エキセントリックコンパクトバイナリの軌道力学

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Averaged redshift

First law and applications 0000000

Outline

① Gravitational wave source modelling

^② Averaged redshift for eccentric orbits

③ First law of mechanics and applications

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Source modelling of compact binaries



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Why?

- Independent checks of long and complicated calculations
- Identify domains of validity of approximation schemes
- Extract information inaccessible to other methods
- Develop a universal model for compact binaries

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What?

- Gravitational waveforms at future null infinity
- Conservative effects on the orbital dynamics

Paper	Year	Methods	Observable	Orbit	Spin
Detweiler	2008	SF/PN	redshift observable		
Blanchet et al.	2010	SF/PN	redshift observable		
Damour	2010	SF/EOB	ISCO frequency		
Mroué et al.	2010	NR/PN	periastron advance		
Barack et al.	2010	SF/EOB	periastron advance		
Favata	2011	SF/PN/EOB	ISCO frequency		
Le Tiec et al.	2011	NR/SF/PN/EOB	periastron advance		
Damour et al.	2012	NR/EOB	binding energy		
Le Tiec et al.	2012	NR/SF/PN/EOB	binding energy		
Akcay et al.	2012	SF/EOB	redshift observable		
Hinderer et al.	2013	NR/EOB	periastron advance		1
Le Tiec et al.	2013	NR/SF/PN	periastron advance		1
Damour et al.	2014	NR/PN/EOB	scattering angle	hyperbolic	
Bini, Damour Shah et al. Blanchet et al }	2014	SF/PN	redshift observable		
Dolan et al. Bini, Damour }	2014	SF/PN	precession angle		1
lsoyama et al. Akcay et al.	2014 2015	SF/PN/EOB SF/PN	ISCO frequency averaged redshift	eccentric	1

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First law and applications

Post-Newtonian expansions and black hole perturbations



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Redshift invariant for circular orbits

• It measures the redshift of light emitted from the point particle [Detweiler 2008]

$$\frac{\mathcal{E}_{obs}}{\mathcal{E}_{em}} = \frac{(p^a u_a)_{obs}}{(p^a u_a)_{em}} = z$$

 It is a constant of the motion associated with the helical Killing field k^a:

$$z = -k^a u_a$$

In coordinates adapted to the symmetry:

$$z = \frac{\mathrm{d}\tau}{\mathrm{d}t} = \frac{1}{u^t}$$





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Averaged redshift for eccentric orbits

• Generic eccentric orbit parameterized by the two invariant frequencies

$$m = \frac{2\pi}{P}, \quad \omega = \frac{\Phi}{P}$$

• Time average of $z = d\tau/dt$ over one radial period [Barack & Sago 2010]

$$\langle z \rangle \equiv \frac{1}{P} \int_0^P z(t) \, \mathrm{d}t = \frac{T}{P}$$

• Coordinate-invariant relation $\langle z \rangle (n, \omega)$ is well defined in GSF and PN frameworks



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	Coeff.	Exact value	Fitted value	Fitted value
		[Akcay et al. 2015]	[Akcay et al. 2015]	[Meent, Shah 2015]
	e^2	4	4.0002(8)	$4\pm6\times10^{-12}$
1PN	e^4	-2	-2.00(1)	$-2\pm4\times10^{-10}$
	e^{6}	0		$0 \pm 4 \times 10^{-9}$

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2PN	e^4	1/4		$^{1/4} \pm 4 \times 10^{-7}$
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	e^2	$-14.312097\cdots$	-14.5(4)	-14.3120980(5)
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New coefficients at 4PN and 5PN orders [van de Meent, Shah 2015]

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First law of binary mechanics

- Canonical ADM Hamiltonian H of two point masses m_a
- Variation δH + Hamilton's equation + orbital averaging:

$$\delta M = \omega \, \delta L + n \, \delta R + \sum_{a} \langle z_a \rangle \, \delta m_a$$

• First integral associated with the variational first law:

$$M = 2\left(\omega L + nR\right) + \sum_{a} \langle z_{a} \rangle m_{a}$$

These relations are satisfied up to at least 3PN order

Applications of the first law

- Conservative dynamics beyond the geodesic approximation
- Shift of the Schwarzschild separatrix and singular curve
- Calibration of EOB potentials for generic bound orbits

$$\frac{\partial M}{\partial m_1} = \langle z \rangle - \omega \frac{\partial \langle z \rangle}{\partial \omega} - n \frac{\partial \langle z \rangle}{\partial n}$$
$$\frac{\partial L}{\partial m_1} = -\frac{\partial \langle z \rangle}{\partial \omega}$$
$$\frac{\partial R}{\partial m_1} = -\frac{\partial \langle z \rangle}{\partial n}$$

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Schwarzschild separatrix



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Schwarzschild separatrix



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Shift of the Schwarzschild separatrix

• Separatrix $\omega = \omega_{sep}(e)$ characterized by the condition

n = 0

 GSF-induced shift of Schwarzschild ISCO frequency [Barack & Sago 2009; Le Tiec et al. 2012; Akcay et al. 2012]

$$\frac{\Delta\omega_{\rm isco}}{\omega_{\rm isco}} = 1.2101539(4) \, q$$

- GSF-induced shift of Schwarzschild IBSO frequency ?
- $\mathcal{O}(q)$ shift in $\omega = \omega_{sep}(e)$ controlled by $\langle z \rangle_{GSF}(n, \omega)$

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Schwarzschild singular curve



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Schwarzschild singular curve



Shift of the Schwarzschild singular curve

• Singular curve $\omega = \omega_{sing}(n)$ characterized by condition

$$\left|\frac{\partial(n,\omega)}{\partial(M,L)}\right| = 0$$

• In the test-particle limit $q \rightarrow 0$ this is equivalent to

$$\left[\left(\partial_{n\omega}^{2}\langle z\rangle\right)^{2}-\partial_{n}^{2}\langle z\rangle\,\partial_{\omega}^{2}\langle z\rangle\right]^{-1}=0$$

• $\mathcal{O}(q)$ shift in $\omega = \omega_{sing}(n)$ controlled by $\langle z \rangle_{GSF}(n, \omega)$

First law and applications

EOB dynamics beyond circular motion



Conservative EOB dynamics determined by "potentials"

$$A = 1 - 2u + \nu a(u) + \mathcal{O}(\nu^2)$$
$$\bar{D} = 1 + \nu \bar{d}(u) + \mathcal{O}(\nu^2)$$
$$Q = \nu q(u) p_r^4 + \mathcal{O}(\nu^2)$$

• Functions a(u), $\overline{d}(u)$ and q(u) controlled by $\langle z \rangle_{GSF}(n, \omega)$

Summary

- GSF/PN comparison for eccentric orbits relying on $\langle z \rangle (n, \omega)$
- First law of mechanics can be extended to eccentric orbits
- Numerous applications of the first law:
 - Conservative dynamics beyond the geodesic approximation
 - Shift of the Schwarzschild separatrix and singular curve
 - Calibration of EOB potentials for generic bound orbits

o ...

Prospects

- GSF/PN comparison for eccentric orbits relying on $\langle \psi \rangle(n,\omega)$
- Extension of the first law to precessing spinning binaries