

Adiabatic radiation reaction

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Collaborators

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Hughes, Drasco, Flanagan, Franklin, PRL **94**, 221101 (2005)

Outline

- Adiabatic?
- Generic orbits
- Snapshots
- Future

Adiabatic?

Types of waveforms

- **Kludge** (available)

- Gair et al.
- explore data analysis & event rates

- **Adiabatic** (our goal)

- Inspiral is sequence of geodesics
- detection waveforms

- **Self-force** (Capra)

- measurement waveforms

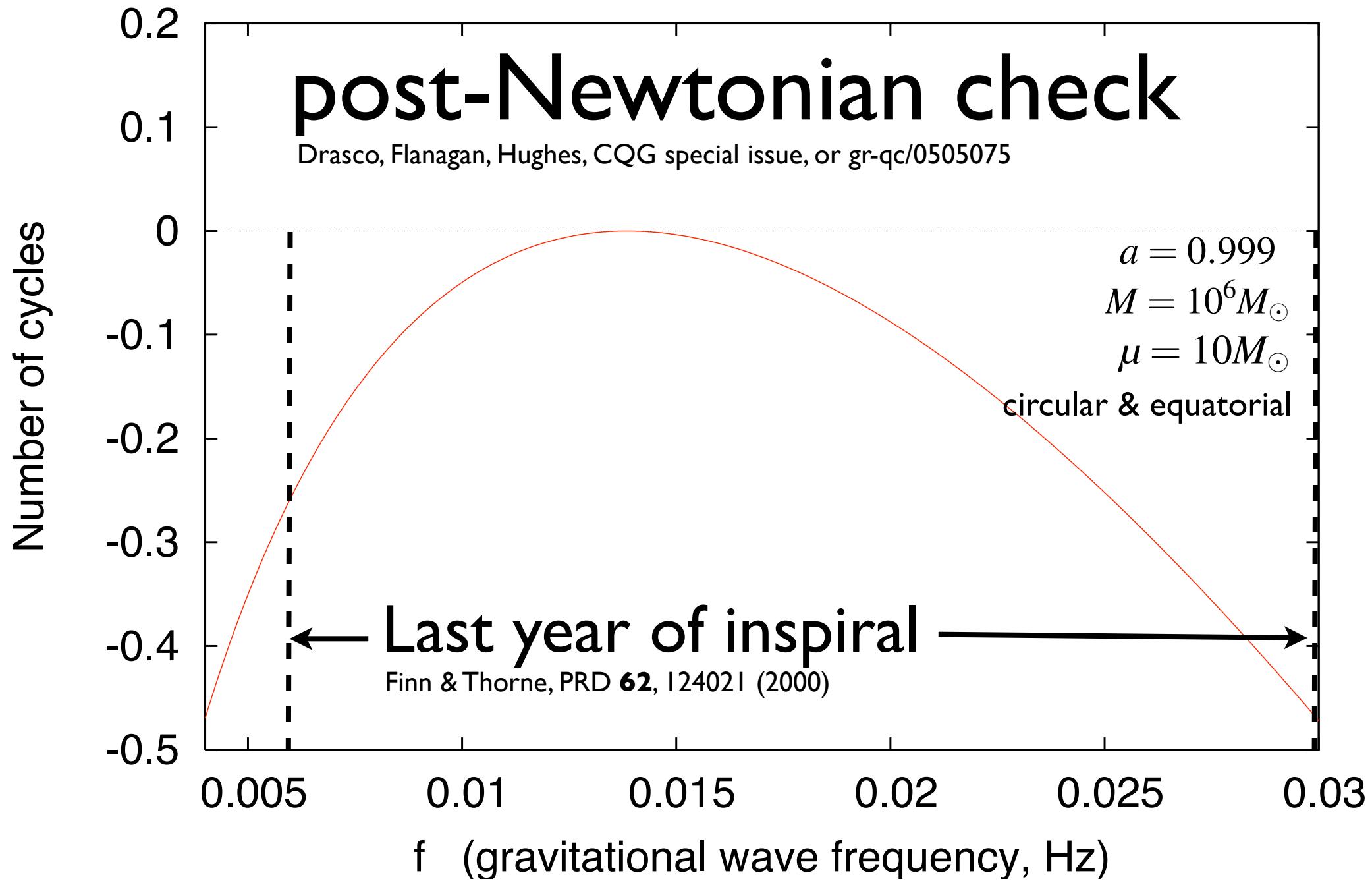
Accuracy

$$\vec{a} = \frac{\mu}{M} \left[\vec{a}_{1,\text{diss}} + \vec{a}_{1,\text{cons}} + \frac{\mu}{M} (\vec{a}_{2,\text{diss}} + \vec{a}_{2,\text{cons}}) + O\left(\frac{\mu^2}{M^2}\right) \right]$$

$$\phi(t) = \frac{M}{\mu} \left[\phi_1(t, t_1) + \frac{\mu}{M} \phi_2(t, t_1) + O\left(\frac{\mu^2}{M^2}\right) \right]$$

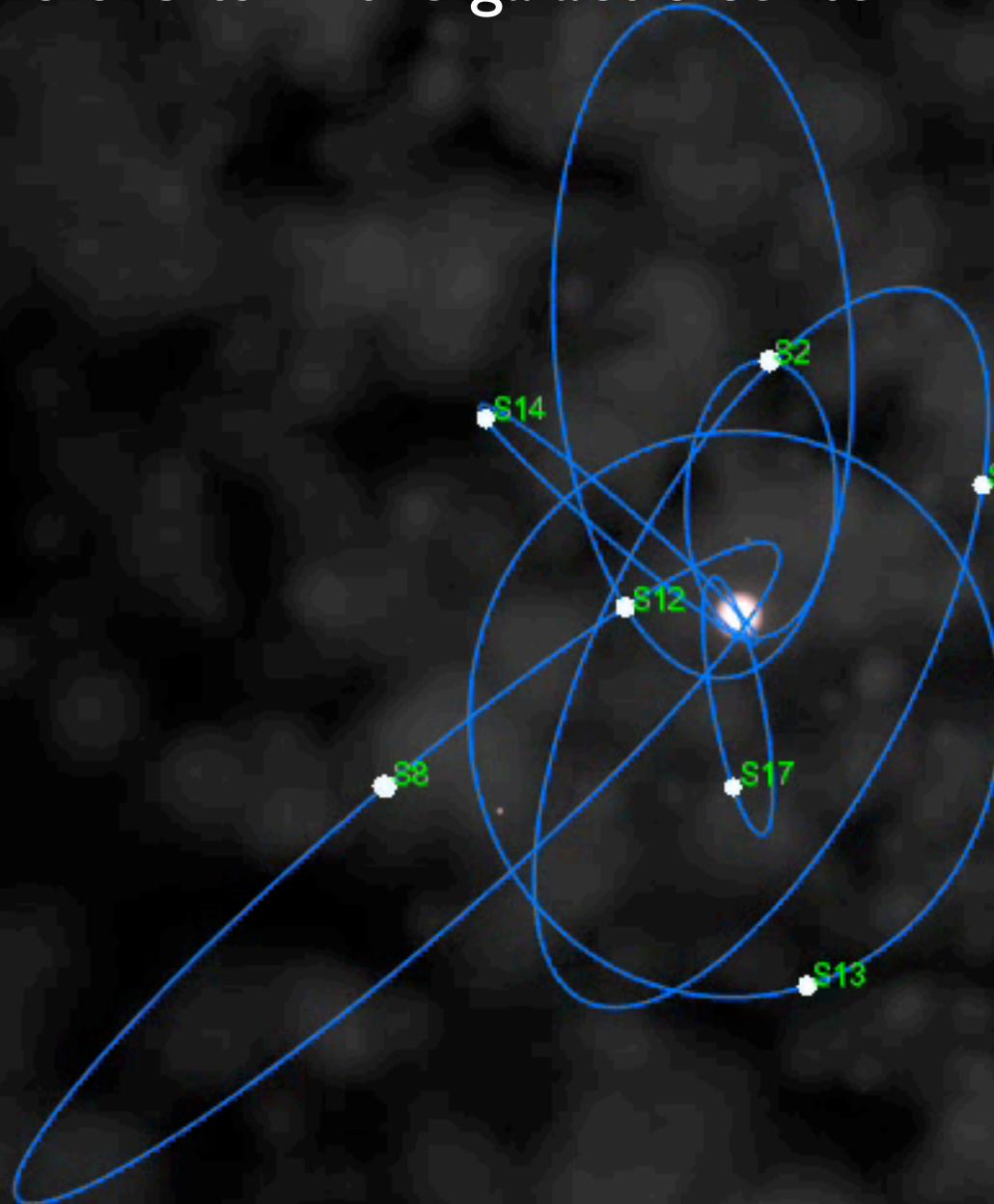
Dissipative effects accumulate secularly.
Conservative effects oscillate, and are **reduced in order**

