

Eccentric motion on a Schwarzschild background:

Self-force in a modified Regge-Wheeler gauge

Seth Hopper - Albert Einstein Institute

Capra 16 - July 16, 2013

Outline

- Gauge freedom on Schwarzschild
- Infinitesimal gauge transformations
- Lorenz gauge
- Modified Regge-Wheeler gauge

Gauge freedom

- In GR, gauge freedom is coordinate freedom
- Zeroth order: use Schwarzschild coordinates
- First-order options:
 - Lorenz
 - Regge-Wheeler
 - Modified Regge-Wheeler?

Lorenz gauge

$$g_{\mu\nu} = g_{\mu\nu}^{\text{Schw}} + p_{\mu\nu}$$

Lorenz gauge

$$g_{\mu\nu} = g_{\mu\nu}^{\text{Schw}} + p_{\mu\nu}$$

$$\square \bar{p}_{\mu\nu} + 2R_{\alpha\mu\beta\nu} \bar{p}^{\alpha\beta} = -16\pi T_{\mu\nu}, \quad \bar{p}^{\mu\nu}{}_{|\nu} = 0$$

Lorenz gauge

$$g_{\mu\nu} = g_{\mu\nu}^{\text{Schw}} + p_{\mu\nu}$$

$$\square \bar{p}_{\mu\nu} + 2R_{\alpha\mu\beta\nu} \bar{p}^{\alpha\beta} = -16\pi T_{\mu\nu}, \quad \bar{p}^{\mu\nu}{}_{|\nu} = 0$$

- 10 coupled wave equations
- Locally isotropic solutions
- Regularization procedure in Lorenz gauge
 - Other gauges may be possible ...

Regge-Wheeler gauge

- Schematically: $p_{\mu\nu} \rightarrow \sum_{\ell,m} h_{\mu\nu}^{\ell m} Y^{\ell m}$
- Set four components of $h_{\mu\nu}^{\ell m}$ to zero
- Field equations simplify greatly:

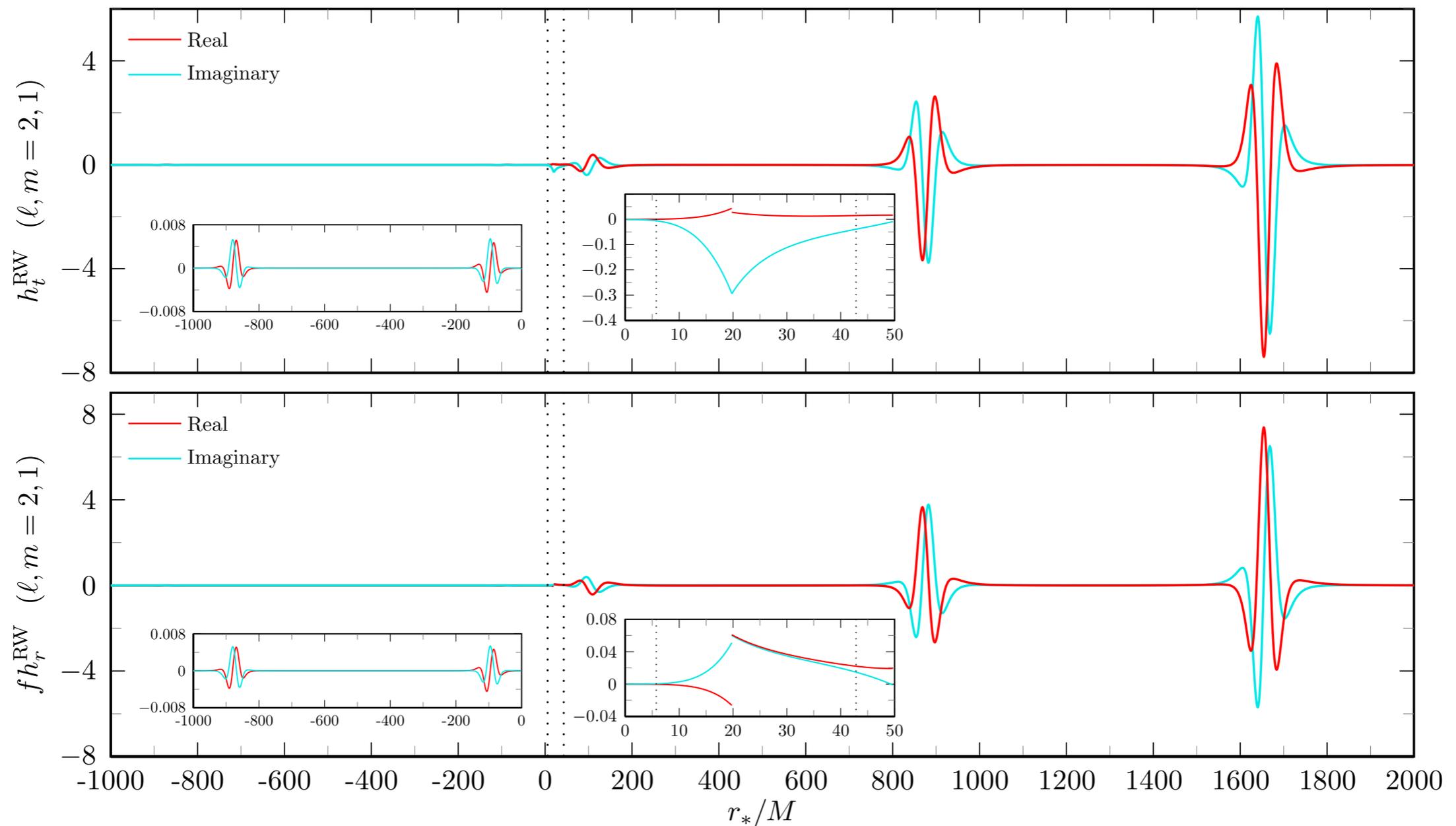
$$\left[-\frac{\partial^2}{\partial t^2} + \frac{\partial^2}{\partial r_*^2} - V_\ell(r) \right] \Psi_{\ell m}(t, r) = S_{\ell m}(t)$$

- Reconstruct metric perturbation:

