

Gauge invariant perturbations of Petrov type D spacetimes - II

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Schwarzschild spacetime

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Introduction

- Regge-Wheeler, Moncrief, Zerilli - they all used harmonic decomposition to find the equations (2nd order in 't' and 'r') their invariants/variables solve. All this is nicely summarized in Martel-Poisson.
- This will not be possible* to do in Kerr spacetime.
- A warm-up exercise to tackle Kerr - lets not do the harmonic decomposition in Schwarzschild, and derive the 2nd order equation (now a PDE, but, ofcourse, separable) that these or new invariants solve.

* to the best of my knowledge.

Introduction

- The invariants are divided into two categories - usually, spin odd-type, and spin even-type. We do the same, too.
- Odd-type invariants are relatively easy to calculate and are fewer in number. The equations they solve are also easier to derive.
- One of the odd-type invariants is the imaginary part of ψ_2
- The other one we know is its time derivative.
- Nothing simple exists for the even-parity Zerilli invariant.

Odd-type invariants

Imaginary part of ψ_2

$$\mathcal{I} = \mathbf{p} \bar{\partial}' h_{23} - \mathbf{p}' \bar{\partial}' h_{13} - \mathbf{p} \bar{\partial} h_{24} + \mathbf{p}' \bar{\partial} h_{14}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho\mathbf{p}' + \rho'\mathbf{p} - 6\rho\rho')(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} + \rho\bar{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho\mathbf{p}' + \rho'\mathbf{p})(\rho\mathbf{p}' + \rho'\mathbf{p}) - \frac{1}{2}(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho\mathbf{p}' - \rho'\mathbf{p}) \right. \\ & \quad \left. - 7\rho\rho'(\rho\mathbf{p}' + \rho'\mathbf{p}) - \rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = \mathbf{p} \vec{\partial}' h_{23} - \mathbf{p}' \vec{\partial}' h_{13} - \mathbf{p} \vec{\partial} h_{24} + \mathbf{p}' \vec{\partial} h_{14}$$

