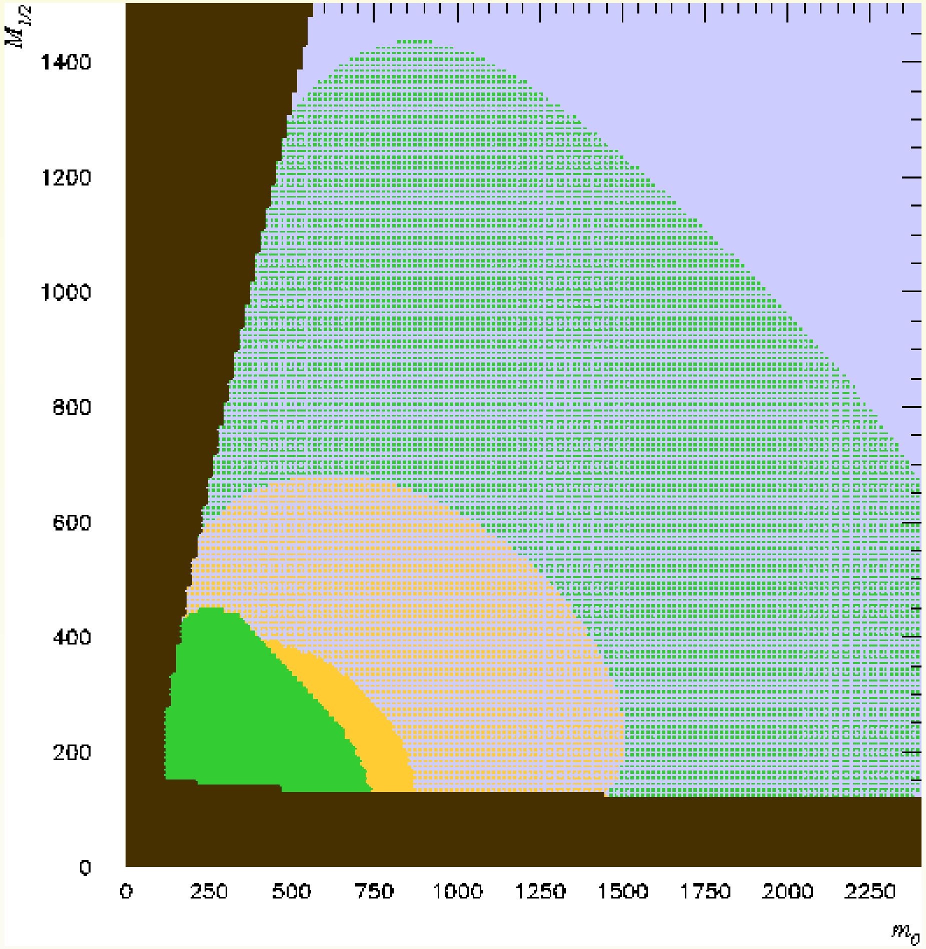
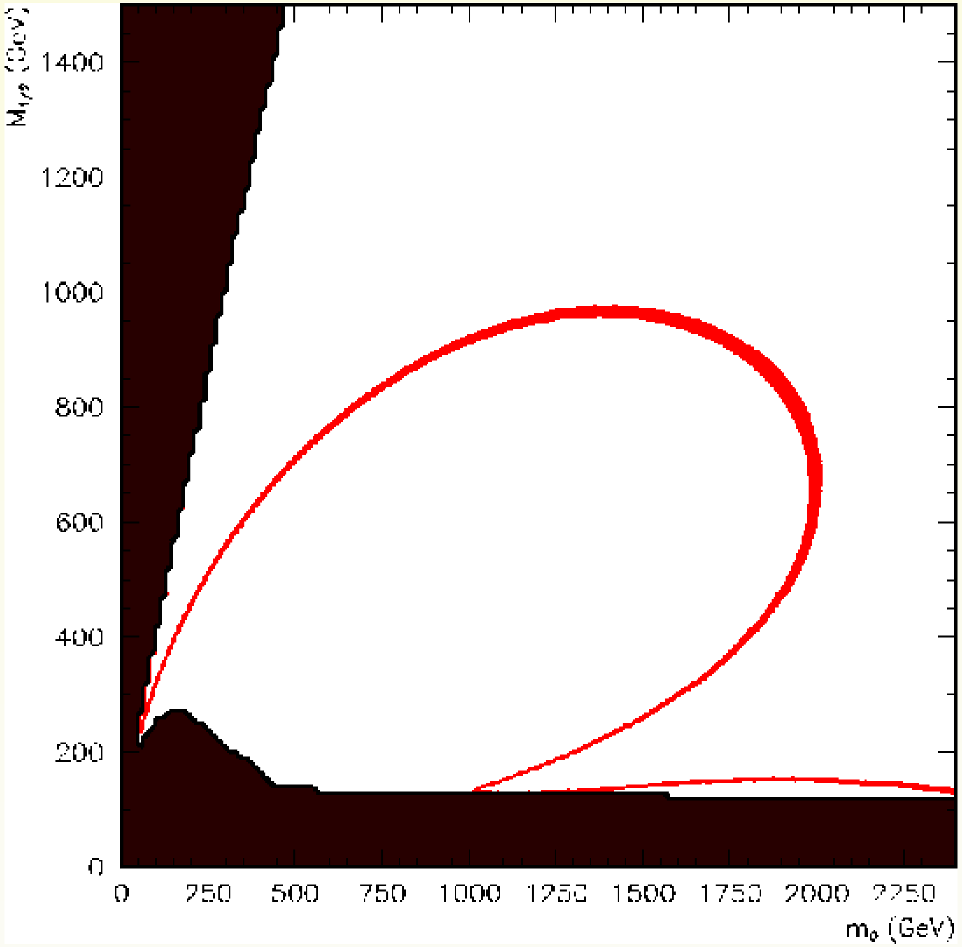