$$\Psi_{\ell m}(t, r) \rightarrow p_{\mu\nu}$$

RW gauge fields: odd-parity

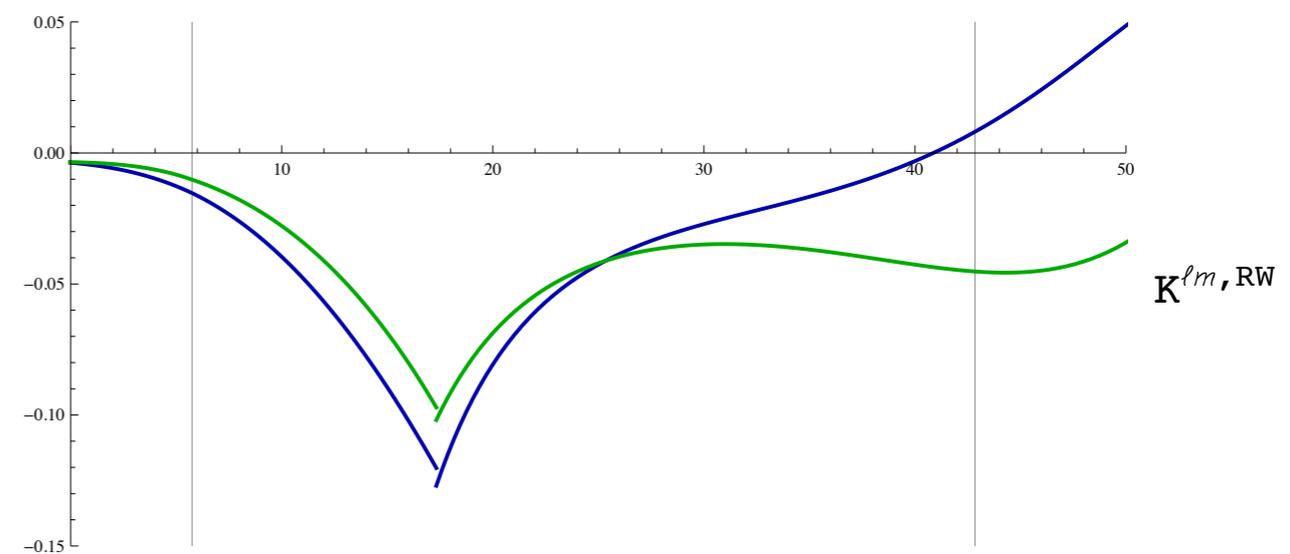
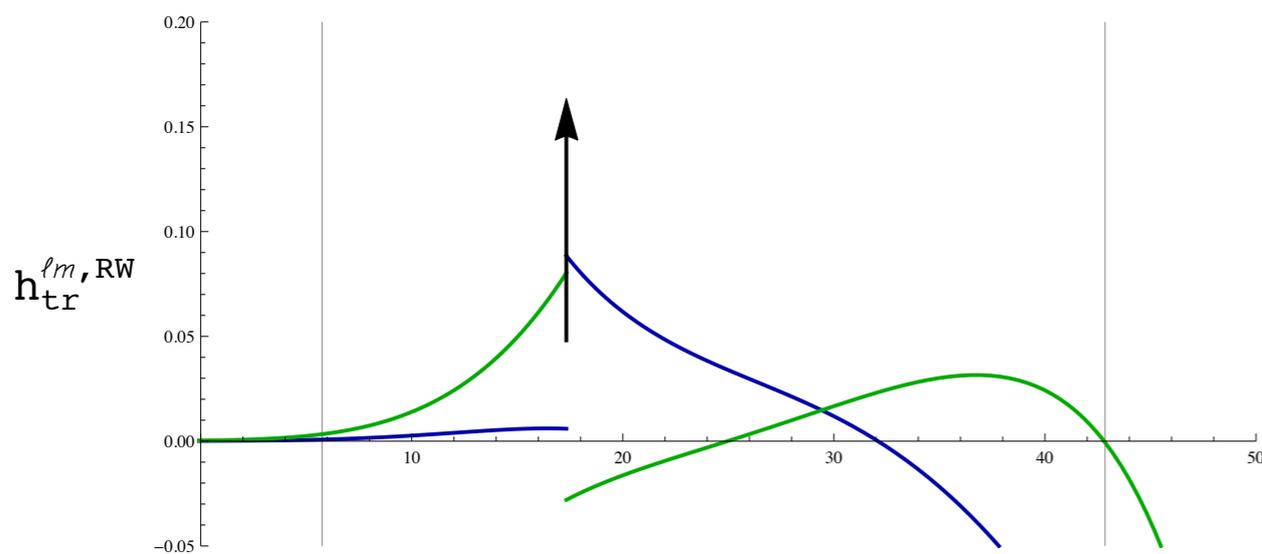
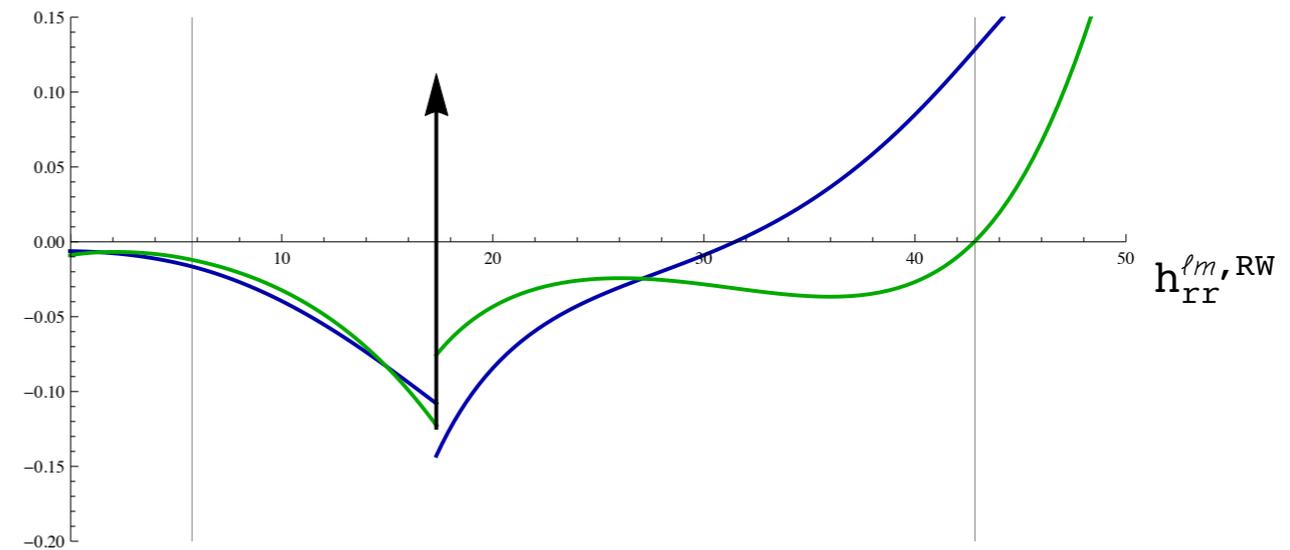
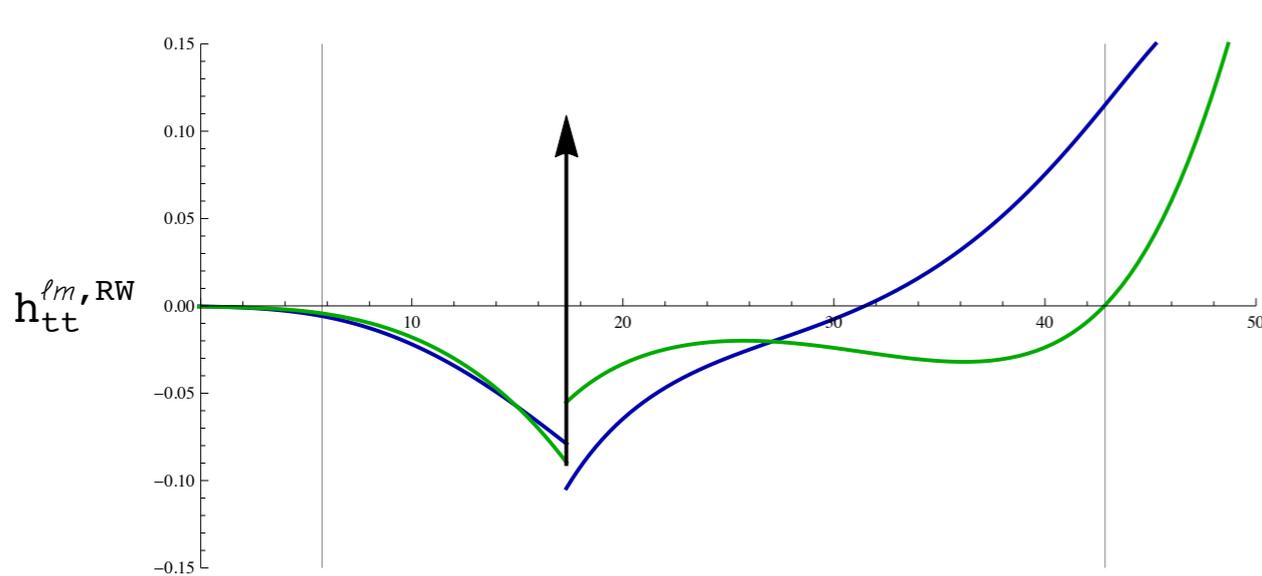
- Two non-vanishing amplitudes: $h_t^{\ell m}$, $h_r^{\ell m}$



RW gauge fields: even-parity

- Four non-vanishing amplitudes: $h_{tt}^{\ell m}$, $h_{tr}^{\ell m}$, $h_{rr}^{\ell m}$, $K^{\ell m}$

$$(\ell, m) = (2, 2)$$



Benefits and drawbacks of RW gauge

- Benefits:
 - Simple field equations
 - Computationally efficient

Benefits and drawbacks of RW gauge

- Benefits:
 - Simple field equations
 - Computationally efficient
- Drawbacks:
 - Only valid for radiative modes, $\ell \geq 2$
 - Singularities/discontinuities at the particle
 - Self-force not well-defined

A regular, well-defined self-force

- Barack and Ori, 2001
 - “Gravitational self force and gauge transformations”

A regular, well-defined self-force

- Barack and Ori, 2001
 - “Gravitational self force and gauge transformations”
- Lorenz gauge yields a regular, well-defined SF

A regular, well-defined self-force

- Barack and Ori, 2001
 - “Gravitational self force and gauge transformations”
- Lorenz gauge yields a regular, well-defined SF
- SF transformation with gauge vector Ξ^μ :

$$\delta F_{\text{self}}^\alpha = -\mu \left[(g^{\alpha\beta} + u^\alpha u^\beta) \ddot{\Xi}_\beta + R^\alpha{}_{\mu\beta\nu} u^\mu \Xi^\beta u^\nu \right]$$

A regular, well-defined self-force

- Barack and Ori, 2001
 - “Gravitational self force and gauge transformations”
- Lorenz gauge yields a regular, well-defined SF
- SF transformation with gauge vector Ξ^μ :

$$\delta F_{\text{self}}^\alpha = -\mu \left[(g^{\alpha\beta} + u^\alpha u^\beta) \ddot{\Xi}_\beta + R^\alpha{}_{\mu\beta\nu} u^\mu \Xi^\beta u^\nu \right]$$

- SF is well-defined if and only if $\delta F_{\text{self}}^\alpha$ relative to Lorenz gauge is

A regular, well-defined self-force

- Barack and Ori, 2001

- “Gravitational self force and gauge transformations”

- Lorenz gauge yields a regular, well-defined SF

- SF transformation with gauge vector Ξ^μ :

$$\delta F_{\text{self}}^\alpha = -\mu \left[(g^{\alpha\beta} + u^\alpha u^\beta) \ddot{\Xi}_\beta + R^\alpha{}_{\mu\beta\nu} u^\mu \Xi^\beta u^\nu \right]$$

- SF is well-defined if and only if $\delta F_{\text{self}}^\alpha$ relative to Lorenz gauge is

- If vector Ξ^μ is well-defined, the SF will be also

A regular, well-defined self-force

- Barack and Ori, 2001

- “Gravitational self force and gauge transformations”

- Lorenz gauge yields a regular, well-defined SF

- SF transformation with gauge vector Ξ^μ :

$$\delta F_{\text{self}}^\alpha = -\mu \left[(g^{\alpha\beta} + u^\alpha u^\beta) \ddot{\Xi}_\beta + R^\alpha{}_{\mu\beta\nu} u^\mu \Xi^\beta u^\nu \right]$$

- SF is well-defined if and only if $\delta F_{\text{self}}^\alpha$ relative to Lorenz gauge is

- If vector Ξ^μ is well-defined, the SF will be also

- Then, regularization is done with Lorenz gauge parameters $A^\alpha, B^\alpha, C^\alpha, D^\alpha$

First-order gauge transformations

- Transform from RW to Lorenz gauge:

$$x_{\text{L}}^{\mu} = x_{\text{RW}}^{\mu} + \Xi_{\text{RW} \rightarrow \text{L}}^{\mu}, \quad |\Xi_{\text{RW} \rightarrow \text{L}}^{\mu}| \sim |p_{\mu\nu}| \ll |g_{\mu\nu}^{\text{Schw}}|$$