$$X_{(i)(j)} = X_{\alpha\beta} e_{(i)}^\alpha e_{(j)}^\beta$$

$$\begin{aligned} & 8\pi(\rho\mathbf{p}' + \\ & + 8\pi(\rho\mathbf{p}' + \\ & = \left(\frac{1}{2}(\rho\mathbf{p}' + \right. \\ & - 7\rho\rho'(\rho \left. \begin{array}{l} e_{(1)}^\alpha = l^\alpha \\ e_{(2)}^\alpha = n^\alpha \\ e_{(3)}^\alpha = m^\alpha \\ e_{(4)}^\alpha = \bar{m}^\alpha \end{array} \right) \rho\vec{\partial}' \mathcal{T}_{23} + \rho\vec{\partial} \mathcal{T}_{24}) \\ & \quad \left. \mathcal{T}_{23} - \rho\vec{\partial} \mathcal{T}_{24} \right) \\ & \quad \rho\mathbf{p}' - \rho'\mathbf{p}) \\ & \quad \left. \rho' - 12\psi_2 \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = \boxed{b} \bar{\partial}' h_{23} - b' \bar{\partial}' h_{13} - \boxed{b} \bar{\partial} h_{24} + b' \bar{\partial} h_{14}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho b' + \rho'b - 6\rho\rho')(\rho \sim l^\alpha \partial_\alpha - \rho' \bar{\partial} \mathcal{T}_{14} - \rho \bar{\partial}' \mathcal{T}_{23} + \rho \bar{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho b' - \rho'b)(\rho' \bar{\partial}' \mathcal{T}_{13} - \rho \bar{\partial}' \mathcal{T}_{14} + \rho \bar{\partial}' \mathcal{T}_{23} - \rho \bar{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho b' + \rho'b)(\rho b' + \rho'b) - \frac{1}{2}(\rho b' - \rho'b)(\rho b' - \rho'b) \right. \\ & \quad \left. - 7\rho\rho'(\rho b' + \rho'b) - \rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = b \bar{\partial}' h_{23} - \boxed{b'} \bar{\partial}' h_{13} - b \bar{\partial} h_{24} + \boxed{b'} \bar{\partial} h_{14}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho b' + \rho' b - 6\rho\rho')(\rho' \bar{\partial}' \boxed{n^\alpha \partial_\alpha} \mathcal{T}_{14} - \rho \bar{\partial}' \mathcal{T}_{23} + \rho \bar{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho b' - \rho' b)(\rho' \bar{\partial}' \mathcal{T}_{13} + \rho \bar{\partial}' \mathcal{T}_{23} - \rho \bar{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho b' + \rho' b)(\rho b' + \rho' b) - \frac{1}{2}(\rho b' - \rho' b)(\rho b' - \rho' b) \right. \\ & \quad \left. - 7\rho\rho'(\rho b' + \rho' b) - \rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = b \bar{\partial}' h_{23} - b' \bar{\partial}' h_{13} - b \boxed{\bar{\partial}} h_{24} + b' \boxed{\bar{\partial}} h_{14}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho b' + \rho' b - 6\rho\rho')(\rho' \bar{\partial}' \mathcal{T}_{13} - \rho' \bar{\partial} \mathcal{T}_{23} + \rho \bar{\partial}' \mathcal{T}_{24}) \\ & + 8\pi(\rho b' - \rho' b)(\rho' \bar{\partial}' \mathcal{T}_{13} - \rho' \bar{\partial} \mathcal{T}_{14} + \rho \bar{\partial}' \mathcal{T}_{23} - \rho \bar{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho b' + \rho' b)(\rho b' + \rho' b) - \frac{1}{2}(\rho b' - \rho' b)(\rho b' - \rho' b) \right. \\ & \quad \left. - 7\rho\rho'(\rho b' + \rho' b) - \rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = b \boxed{\partial' h_{23}} - b' \boxed{\partial' h_{13}} - b \partial h_{24} + b' \partial h_{14}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho b' + \rho' b - \bar{m}^\alpha \partial_\alpha) (\mathcal{T}_{13} - \rho' \partial \mathcal{T}_{14} - \rho \partial' \mathcal{T}_{23} + \rho \partial \mathcal{T}_{24}) \\ & + 8\pi(\rho b' - \rho' b) (\rho' \partial' \mathcal{T}_{13} - \rho' \partial \mathcal{T}_{14} + \rho \partial' \mathcal{T}_{23} - \rho \partial \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho b' + \rho' b)(\rho b' + \rho' b) - \frac{1}{2}(\rho b' - \rho' b)(\rho b' - \rho' b) \right. \\ & \quad \left. - 7\rho\rho'(\rho b' + \rho' b) - \rho\rho'(\partial \partial' + \partial' \partial - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = \mathbf{p} \vec{\partial}' h_{23} - \mathbf{p}' \vec{\partial}' h_{13} - \mathbf{p} \vec{\partial} h_{24} + \mathbf{p}' \vec{\partial} h_{14}$$

Eq. $\sim \partial_{r_*} \partial_{r_*}$ satisfies is

$$\begin{aligned} & 8\pi(\rho\mathbf{p}' + \rho'\mathbf{p} - 6\rho\rho')(\rho'\vec{\partial}' \mathcal{T}_{13} - \rho'\vec{\partial} \mathcal{T}_{14} - \rho\vec{\partial}' \mathcal{T}_{23} + \rho\vec{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho'\vec{\partial}' \mathcal{T}_{13} - \rho'\vec{\partial} \mathcal{T}_{14} + \rho\vec{\partial}' \mathcal{T}_{23} - \rho\vec{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}[(\rho\mathbf{p}' + \rho'\mathbf{p})(\rho\mathbf{p}' + \rho'\mathbf{p})] - \frac{1}{2}(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho\mathbf{p}' - \rho'\mathbf{p}) \right. \\ & \quad \left. - 7\rho\rho'(\rho\mathbf{p}' + \rho'\mathbf{p}) - \rho\rho'(\vec{\partial}\vec{\partial}' + \vec{\partial}'\vec{\partial} - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = \mathbf{b}' \bar{\partial}' h_{23} - \mathbf{b}' \bar{\partial}' h_{13} - \mathbf{b} \bar{\partial} h_{24} + \mathbf{b}' \bar{\partial} h_{14}$$

