# Self-force meets post-Newtonian theory and more...

Abhay Shah (University of Southampton)

with J Friedman, B Whiting, N Johnson McDaniel, A Nagar, A Pound, A Le Tiec

\*Special thanks to Cesar, Adam and Marta

# Overlap between self-force, post-Newtonian and EOB theories

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# PLAN OF THE TALK

## Part I

#### • PN coefficients of $\Delta U$ in Kerr

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Tuesday, June 30, 2015

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#### • PN coefficients of $\Delta U$ in Kerr Detweiler's redshift invariant

## Part II

### • Resummation of pN flux or waveforms

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•  $\Delta \lambda_{B/E}$  transcribe to the dynamically significant EOB description of tidal interaction energy.

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# Brief summary

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Perturbed Weyl scalars,  $\psi_0$  or  $\psi_4$ 

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Homogenous solutions of Teukolsky equation are written as a sum over known analytical functions Mano-Suzuki-Takasugi

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Circular orbits in Kerr

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$$\Delta U = \frac{-1}{R} + \frac{-2}{R^2} + \frac{-5}{R^3} + \frac{\left(\frac{-121}{3} + \frac{41}{32}\pi^2\right)}{R^4} + \cdots$$

If we subtracted the known PN-series from it

 $\Delta U - \left(\frac{-1}{R} + \frac{-2}{R^2} + \frac{-5}{R^3} + \frac{\left(\frac{-121}{3} + \frac{41}{32}\pi^2\right)}{R^4} + \frac{\frac{64}{5}\log(R)}{R^5} + \frac{\frac{956}{105}\log(R)}{R^6}\right)$ 

= -114.348951367572602952040002444836538764412 8652844070388692348480929255963692827665976343 7619372125523054160542189937045693826002042714 825538690979057075189... X 10<sup>-150</sup>

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# $$\begin{split} \gamma_7 &= 52.17523809523809523809523809523809523809\\ &\quad 523809523809523809523809523809523809\\ &\quad 5238095237538043489164331\cdots \end{split}$$

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$$523809523809523809523809523809523809$$
  
$$5238095237538043489164331\cdots$$

#### More than 11 repetition cycles

$$\gamma_7 = \frac{27392}{525}$$

# Extracting spin-dependent pN parameters by just staring at $\Delta U$ calculated for a given R and a.

For a circular, equatorial orbit in Kerr,  $\Delta U$  for  $R = 10^{20} \& a = 10^{-3}$  is

If we subtract the known *a*-independent PN-terms from it, we get

2.33333333333333333333348666666666666657111118811111211110470000333 61101449462298052019375307453520567434360997918586330437827422 85790789258060941125270718078251256276684354908208083322490189 27613456043301239146966332282381109537868453512458593272282243 42962251368467464109905674383013754100571040568288348387074195 38774336951906007694268879889268637496375852038699098746116349 07608445730228968534690145005883609409871989370801221523747813 79214778427338448055324558181161062657272652196860626477197489 998371250766911017733058814377904976565274243148 X 10<sup>-53</sup>

 $\Delta U - \Delta U_{\rm Schw PN}$ 

2.33333333333333333333348666666666666657111118811111211110470000333 61101449462298052019375307453520567434360997918586330437827422 85790789258060941125270718078251256276684354908208083322490189 27613456043301239146966332282381109537868453512458593272282243 42962251368467464109905674383013754100571040568288348387074195 38774336951906007694268879889268637496375852038699098746116349 07608445730228968534690145005883609409871989370801221523747813 79214778427338448055324558181161062657272652196860626477197489 998371250766911017733058814377904976565274243148 X 10<sup>-53</sup>