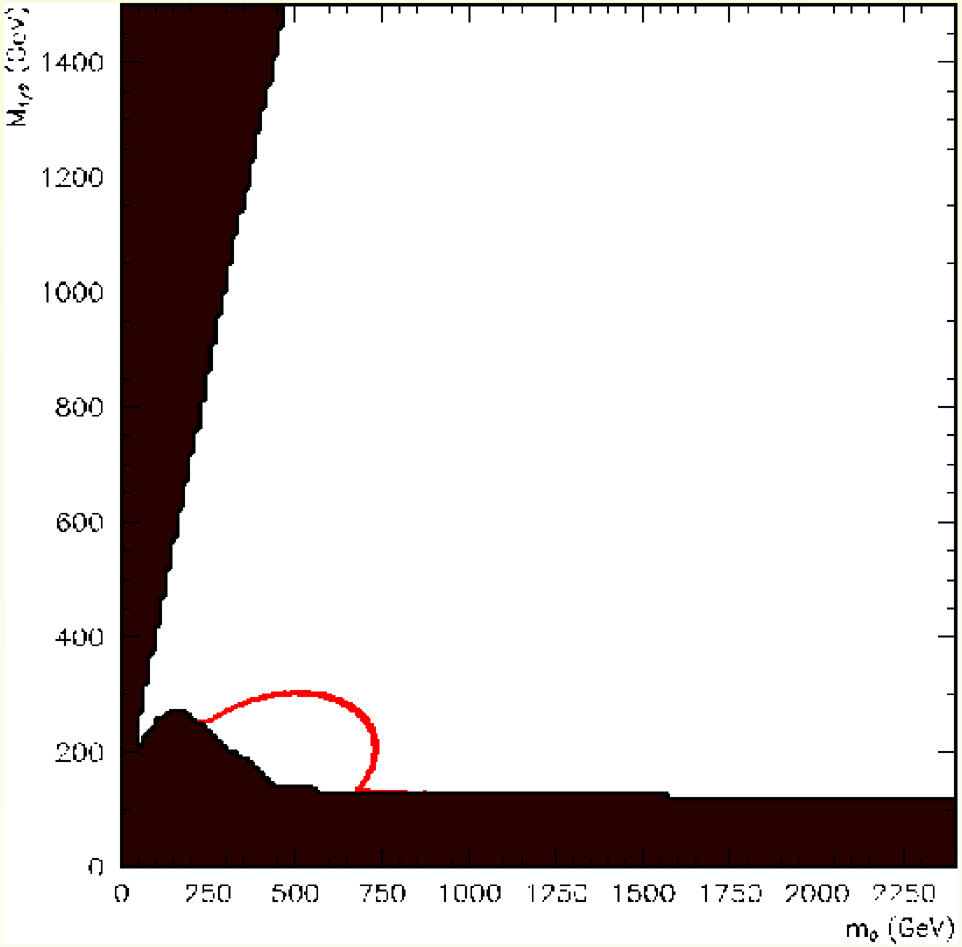
Green (hatched): Present (fut.) $\mu \to e\gamma$ bounds, Yellow: Present (fut.) $\tau \to \mu\gamma$ bounds.