Phase error from adiabatic waveforms



Generic orbits

Generic orbits in the galactic center



|-10 light days|

Speed: 0.000 m/s

Follow GC

FOV: 13° 59' 60.0" (1.00×)

Mino time

Mino, PRD **67**, 084027 (2003)

$$\frac{d\tau}{d\lambda} = r^2 + a^2 \cos^2 \theta = \Sigma$$

- Decouples radial and polar geodesic equations

$$\left(\frac{dr}{d\lambda} \right)^2 = V_r(r) \quad \left(\frac{d\theta}{d\lambda} \right)^2 = V_\theta(\theta)$$

$$\frac{dt}{d\lambda} = V_t(r, \theta) \quad \frac{d\phi}{d\lambda} = V_\phi(r, \theta)$$

Harmonic Structure

Schmidt, CQG **19**, 2743 (2002)

Drasco & Hughes, PRD **69**, 044015 (2004)

- Radial and polar motion are periodic

$$r(\lambda) = \sum_{n=-\infty}^{\infty} r_n e^{-in\Upsilon_r \lambda} \quad \theta(\lambda) = \sum_{k=-\infty}^{\infty} \theta_k e^{-ik\Upsilon_\theta \lambda}$$

- Axial and time motion are linear & doubly-periodic

$$t(\lambda) = t_0 + \Gamma \lambda + \Delta t(r, \theta) \quad \phi(\lambda) = \phi_0 + \Upsilon_\phi \lambda + \Delta \phi(r, \theta)$$

$$\Delta t = \sum_{k=1}^{\infty} \Delta t_k^\theta \sin(k\Upsilon_\theta \lambda) \quad \Delta \phi = \sum_{k=1}^{\infty} \Delta \phi_k^\theta \sin(k\Upsilon_\theta \lambda)$$

$$+ \sum_{n=1}^{\infty} \Delta t_n^r \sin(n\Upsilon_r \lambda) \quad + \sum_{n=1}^{\infty} \Delta \phi_n^r \sin(n\Upsilon_r \lambda)$$

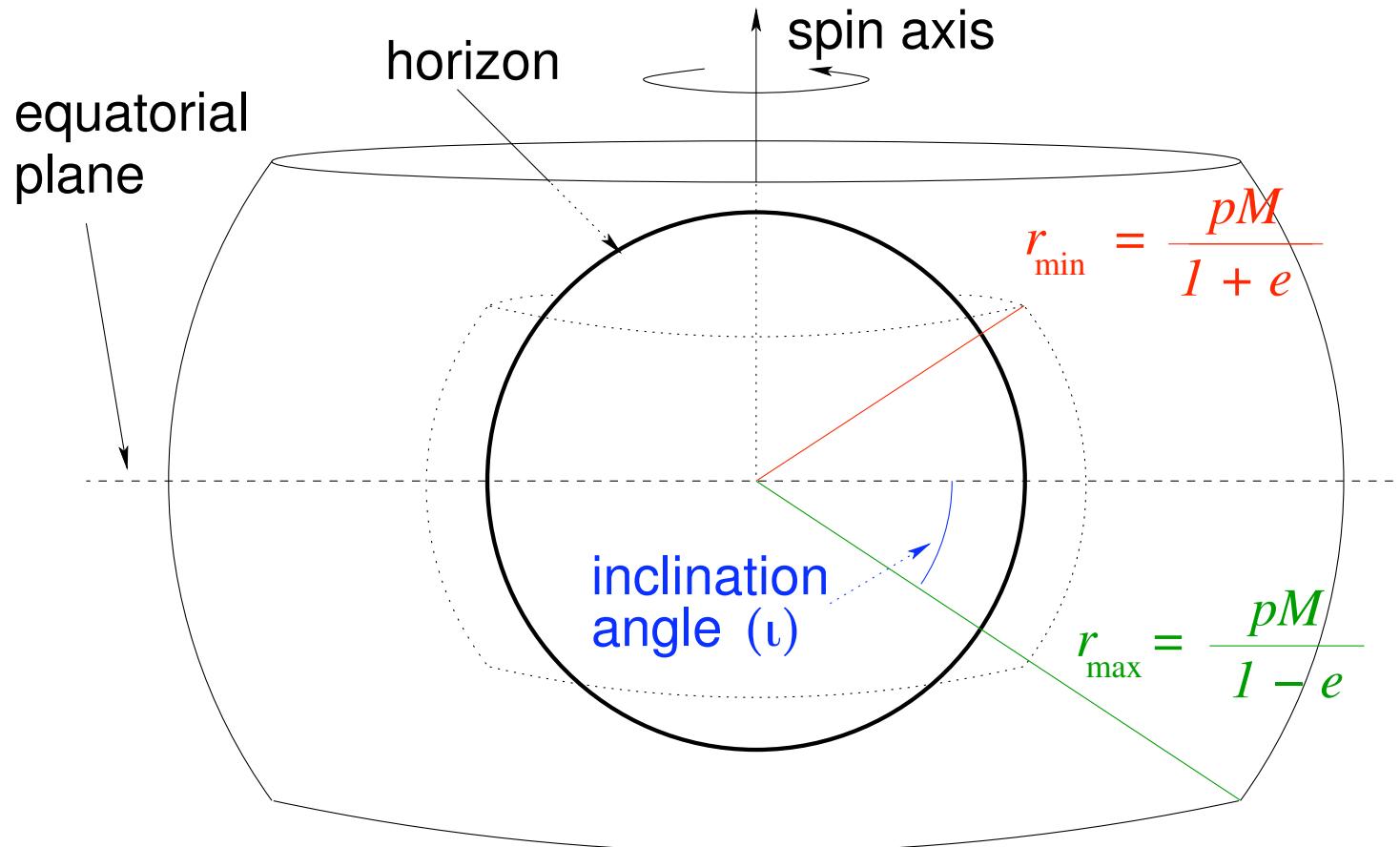
Coordinate frequencies

- Waveforms oscillate at integer-linear combinations of coordinate-time frequencies