 $\Delta U - \Delta U_{\rm Schw PN}$ 



$$\Delta U - \Delta U_{\rm Schw PN} - \frac{\frac{7}{3}a}{R^{5/2}}$$

-0.9999984666666666666666676222214522222122222863332999722318838710 422315600274518955428356768059090185106991 X 10-66



$$\Delta U - \Delta U_{\rm Schw PN} - \frac{\frac{7}{3}a}{R^{5/2}}$$

-0.99999846666666666666666676222214522222122222863332999722318838710 35281313958025879812765898972335414747002895505910475425440752 72392208062615255082077056648978425125250010843144057198772900 32094186367001050952223795464879820874740061051089903710819648 65869223427658950319579232762292765044984946259137945589963814 27325639064453444064695836957481294634234587216984257248876031 04364798643188327449723923461343962532111809585519541185549059 94885278008775152172270676060681136472706856135843334962082566 422315600274518955428356768059090185106991 X 10<sup>-66</sup>

 $\frac{-a^2}{R^3}$ 

$$\Delta U - \Delta U_{\rm Schw PN} - \frac{\frac{7}{3}a}{R^{5/2}} - \frac{-a^2}{R^3}$$



$$\Delta U - \Delta U_{\rm Schw PN} - \frac{\frac{7}{3}a}{R^{5/2}} - \frac{-a^2}{R^3}$$

$$\frac{\frac{46}{3}a}{R^{7/2}}$$

$$\Delta U - \Delta U_{\text{Schw PN}} - \frac{\frac{7}{3}a}{R^{5/2}} - \frac{-a^2}{R^3} - \frac{\frac{46}{3}a}{R^{7/2}}$$

-9.55554785555545555619666633305565217204368614647291359213146 169010139242351844032438954958070 X 10<sup>-86</sup>



$$\Delta U - \Delta U_{\text{Schw PN}} - \frac{\frac{7}{3}a}{R^{5/2}} - \frac{-a^2}{R^3} - \frac{\frac{46}{3}a}{R^{7/2}}$$

-9.55554785555545555619666633305565217204368614647291359213146 169010139242351844032438954958070 X 10<sup>-86</sup>

$$\frac{\frac{-86}{9}a^2}{R^4}$$

Careful:

- Coefficient of  $\frac{a^i}{R^j}$  has to be of order unity or "so", and grow consistently with each pN order for a given choice of *i*.
- The *i* in  $\frac{a^i}{R^j}$  has to be a positive integer.
- After subtracting  $\frac{c_{i,j}a^i}{R^j}$ , one should see a "consistent" reduction in magnitude that would help you in predicting the next higher-order pN coefficient.



77.00000099999358889222499903383511869409082641963424094563232 163132037115231166005974855327176 X 10<sup>-93</sup>





77.00000099999358889222499903383511869409082641963424094563232 49886807475219326716311746796781469498300141596069671401451655 73243797096972211379078165023449321901280358552211712699984267 57342401347482161171132318511402573563529987945632719026429894 59929457177237275963084276632258407948965831577687781575263852 64740927587987635005237964973346191178574235790338947724982987 60878259690110412636702681036673162273369442134470700499515461 61541085749515366086378887260139655799906621947703266793865454 163132037115231166005974855327176 X 10<sup>-93</sup>

 $\frac{77\,a}{R^{9/2}}$ 

$$\Delta U - \Delta U_{\text{Schw PN}} - \frac{\frac{7}{3}a}{R^{5/2}} - \frac{-a^2}{R^3} - \frac{\frac{46}{3}a}{R^{7/2}} - \frac{\frac{-86}{9}a^2}{R^4} - \frac{77a}{R^{9/2}}$$

0.999993588892224999033835118694090826419634240945632324988680 371152311660059748553271763421 X 10<sup>-99</sup>



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$$\Delta U - \Delta U_{\text{Schw PN}} - \frac{\frac{7}{3}a}{R^{5/2}} - \frac{-a^2}{R^3} - \frac{\frac{46}{3}a}{R^{7/2}} - \frac{\frac{-86}{9}a^2}{R^4} - \frac{77a}{R^{9/2}} - \frac{a^3}{R^{9/2}}$$

-64.1110777500096616488130590917358036575905436767501131925247 7688339940251446728236578214 X 10<sup>-106</sup>