First-order gauge transformations

- Transform from RW to Lorenz gauge:

$$x_{\text{L}}^{\mu} = x_{\text{RW}}^{\mu} + \Xi_{\text{RW} \rightarrow \text{L}}^{\mu}, \quad |\Xi_{\text{RW} \rightarrow \text{L}}^{\mu}| \sim |p_{\mu\nu}| \ll |g_{\mu\nu}^{\text{Schw}}|$$

- Metric perturbation transforms:

$$p_{\mu\nu}^{\text{L}} = p_{\mu\nu}^{\text{RW}} - \Xi_{\mu|\nu}^{\text{RW} \rightarrow \text{L}} - \Xi_{\nu|\mu}^{\text{RW} \rightarrow \text{L}}$$

First-order gauge transformations

- Transform from RW to Lorenz gauge:

$$x_{\text{L}}^{\mu} = x_{\text{RW}}^{\mu} + \Xi_{\text{RW} \rightarrow \text{L}}^{\mu}, \quad |\Xi_{\text{RW} \rightarrow \text{L}}^{\mu}| \sim |p_{\mu\nu}| \ll |g_{\mu\nu}^{\text{Schw}}|$$

- Metric perturbation transforms:

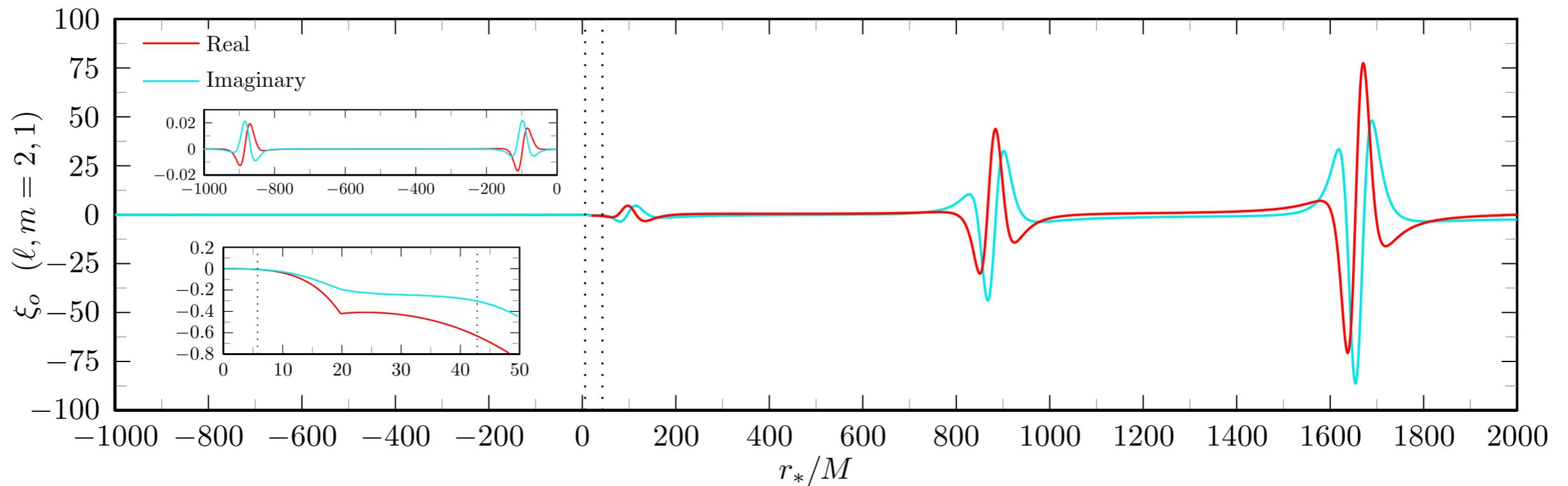
$$p_{\mu\nu}^{\text{L}} = p_{\mu\nu}^{\text{RW}} - \Xi_{\mu|\nu}^{\text{RW} \rightarrow \text{L}} - \Xi_{\nu|\mu}^{\text{RW} \rightarrow \text{L}}$$

- Gauge vector satisfies a wave equation:

$$\square \Xi_{\text{RW} \rightarrow \text{L}}^{\mu} = \bar{p}_{\text{RW}|\nu}^{\mu\nu}$$

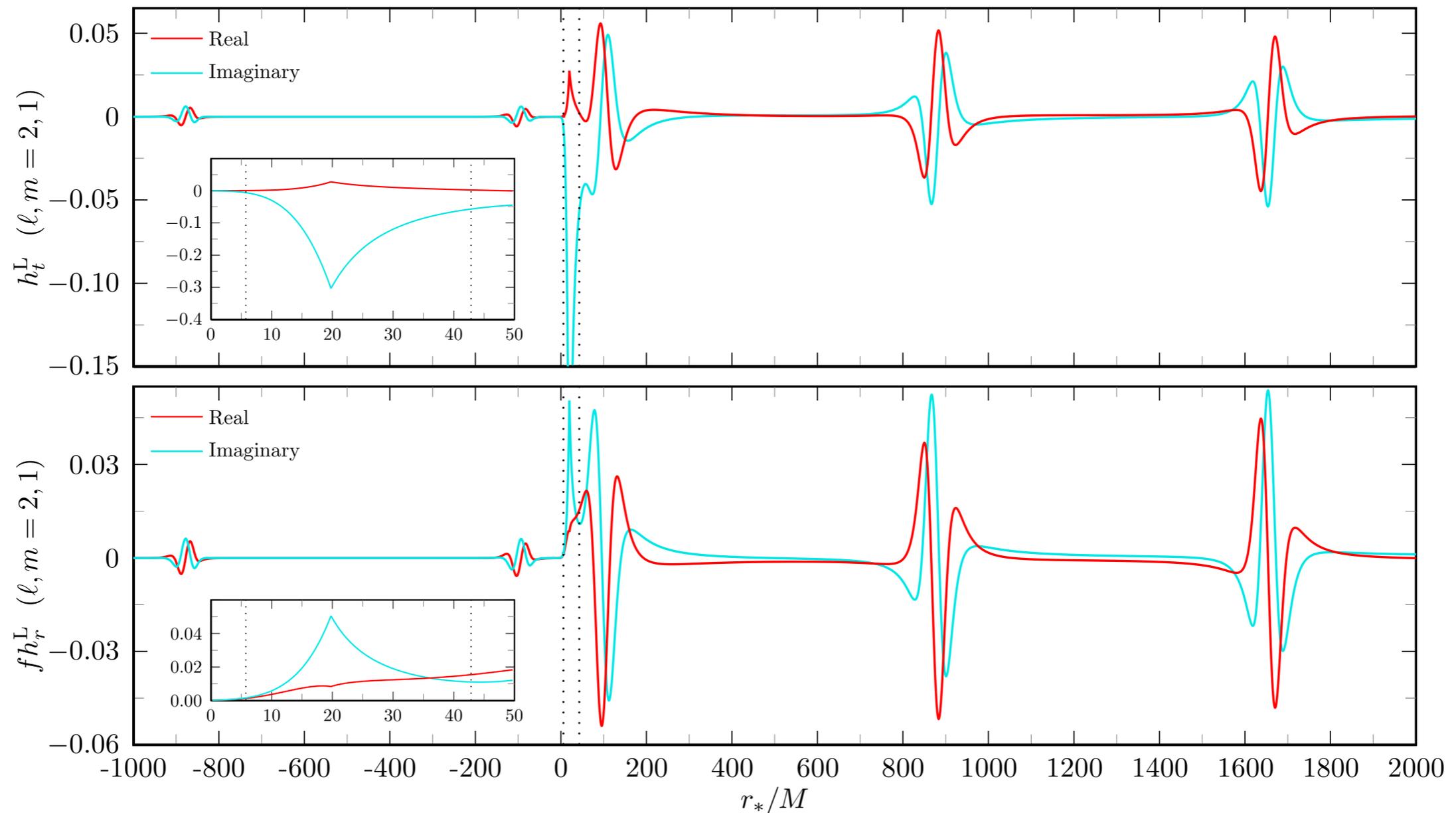
RW->Lorenz gauge vector: odd-parity

- Transform the global solution, mode-by-mode



Lorenz gauge fields: odd-parity

- Amplitudes are now C^0 and asymptotically flat



Benefits/drawbacks of global gauge transf.

$$\square \Xi_{RW \rightarrow L}^{\mu} = \bar{p}_{RW|\nu}^{\mu\nu}$$

Benefits/drawbacks of global gauge transf.

$$\square \Xi_{\text{RW} \rightarrow \text{L}}^{\mu} = \bar{p}_{\text{RW}|\nu}^{\mu\nu}$$

- Benefits:
 - Gives the solution, everywhere in Lorenz gauge
 - Gives solution to low-order modes

Benefits/drawbacks of global gauge transf.

$$\square \Xi_{\text{RW} \rightarrow \text{L}}^{\mu} = \bar{p}_{\text{RW}|\nu}^{\mu\nu}$$

- Benefits:
 - Gives the solution, everywhere in Lorenz gauge
 - Gives solution to low-order modes
- Drawbacks:
 - Computationally difficult and expensive
 - Discontinuous, extended source terms
 - Excessive, if you just want the self-force

A modified RW gauge

- Gralla, 2011 - Simple gauge transf. to reach “parity-regular” gauge

A modified RW gauge

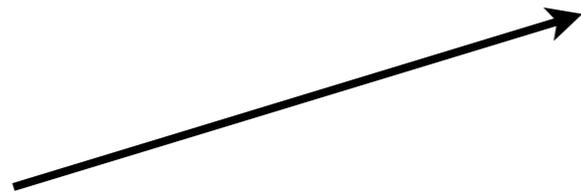
- Gralla, 2011 - Simple gauge transf. to reach “parity-regular” gauge
- Split gauge transformation into two steps

$$\Xi_{RW \rightarrow L}^{\mu} = \Xi_{RW \rightarrow MRW}^{\mu} + \Xi_{MRW \rightarrow L}^{\mu}$$