ε_K and B_s mixing

$\tan\beta = 10$, $\varepsilon_K^{\mathrm{SUSY}} = 0.2\varepsilon_K^{\mathrm{exp}}$



$\tan\beta = 10$, $\varepsilon_K^{\mathrm{SUSY}} = -2\varepsilon_K^{\mathrm{exp}}$,



However, things are difficult in B system...

- SM phase in B_s small: $\beta_S = 0.035$, where the SM contribution to mixing:

$$M_{12}^{\rm SM} \simeq \frac{\alpha_{em}^2}{8 M_W^2 \sin^2\theta_W}\, \frac{m_t^2}{M_W^2}\, \tfrac{1}{3} f_B^2 B_B \left(V_{tb}^* V_{ts}\right)^2$$

while SUSY contribution:

$$M_{12}^{\rm SUSY} \simeq \frac{\alpha_s^2}{216 M_{\rm SUSY}^2}\, f(x)\, \tfrac{1}{3} f_B^2 B_B \left(\delta_{LL}^d\right)_{12}^2$$

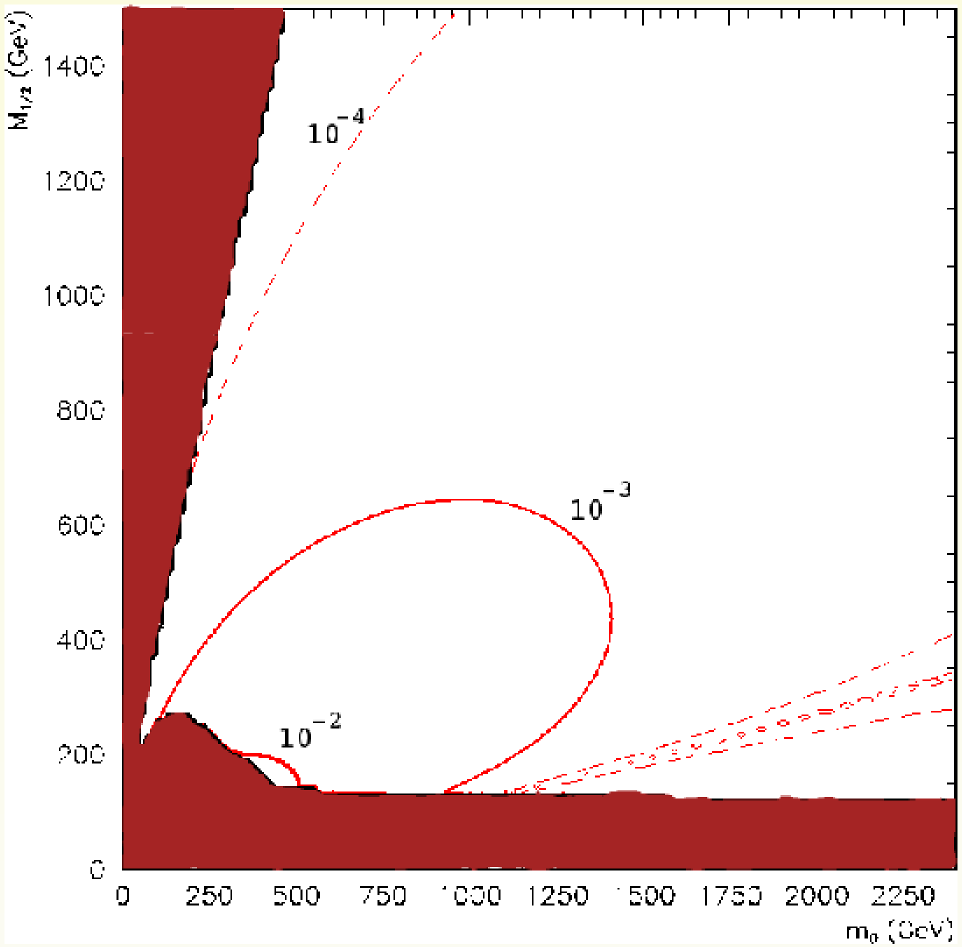
- To have a large phase in mixing $M_{12}^{B_s} = M_{12}^{\rm SM} + M_{12}^{\rm SUSY}$, we need,

$$1 \simeq \frac{M_{12}^{\rm SUSY}}{M_{12}^{\rm SM}} = \frac{\alpha_s^2 \sin^2\theta_W}{\alpha_{em}^2}\, \frac{M_W^2}{m_t^2 M_{\rm SUSY}^2}\, \frac{8 f(x)}{216}\, \frac{\left(\delta_{LL}^d\right)_{12}^2}{\left(V_{tb}^* V_{ts}\right)^2} = 12.5 \times 0.005 \times 0.04 \times \frac{\left(\delta_{LL}^d\right)_{12}^2}{(0.008)^2}$$