$$\omega_r = \frac{\Upsilon_r}{\Gamma} \quad \omega_\theta = \frac{\Upsilon_\theta}{\Gamma} \quad \omega_\phi = \frac{\Upsilon_\phi}{\Gamma}$$

$$\omega_{mkn} = m\omega_\phi + k\omega_\theta + n\omega_r$$

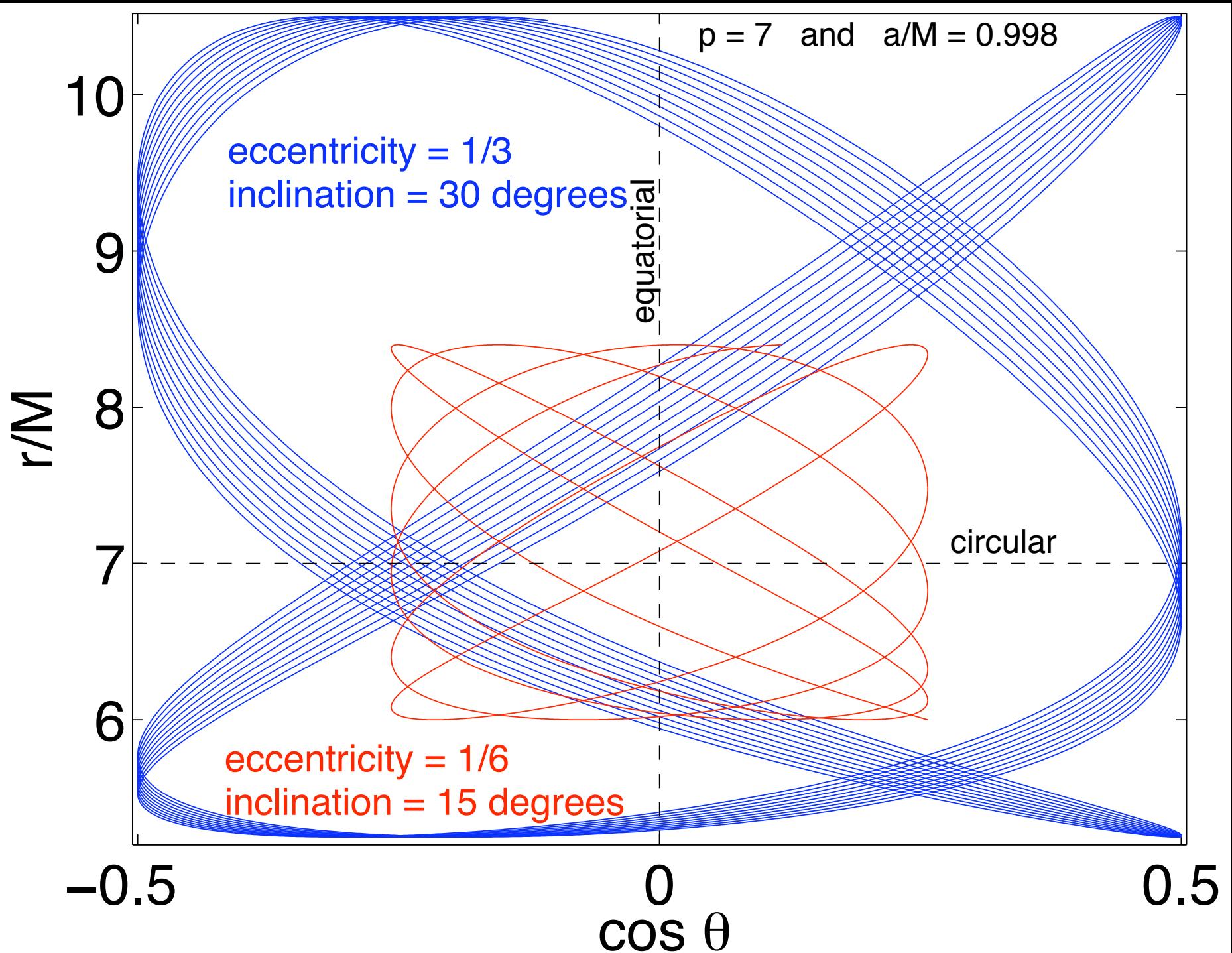
The orbital torus



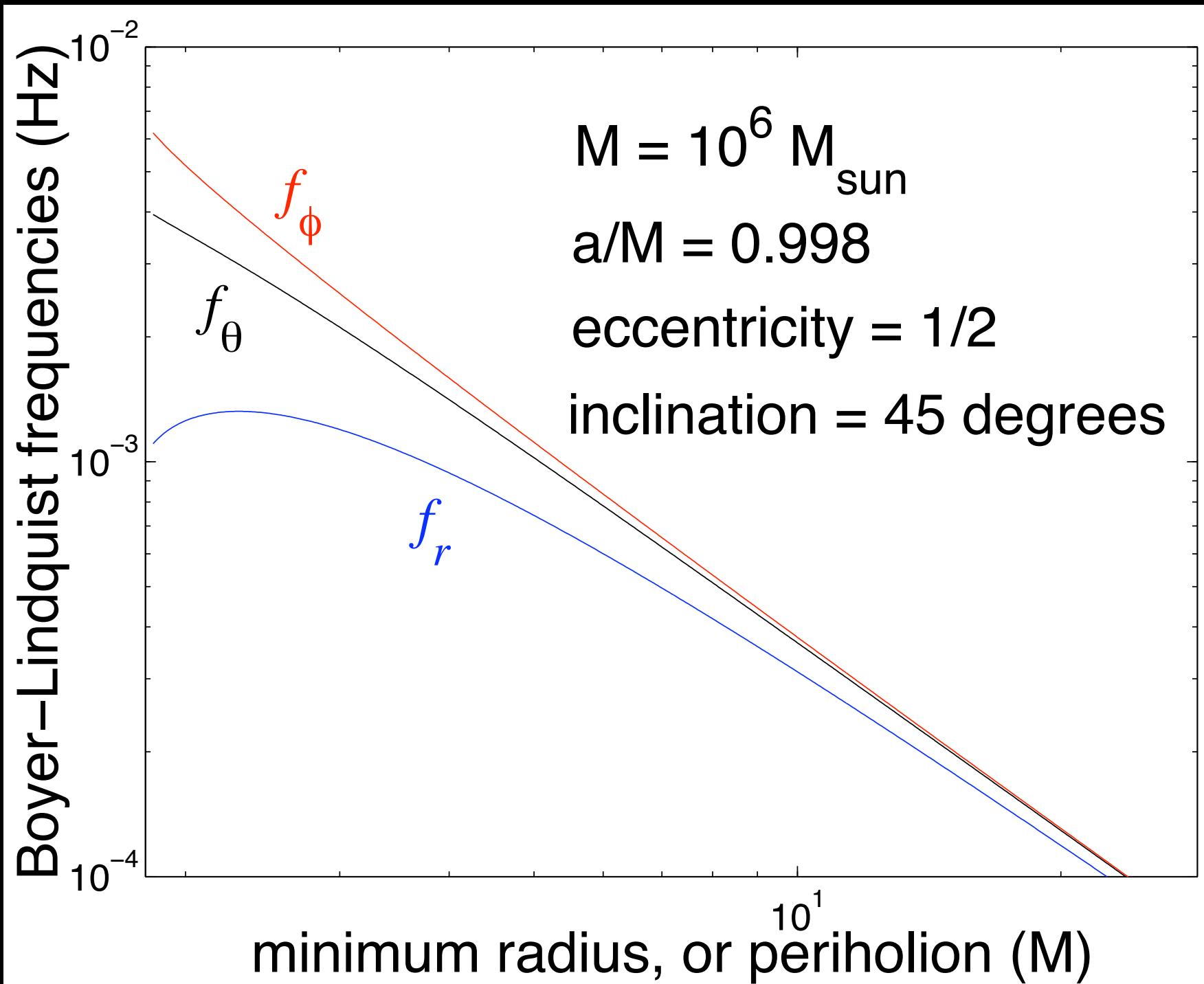
Orbit “shape” determined by three constants:

$(r_{\min}, r_{\max}, \iota)$ or (e, p, ι) or (E, L_z, Q) or...

Comparing a pair of generic orbits



Frequencies of generic orbits



Snapshots

Teukolsky-Sasaki-Nakamura formalism

- Source of perturbation: generic orbiting test mass
- Solution: discrete set of mode amplitudes

$$Z_{lmkn}^{\text{H},\infty} = \frac{2\pi}{\Gamma} \left\langle J^{\text{H},\infty}(r, \theta, \omega_{mkn}) \right\rangle$$

$$\langle f(r, \theta) \rangle = \frac{1}{\Lambda_\theta \Lambda_r} \int_0^{\Lambda_\theta} d\lambda_\theta \int_0^{\Lambda_r} d\lambda_r f[r(\lambda_r), \theta(\lambda_\theta)] e^{ik\Upsilon_\theta \lambda_\theta + in\Upsilon_r \lambda_r}$$

$$\Lambda_{r,\theta} = \frac{2\pi}{\Upsilon_{r,\theta}}$$

Waveforms and fluxes

$$h_+ - i h_\times = -\frac{2}{r} \sum_{lmkn} \frac{Z_{lmkn}^H}{\underline{\omega_{mkn}^2}} S_{lmkn}(\theta) e^{-i\omega_{mkn}(t-r)+im\phi}$$

$$\left\langle \frac{dE}{dt} \right\rangle = \sum_{lmkn} \frac{1}{4\pi \underline{\omega_{mkn}^2}} \left(|Z_{lmkn}^H|^2 + \alpha_{lmkn} |Z_{lmkn}^\infty|^2 \right)$$

$$\left\langle \frac{dL_z}{dt} \right\rangle = \sum_{lmkn} \frac{m}{4\pi \underline{\omega_{mkn}^3}} \left(|Z_{lmkn}^H|^2 + \alpha_{lmkn} |Z_{lmkn}^\infty|^2 \right)$$

Where α_{lmkn} is a simple quantity.

“Flux” of Carter

- Sago, Tanaka, Hikida, Nakano, to appear in PTP, or gr-qc/0506092

$$\left\langle \frac{dQ}{dt} \right\rangle^\infty = (\text{constants}) + 2 \sum_{lmkn} n \omega_r \left| \underline{\underline{Z}}_{lmkn}^H \right|^2$$

due to gravitational perturbations.

- Drasco, Flanagan, Hughes, CQG special issue, or gr-qc/0505075

$$\left\langle \frac{dQ}{dt} \right\rangle = A \sum_{lmkn} \frac{\omega_{lmkn}}{\left| \underline{\omega}_{lmkn} \right|} \left[(\underline{\underline{\tilde{Z}}}_{lmkn}^H)^* \underline{\underline{Z}}_{lmkn}^H + \alpha_{lmkn} (\underline{\underline{\tilde{Z}}}_{lmkn}^\infty)^* \underline{\underline{Z}}_{lmkn}^\infty \right]$$

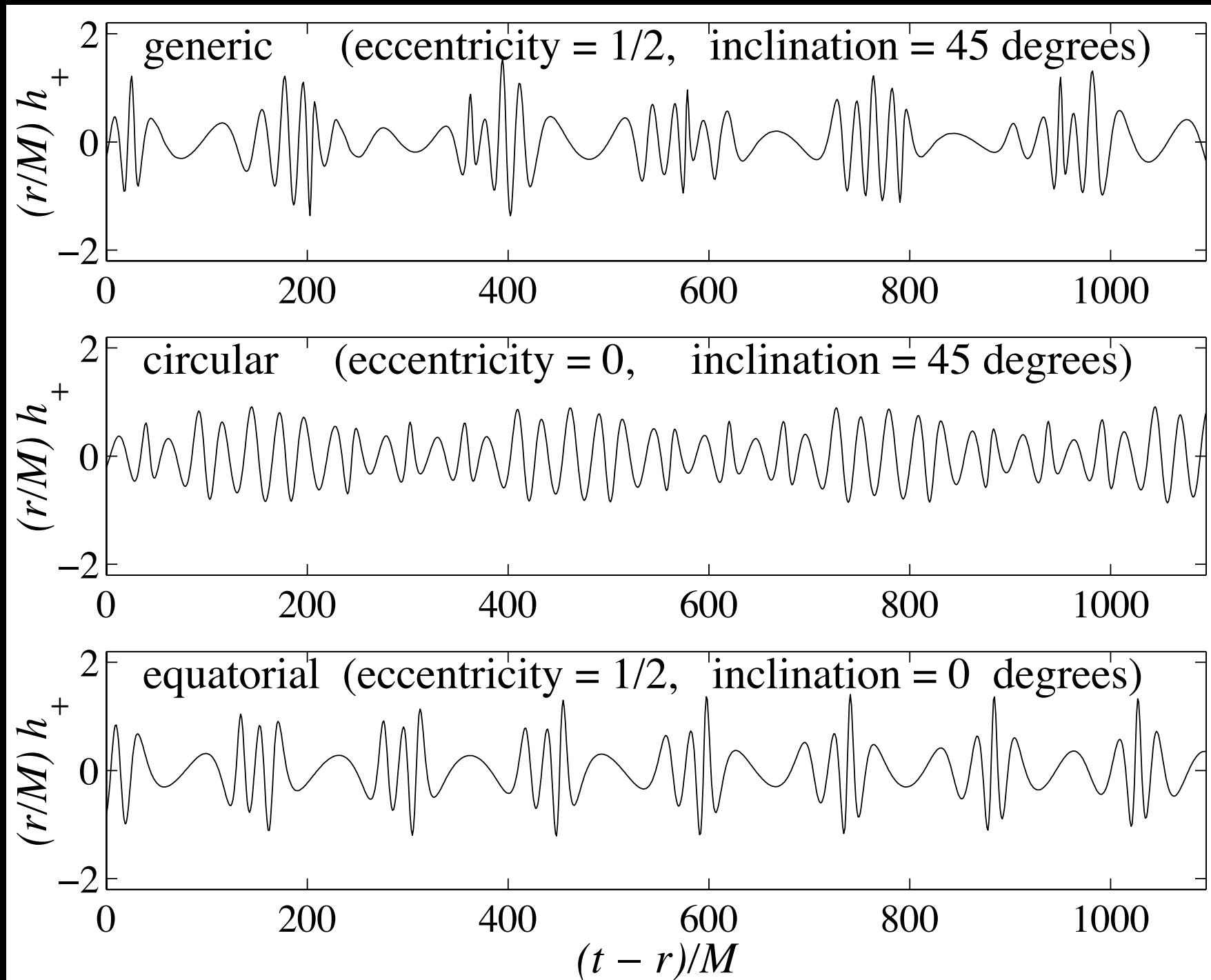
due to scalar perturbations.