A modified RW gauge

- Gralla, 2011 - Simple gauge transf. to reach “parity-regular” gauge
- Split gauge transformation into two steps

$$\Xi_{RW \rightarrow L}^{\mu} = \underline{\Xi_{RW \rightarrow MRW}^{\mu}} + \Xi_{MRW \rightarrow L}^{\mu}$$



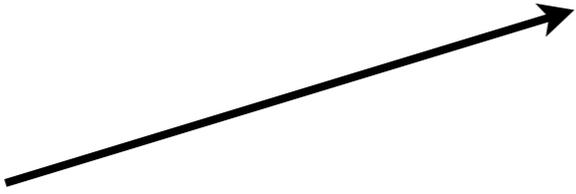
Remove major discontinuities

A modified RW gauge

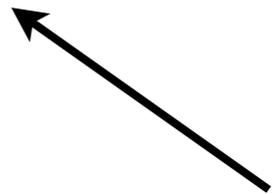
- Gralla, 2011 - Simple gauge transf. to reach “parity-regular” gauge
- Split gauge transformation into two steps

$$\Xi_{\text{RW} \rightarrow \text{L}}^{\mu} = \underline{\Xi_{\text{RW} \rightarrow \text{MRW}}^{\mu}} + \underline{\Xi_{\text{MRW} \rightarrow \text{L}}^{\mu}}$$

Remove major discontinuities



Smooth enough to ignore



A modified RW gauge

- The metric perturbation in MRW gauge is

$$p_{\mu\nu}^{\text{MRW}} = p_{\mu\nu}^{\text{RW}} - 2\Xi_{(\mu|\nu)}^{\text{RW}\rightarrow\text{MRW}}$$

A modified RW gauge

- The metric perturbation in MRW gauge is

$$p_{\mu\nu}^{\text{MRW}} = p_{\mu\nu}^{\text{RW}} - 2\Xi_{(\mu|\nu)}^{\text{RW}\rightarrow\text{MRW}}$$

- Decompose into spherical harmonics, e.g.

$$h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$$

A modified RW gauge

- We demand $[[h_t^{\ell m, \text{MRW}}]] = [[h_t^{\ell m, \text{L}}]] = 0$

A modified RW gauge

- We demand $[[h_t^{\ell m, \text{MRW}}]] = [[h_t^{\ell m, \text{L}}]] = 0$
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$

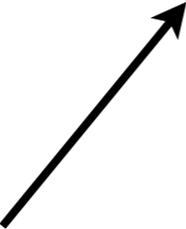
A modified RW gauge

- We demand $[[h_t^{\ell m, \text{MRW}}]] = [[h_t^{\ell m, \text{L}}]] = 0$
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$
- Therefore $[[h_t^{\ell m, \text{RW}}]] = [[\partial_t \xi_{\text{odd}}^{\ell m}]]$

A modified RW gauge

- We demand $[[h_t^{\ell m, \text{MRW}}]] = [[h_t^{\ell m, \text{L}}]] = 0$
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$
- Therefore $[[h_t^{\ell m, \text{RW}}]] = [[\partial_t \xi_{\text{odd}}^{\ell m}]]$

We know this



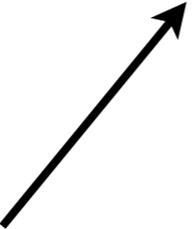
A modified RW gauge

- We demand $[[h_t^{\ell m, \text{MRW}}]] = [[h_t^{\ell m, \text{L}}]] = 0$

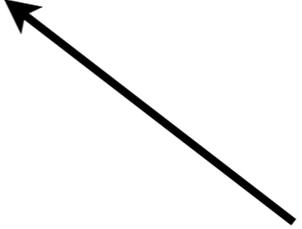
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$

- Therefore $\underline{[[h_t^{\ell m, \text{RW}}]]} = \underline{[[\partial_t \xi_{\text{odd}}^{\ell m}]]}$

We know this



Restriction on the gauge vector



A modified RW gauge

- Restrictions on $[\xi_{\text{odd}}^{\ell m}]$, $[\partial_t \xi_{\text{odd}}^{\ell m}]$, $[\partial_r \xi_{\text{odd}}^{\ell m}]$
- Away from the particle, no restrictions

A modified RW gauge

- Restrictions on $[\xi_{\text{odd}}^{\ell m}]$, $[\partial_t \xi_{\text{odd}}^{\ell m}]$, $[\partial_r \xi_{\text{odd}}^{\ell m}]$
- Away from the particle, no restrictions
- A possible vector:

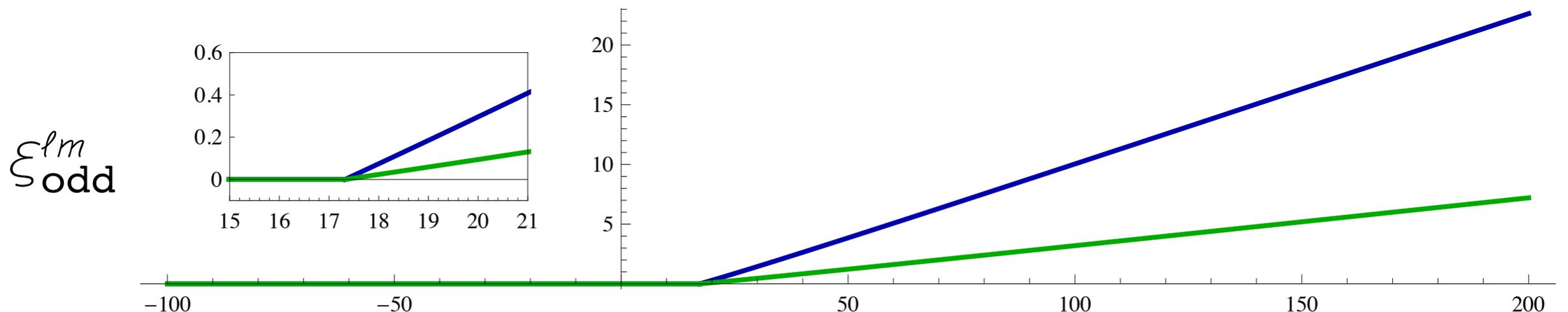
$$\xi_{\text{odd}}^{\ell m}(t, r) = (r - r_p) [h_r^{\ell m, \text{RW}}] \theta [r - r_p]$$

A modified RW gauge

- Restrictions on $[\xi_{\text{Sodd}}^{\ell m}]$, $[\partial_t \xi_{\text{Sodd}}^{\ell m}]$, $[\partial_r \xi_{\text{Sodd}}^{\ell m}]$
- Away from the particle, no restrictions
- A possible vector:

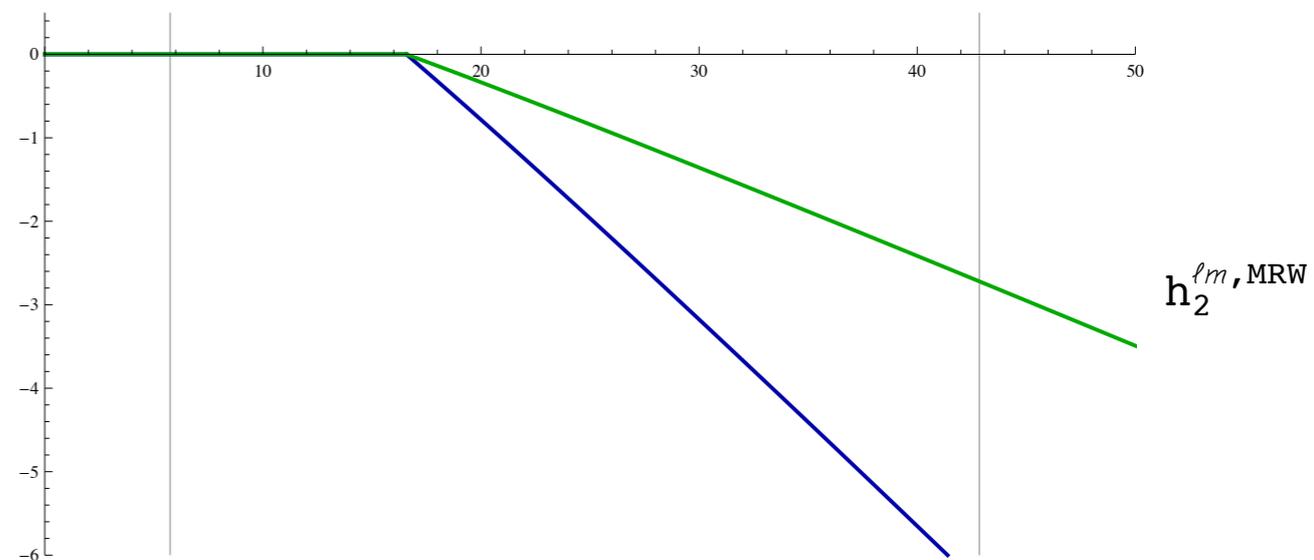
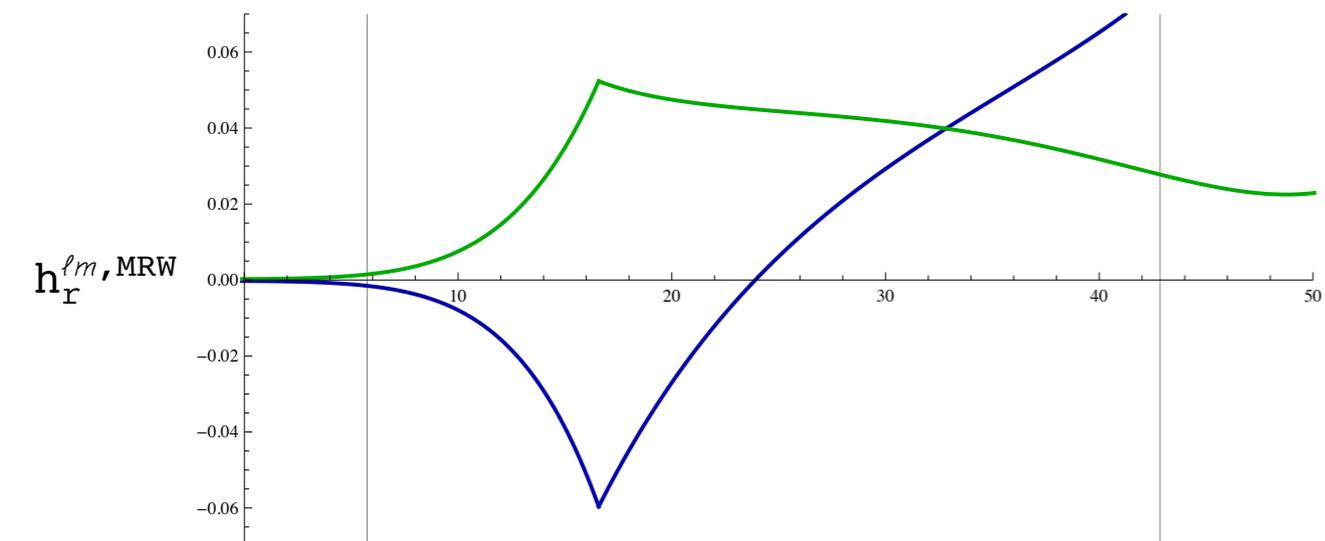
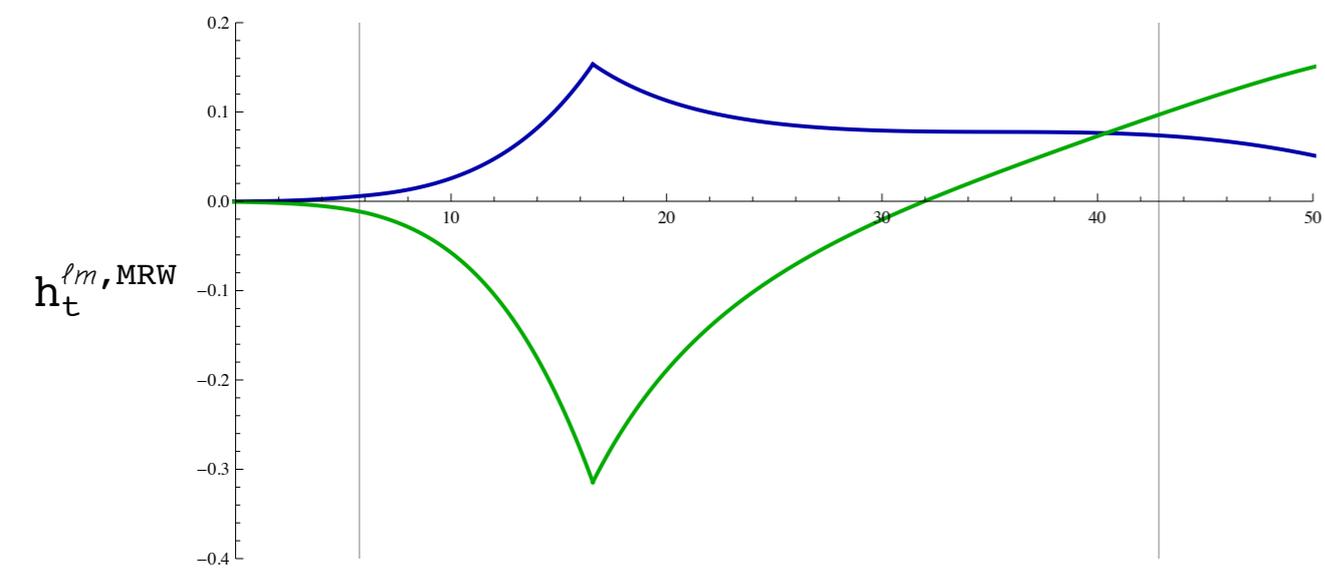
$$\xi_{\text{Sodd}}^{\ell m}(t, r) = (r - r_p) [h_r^{\ell m, \text{RW}}] \theta [r - r_p]$$

$$(p, e, t_p) = (8.75455, 0.764124, 80.17) \quad (\ell, m) = (2, 1)$$



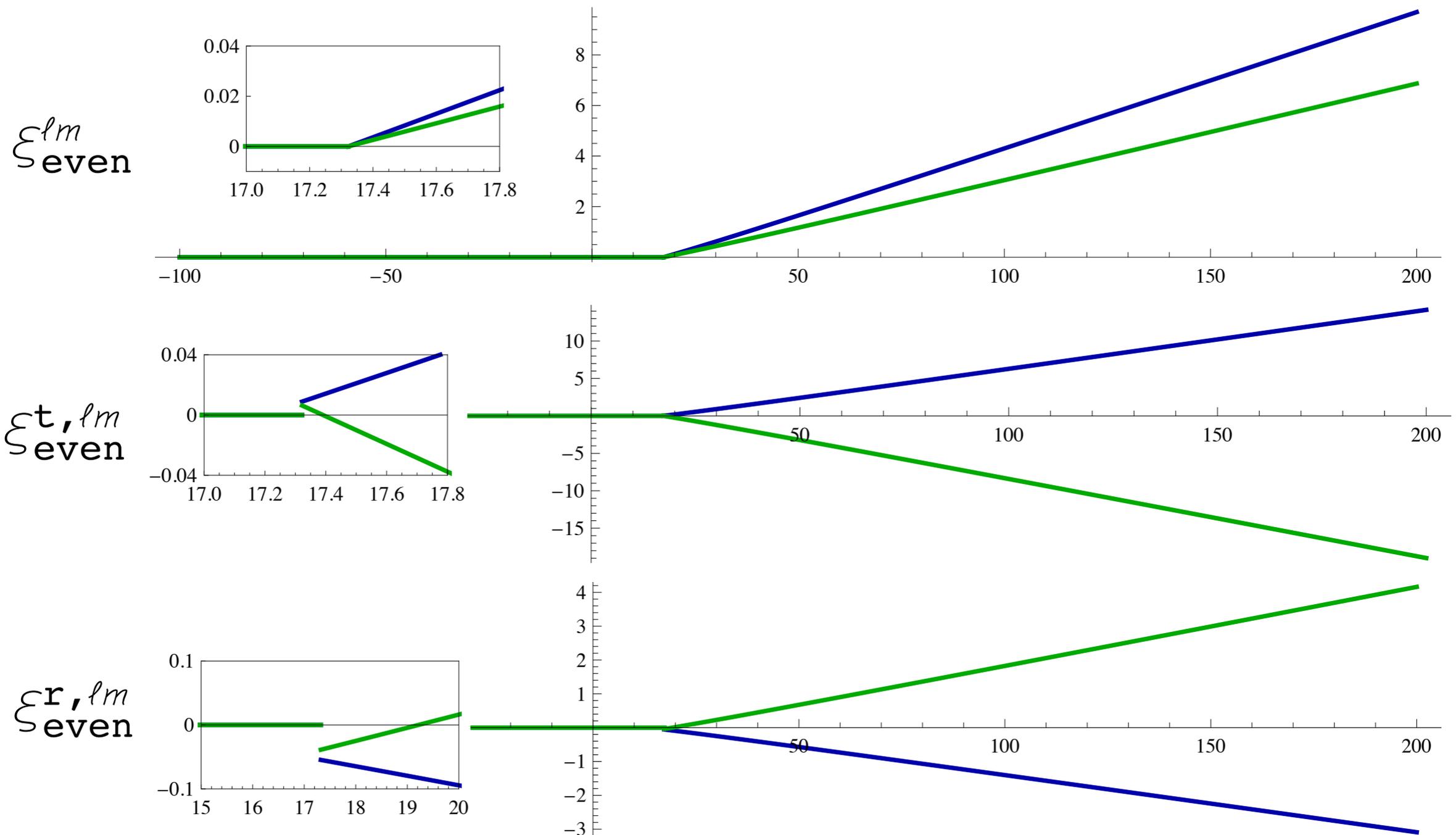
Metric perturbation in modified RW gauge