Eq $\sim \partial_t \partial_t$ satisfies is

$$\begin{aligned} & 8\pi(\rho\mathbf{b}' + \rho'\mathbf{b} - 6\rho\rho')(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} + \rho\bar{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho\mathbf{b}' - \rho'\mathbf{b})(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho\mathbf{b}' + \rho'\mathbf{b})(\rho\mathbf{b}' + \rho'\mathbf{b}) - \frac{1}{2}(\rho\mathbf{b}' - \rho'\mathbf{b})(\rho\mathbf{b}' - \rho'\mathbf{b}) \right. \\ & \quad \left. - 7\rho\rho'(\rho\mathbf{b}' + \rho'\mathbf{b}) - \rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Preliminaries

Imaginary part of ψ_2

$$\mathcal{I} = \mathbf{p} \vec{\partial}' h_{23} - \mathbf{p}' \vec{\partial}' h_{13} - \mathbf{p} \vec{\partial} h_{24} + \mathbf{p}' \vec{\partial} h_{14}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho\mathbf{p}' + \boxed{\sim \ell(\ell+1)} \mathcal{T}_{13} - \rho'\vec{\partial} \mathcal{T}_{14} - \rho\vec{\partial}' \mathcal{T}_{23} + \rho\vec{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho'\vec{\partial}' \mathcal{T}_{13} - \rho'\vec{\partial} \mathcal{T}_{14} + \rho\vec{\partial}' \mathcal{T}_{23} - \rho\vec{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho\mathbf{p}' + \rho'\mathbf{p})(\rho\mathbf{p}' + \rho'\mathbf{p}) - \frac{1}{2}(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho\mathbf{p}' - \rho'\mathbf{p}) \right. \\ & \quad \left. - 7\rho\rho'(\rho\mathbf{p}' + \rho'\mathbf{p}) - \rho\rho' \boxed{(\vec{\partial}\vec{\partial}' + \vec{\partial}'\vec{\partial} - 12\rho\rho' - 12\psi_2)} \right) \mathcal{I} \end{aligned}$$

Odd-type invariants

Imaginary part of ψ_2

$$\mathcal{I} = \mathbf{p} \bar{\partial}' h_{23} - \mathbf{p}' \bar{\partial}' h_{13} - \mathbf{p} \bar{\partial} h_{24} + \mathbf{p}' \bar{\partial} h_{14}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho\mathbf{p}' + \rho'\mathbf{p} - 6\rho\rho')(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} + \rho\bar{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\ & = \left(\frac{1}{2}(\rho\mathbf{p}' + \rho'\mathbf{p})(\rho\mathbf{p}' + \rho'\mathbf{p}) - \frac{1}{2}(\rho\mathbf{p}' - \rho'\mathbf{p})(\rho\mathbf{p}' - \rho'\mathbf{p}) \right. \\ & \quad \left. - 7\rho\rho'(\rho\mathbf{p}' + \rho'\mathbf{p}) - \rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 12\rho\rho' - 12\psi_2) \right) \mathcal{I} \end{aligned}$$

Odd-type invariants

Time derivative of the imaginary part of ψ_2

$$\boxed{\mathcal{I}_t} = (\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2) (\rho'\bar{\partial}'h_{13} - \rho'\bar{\partial}h_{14} + \rho\bar{\partial}'h_{23} - \rho\bar{\partial}h_{24}) \\ - (\rho\bar{p}' + \rho'\bar{p} - 4\rho\rho') (\bar{\partial}'\bar{\partial}'h_{33} - \bar{\partial}\bar{\partial}h_{44})$$

ψ_{RW}
in Bernard's talk

The relation between the two

$$\mathcal{I}_t = (\rho\bar{p}' - \rho'\bar{p})\mathcal{I} + 2\rho\bar{\partial}' \mathcal{E}_{23} - 2\rho\bar{\partial} \mathcal{E}_{24} + 2\rho'\bar{\partial}' \mathcal{E}_{13} - 2\rho'\bar{\partial} \mathcal{E}_{14}$$

