

Periastron Advance in Black Hole Binaries

Alexandre Le Tiec

University of Maryland

Collaborators:

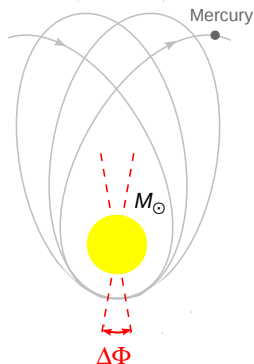
L. Barack, A. Buonanno, A. H. Mroué,
H. P. Pfeiffer, N. Sago & A. Taracchini

Relativistic perihelion advance of Mercury

- Observed anomalous precession of Mercury's perihelion of $\sim 43''$ /century
- Accounted for by the leading order relativistic angular advance per orbit

$$\Delta\Phi_{\text{GR}} = \frac{6\pi GM_{\odot}}{c^2 a (1 - e^2)}$$

- One of the first **successes** of Einstein's theory of general relativity
- Relativistic periastron advance of $\sim \text{''}/\text{year}$ now measured in **binary pulsars**



Periastron advance in black hole binaries

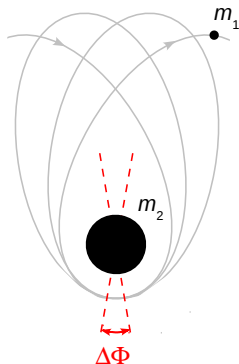
- **Conservative** part of the dynamics only
- Generic non-circular orbit parametrized by the two **invariant** frequencies

$$\Omega_r = \frac{2\pi}{P}, \quad \Omega_\varphi = \frac{1}{P} \int_0^P \dot{\varphi}(t) dt$$

- Periastron advance per radial period

$$K \equiv \frac{\Omega_\varphi}{\Omega_r} = 1 + \frac{\Delta\Phi}{2\pi}$$

- In the **circular** orbit limit $e \rightarrow 0$, the relation $K(\Omega_\varphi)$ is coordinate **invariant**



Post-Newtonian and EOB results

- Third post-Newtonian result [Damour, Jaranowski & Schäfer 2000]

$$K_{3\text{PN}} = 1 + 3x + \left(\frac{27}{2} - 7\nu \right) x^2 + \left(\frac{135}{2} - \left[\frac{649}{4} - \frac{123}{32} \pi^2 \right] \nu + 7\nu^2 \right) x^3 + \mathcal{O}(x^4),$$

where $\nu \equiv m_1 m_2 / m^2$ and $x \equiv (m \Omega_\varphi)^{2/3} \sim v^2$

- Effective-one-body result [Buonanno & Damour 1999; Damour 2010]

$$K_{\text{EOB}} = \sqrt{\frac{AA'B}{AA' + 2u(A')^2 - uAA''}},$$

where A and B are the EOB potentials

Gravitational self-force result(s)

- The conservative SF correction ρ to the Schwarzschild result has recently been computed numerically [Barack & Sago 2010]

$$W \equiv \frac{1}{K^2} = \underbrace{1 - 6x}_{\text{Schw.}} + \underbrace{q \rho(x)}_{\text{SF effect}} + \mathcal{O}(q^2)$$

- Usual SF result expressed in terms of mass ratio $q \equiv m_1/m_2$:

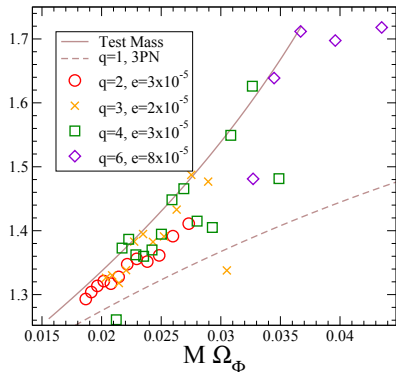
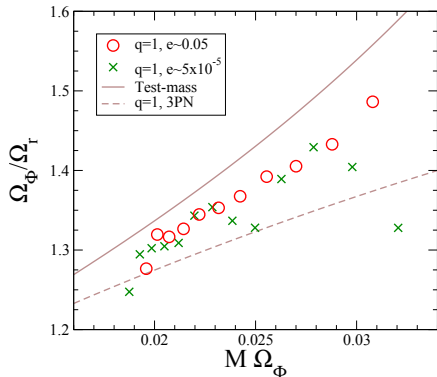
$$K_{\text{GSF}}^q = \frac{1}{\sqrt{1 - 6x}} \left[1 - \frac{q}{2} \frac{\rho(x)}{1 - 6x} + \mathcal{O}(q^2) \right]$$

- Since at first order $q = \nu + \mathcal{O}(\nu^2)$, we also have

$$K_{\text{GSF}}^\nu = \frac{1}{\sqrt{1 - 6x}} \left[1 - \frac{\nu}{2} \frac{\rho(x)}{1 - 6x} + \mathcal{O}(\nu^2) \right]$$

“Early” numerical relativity results for $K(\Omega_\varphi)$

[Mroué, Pfeiffer, Kidder & Teukolsky 2010]



Improved measurement of $K(\Omega_\varphi)$ in NR (1/2)

- Extract the orbital frequency $\Omega(t)$ from the coordinate separation $\mathbf{r}(t)$ *via*