 $\Delta U - \Delta U_{\text{Schw PN}} - \frac{\frac{7}{3}a}{R^{5/2}} - \frac{-a^2}{R^3} - \frac{\frac{46}{3}a}{R^{7/2}} - \frac{\frac{-86}{9}a^2}{R^4} - \frac{77a}{R^{9/2}} - \frac{a^3}{R^{9/2}}$ 

-64.1110777500096616488130590917358036575905436767501131925247 80673283688253203218530501699858403930328598548344267562029030 27788620921834976550678098719641447788287300015732426575986525 17838828867681488597426436470012054367280973570105400705428227 62724036915723367741592051034168422312218424736147352590724120 12364994762035026653808821425764209661052275017012391217403098 89587363297318963326837726630557865529299500484538384589142504 84633913621112739860344200093378052296733206134545836867962884 7688339940251446728236578214 X 10<sup>-106</sup>

$$\frac{\frac{-577}{9}a^2}{R^5}$$



Polygamma functions enter at 6.5 PN order, i.e., at  $O\left(\frac{1}{R^{15/2}}\right)$ .

#### Resummation of fluxes or waveforms

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Find an analytical form for linear-in-mass-ratio part of the flux or waveform valid in all region.

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Resummation introduced by Damour-Iyer-Nagar:

Factorized multipolar resummation where each  $h_{l,m}$  is factored into five parts and each of these parts is resummed individually.

$$h_{\ell,m} = h_{\ell,m}^{N} \times S_{\ell,m} \times T_{\ell,m} \times e^{i\delta_{\ell,m}} \times \rho_{\ell,m}^{\ell}$$

$$\uparrow$$
Newtonian
contribution

$$= \frac{1}{R} n_{\ell,m}^{(p)} (-1)^{\ell+p} v^{\ell+p} Y_{\ell+p,m}$$

 $h_{\ell,m} = h_{\ell,m}^{\mathcal{N}} \times S_{\ell,m} \times T_{\ell,m} \times e^{i\delta_{\ell,m}} \times \rho_{\ell,m}^{\ell}$ Source,  $\mathcal{E}$  or  $\mathcal{J}$ 

 $h_{\ell,m} = h_{\ell,m}^{\mathcal{N}} \times S_{\ell,m} \times T_{\ell,m} \times e^{i\delta_{\ell,m}} \times \rho_{\ell,m}^{\ell}$ "Tail term"

$$\frac{\Gamma(\ell+1-2i\hat{k})}{\Gamma(\ell+1)} e^{\pi\hat{k}} e^{2i\hat{k}\ln(2\hat{k}c)} (\hat{k}=m\Omega)$$

 $h_{\ell,m} = h_{\ell,m}^{\mathcal{N}} \times S_{\ell,m} \times T_{\ell,m} \times e^{i\delta_{\ell,m}} \times \rho_{\ell,m}^{\ell}$   $\uparrow$ Dephasing term

 $h_{\ell,m} = h_{\ell,m}^{\mathcal{N}} \times S_{\ell,m} \times T_{\ell,m} \times e^{i\delta_{\ell,m}} \times \rho_{\ell,m}^{\ell}$ Whatever left

Motive  $\rightarrow$  Analytical  $\rho_{\ell,m}$  is valid everywhere.

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#### 



Proposal  $\rightarrow$  Hybridize it and impose agreement with pN.

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$$\rho_{22}^{\text{Resum}} = \frac{1 + \sum_{i=1}^{i_{\text{max}}} \frac{n_i}{R^i}}{1 + \sum_{j=1}^{j_{\text{max}}} \frac{d_j}{R^j}}$$

Choose  $n_i$  and  $d_j$  such that the Taylor-expanded  $\rho_{22}^{\text{Resum}}$  agree with pN-expansion of  $\rho_{22}$ . Left with some undetermined coefficients.

Fix these remaining coefficients,  $n_i$  and  $d_j$ , by imposing agreement with numerical  $\rho_{22}$ .





#### Doable work in the not-so-distant-future

- Calculate PN parameters of spin and tidal invariants in Kerr.
- Resummation of ingoing and outgoing fluxes in Kerr.