$$(p, e, t_p) = (8.75455, 0.764124, 80.17) \quad (\ell, m) = (2, 1)$$



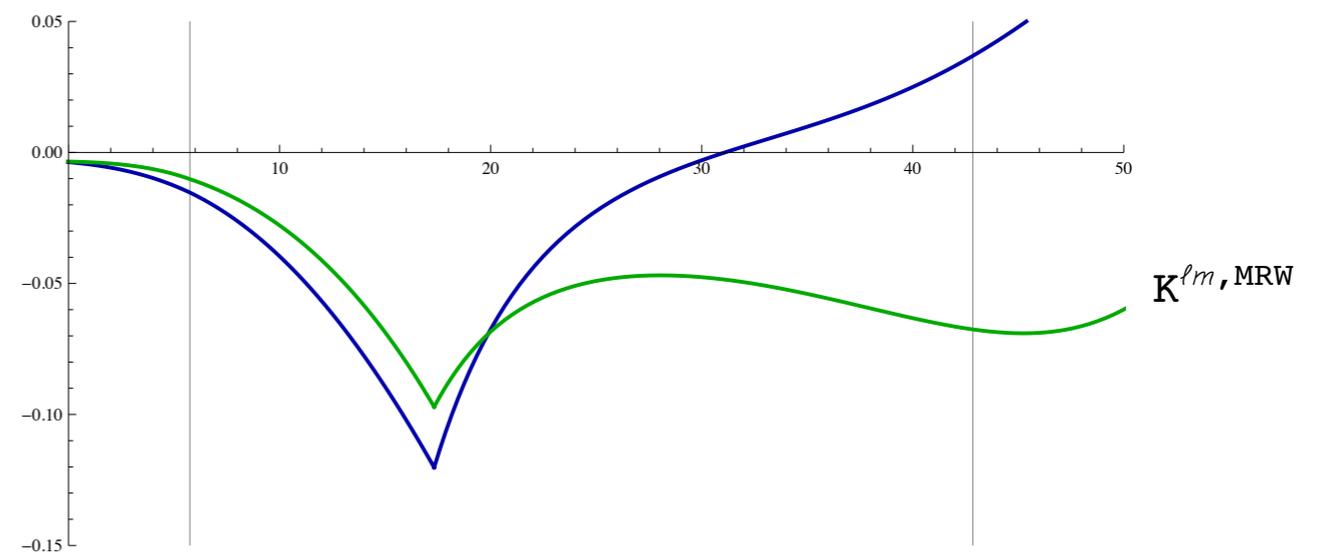
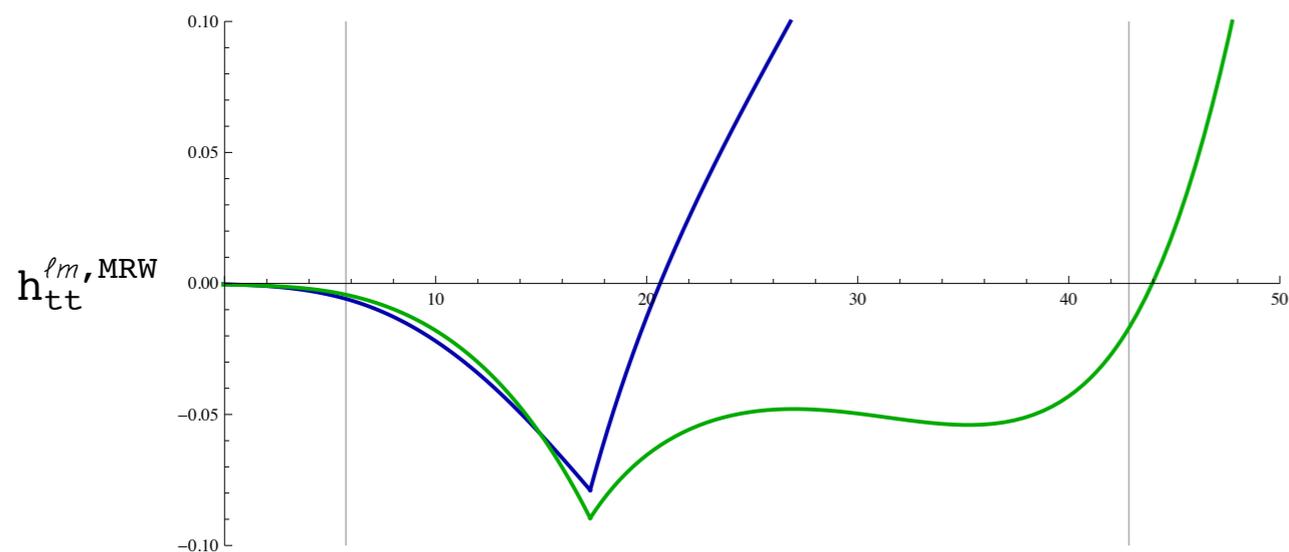
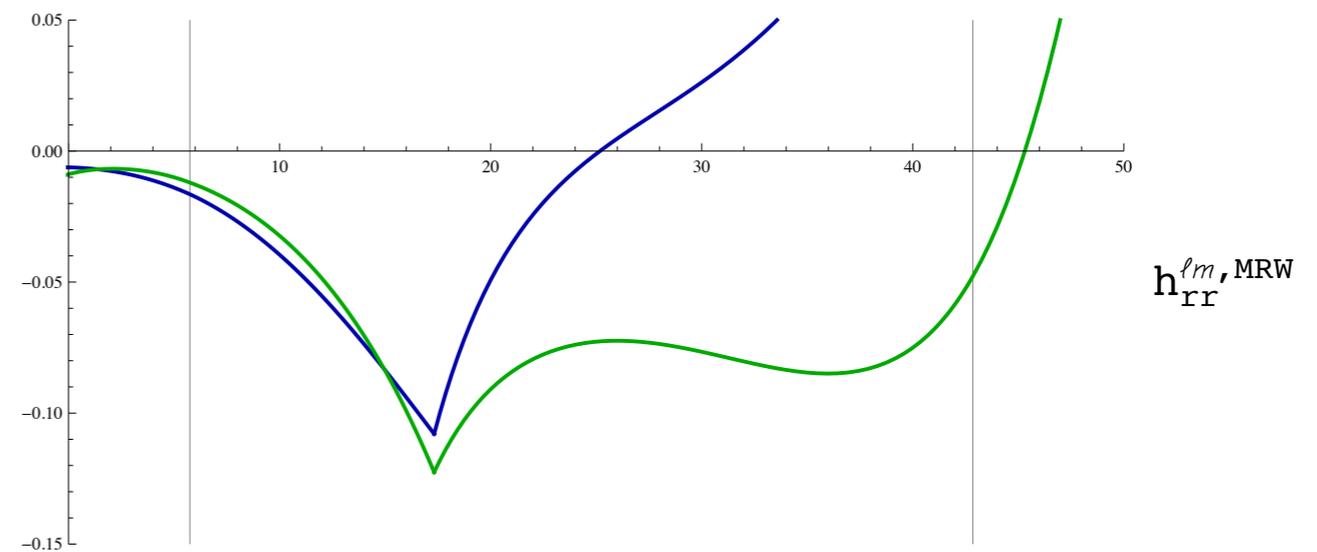
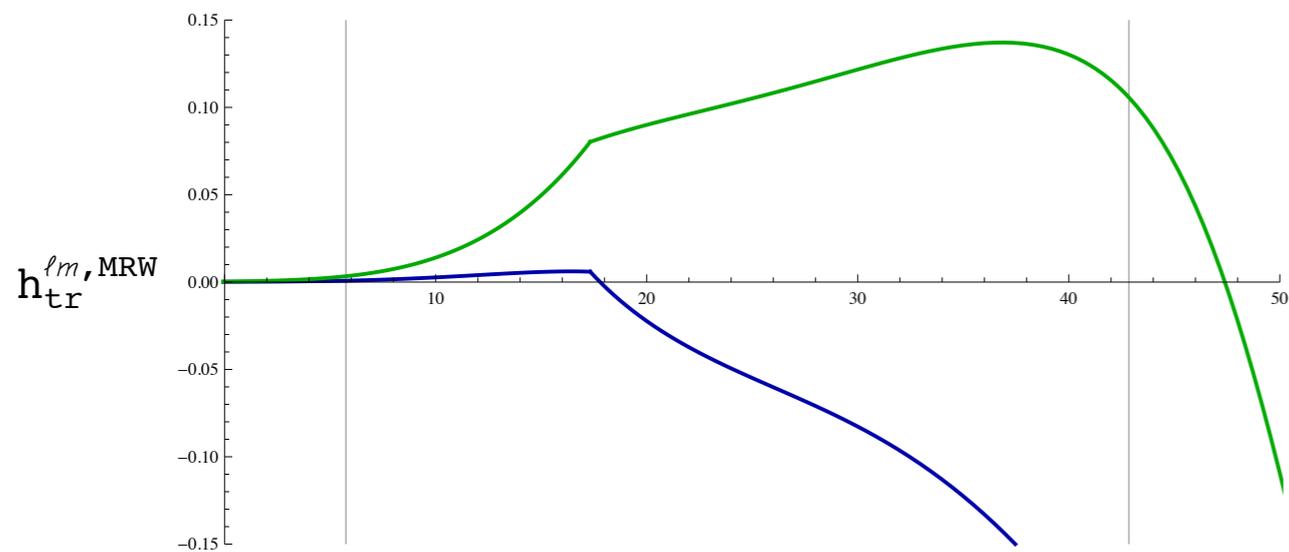
Even-parity gauge vector