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- Fit the NR data for $\Omega(t)$ to a model of the form

$$\Omega_\varphi(t) = p_0 (p_1 - t)^{p_2}$$

$$\delta\Omega(t) = p_3 \cos [p_4 + p_5(t - T) + p_6(t - T)^2]$$

for a time interval of width $W \times 2\pi/\Omega(T)$, centered on $t = T$

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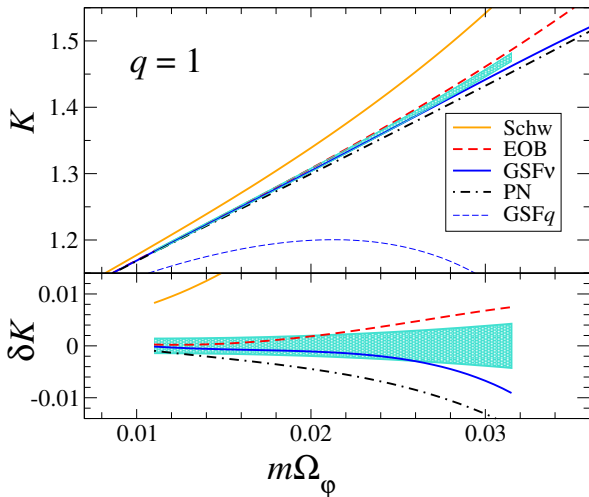
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- An error estimate is provided by changing the window size W
- Error introduced by non-zero eccentricity is subdominant
- Motion well approximated by a sequence of quasi-circular orbits because $0.3\% \lesssim \dot{\Omega}_\varphi/\Omega_\varphi^2 \lesssim 1.7\%$

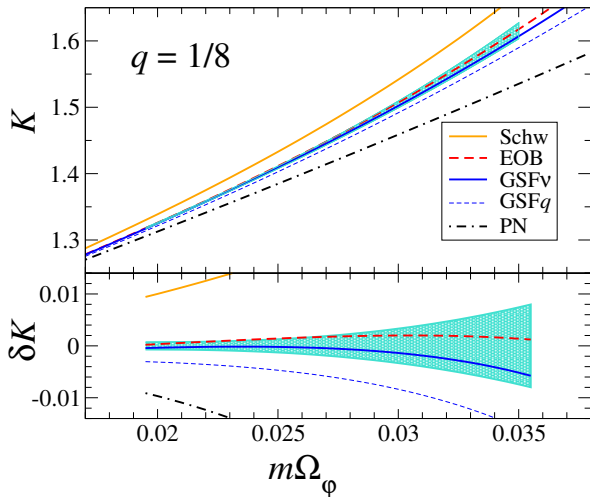
Extensive comparison for a mass ratio 1:1

[Le Tiec, Mroué, Barack, Buonanno, Pfeiffer, Sago & Taracchini 2011]



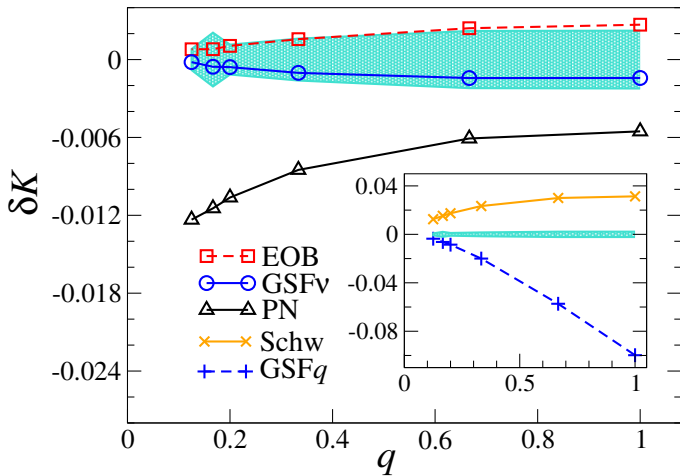
Extensive comparison for a mass ratio 1:8

[Le Tiec, Mroué, Barack, Buonanno, Pfeiffer, Sago & Taracchini 2011]



Variation with respect to the mass ratio

[Le Tiec, Mroué, Barack, Buonanno, Pfeiffer, Sago & Taracchini 2011]



Why does the GSF ν result perform so well?

A few *a posteriori* heuristic explanations:

- K is **linear in ν** up to 3PN order
- The 3PN term $\propto \nu^2$ contributes **less than 1%** to $K_{3\text{PN}}$
- The exact $K(\Omega_\varphi; m_1, m_2)$ must be **symmetric** by exchange $1 \leftrightarrow 2$ of the black holes
- **Assuming** the coefficients a_n in the series $K = \sum_n a_n(x) \nu^n$ do not increase quickly with n , this expansion will exhibit a **fast convergence** since $0 < \nu \leq 1/4$

Summary and outlook

- **Good convergence** of the PN series in the comparable mass regime $q \lesssim 1$, unlike in the extreme mass ratio limit $q \ll 1$
- The EOB (3PN) prediction is in **very good agreement** with the NR results for all mass ratios and frequencies considered
- The standard GSF result in terms of the mass ratio q agrees with the NR data up to a difference $\mathcal{O}(q^2)$
- The **GSF result with $q \rightarrow \nu$** compares remarkably well to the NR results, even for binaries with $m_1 \simeq m_2$
- First order BH perturbation theory may be good enough to model the GW emission from **IMRIs with $q \sim 10^{-4} - 10^{-2}$**