## Kepler's laws

The explanation on Kepler's laws in the textbook is somewhat insufficient. I will show how these laws arise from the equation of motion.

#### 1 Ellipses

We consider an ellipse of semimajor axis a, semiminor axis b, half-distance between the focuses c. Obviously,  $a^2 = b^2 + c^2$ . The eccentricity is e = c/a. If we take one of the focuses as the origin, the equation of the ellipse in the cartesian coordinates can be written as

$$\frac{(x+c)^2}{a^2} + \frac{y^2}{b^2} = 1.$$
 (1)

In polar coordinates, the ellipse is

$$r = \frac{\ell}{1 + e\cos\theta},\tag{2}$$

where  $\ell = b^2/a$ . These are related, as usual, by  $x = r \cos \theta$  and  $y = r \sin \theta$ .

### 2 The equation of motion

Let the mass of the heavier star ('the sun') M and that of the lighter star ('the earth') m. The gravitational force exerted on the earth is

$$\vec{F} = -G\frac{Mm}{r^2}\vec{e_r},\tag{3}$$

where

$$\vec{e}_r = \begin{pmatrix} \cos\theta\\ \sin\theta \end{pmatrix} \tag{4}$$

is a unit vector on the earth, pointing in the opposite direction to the sun. Also defining an orthogonal unit vector

$$\vec{e}_{\theta} = \begin{pmatrix} -\sin\theta \\ \cos\theta \end{pmatrix},\tag{5}$$

it is easy to show that the differentials of these unit vectors are

$$d\vec{e}_r = \vec{e}_\theta d\theta, \qquad d\vec{e}_\theta = -\vec{e}_r d\theta. \tag{6}$$

As the position of the earth is

$$\begin{pmatrix} x\\ y \end{pmatrix} = r\vec{e_r},\tag{7}$$

the second law of Newton, with the force (3), is

$$-G\frac{Mm}{r^2}\vec{e_r} = m\frac{d^2}{dt^2}(r\vec{e_r}).$$
(8)

Using (6), we have (the overdot denoting the time derivative  $\frac{d}{dt}$ )

$$\frac{d^2}{dt^2}r\vec{e}_r = (\ddot{r} - r\dot{\theta}^2)\vec{e}_r + (2\dot{r}\dot{\theta} + r\ddot{\theta})\vec{e}_{\theta}.$$
(9)

Then the equation of motion (8) becomes two equations,

$$-G\frac{M}{r^2} = \ddot{r} - r\dot{\theta}^2,\tag{10}$$

$$2\dot{r}\dot{\theta} + r\ddot{\theta} = 0. \tag{11}$$

## 3 Kepler's 2nd law

Kepler's 2nd law states that the radius vector of a planet  $(r\vec{e}_r)$  sweeps equal area in unit time. The equation (11) can be written as  $\frac{d}{dt}(r^2\dot{\theta}) = 0$ , meaning that

$$A = r^2 \dot{\theta},\tag{12}$$

is a constant. Since  $A/2 = \frac{1}{2}r^2\dot{\theta}$  is the area velocity, namely the area swept by the radius vector of the planet in unit time, we have shown the 2nd law of Kepler.

#### 4 Kepler's 1st law

The nasty-looking differential equation (10) can actually be solved easily if you introduce a new variable  $q \equiv 1/r$ . Using (12),

$$\frac{dq}{d\theta} = \frac{d}{d\theta} \left(\frac{1}{r}\right) = -\frac{1}{r^2} \frac{dr}{d\theta} = -\frac{1}{r^2} \frac{dr}{dt} \frac{dt}{d\theta} = -\frac{1}{A} \dot{r},\tag{13}$$

and then

$$\frac{d^2q}{d\theta^2} = -\frac{1}{A}\frac{d\dot{r}}{d\theta} = -\frac{1}{A}\ddot{r}\frac{dt}{d\theta} = -\frac{r^2}{A^2}\ddot{r}.$$
(14)

Using this equation and (12), we can rewrite (10) as

$$\frac{d^2q}{d\theta^2} + q - \frac{GM}{A^2} = 0. \tag{15}$$

This is a 2nd order differential equation for a harmonic oscillator. The solution is

$$q - \frac{GM}{A^2} = B\cos(\theta + \phi), \tag{16}$$

where B is the amplitude of the oscillator and  $\phi$  is a phase (B and  $\phi$  are the integration constants). Choosing  $\phi = 0$  by shifting the origin of time t, we find

$$r = \frac{1}{q} = \frac{\frac{A^2}{GM}}{1 + \frac{A^2B}{GM}\cos\theta}.$$
(17)

This is an ellipse in the polar coordinates (2), with

$$\ell = \frac{b^2}{a} = \frac{A^2}{GM}, \qquad e = \frac{c}{a} = \frac{A^2B}{GM}.$$
(18)

This shows Kepler's 1st law.

#### 5 Kepler's 3rd law

The period of the orbital motion T, is equal to the area of the ellipse divided by the area swept by the radius vector in unit time, i.e.

$$T = \frac{\pi ab}{A/2}.\tag{19}$$

Using the first equation of (18), we have

$$T^{2} = \frac{4\pi^{2}a^{2}b^{2}}{A^{2}} = \frac{4\pi^{2}a^{3}}{GM} \propto a^{3}.$$
 (20)

This shows Kepler's 3rd law.

# 6 The mechanical energy of a planet

The mechanical energy (the sum of the kinetic and potential energy) of the planet can be found as follows. The position of the planet is  $\vec{r} = r\vec{e}_r$  and hence the velocity is

$$\vec{v} = \frac{d\vec{r}}{dt} = \frac{dr}{dt}\vec{e}_r + r\frac{d\vec{e}_r}{dt} = \dot{r}\vec{e}_r + r\dot{\theta}\vec{e}_\theta, \qquad (21)$$

where (6) has been used. Noticing that  $\vec{e}_r$  and  $\vec{e}_{\theta}$  are unit vectors which are orthogonal to each other, we have  $v^2 = \dot{r}^2 + r^2 \dot{\theta}^2$ . Thus the sum of the kinetic and the potential energies is

$$E = \frac{1}{2}mv^2 - \frac{GMm}{r} = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\theta}^2) - \frac{GMm}{r}.$$
 (22)

This expression becomes simple when the planet is at the perihelion (r = a - c) or at the aphelion (r = a + c), because there  $\dot{r} = 0$ . Let us consider the perihelion case below. Using  $A = r^2 \dot{\theta}$  (12) and  $A^2 = b^2 GM/a$  (18), we find

$$E = \frac{mA^2}{2r^2} - \frac{GMm}{r} = \frac{GMm}{2r^2a}(b^2 - 2ra).$$
 (23)

Now using the geometry of the ellipse,

$$b^{2} - 2ra = a^{2} - c^{2} - 2(a - c)a$$
  
=  $-(a - c)^{2}$   
=  $-r^{2}$ . (24)

Hence

$$E = -\frac{GMm}{2a}.$$
(25)