Here $\underline{\underline{\tilde{Z}}}_{lmkn}^{H,\infty}$ is as difficult to calculate as $\underline{\underline{Z}}_{lmkn}^{H,\infty}$.

History (frequency domain)

Ref.	$a \neq 0$	$e > 0$	$\iota \neq 0$	evolve
Cutler et al. (1994)		✓	(NA)	✓
Finn & Thorne (2000)	✓			✓
Shibata (1993)	✓		✓	
Shibata (1994)	✓	✓		
Hughes (2001)	✓		✓	✓
Glampedakis & Kennefick (2002)	✓	✓		✓
Drasco & Hughes (2005)	✓	✓	✓	

Comparing waveforms: $a = 0.9M$, $p = 4$



Radiation and orbit geometry

$$H = h_+ - ih_\times = \sum_{kn} H_{kn} e^{-i\omega_{kn}(t-r)}$$
$$\omega_{kn} = k\omega_\theta + n\omega_r$$

$H_{kn} \propto \delta_{n0}$ (circular orbits)

$H_{kn} \propto \delta_{0k}$ (equatorial orbits)

$H_{kn} \propto \delta_{0k}\delta_{n0}$ (circular-equatorial orbits)

Radial and polar voices

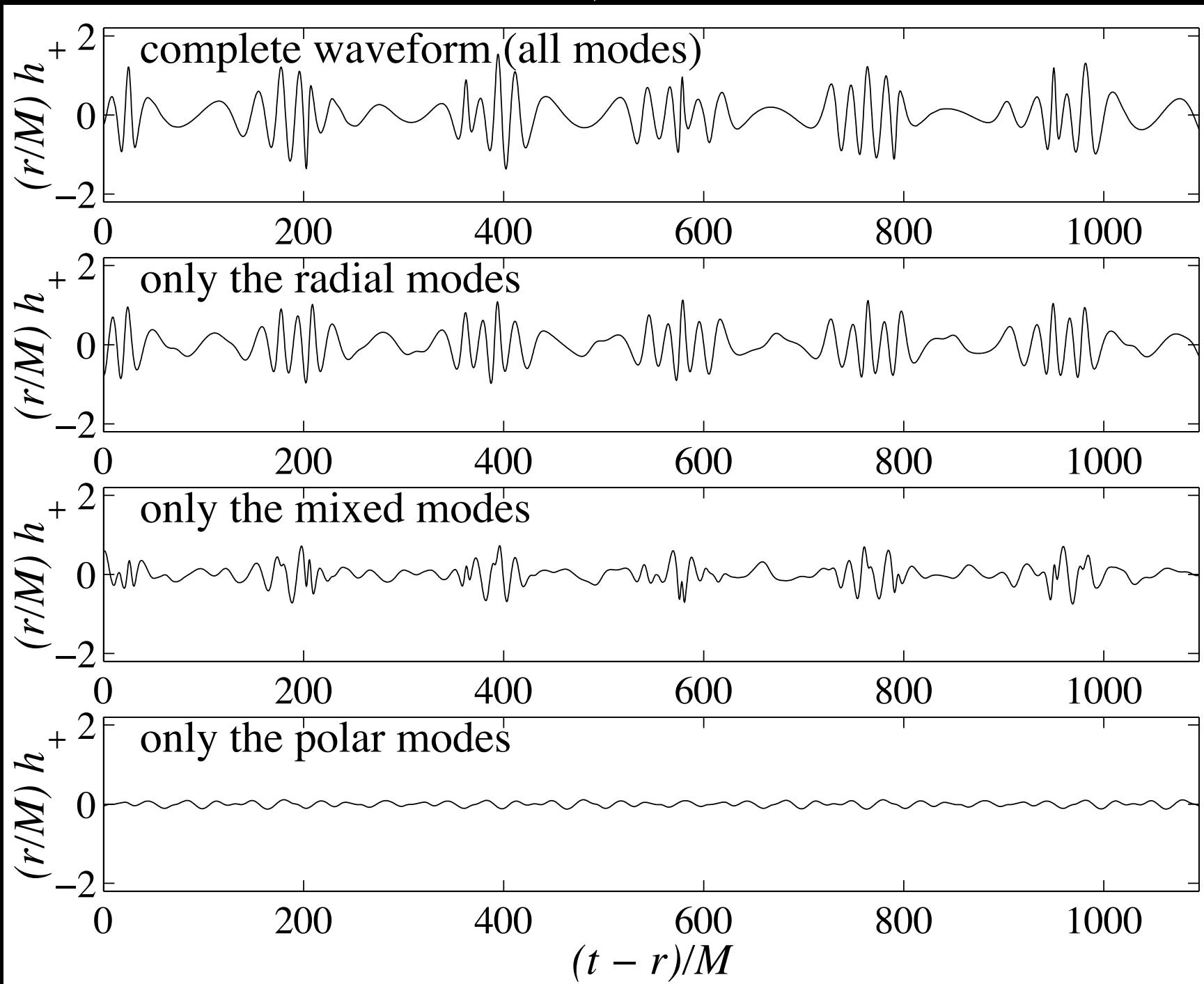
$$H = H_{\text{polar}} + H_{\text{radial}} + H_{\text{mixed}},$$

$$H_{\text{polar}} = \sum_{k \neq 0} H_{k0} e^{-i\omega_{k0}t}$$

$$H_{\text{radial}} = \sum_{n \neq 0} H_{0n} e^{-i\omega_{0n}t}$$

$$H_{\text{mixed}} = H_{00} + \sum_{k \neq 0} \sum_{n \neq 0} H_{kn} e^{-i\omega_{kn}t}$$

$$e = 0.5, \iota = 45^\circ$$



Radial and polar voices

$$e = 0.5, \ i = 45^\circ$$

	polar/total	radial/total	mixed/total
$\langle dE/dt \rangle$	1.2%	62%	37%
$\langle dL_z/dt \rangle$	0.77%	73%	26%

