Adaptive Mesh Refinement for Self-Force Calculations

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in collaboration with

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As discussed by Leor Barack & Norichika Sago in their talks, we (the Southampton group) do self-force calculations by finding the **metric perturbation** in the **Lorenz gauge**: [details in Barack & Sago, *Phys. Rev. D* **75**, 064021 (2007)]

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- decompose the metric perturbation into (ℓ, m) multipole modes
- for each (ℓ, m) multipole mode, numerically integrate the metric-perturbation equations in the time domain
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Most of the material in this talk would be equally applicable to other self-force calculation schemes, or indeed to other problems involving the numerical solution of time-evolution PDEs.

After the (ℓ, m) multipole decomposition, our typical metric perturbation equation (to be numerically integrated) looks like

where $V_{\ell}(r)$, $S_{\ell m}(t)$, and $\vec{x}_{\text{particle}}(t)$ are known, and ϕ is a complex field on a Schwarzschild or Kerr background (1 or 2 space dimensions × time)

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 \Rightarrow Want mesh refinement!



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- relatively hard to implement

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- coarsest grid covers entire problem domain



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- \bullet coarsest grid covers entire problem domain
- refine in both space and time
- fine grids **overlay** coarser grids
- fine-grid initial & boundary data is **interpolated** from coarse grids
- integrate each grid independently
- inject fine results back into coarse grid when/where points coincide (keeps coarse grid accurate)



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Some people use generic toolkits:

- AMRD/PAMR (Choptuik/Pretorius)
- DAGH/GRACE (Parashar)
- PARAMESH (MacNeice et al., NASA/Goddard 2BH project)
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Another alternative is to write a Berger-Oliger code from scratch. \Rightarrow takes ~ 5K-10K lines of code (including test drivers, comments, etc.)

Berger-Oliger in Cauchy (Numerical) Relativity

Choptuik pioneered the use of Berger-Oliger mesh refinement in numerical relativity in his discovery of self-similarity and critical phenomena in gravitational collapse. [*Phys. Rev. Lett.* **70**, 9 (1993)]

Berger-Oliger methods are now widely used in 3 + 1 numerical relativity. Some good references include:

- Berger & Oliger, J. Comp. Phys. 53, 484 (1984)
- Berger, SIAM J. Sci. Stat. Comput. 7, 904 (1986)
- Choptuik, "Experiences with an Adaptive Mesh Refinement Algorithm in Numerical Relativity", pages 206–221 in Evans, Finn, and Hobill, *Frontiers in Numerical Relativity*, Cambridge U.P., 1989
- Schnetter, Hawley, and Hawke, *Class. Quant. Grav.* **21**, 1465 (2004) [Nicely explains the complications which arise when equations] contain 1st time derivatives but <u>2nd</u> spatial derivatives.

Characteristic Coordinates

Motivation: null boundaries are <u>much</u> nicer than timelike boundaries (both analytically and numerically) $t_{a,i} \neq t_{a,i}$

Null coordinates: $u = t_{\text{Schw}} - r_*$

$$v = t_{\text{Schw}} + r_*$$



Characteristic Coordinates

Motivation: null boundaries are <u>much</u> nicer than timelike boundaries (both analytically and numerically) Null coordinates: $u = t_{\text{Schw}} - r_*$ $v = t_{\text{Schw}} + r_*$ Fundamental discretization uses double-null "diamond" cells

south

Characteristic Coordinates



Self-Force Calculations (Characteristic Coordinates)

As discussed by Leor Barack & Norichika Sago in their talks, for self-force calculations we integrate the (discretized) metric perturbation equations for each (ℓ, m) multipole mode, in a square "grid box" in (u, v) space, chosen

to be big enough for the initial-data field perturbations to have decayed below our numerical error levels by the end of the integration.

The integration proceeds one v = constant slice at a time;each slice is integrated one diamond cell at a time.



