THE ASTROPHYSICS OF EMRIS

Capture of compact objects by SMBHs

Pau Amaro Seoane

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Max-Planck Institute for Gravitational Physics (Albert Einstein Institute)

CAPTURE OF COMPACT OBJECTS

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- \cdot Stellar mass object spiraling into $10^4 10^6 \, M_{ullet}$
- · This range of masses corresponds to relaxed nuclei (!)
- · Only compact objects (extended stars disrupted early)
- $\cdot\,$ With eLISA stellar BH z $\gtrsim 0.7$

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- ▷ Bridge between astrophysics and GR

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- Many years before launch we're making new discoveries
- In this talk we'll see some of these difficulties, and how we've made progress: Microphysics around SMBHs



[PAS, LRR 2015]

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DISTRIBUTION OF STARS AROUND SMBHS

THE THREE REALMS OF STELLAR DYNAMICS



[PAS, LRR 2015]



Oth question: How many stars? How do they distribute?

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- Oth question: How many stars? How do they distribute?
- Very few observations R_h difficult to resolve
- To study inner region have to assume underlying population, deproject observation, assume observed star is tracing invisible population
 - Considerable amount of modelling: Are these profiles a coincidence?

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- ▷ If single-mass: quasi-steady solution takes power-law form (isotropic DF) $f(E) \sim E^p$, $\rho(r) \sim r^{-\gamma}$, with $\gamma = 3/2 + p$

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- ▷ If single-mass: quasi-steady solution takes power-law form (isotropic DF) $f(E) \sim E^p$, $\rho(r) \sim r^{-\gamma}$, with $\gamma = 3/2 + p$
- ▷ Confirmed later with a detailed kinematic treatment for single-mass [Bahcall & Wolf 1976]: $\gamma = 7/4$ and $p = \gamma 3/2 = 1/4$

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- More realistic models: Properties of multi-mass systems poorly reproduced by single-mass models
- ▷ Initial Mass Functions ∈ [0.1, ~ 120]M_☉ to first order by two (well-separated) mass scales: O(1M_☉) (Main Sequence, White Dwarfs, Neutron Stars) and O(10M_☉) (Stellar Black Holes)

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- Two branches for the solution: A "weak" (unrealistic) branch and a "strong" branch

[Hopman & Alexander 2009, Preto & Amaro-Seoane 2010, Amaro-Seoane & Preto 2011]

CUSPS IN DISTRESS



Deficit of old stars based on number counts of spectroscopically identified, old stars in sub-parsec SgrA* (down to magnitude K = 15.5)

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- ▷ Possibility of a core with ρ_{\star} decreasing, $\gamma < 0$
- Observers only see essentially late-type giants: Detectable stars are still a small fraction

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- Must invoke unlikely events to get rid of it
- Let's play the game What is the time necessary for cusp growth if at some point a central core is carved?
- We have now the correct, more efficient, solution of mass segregation



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- Our results confirmed later

[Gualandris & Merritt 2011]

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- The Milky Way nucleus is not necessarily the prototype of the nucleus from which e-LISA detections will be more frequent
- We still expect that a substantial fraction of EMRI events will originate from segregated stellar cusps, in particular with our new solution of mass segregation

EVENT RATES



DISGUISED CAPTURES

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- Plunges are more frequent than "adiabatic" EMRIS A common result to all event rate estimates
- What if these stars did not plunge? We'd have extremely eccentric sources, and event rates orders of magnitude larger

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- ▷ Calculate time to go from apo to periapsis and back (radial periode) and thus the change in (E, L_z, C) and so the new constants of motion, therefore: (p_{new}, e_{new}, i_{new})

\mathcal{M}_{ullet}	Spin (a/M)	a_0 (pc)	e_i	i (rad)	$\tau_{mrg~(yrs)}$	$\tau_{\mathrm{e-LISA}}$	Peri (e-LISA)
3E6	0.990	8.6182E-4	0.9990	0.6	2.6755E3	6.8409E2	432503
1E6	0.990	2.8727E-4	0.9990	0.6	2.9743E2	1.1915E2	146074
1E6	0.500	2.8727E-4	0.9990	0.6	2.4714E2	9.8328E1	97715
3E6	0.500	8.6182E-4	0.9990	0.6	2.2229E3	5.6105E2	288372
1E6	0.900	2.3939E-4	0.9990	0.2	1.5328E2	6.8038E1	90555
3E6	0.900	7.1818E-4	0.9990	0.2	1.3785E3	3.9237E2	268423
3E6	0.900	7.1786E-3	0.9999	0.2	4.6101E3	3.9131E2	267802
3E6	0.900	5.7429E-3	0.9999	0.2	2.0757E3	1.9956E2	149747
3E6	0.900	5.0250E-3	0.9999	0.2	1.3164E3	1.3607E2	106563
1E6	0.900	1.6750E-3	0.9999	0.2	1.4843E2	2.3449E1	35889
1E6	0.900	1.4357E-3	0.9999	0.2	9.1260E1	1.5533E1	24593
1E6	0.900	1.4357E-3	0.9999	0.1	9.2711E1	1.5769E1	25038
3E6	0.900	4.3071E-3	0.9999	0.1	8.1857E2	9.1641E1	74371
5E6	0.900	7.1786E-3	0.9999	0.1	2.2652E3	2.0548E2	122993
1E6	0.900	1.4357E-3	0.9999	0.1	1.8272E2	3.1556E1	50075
4E6	0.700	6.7000E-3	0.9999	0	1.8937E3	1.7207E2	96284
4E6	0.998	6.7000E-3	0.9999	0	2.6993E3	2.4753E2	170494
4E6	0.998	9.5714E-3	0.9999	0	8.7952E3	6.6162E2	395248
4E6	0.998	7.6571E-3	0.9999	0	4.1097E3	3.5062E2	230973
4E6	0.998	6.7000E-3	0.9999	0	2.6993E3	2.4753E2	170494
4E6	0.998	5.7429E-3	0.9999	0	1.7598E3	1.7468E2	123868
4E6	0.998	5.7429E-3	0.9999	0.3	1.6574E3	1.6506E2	117974