$$(p, e, t_p) = (8.75455, 0.764124, 80.17) \quad (\ell, m) = (2, 2)$$



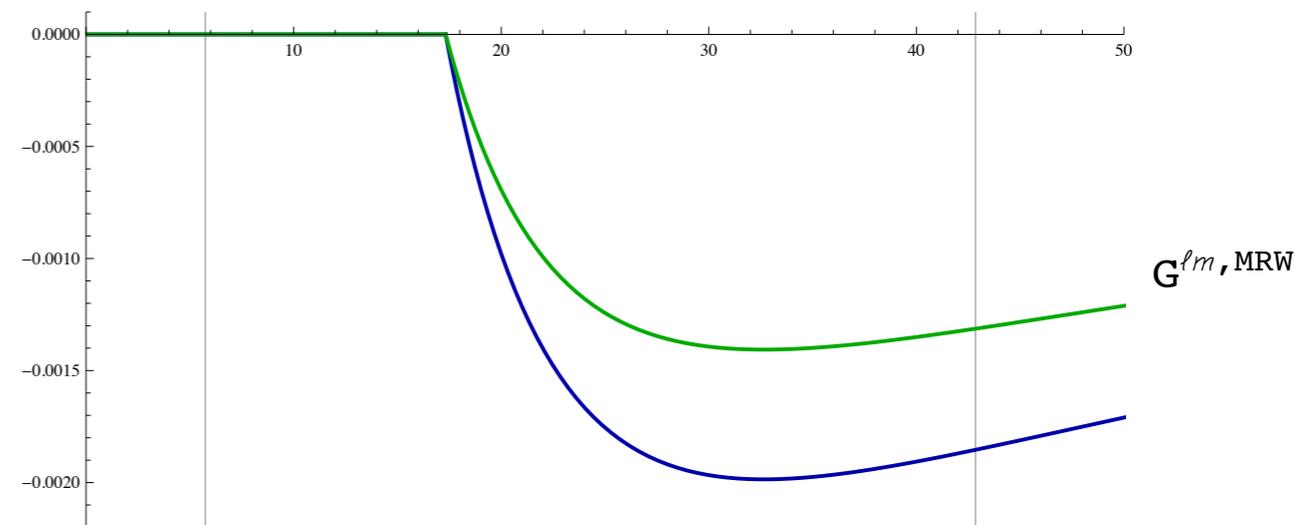
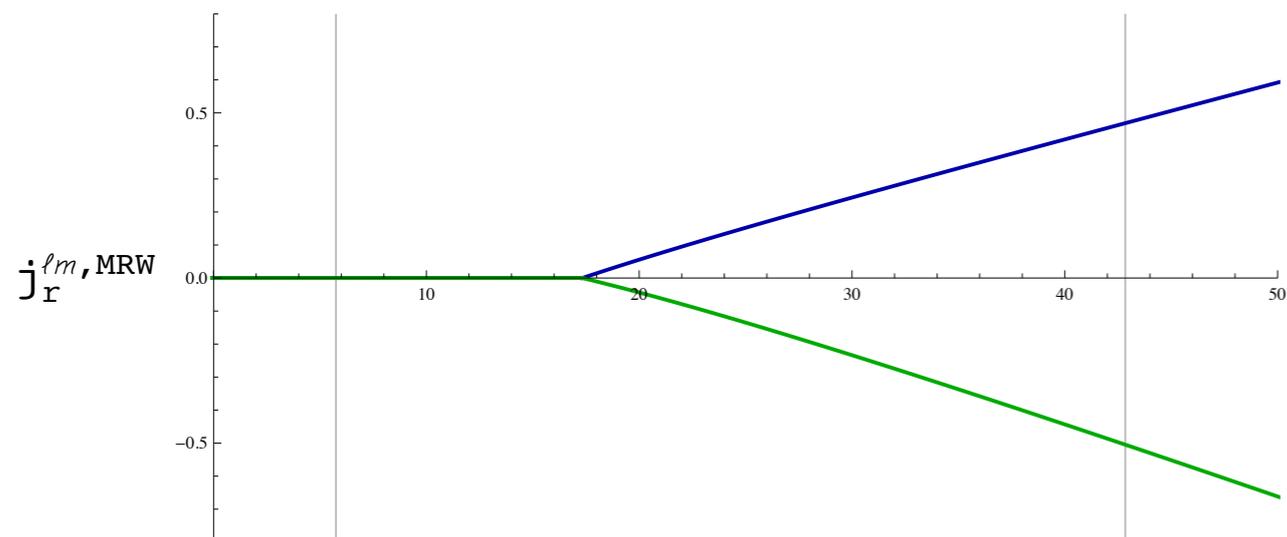
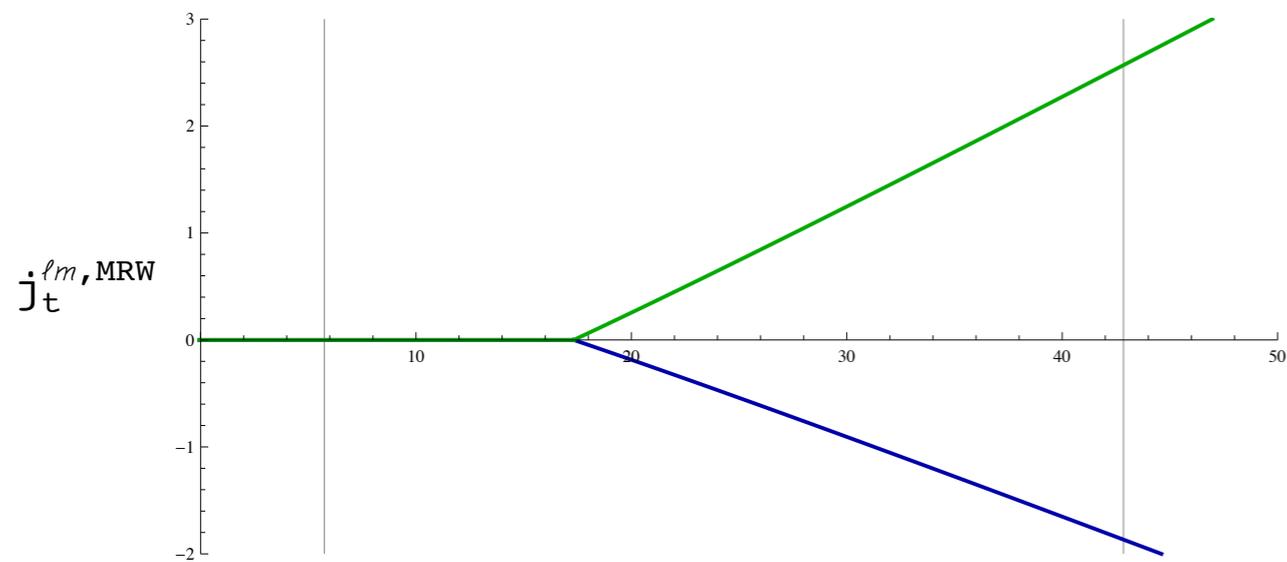
Metric perturbation in modified RW gauge

$$(p, e, t_p) = (8.75455, 0.764124, 80.17) \quad (\ell, m) = (2, 2)$$



Metric perturbation in modified RW gauge

$$(p, e, t_p) = (8.75455, 0.764124, 80.17) \quad (\ell, m) = (2, 2)$$

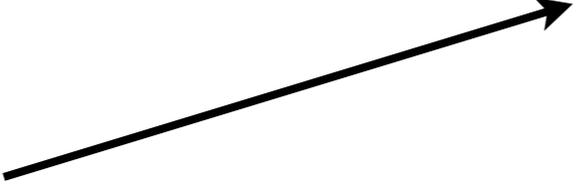


Is this good enough?

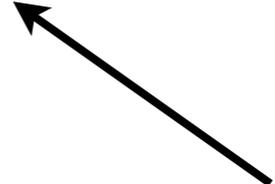
- Split gauge transformation into two steps

$$\Xi_{RW \rightarrow L}^\mu = \underline{\Xi_{RW \rightarrow MRW}^\mu} + \underline{\Xi_{MRW \rightarrow L}^\mu}$$

Remove major discontinuities



Smooth enough to ignore



Is this good enough?

- Split gauge transformation into two steps

$$\Xi_{RW \rightarrow L}^{\mu} = \underline{\Xi_{RW \rightarrow MRW}^{\mu}} + \underline{\Xi_{MRW \rightarrow L}^{\mu}}$$

Remove major discontinuities

Smooth enough to ignore

- The remaining transformation comes from

$$\square \Xi_{MRW \rightarrow L}^{\mu} = \bar{p}_{MRW}^{\mu\nu} |_{\nu}$$

Is this good enough?

- Split gauge transformation into two steps

$$\Xi_{RW \rightarrow L}^\mu = \underline{\Xi_{RW \rightarrow MRW}^\mu} + \underline{\Xi_{MRW \rightarrow L}^\mu}$$

Remove major discontinuities

Smooth enough to ignore

- The remaining transformation comes from

$$\square \Xi_{MRW \rightarrow L}^\mu = \underline{\bar{p}_{MRW}^{\mu\nu}} |_\nu$$

This is now C^0

Is this good enough?

- Split gauge transformation into two steps

$$\Xi_{RW \rightarrow L}^\mu = \underline{\Xi_{RW \rightarrow MRW}^\mu} + \underline{\Xi_{MRW \rightarrow L}^\mu}$$

Remove major discontinuities

Smooth enough to ignore

- The remaining transformation comes from

$$\square \Xi_{MRW \rightarrow L}^\mu = \underline{\bar{p}_{MRW}^{\mu\nu}} |_\nu$$

This is now C^0

- But we can do better

Extending the modified RW gauge

- We demand $[[\partial_r h_t^{\ell m, \text{MRW}}]] = [[\partial_r h_t^{\ell m, \text{L}}]]$

Extending the modified RW gauge

- We demand $[[\partial_r h_t^{\ell m, \text{MRW}}]] = [[\partial_r h_t^{\ell m, \text{L}}]]$
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$

Extending the modified RW gauge

- We demand $[[\partial_r h_t^{\ell m, \text{MRW}}]] = [[\partial_r h_t^{\ell m, \text{L}}]]$
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$
- Therefore $[[\partial_t \partial_r \xi_{\text{odd}}^{\ell m}]] = [[\partial_r h_t^{\ell m, \text{RW}}]] - [[\partial_r h_t^{\ell m, \text{L}}]]$

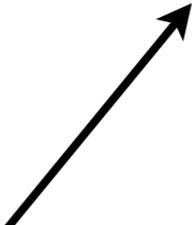
Extending the modified RW gauge

- We demand $[[\partial_r h_t^{\ell m, \text{MRW}}]] = [[\partial_r h_t^{\ell m, \text{L}}]]$
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$
- Therefore $[[\partial_t \partial_r \xi_{\text{odd}}^{\ell m}]] = \underbrace{[[\partial_r h_t^{\ell m, \text{RW}}]]}_{\substack{\swarrow \\ \text{We know these}}} - \underbrace{[[\partial_r h_t^{\ell m, \text{L}}]]}_{\substack{\nearrow \\ \text{We know these}}}$