Thus, to have a large phase in B_s $\Longrightarrow$ $\left(\delta_{LL}^d\right)_{12}^2 \geq 0.16$

Also, $|M_{12}^{B_s}|$ roughly agrees with SM prediction ...

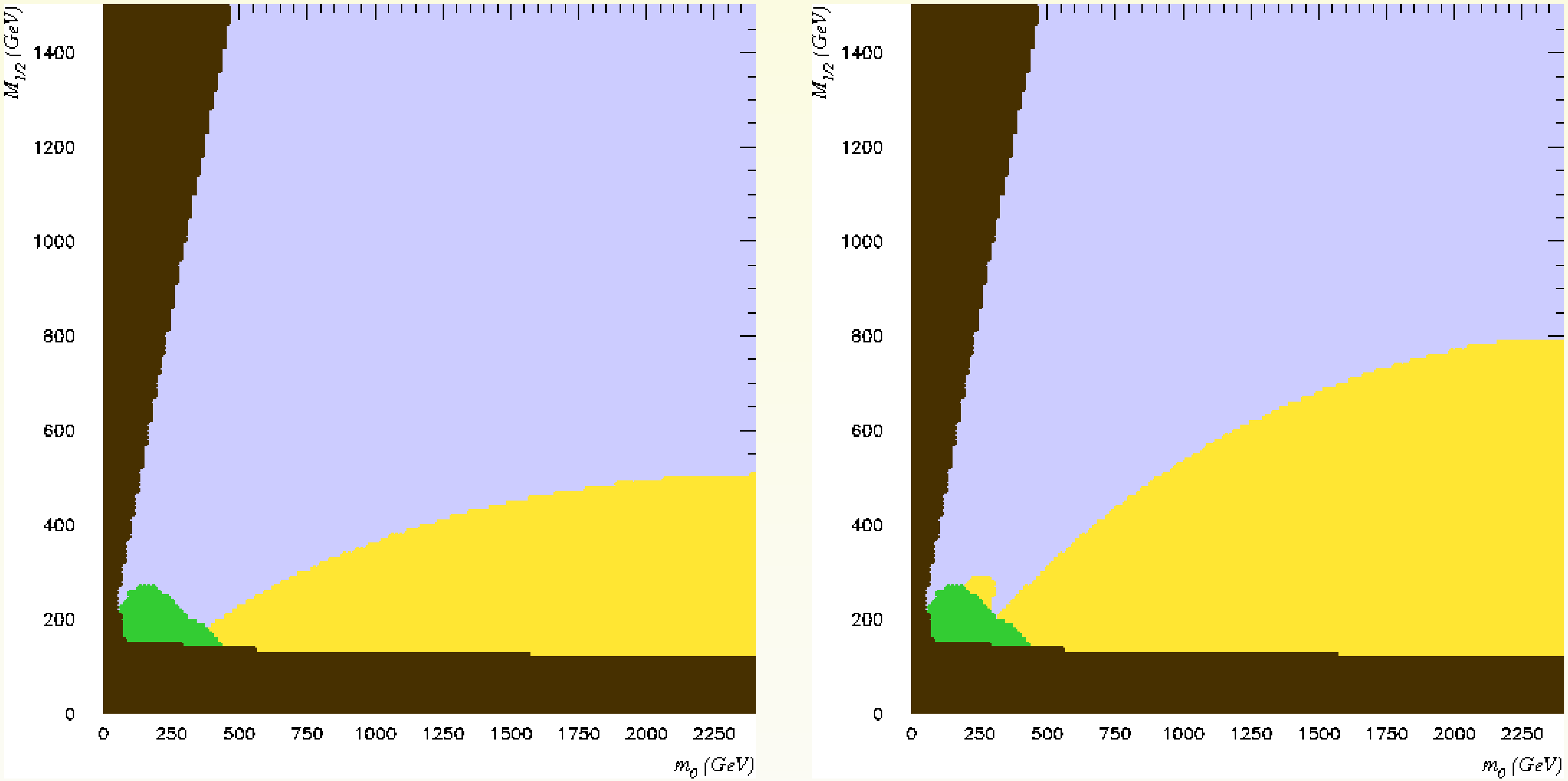
$$\phi_{B_s} = \arg\left[M_{12}^{\rm SUSY}\right]$$



to be compared with $\beta_s^{\rm SM} = 0.034$

Neutron EDM

Chiral quark model, $\tan\beta = 10$ Quark parton model, $\tan\beta = 10$



Yellow: $d_n = 1 \times 10^{-28}$ e cm (Future bound)