Increasing eccentricity excites radial harmonics

Increasing inclination excites polar harmonics (less so)

$$a = 0.9M, p = 6, \theta_{\text{inc}} = 1$$

e	θ_{inc}	$\frac{\langle dE/dt \rangle_{\text{radial}}}{\langle dE/dt \rangle}$	$\frac{\langle dE/dt \rangle_{\text{polar}}}{\langle dE/dt \rangle}$	$\frac{\langle dL_z/dt \rangle_{\text{radial}}}{\langle dL_z/dt \rangle}$	$\frac{\langle dL_z/dt \rangle_{\text{polar}}}{\langle dL_z/dt \rangle}$
0.1	20°	16%	8.3%	14%	5.0%
	40°	13%	28%	14%	17%
	60°	9.4%	48%	14%	29%
	80°	6.7%	53%	20%	24%
0.3	20°	76%	1.9%	75%	1.3%
	40°	58%	5.4%	67%	3.6%
	60°	37%	5.9%	58%	2.7%
	80°	17%	1.9%	51%	-2.4%
0.5	20°	90%	0.077%	93%	0.083%
	40°	66%	0.24%	77%	0.18%
	60°	39%	0.93%	58%	0.68%
	80°	19%	1.6%	53%	-2.0%
0.7	20°	90%	0.057%	92%	0.10%
	40°	67%	0.22%	77%	0.35%
	60°	40%	0.19%	61%	-0.17%
	80°	19%	0.16%	53%	-0.47%

Radial **Polar**

Summed-spectra

$$\left\langle \frac{dX}{dt} \right\rangle_l^{H,\infty} = \sum_{mkn} \left\langle \frac{dX}{dt} \right\rangle_{lmkn}^{H,\infty} \quad X = E, L_z$$

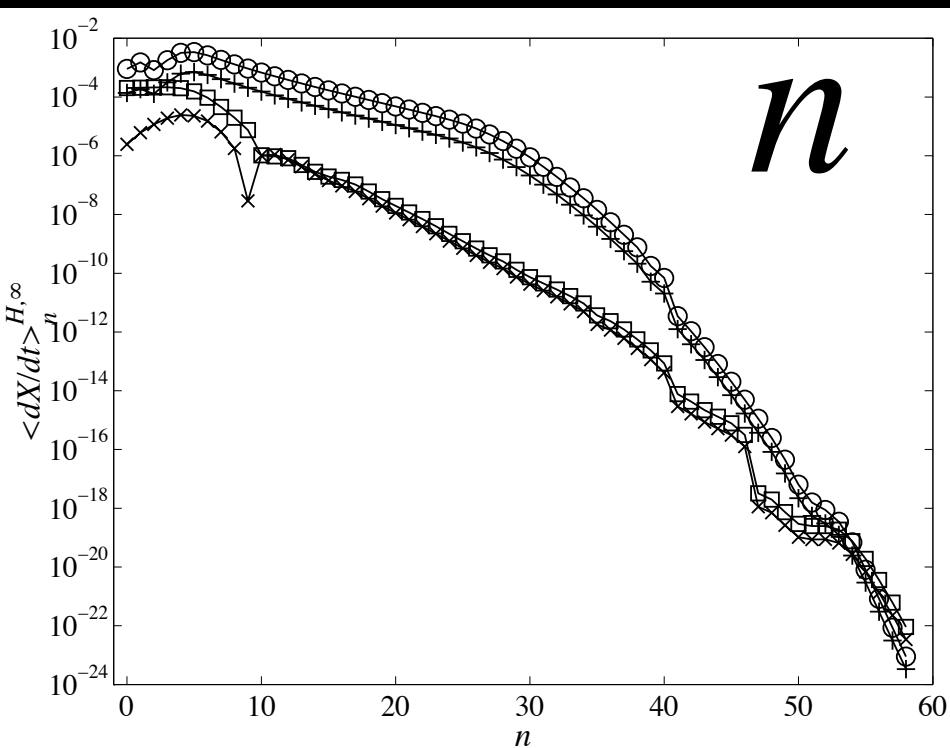
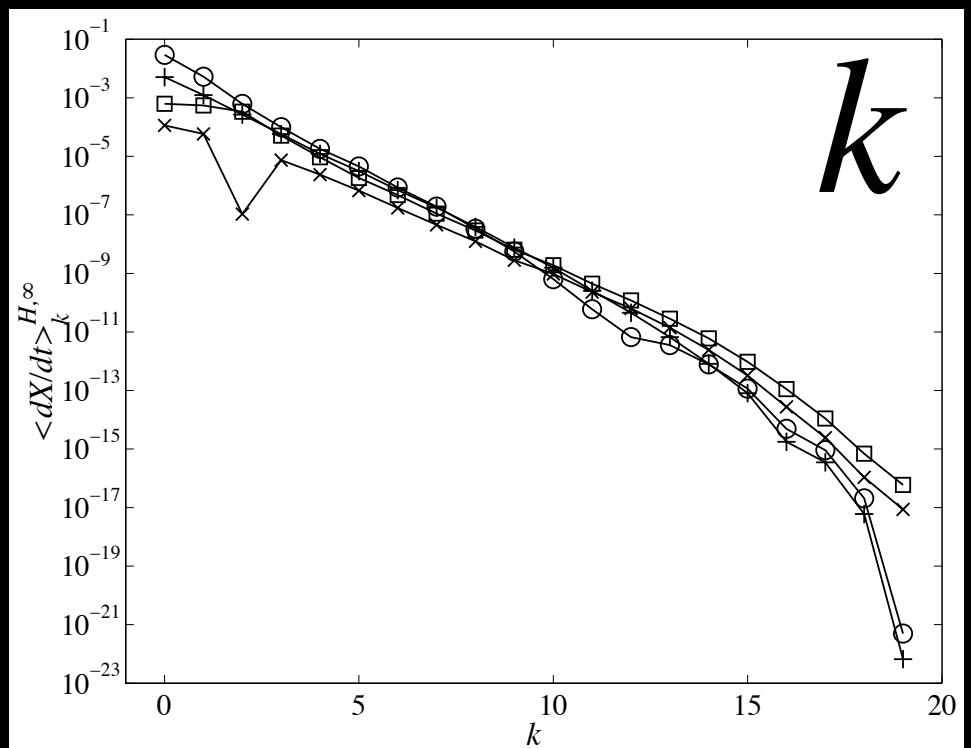
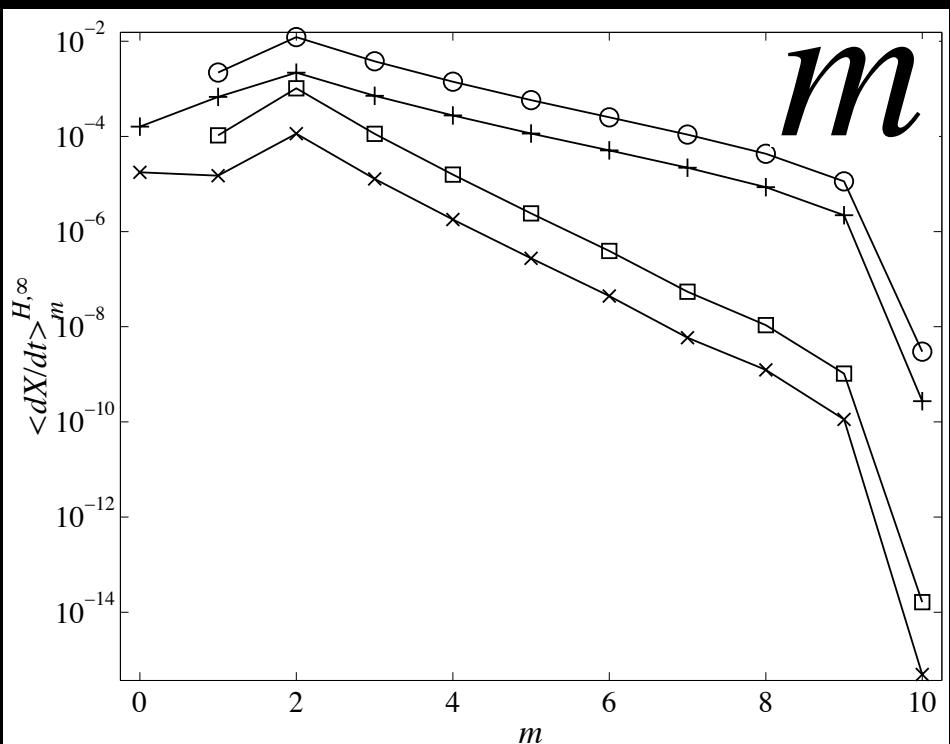
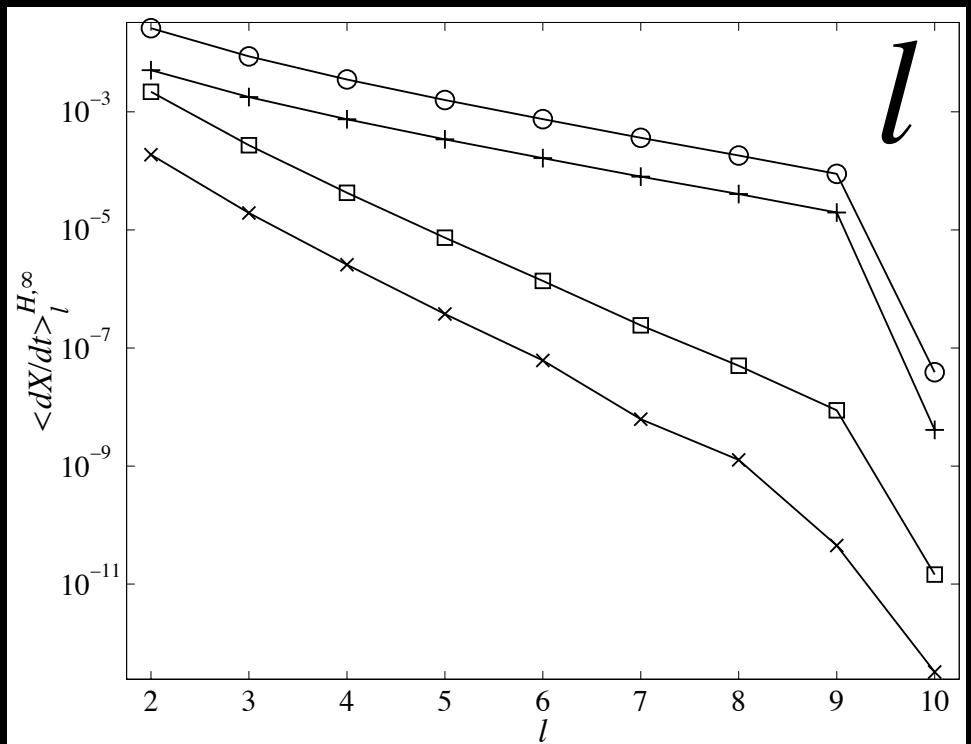
$$\left\langle \frac{dX}{dt} \right\rangle_m^{H,\infty} = \sum_{lkn} \left\langle \frac{dX}{dt} \right\rangle_{lmkn}^{H,\infty}$$