Odd-type invariants

Time derivative of the imaginary part of ψ_2

$$\begin{aligned} \mathcal{I}_t = & (\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2) (\rho'\bar{\partial}'h_{13} - \rho'\bar{\partial}h_{14} + \rho\bar{\partial}'h_{23} - \rho\bar{\partial}h_{24}) \\ & - (\rho\bar{b}' + \rho'\bar{b} - 4\rho\rho') (\bar{\partial}'\bar{\partial}'h_{33} - \bar{\partial}\bar{\partial}h_{44}) \end{aligned}$$

Equation it satisfies is

$$\begin{aligned} & 8\pi(\rho\bar{b}' + \rho'\bar{b} - 8\rho\rho')(\rho\bar{b}' - \rho'\bar{b})(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} + \rho\bar{\partial} \mathcal{T}_{24}) \\ & + 8\pi(\rho\bar{b}' + \rho'\bar{b} - 18\rho\rho')(\rho\bar{b}' + \rho'\bar{b})(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\ & - 16\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 24\rho\rho' - 14\psi_2)(\rho'\bar{\partial}' \mathcal{T}_{13} - \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\ = & \left(\frac{1}{2}(\rho\bar{b}' + \rho'\bar{b})(\rho\bar{b}' + \rho'\bar{b}) - \frac{1}{2}(\rho\bar{b}' - \rho'\bar{b})(\rho\bar{b}' - \rho'\bar{b}) \right. \\ & \left. - 9\rho\rho'(\rho\bar{b}' + \rho'\bar{b}) - \rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 24\rho\rho' - 14\psi_2) \right) \mathcal{I}_t \end{aligned}$$

Even-type invariants

Zerilli's even-parity invariant

$$\begin{aligned}\mathcal{I}_z = & (\eth\eth' + \eth'\eth - 4\rho\rho' - 4\psi_2) (\eth\eth' + \eth'\eth) (\rho b' - \rho'b) h_{34} \\ & + (\eth\eth' + \eth'\eth - 4\rho\rho' - 4\psi_2) (\eth\eth' + \eth'\eth) \left(\rho'^2 h_{11} - \rho^2 h_{22}\right) \\ & - (\eth\eth' + \eth'\eth - 4\rho\rho' + 2\psi_2) (\rho b' - \rho'b) (\eth'\eth'h_{33} + \eth\eth'h_{44}) \\ & + 2(2\rho\rho' - \psi_2) (\eth\eth' + \eth'\eth - 4\rho\rho' - 4\psi_2) (\rho'\eth'h_{13} + \rho'\eth'h_{14} - \rho\eth'h_{23} - \rho\eth'h_{24}) \\ & - 2\rho\rho' (\eth\eth' + \eth'\eth - 4\rho\rho' - 4\psi_2) (b'\eth'h_{13} + b'\eth'h_{14} - b\eth'h_{23} - b\eth'h_{24})\end{aligned}$$

5th order

Even-type invariants

$$\begin{aligned}
& 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' + \rho'\bar{b}) \\
& \times (\rho'\bar{\partial}' \mathcal{T}_{13} + \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\
& - 2\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' + \rho'\bar{b}) \\
& \times (\rho'^2 \mathcal{T}_{11} - \rho^2 \mathcal{T}_{22}) \\
& - 2\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' + \rho'\bar{b}) \\
& \times (\rho'^2 \mathcal{T}_{11} + 2\rho\rho' \mathcal{T}_{12} + \rho^2 \mathcal{T}_{22}) \\
& + 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^3(\rho\bar{b}' + \rho'\bar{b})(\bar{\partial}'\bar{\partial}' \mathcal{T}_{33} + \bar{\partial}\bar{\partial} \mathcal{T}_{44}) \\
& + 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' - \rho'\bar{b}) \\
& \times (\rho'\bar{\partial}' \mathcal{T}_{13} + \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} + \rho\bar{\partial} \mathcal{T}_{24}) \\
& + 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2) \\
& \times (\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 8\psi_2)(\rho'^2 \mathcal{T}_{11} - \rho^2 \mathcal{T}_{22}) \\
& + 4\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2) \\
& \times \left((\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})^2 + 2(\psi_2 + 8\rho\rho')(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial}) - 8(\psi_2^2 + 7\rho\rho'\psi_2 - 6\rho^2\rho'^2) \right) \\
& \times (\rho'\bar{\partial}' \mathcal{T}_{13} + \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\
& = \\
& \frac{1}{4}(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2[(\rho\bar{b}' + \rho'\bar{b})^2 - (\rho\bar{b}' - \rho'\bar{b})^2] \mathcal{I}_z \\
& - \frac{1}{2}\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)(13(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial}) + 38\psi_2 - 52\rho\rho') \mathcal{I}_z \\
& - \frac{1}{2}\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2) \\
& \times ((\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})^2 - 16(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\psi_2 + 4\rho\rho') - 24(\psi_2 - \rho\rho')(\psi_2 + 10\rho\rho')) \mathcal{I}_z
\end{aligned}$$