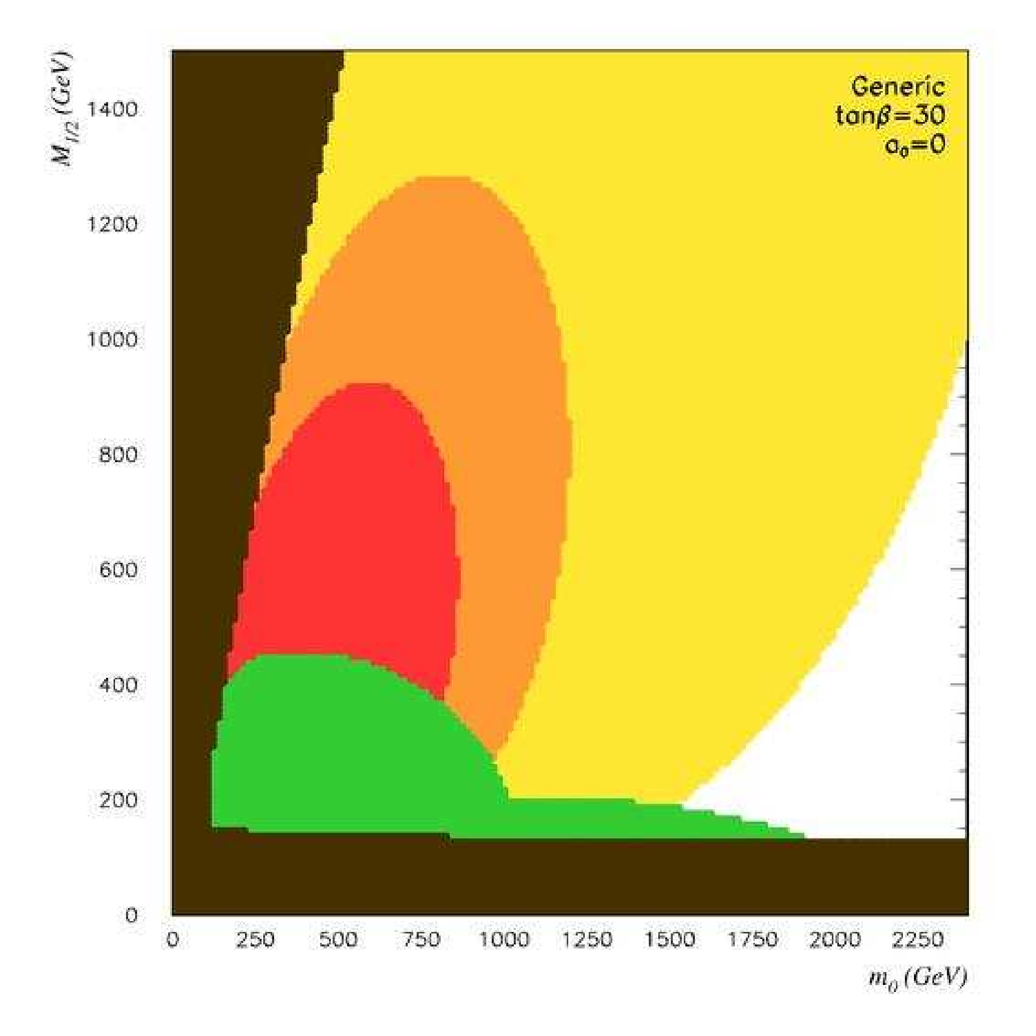
Conclusions

Flavour symmetries solve the Flavour and CP problems
both in SUSY and in the SM!