$$\left\langle \frac{dX}{dt} \right\rangle_n^{H,\infty} = \sum_{lmk} \left\langle \frac{dX}{dt} \right\rangle_{lmkn}^{H,\infty}$$

$$\left\langle \frac{dX}{dt} \right\rangle_k^{H,\infty} = \sum_{lmn} \left\langle \frac{dX}{dt} \right\rangle_{lmkn}^{H,\infty}$$

$$\left\langle \frac{dX}{dt} \right\rangle_v^{H,\infty} \propto \exp(-|v|\gamma_v^{H,\infty})$$

$$v = l, m, k, n$$



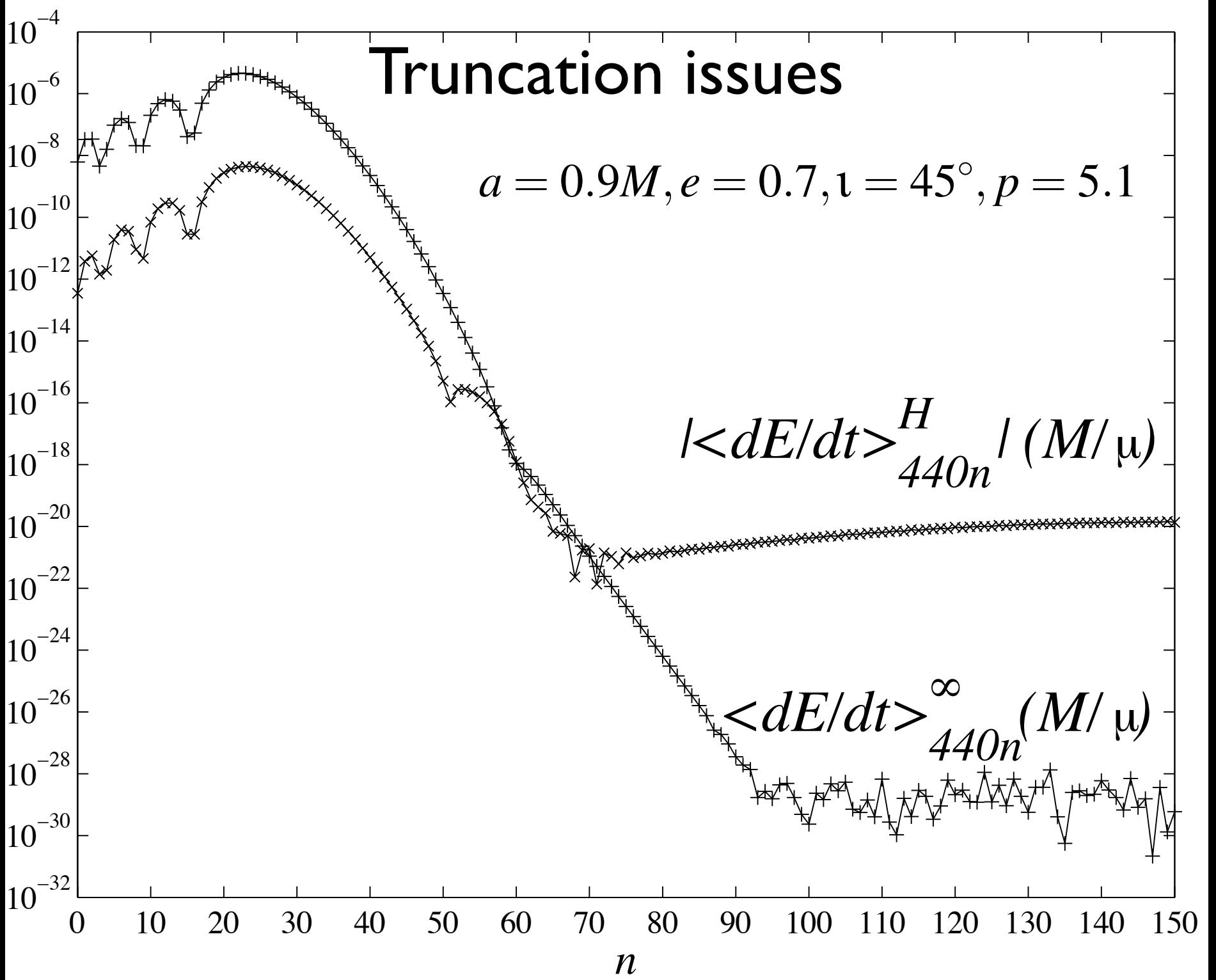
$$a = 0.9M, p = 6, \theta_{\text{inc}} = 1$$