Even-type invariants

$$\begin{aligned}
& 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' + \rho'\bar{b}) \\
& \times (\rho'\bar{\partial}' \mathcal{T}_{13} + \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\
& - 2\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' + \rho'\bar{b}) \\
& \times (\rho'^2 \mathcal{T}_{11} - \rho^2 \mathcal{T}_{22}) \\
& - 2\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' + \rho'\bar{b}) \\
& \times (\rho'^2 \mathcal{T}_{11} + 2\rho\rho' \mathcal{T}_{12} + \rho^2 \mathcal{T}_{22}) \\
& + 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^3(\rho\bar{b}' + \rho'\bar{b})(\bar{\partial}'\bar{\partial}' \mathcal{T}_{33} + \bar{\partial}\bar{\partial} \mathcal{T}_{44}) \\
& + 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\rho\bar{b}' - \rho'\bar{b}) \\
& \times (\rho'\bar{\partial}' \mathcal{T}_{13} + \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} + \rho\bar{\partial} \mathcal{T}_{24}) \\
& + 8\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2) \\
& \times (\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 8\psi_2)(\rho'^2 \mathcal{T}_{11} - \rho^2 \mathcal{T}_{22}) \\
& + 4\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2) \\
& \times \left((\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})^2 + 2(\psi_2 + 8\rho\rho')(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial}) - 8(\psi_2^2 + 7\rho\rho'\psi_2 - 6\rho^2\rho'^2) \right) \\
& \times (\rho'\bar{\partial}' \mathcal{T}_{13} + \rho'\bar{\partial} \mathcal{T}_{14} - \rho\bar{\partial}' \mathcal{T}_{23} - \rho\bar{\partial} \mathcal{T}_{24}) \\
& = \\
& \frac{1}{4}(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2 \left[(\rho\bar{b}' + \rho'\bar{b})^2 - (\rho\bar{b}' - \rho'\bar{b})^2 \right] \mathcal{I}_z \\
& - \frac{1}{2}\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)(13(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial}) + 38\psi_2 - 52\rho\rho') \mathcal{I}_z \\
& - \frac{1}{2}\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2) \\
& \times ((\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})^2 - 16(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(\psi_2 + 4\rho\rho') - 24(\psi_2 - \rho\rho')(\psi_2 + 10\rho\rho')) \mathcal{I}_z
\end{aligned}$$

Even-type invariants

New even-parity invariant

$$\begin{aligned}\mathcal{I}_{E1} = & (\eth\bar{\eth}' + \bar{\eth}'\eth - 4\rho\rho' - 4\psi_2) \left(\rho'^2 h_{11} + 2\rho\rho' h_{12} + \rho^2 h_{22} + \rho'\bar{\eth}'h_{13} + \rho'\eth h_{14} + \rho\bar{\eth}'h_{23} + \rho\bar{\eth}h_{24} \right. \\ & \left. - \left[\rho\mathbb{P}' + \rho'\mathbb{P} + \frac{1}{2}\eth\bar{\eth}' + \frac{1}{2}\bar{\eth}'\eth - 2\rho\rho' + \psi_2 \right] h_{34} + \frac{1}{2}\bar{\eth}'\bar{\eth}'h_{33} + \frac{1}{2}\bar{\eth}\bar{\eth}h_{44} \right) + 3\psi_2 (\bar{\eth}'\bar{\eth}'h_{33} + \bar{\eth}\bar{\eth}h_{44})\end{aligned}$$