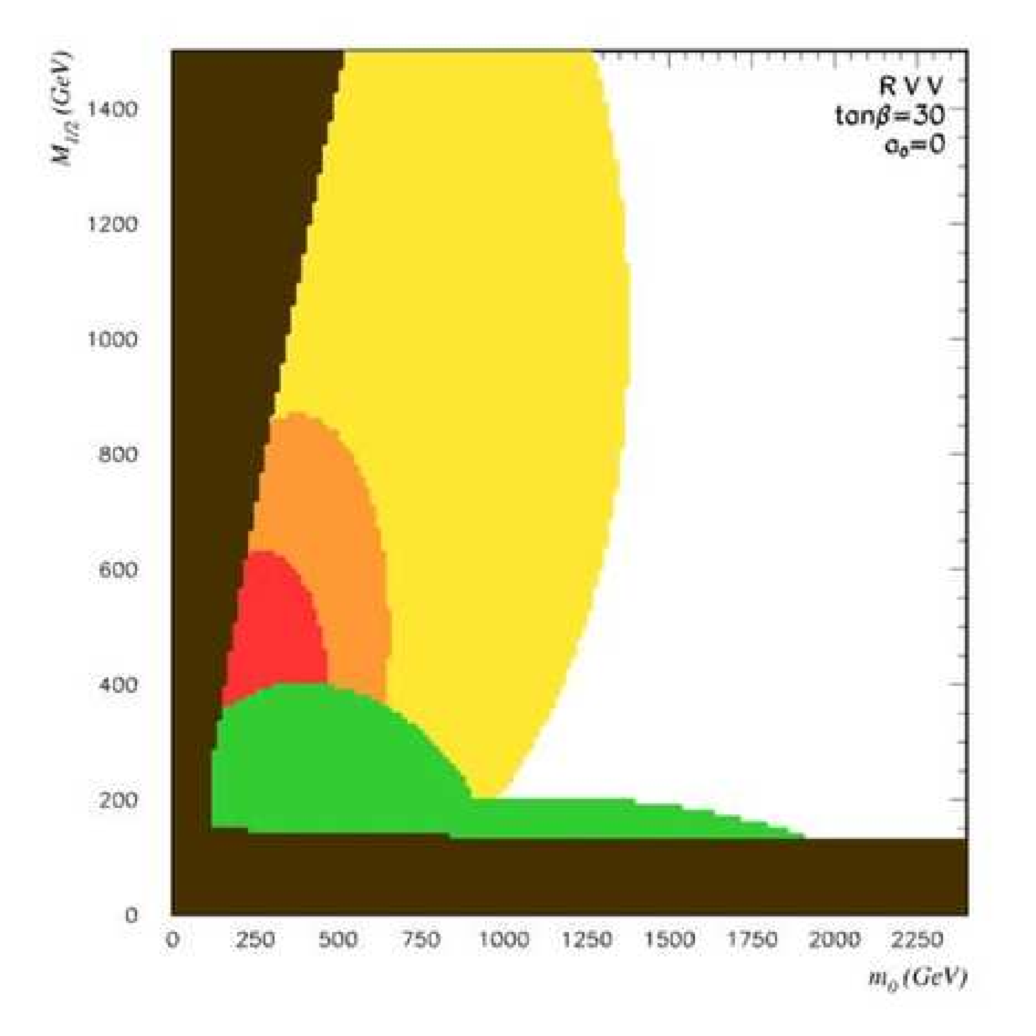
$\Downarrow$

- Flavour phases (already obs. in Yukawas) contribute to EDMs.
- d_e and d_n in reach of proposed experiments for LHC sfermions.
- LFV processes ($\mu \to e\gamma$) close to present exp. limits.
- Sizeable contribution in the Kaon sector natural.
- LFV and EDMs can explore large areas of flavour MSSM in near future.