Note: Prograde orbits, $m_{ullet}~=~10~M_{\odot}$

A FAMILY OF SEPARATRICES: s = 0.1





a (pc)

22

A FAMILY OF SEPARATRICES: s = 0.4





A family of separatrices: $\ensuremath{\mathsf{s}}=0.7$

s = 0.7



A family of separatrices: s = 0.99

s = 0.99



A family of separatrices: s = 0.999

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IMPACT OF THE SPIN ON THE RATES?



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IT'S ALL ABOUT AN UPPER LIMIT



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 $\boldsymbol{\Omega}_{\rm EMRI}^{\rm Kerr} = \boldsymbol{\Omega}_{\rm EMRI}^{\rm Schw} \times \boldsymbol{\mathcal{W}}^{\frac{-5}{6-2\gamma}}(\boldsymbol{\iota},\,\boldsymbol{S})$

$$\dot{N}_{\mathrm{EMRI}}^{\mathrm{Kerr}} = \dot{N}_{\mathrm{EMRI}}^{\mathrm{Schw}} imes \mathcal{W}^{rac{20\gamma - 45}{12 - 4\gamma}}(\iota, S)$$

 \triangleright Take a typical value of a prograde orbit with high spin: $\mathcal{W} = 0.15$; then for a modest $\gamma = 1.5$

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 \triangleright When taking into account spinning MBHs EMRI rates are boosted

THE BUTTERFLY EFFECT

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- \triangleright Assume power-law of $R^{-\gamma}$
- ⊳ Then

$$M(R) = \int_0^R 4\pi r^2 \rho(r) \, dr \propto \int_0^R r^{-\gamma+2} dr \propto R^{3-\gamma}$$
$$N(R) \simeq 8.6 \times 10^4 \left(\frac{R}{6 \times 10^{-4} \text{ pc}}\right)^{3-\gamma}$$
$$R_1 \simeq 6 \times 10^{-4} \text{ pc} \times \left(\frac{1}{8.6 \times 10^4}\right)^{\frac{1}{3-\gamma}}$$

 \triangleright $R_1 \simeq 3 \times 10^{-7}$ pc for $\gamma = 1.5$



Fig. 11.—Evolution of the profiles of enclosed mass for GN25. The solid lines are the results of the MC simulation. For reference, the dashed lines show $\eta = 1.5$ profiles adjusted on the total mass and half-mass radius of each component. The top thin line is the total mass, including the central MBH; it is compared to the observational constrains for the MV center (see Fig. 3). [See the electronic edition of the Journal for a color version of this figure.]

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- $ightarrow \sim 15 \, M_{\odot}$ within $3 \times 10^{-4} \, \mathrm{pc}$ in our Milky Way G25 model

[Freitag, Amaro-Seoane & Kalogera 2006]



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▷ Watch out: I am cheating



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▷ $\mathcal{M}_{\bullet} = 10^{6} M_{\odot}$ ▷ $a_{\bullet,i} \simeq 1.45 \times 10^{-6} \text{ pc}$ (well within e-LISA) ▷ $m_{\bullet} = 10 M_{\odot}$ (also successfully tested 5, 1.44 M_{\odot})



 $\succ \mathcal{M}_{\bullet} = 10^{6} M_{\odot}$ $\succ a_{\bullet,i} \simeq 1.45 \times 10^{-6} \text{ pc (well within e-LISA)}$ $\succ m_{\bullet} = 10 M_{\odot} \text{ (also successfully tested 5, } 1.44 M_{\odot} \text{)}$ $\succ m_{\star} = 10 M_{\odot}, a_{\star,i} \simeq 4.1 \times 10^{-6} \text{ pc, } e_{\star,i} = 0.5, i_{\bullet,\star} = 30^{\circ}$



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- Evolution of the eccentricity when taking energy loss, i.e. 2.5 PN into account?







в,



Red $i_{\star} = 30^{\circ}$, green $i_{\star} = 30.001^{\circ}$, blue fiducial plus a *ten billionth* of a degree, $i_{\star} = 30.000000001^{\circ}$ and magenta plus a *ten trillionth* of a degree, $i_{\star} = 30.000000000001^{\circ}$

NO, IT'S NOT A BUG



 $a_{\star} = 4 \times 10^{-6}$ pc, 6×10^{-6} pc, 9×10^{-6} pc and 4.07243×10^{-5} pc

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- ▷ This is not a classical system
- ▶ How to characterise the chaos?



Start with a fiducial case and another one a *bit* different in phase space. Let them evolve. Calculate time for the "distance" to be $2 \times a_{\bullet}$ and divide it by the isolated inspiral time : Characteristic time

CONCLUSIONS

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- Originating in the bulk of the stellar system: Enhanced rates
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QUESTIONS?