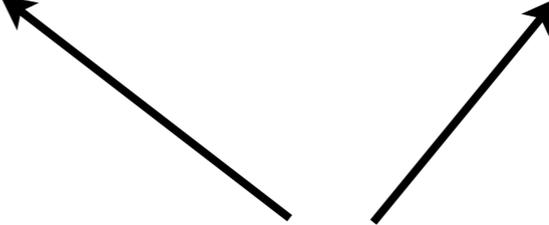
Extending the modified RW gauge

- We demand $[[\partial_r h_t^{\ell m, \text{MRW}}]] = [[\partial_r h_t^{\ell m, \text{L}}]]$
- Given $h_t^{\ell m, \text{MRW}} = h_t^{\ell m, \text{RW}} - \partial_t \xi_{\text{odd}}^{\ell m}$
- Therefore $\underline{[[\partial_t \partial_r \xi_{\text{odd}}^{\ell m}]]} = \underline{[[\partial_r h_t^{\ell m, \text{RW}}]]} - \underline{[[\partial_r h_t^{\ell m, \text{L}}]]}$

Restriction on the gauge vector

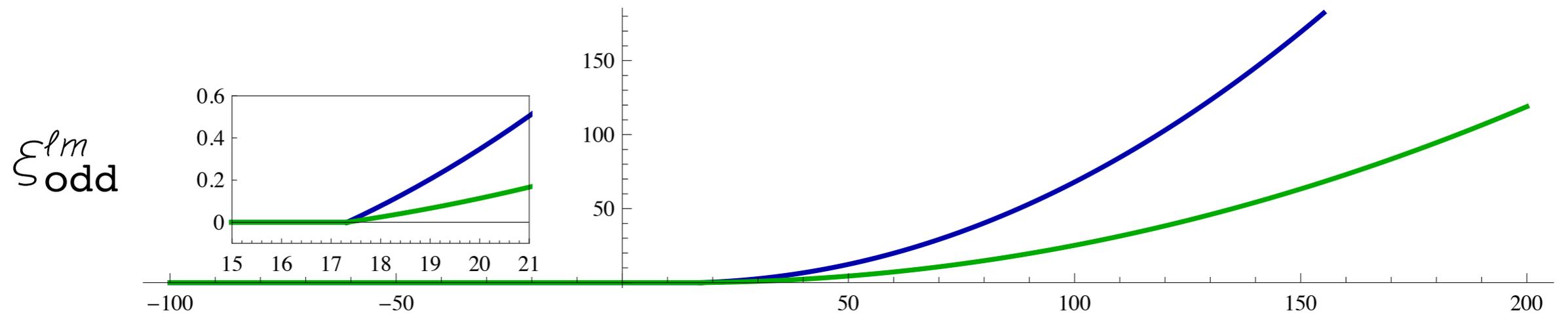


We know these



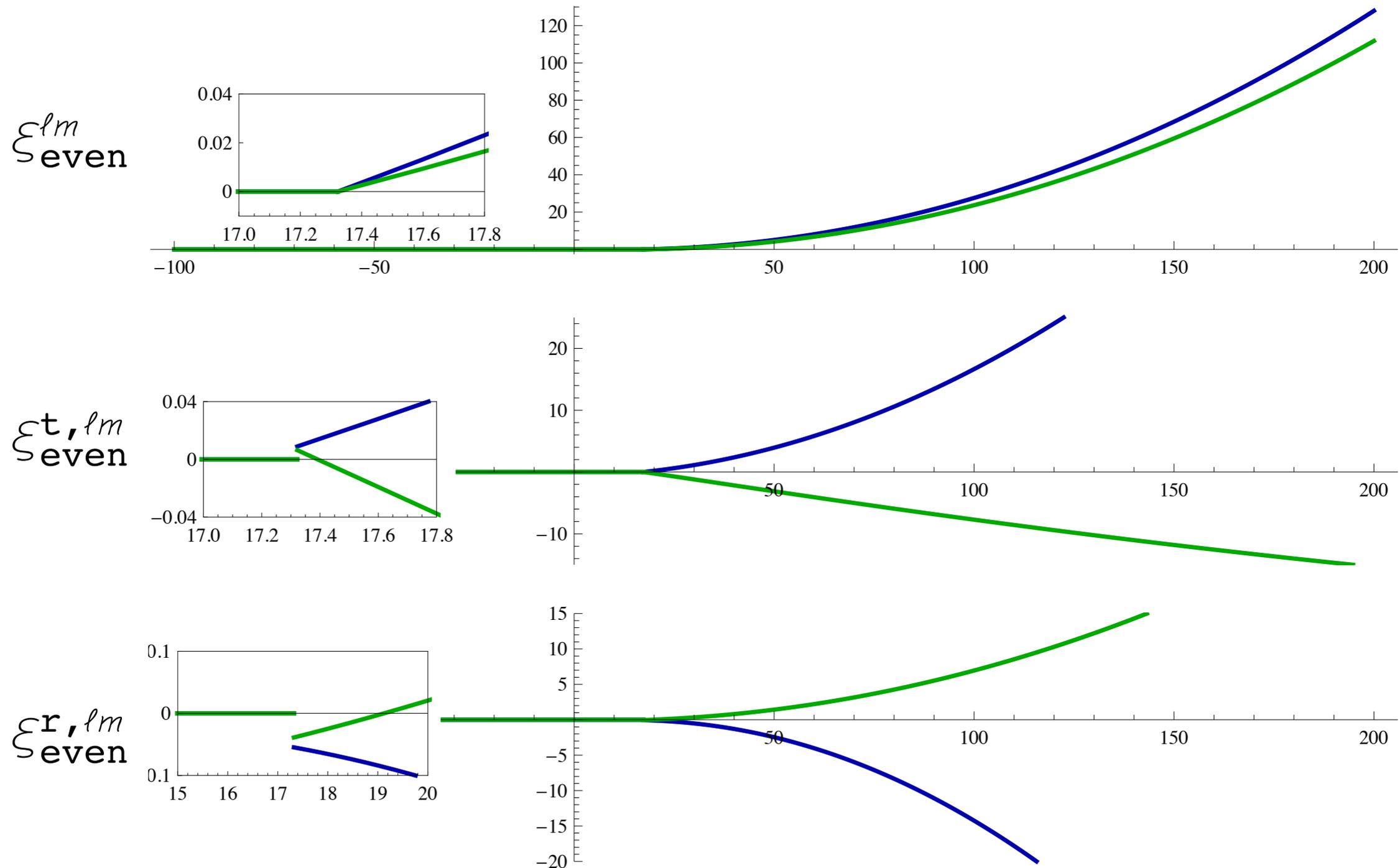
Updated gauge vector: odd-parity

$$(p, e, t_p) = (8.75455, 0.764124, 80.17) \quad (\ell, m) = (2, 1)$$



Updated gauge vector: even-parity

$$(p, e, t_p) = (8.75455, 0.764124, 80.17) \quad (\ell, m) = (2, 2)$$



And so on ...

- Jumps in RW amplitudes and jumps in Lorenz amplitudes yield new restrictions on gauge vector

And so on ...

- Jumps in RW amplitudes and jumps in Lorenz amplitudes yield new restrictions on gauge vector
- Will always disagree with Lorenz by:

$$p_{\mu\nu}^L = p_{\mu\nu}^{\text{MRW}} - 2\Xi_{(\mu|\nu)}^{\text{MRW} \rightarrow \text{L}}$$

And so on ...

- Jumps in RW amplitudes and jumps in Lorenz amplitudes yield new restrictions on gauge vector
- Will always disagree with Lorenz by:

$$p_{\mu\nu}^L = p_{\mu\nu}^{\text{MRW}} - 2\Xi_{(\mu|\nu)}^{\text{MRW}\rightarrow\text{L}}$$

- Modified RW gauge will have the same discontinuities as Lorenz gauge, to arbitrary orders of discontinuity

Where does this leave us?

- SF is well-defined if and only if $\delta F_{\text{self}}^{\alpha}$ is also
- If vector Ξ^{μ} is well-defined, the SF will be also

Where does this leave us?

- SF is well-defined if and only if $\delta F_{\text{self}}^\alpha$ is also
- If vector Ξ^μ is well-defined, the SF will be also
- We can make the gauge vector $\Xi_{\text{MRW} \rightarrow \text{L}}^\mu$ as smooth as necessary

Preliminary results

Preliminary results

- Dissipative SF yields local E and J losses which agree with fluxes

Preliminary results

- Dissipative SF yields local E and J losses which agree with fluxes
- Working on the conservative SF
 - Want same value on both sides of particle
 - If not, why not?

Preliminary results

- Dissipative SF yields local E and J losses which agree with fluxes
- Working on the conservative SF
 - Want same value on both sides of particle
 - If not, why not?
- Non-radiative modes should follow from a “Modified Zerilli gauge”

Preliminary results

- Dissipative SF yields local E and J losses which agree with fluxes
- Working on the conservative SF
 - Want same value on both sides of particle
 - If not, why not?
- Non-radiative modes should follow from a “Modified Zerilli gauge”
- Solving field equations in Mathematica yields high (theoretically arbitrary) accuracy

Conclusions

Conclusions

- Regge-Wheeler gauge is very convenient for solving the field equations on Schwarzschild

Conclusions

- Regge-Wheeler gauge is very convenient for solving the field equations on Schwarzschild
- Local singularities make a “well-defined” self-force impossible in this gauge

Conclusions

- Regge-Wheeler gauge is very convenient for solving the field equations on Schwarzschild
- Local singularities make a “well-defined” self-force impossible in this gauge
- Global gauge transformations to Lorenz are possible but difficult

Conclusions

- Regge-Wheeler gauge is very convenient for solving the field equations on Schwarzschild
- Local singularities make a “well-defined” self-force impossible in this gauge
- Global gauge transformations to Lorenz are possible but difficult
- Modified RW gauge (hopefully) yields a way to find the self-force with no extra computational cost