Generic $\tan\beta = 30,\ A_0 = 0$

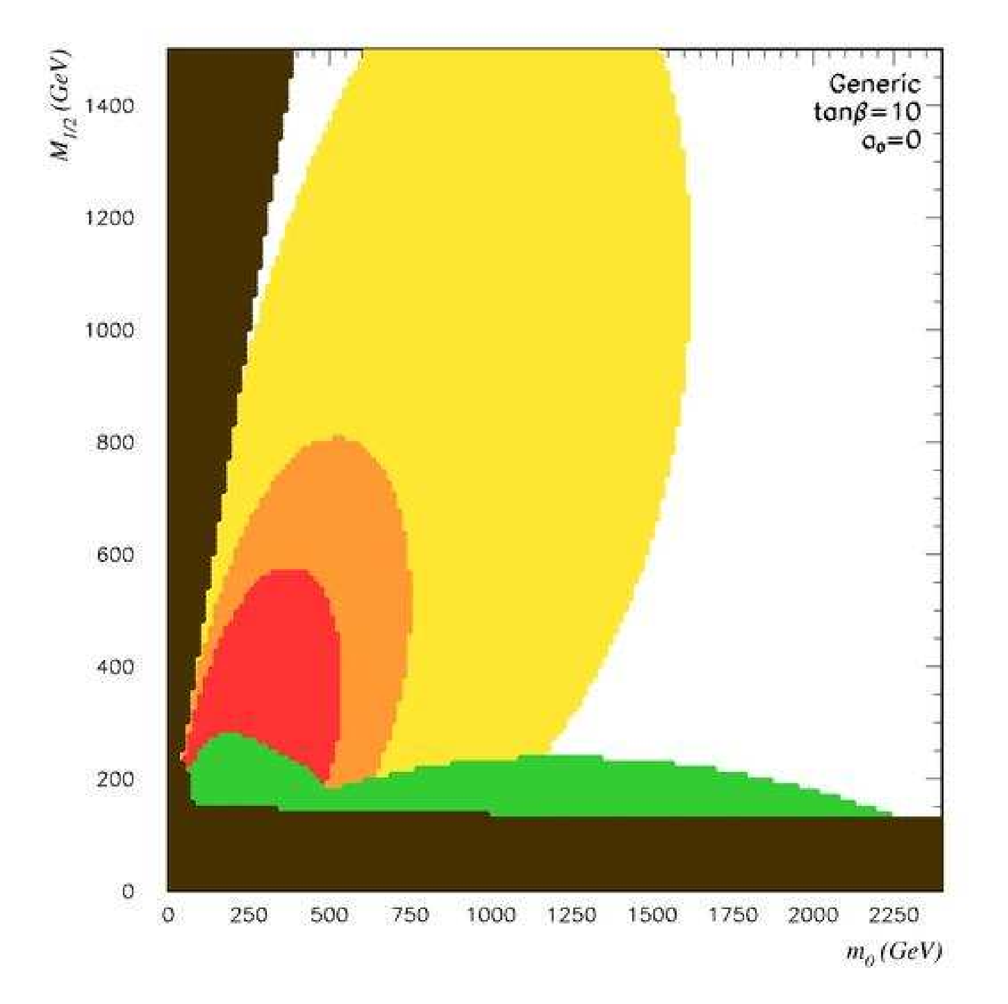


RVV $\tan\beta = 30,\ A_0 = 0$

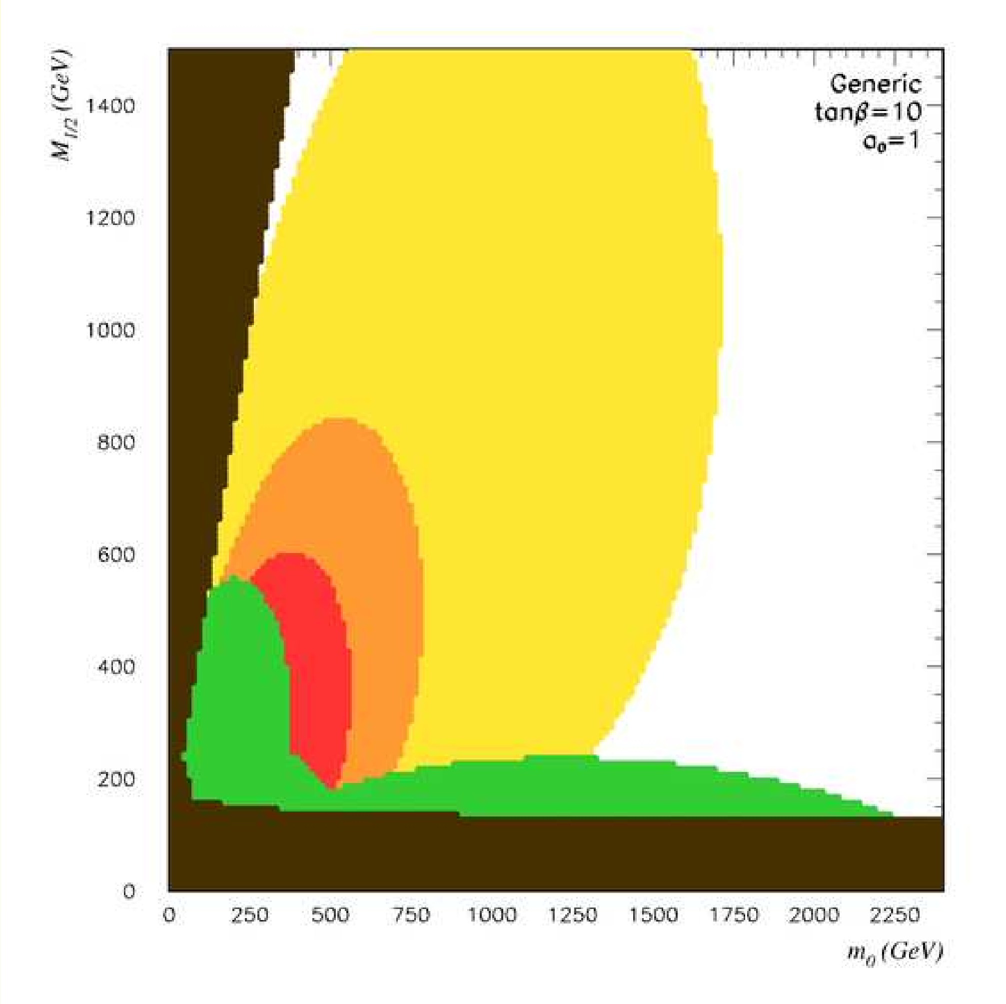


$$d_e \propto m_\tau \mu \tan\beta \cdot \mathrm{Im}\left[\delta^{e_R}_{13} \cdot \delta^{e_L}_{31}\right]$$

generic phases $\tan\beta = 10,\ A_0 = 0$



generic phases $\tan\beta = 10,\ A_0 = m_0$



Yellow: $d_e = 10^{-29}$ e cm, Orange: $d_e = 5 \times 10^{-29}$ e cm, green: LFV constraints

Field	θ_3	θ_{23}	$\overline{\theta_3}$	$\overline{\theta_{23}}$	Σ	H	ψ	W
SU(3)	$\bar{3}$	$\bar{3}$	3	1	1	3	1	1
U(1)	1	0	0	2	1	0	-1	0
$\mathbf{Z_9}$	0	0	1	-1	-2	0	0	0
$\mathbf{Z_6}$	0	1	0	-2	-2	0	0	0

with $\theta_3, \overline{\theta_3} \sim \mathbf{3} \oplus \mathbf{1}$, under $SU(2)_R$.

Yukawa Superpotential

$$W_Y = H\psi_i\psi_j^c \left[\theta_3^i\theta_3^j + \theta_{23}^i\theta_{23}^j\Sigma + \left(\epsilon^{ikl}\overline{\theta}_{23,k}\overline{\theta}_{3,l}\theta_{23}^j \left(\theta_{23}\overline{\theta_3}\right) + (i \leftrightarrow j)\right)\right.$$