4th order

Even-type invariants

$$\begin{aligned}
& 8\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial}) \\
& \quad \times (\rho'^2 \mathcal{T}_{11} + 2\rho\rho' \mathcal{T}_{12} + \rho^2 \mathcal{T}_{22}) \\
& - 32\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)^2(\bar{\partial}'\bar{\partial}' \mathcal{T}_{33} + \bar{\partial}\bar{\partial} \mathcal{T}_{44}) \\
& - 16\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\rho\bar{b}' - \rho'\bar{b}) \\
& \quad \times (\rho'^2 \mathcal{T}_{11} - \rho^2 \mathcal{T}_{22}) \\
& - 32\pi\rho\rho'(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2) \\
& \quad \times (\rho'\bar{\partial}' \mathcal{T}_{13} + \rho'\bar{\partial} \mathcal{T}_{14} + \rho\bar{\partial}' \mathcal{T}_{23} + \rho\bar{\partial} \mathcal{T}_{24}) \\
& - 16\pi(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\rho\bar{b}' + \rho'\bar{b}) \\
& \quad \times (\rho'^2 \mathcal{T}_{11} - 2\rho\rho' \mathcal{T}_{12} - \rho^2 \mathcal{T}_{22}) \\
& - 8\pi[(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})^2 + (\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(2\psi_2 - 20\rho\rho') + 64\rho^2\rho'^2 - 80\psi_2\rho\rho'] \\
& \quad \times (\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' - 4\psi_2)(\rho'^2 \mathcal{T}_{11} - 2\rho\rho' \mathcal{T}_{12} - \rho^2 \mathcal{T}_{22}) \\
& = \\
& 2(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial} - 4\rho\rho' + 2\psi_2)[(\rho\bar{b}' + \rho'\bar{b})^2 - (\rho\bar{b}' - \rho'\bar{b})^2] \mathcal{I}_{E1} \\
& - 4\rho\rho'(11\bar{\partial}\bar{\partial}' + 11\bar{\partial}'\bar{\partial} - 44\rho\rho' + 34\psi_2)(\rho\bar{b}' + \rho'\bar{b}) \mathcal{I}_{E1} \\
& - 4\rho\rho'[(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})^2 - 2(\bar{\partial}\bar{\partial}' + \bar{\partial}'\bar{\partial})(7\psi_2 + 22\rho\rho') - 20(\psi_2^2 + 8\rho\rho'\psi_2 - 8\rho^2\rho'^2) \mathcal{I}_{E1}]
\end{aligned}$$

Invariants

- Until now we had 2 odd-parity invariants, and an even parity invariant.
- We also had the “2nd order” PDE these three solved.
- And now, we have a new even parity invariant, and the “2nd order” PDE it solves.
- This brings the list to 2 odd-parity and 2 even-parity invariants in Schwarzschild spacetime.
- We later had another odd-parity invariant but didn’t look for the equation it solves yet.

Invariants

- One then wonders - are there more invariants?
- If so, how does one find them? Is there a recipe?
- Yes, there is... We use a combination of Regge-Wheeler's and Steve Detweiler's ideas.
- Just like RW, we write the different “components” of the metric perturbation and look at their gauge-contribution.
- Just like SD, we look for different possible combinations to eliminate the gauge-contributions.
- All this without the separation into radial and angular parts,

GHP-form

Lets look at the odd spin-type case..

The gauge-change for the following “components” are

$$\eth' \eth' h_{33} - \eth \eth h_{44} \longrightarrow (\eth \eth' + \eth' \eth - 4\rho\rho' - 4\psi_2)(\eth' \xi_3 - \eth \xi_4)$$

$$\rho' \eth' h_{13} - \rho' \eth h_{14} + \rho \eth' h_{23} - \rho \eth h_{24} \longrightarrow (\rho \mathbb{P}' + \rho' \mathbb{P})(\eth' \xi_3 - \eth \xi_4)$$

$$\rho' \eth' h_{13} - \rho' \eth h_{14} - \rho \eth' h_{23} + \rho \eth h_{24} \longrightarrow -(\rho \mathbb{P}' - \rho' \mathbb{P})(\eth' \xi_3 - \eth \xi_4)$$

And so you have three possible ways to eliminate the gauge vector, hence three odd invariants.