e	θ_{inc}	γ_l^H	γ_l^∞	γ_m^H	γ_m^∞	γ_k^H	γ_k^∞	γ_n^H	γ_n^∞
0.1	20°	3.2	1.4	2.4	1.0	4.6	4.2	5.8	4.1
0.1	40°	3.2	1.4	2.4	1.3	2.9	2.6	4.9	4.4
0.1	60°	3.1	1.4	2.4	1.6	2.2	2.0	1.6	2.0
0.1	80°	2.7	1.3	2.7	2.0	1.7	1.6	2.6	2.5
0.3	20°	2.9	1.2	2.3	0.92	4.2	3.7	2.6	2.2
0.3	40°	2.9	1.2	2.2	1.3	2.5	2.4	1.9	1.6
0.3	60°	2.7	1.2	2.3	1.4	1.8	1.8	1.5	1.6
0.3	80°	2.0	1.2	2.3	1.8	1.4	1.4	1.3	1.0
0.5	20°	2.7	1.0	2.2	0.81	4.2	3.9	1.1	1.0
0.5	40°	2.6	1.1	2.2	1.0	2.2	2.3	0.78	0.74
0.5	60°	2.6	1.1	2.1	1.3	1.5	1.7	0.72	0.64
0.5	80°	1.8	1.1	1.9	1.7	1.3	1.4	0.60	0.49
0.7	20°	2.9	1.4	2.3	0.92	3.1	3.4	0.40	0.45
0.7	40°	2.4	1.2	2.1	1.1	2.0	2.2	0.30	0.32
0.7	60°	1.8	1.4	2.1	1.4	1.3	1.5	0.27	0.26
0.7	80°	1.5	1.1	1.8	1.8	1.7	1.6	0.23	0.24

Polar

Radial

Future



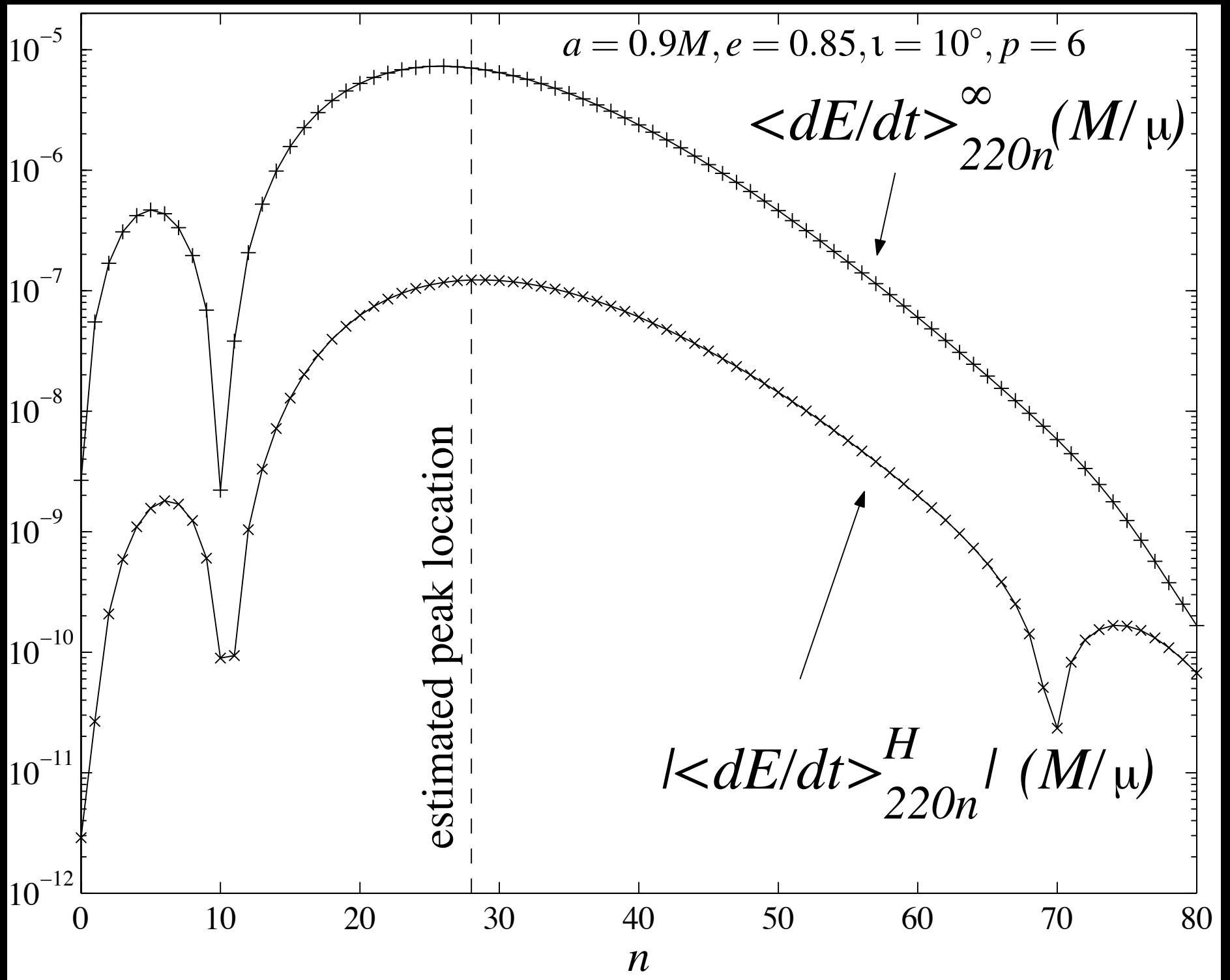
Making good guesses

Gravitational wave spectrum for Keplerian waveforms

Peters & Mathews, Phys. Rev. **131**, 435 (1963)

$$\left\langle \frac{dE}{dt} \right\rangle_n^{\text{PM}} \propto \frac{n^4}{32} \left\{ [J_{n-2}(ne) - 2eJ_{n-1}(ne) + \frac{2}{n}J_n(ne) + 2eJ_{n+1}(ne) - J_{n+2}(ne)]^2 + (1-e^2)[J_{n-2}(ne) - 2J_n(ne) + J_{n+2}(ne)]^2 + \frac{4}{3n^2}[J_n(ne)]^2 \right\}$$

$$n_{\text{peak}} \approx \exp \left[\frac{1}{2} - \frac{3}{2} \ln(1-e) \right]$$



Summary

- Generic snapshots are now available $e < 0.7$
- Resolve & implement Carter evolution
- Evolution will give generic adiabatic inspirals
- Optimization
- Evolution of initial conditions?
- immediate applications to astrophysics?