$$\left. + X\,\epsilon^{ijk}\overline{\theta}_{23,k}\left(\theta_{23}\overline{\theta_3}\right)^2 + X\,\epsilon^{ijk}\overline{\theta}_{3,k}\left(\theta_{23}\overline{\theta_3}\right)\left(\theta_{23}\overline{\theta_{23}}\right) + \ldots\right]$$

Spontaneous Symmetry Breaking

Field	P	S	$\bar{S}$	T	U	$\bar{U}$	V	Y	Z
$\mathbf{SU(3)}$	1	1	1	1	1	1	1	1	1
$\mathbf{U(1)_{PQ}}$	-2	$-9/2$	$9/2$	0	7	-7	-2	-9	-2
$\mathbf{Z_{15}}$	3	1	3	3	-2	-7	8	2	9

Then the flavon Superpotential,

$$W = P\left(T^4 + S^3\bar{S}^3\right) + U\left((\theta_{23}\bar{\theta}_{23}) + S^2\right) + V\left((\theta_3\bar{\theta}_3)^4 + S\bar{S}\right) + \bar{U}\bar{S}^2(\theta_3\bar{\theta}_3)$$
$$+ Y(\theta_3\bar{\theta}_2) + Z\left((\theta_{23}\bar{\theta}_3)(\theta_{23}\bar{\theta}_2) + \bar{S}^2\right) + \mu H_1 H_2\left[1 + (\theta_{23}\bar{\theta}_3)^5 + \ldots\right]$$

Phase determination: 1) T gets a vev radiatively with $\varphi_T = \frac{n\pi}{5}$,

2) $\varphi_{(S\bar{S})} = \frac{n'4}{3}\,\varphi_T$, 3) $\alpha_3 = \frac{m}{4}\,\varphi_{(S\bar{S})}$, 4) $\beta_3 = \frac{m'\pi}{5} - \alpha_3$