GHP-form

The odd-spin type invariants are

$$(\rho \bar{p}' - \rho' \bar{p})(\bar{\partial}' \bar{\partial}' h_{33} - \bar{\partial} \bar{\partial} h_{44}) + (\bar{\partial} \bar{\partial}' + \bar{\partial}' \bar{\partial} - 4\rho\rho' - 4\psi_2)(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} - \rho \bar{\partial}' h_{23} + \rho \bar{\partial} h_{24})$$

$$(\rho \bar{p}' + \rho' \bar{p} - 4\rho\rho')(\bar{\partial}' \bar{\partial}' h_{33} - \bar{\partial} \bar{\partial} h_{44}) - (\bar{\partial} \bar{\partial}' + \bar{\partial}' \bar{\partial} - 4\rho\rho' - 4\psi_2)(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} + \rho \bar{\partial}' h_{23} - \rho \bar{\partial} h_{24})$$

$$(\rho \bar{p}' - \rho' \bar{p})(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} + \rho \bar{\partial}' h_{23} - \rho \bar{\partial} h_{24}) + (\rho \bar{p}' + \rho' \bar{p} - 2\rho\rho')(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} - \rho \bar{\partial}' h_{23} + \rho \bar{\partial} h_{24})$$

Time derivative of the imaginary part of ψ_2

Imaginary part of ψ_2

Radial derivative of the imaginary part of ψ_2

We now have PDEs that each one of these solve!

GHP-form

The odd-spin type invariants are

$$(\rho \bar{p}' - \rho' \bar{p})(\bar{\partial}' \bar{\partial}' h_{33} - \bar{\partial} \bar{\partial} h_{44}) + (\bar{\partial} \bar{\partial}' + \bar{\partial}' \bar{\partial} - 4\rho\rho' - 4\psi_2)(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} - \rho \bar{\partial}' h_{23} + \rho \bar{\partial} h_{24})$$

$$(\rho \bar{p}' + \rho' \bar{p} - 4\rho\rho')(\bar{\partial}' \bar{\partial}' h_{33} - \bar{\partial} \bar{\partial} h_{44}) - (\bar{\partial} \bar{\partial}' + \bar{\partial}' \bar{\partial} - 4\rho\rho' - 4\psi_2)(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} + \rho \bar{\partial}' h_{23} - \rho \bar{\partial} h_{24})$$

$$(\rho \bar{p}' - \rho' \bar{p})(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} + \rho \bar{\partial}' h_{23} - \rho \bar{\partial} h_{24}) + (\rho \bar{p}' + \rho' \bar{p} - 2\rho\rho')(\rho' \bar{\partial}' h_{13} - \rho' \bar{\partial} h_{14} - \rho \bar{\partial}' h_{23} + \rho \bar{\partial} h_{24})$$

One then goes on, and systematically finds the even spin type invariants. Its more complicated to show the components or the invariants here.

Finally, we end up with **thirteen even-type** invariants using this recipe, and **three odd-type** invariants.

Summary

- Used GHP-tools to calculate the already known odd- and even-type invariants, and derived the equation they solve. All this, without separation.
- Calculated a new 4th order even-type invariant, and derived the equation it solves.
- Calculated a new 4th order odd-type invariant.
- Used a systematic approach to find other invariants and ended up with 3 odd-type and 13* even-type invariants.

* though linearly independent a few of them can be related to each other using Einstein tensors and derivatives

Summary

- Taking this to Kerr space-time — even the first steps are very complicated given that operators and variables that “commuted” in Schwarzschild no more commute with each other in Kerr spacetime — every formula has a lot of “unwanted” baggage.
- How many invariants will we get in Kerr (if any possible)?
- Will it be possible to get the PDE those invariants solve?