Basically, use standard Berger-Oliger mesh refinement, treating u as a "spatial" coordinate on v = constant slices, and v as a "time" coordinate



Cell-recursive algorithm:

 $\llbracket Hamadé \& Stewart, Class. Quant. Grav. 13, 497 (1996) \rrbracket$

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 - divide cell into 4 subcells



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- integrate a cell
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 - divide cell into 4 subcells
 - recursively integrate south subcell



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Problem: fine-grained (per-cell) memory management

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- if there are flagged points, recurse:
 - divide slice into 2 subslices
 - recursively integrate smaller-v ("lower") subslice
 - recursively integrate larger-v ("upper") subslice
 - inject upper-slice results back into matching points of original (coarse) slice
 - reintegrate the remainder of the original (coarse) slice starting from the upper-most injected data

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- \Rightarrow Relatively simple bookkeeping & memory management

Error Estimation for <u>Adaptive</u> Mesh Refinement

For **adaptive** mesh refinement (AMR), we need to estimate (during the integration) when the integration is accurate enough, and when it's not. To do this, we use an estimate of the local finite differencing error, also known as the local truncation error (LTE): compare

(a) the result of integrating a standard diamond cell



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- (a) the result of integrating a standard diamond cell
- (b) the result of integrating a double-sized diamond cell using data from the 2nd-to-last and current slices

The difference (b) - (a) is an estimate of the LTE, up to some $\mathcal{O}(1)$ factor; the AMR refinement criterion is simply **if** (|(b) - (a)| > threshold)

then this cell must be redone at higher resolution



Test Cases for Sample Results

Physics:

- Schwarzschild background, scalar particle in circular orbit
- zero initial data on $v = v_{\min}$ and $u = u_{\min}$ grid-box faces

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Numerical Methods:

- 2nd order finite differencing (global accuracy)
- slice-recursive AMR algorithm
- AMR gradually turned starting at 100m (we don't bother resolving the junk radiation at the start of the evolution); AMR fully active by $\approx 200m$

Sample Results (Convergence Tests)

To test convergence, the code writes out a "script" describing the mesh-refinement structure, which can then be "played back" at successively higher resolutions.

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Test case: $(\ell, m) = (10, 10)$, particle at $r_* = 10m$ $(r_{\text{Schw}} = 7.85m)$ script "recorded" at: AMR error threshold 10^{-6} in $|\phi|$ (coarsest grid resolution 0.5m)

The code shows excellent 2nd order convergence, even across the mesh-refinement boundaries and near the particle worldline, (where $\partial_r \phi$ has a jump discontinuity)



Sample Results (ϕ movie)

Test case: same as before, but error threshold 10^{-7} (coarsest grid resolution 0.1m)

Sample frame from **movie** (see also **poster** outside)



phi(u) on relative v=155.0 slice

Sample Results (Self-Force)

Test case:

- particle at $r_{\text{Schw}} = 10m$
- modes computed numerically for $\ell \leq 20$ (121 modes given even/odd symmetry)
- F_{ℓ} for $\ell > 20$ tail estimated via 3-term fit to $\{\ell^{-2}, \ell^{-4}, \ell^{-6}, \dots\}$ series [Detweiler, Messaritaki, & Whiting, *Phys. Rev. D* 67, 104016 (2003)]
- grid box 300m on a side (too small)
- AMR error threshold 10^{-6} in $|\phi|$, Richardson extrap. playback×{6,8}
- error estimate via $|F_{\text{inside}}^r F_{\text{outside}}^r|$

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- $\Rightarrow \text{ this work } F^r = 1.3808 \times 10^{-5} \text{ (43\% numerical, 57\% tail)}$ Detweiler *et al.* $F^r = 1.3784 \times 10^{-5}$ fractional error 0.17%

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Self-Force:

- nice results for Schwarzschild background, scalar particle in circular orbit
- next steps: 4th order, eccentric orbits, Kerr

Self-Force Results $(F_{\ell}(\ell))$



Self-Force Results (scaled $F_{\ell}(\ell)$)



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- recursively integrate **entire** fine slice (A1)
- rotate fine slices
- copy coarse slice (A1) value to get \checkmark starting (u_{\min}) value for fine slice (A1)

Α1

B0

C0

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