

First order Self-force on Kerr generic orbits

Maarten van de Meent

Max Planck Institute for Gravitational Physics

Capra 21, AEI Potsdam, 25 June 2018



Outline

① Generic orbit self-force in 4 easy steps

② Overview of Results



Generic orbit self-force in 4 “easy” steps



Step 1; Solve Teukolsky equation

Separation of variables; frequency domain

$$\psi_4 = \rho^{-4} \sum_{lm\omega} Z_{lm\omega s} R_{lm\omega}(r)_s S_{lm\omega}(\theta) e^{i(m\phi - \omega t)}$$

MST solutions

$${}_s R_{lm\omega}(r) = \mathcal{C} \sum_{n=-\infty}^{\infty} a_n^\nu F_n^\nu(r)$$

$F_n^\nu(r)$: Hypergeometric function, a_n^ν coefficients, n "renormalized angular momentum"

Variation of Parameters + extended homogeneous solutions

$$R_{lm\omega} = \begin{cases} Z_{lm\omega}^+ R_{lm\omega}^+ & r > r_0 \\ Z_{lm\omega}^- R_{lm\omega}^- & r < r_0 \end{cases},$$

$$Z_{lm\omega}^\pm = \int_{r_{\min}}^{r_{\max}} \frac{R_{lm\omega}^\mp(r) T_{lm\omega}(r)}{W[R^+, R^-]}$$



Step 1; Solve Teukolsky equation

Separation of variables; frequency domain

$$\psi_4 = \rho^{-4} \sum_{lm\omega} Z_{lm\omega} {}_s R_{lm\omega}(r) {}_s S_{lm\omega}(\theta) e^{i(m\phi - \omega t)}$$

MST solutions

$${}_s R_{lm\omega}(r) = \mathcal{C} \sum_{n=-\infty}^{\infty} a_n^\nu F_n^\nu(r)$$

$F_n^\nu(r)$: Hypergeometric function, a_n^ν coefficients, n "renormalized angular momentum"

Variation of Parameters + extended homogeneous solutions

$$R_{lm\omega} = \begin{cases} Z_{lm\omega}^+ R_{lm\omega}^+ & r > r_0 \\ Z_{lm\omega}^- R_{lm\omega}^- & r < r_0 \end{cases},$$

$$Z_{lm\omega}^\pm = \int_{r_{\min}}^{r_{\max}} \frac{R_{lm\omega}^\mp(r) T_{lm\omega}(r)}{W[R^+, R^-]}$$



Step 1; Solve Teukolsky equation

Separation of variables; frequency domain

$$\psi_4 = \rho^{-4} \sum_{lm\omega} Z_{lm\omega} s R_{lm\omega}(r) s S_{lm\omega}(\theta) e^{i(m\phi - \omega t)}$$

MST solutions

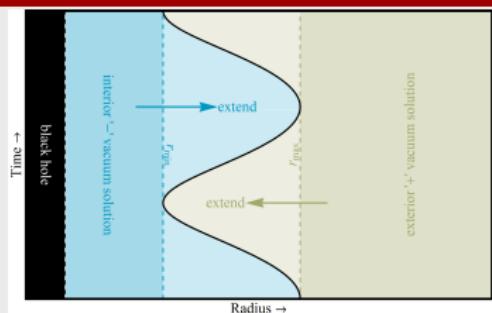
$${}_s R_{lm\omega}(r) = \mathcal{C} \sum_{n=-\infty}^{\infty} a_n^\nu F_n^\nu(r)$$

$F_n^\nu(r)$: Hypergeometric function, a_n^ν coefficients, n "renormalized angular momentum"

Variation of Parameters + extended homogeneous solutions

$$R_{lm\omega} = \begin{cases} Z_{lm\omega}^+ R_{lm\omega}^+ & r > r_0 \\ Z_{lm\omega}^- R_{lm\omega}^- & r < r_0 \end{cases},$$

$$Z_{lm\omega}^\pm = \int_{r_{\min}}^{r_{\max}} \frac{R_{lm\omega}^\mp(r) T_{lm\omega}(r)}{W[R^+, R^-]}$$



Step 2: Metric Reconstruction

A tale of two maps [Chrzanowski, Cohen, Kegeles, Wald, 1970s]

$$\hat{\psi}_4 : h_{\mu\nu} \mapsto \psi_4$$

$\mathcal{S}_{\text{CCK}}^\dagger$: (vacuum solutions of spin-2 teuk. eq.) \rightarrow (vacuum metric perturbations)

"Radial" inversion [Ori, 2001][MvdM, 2014]

$\hat{\psi}_4 \circ \mathcal{S}_{\text{CCK}}^\dagger$ is not the identity map. However, it is diagonal on a basis of frequency domain in/out (up/down) going radial modes. Ergo, easy to invert mode-by-mode.

Metric reconstruction

$\mathcal{S}_{\text{CCK}}^\dagger \circ (\hat{\psi}_4 \circ \mathcal{S}_{\text{CCK}}^\dagger)^{-1}$ is a right inverse for $\hat{\psi}_4$ restricted to vacuum solutions.



Step 2: Metric Reconstruction

A tale of two maps [Chrzanowski, Cohen, Kegeles, Wald, 1970s]

$$\hat{\psi}_4 : h_{\mu\nu} \mapsto \psi_4$$

$\mathcal{S}_{\text{CCK}}^\dagger$: (vacuum solutions of spin-2 teuk. eq.) \rightarrow (vacuum metric perturbations)

“Radial” inversion [Ori, 2001][MvdM, 2014]

$\hat{\psi}_4 \circ \mathcal{S}_{\text{CCK}}^\dagger$ is not the identity map. However, it is diagonal on a basis of frequency domain in/out (up/down) going radial modes. Ergo, easy to invert mode-by-mode.

Metric reconstruction

$\mathcal{S}_{\text{CCK}}^\dagger \circ (\hat{\psi}_4 \circ \mathcal{S}_{\text{CCK}}^\dagger)^{-1}$ is a right inverse for $\hat{\psi}_4$ restricted to vacuum solutions.



Step 2: Metric Reconstruction

A tale of two maps [Chrzanowski, Cohen, Kegeles, Wald, 1970s]

$$\hat{\psi}_4 : h_{\mu\nu} \mapsto \psi_4$$

$\mathcal{S}_{\text{CCK}}^\dagger$: (vacuum solutions of spin-2 teuk. eq.) \rightarrow (vacuum metric perturbations)

“Radial” inversion [Ori, 2001][MvdM, 2014]

$\hat{\psi}_4 \circ \mathcal{S}_{\text{CCK}}^\dagger$ is not the identity map. However, it is diagonal on a basis of frequency domain in/out (up/down) going radial modes. Ergo, easy to invert mode-by-mode.

Metric reconstruction

$\mathcal{S}_{\text{CCK}}^\dagger \circ (\hat{\psi}_4 \circ \mathcal{S}_{\text{CCK}}^\dagger)^{-1}$ is a right inverse for $\hat{\psi}_4$ restricted to vacuum solutions.



Step 3: Completion

[Wald 1968]

$$\ker \hat{\psi}_4 = (\delta M \text{ and } \delta J \text{ perturbations} + \text{gauge modes})$$

δM and δJ perturbations [Merlin et al., 2016][MvdM, 2017]

- Abbott-Deser integrals measure mass/angular momentum inside closed surfaces.
- $\text{im}(\mathcal{S}_{\text{CCK}}^\dagger)$ does not contribute to Abbott-Deser integrals [MvdM, 2017].

Gauge modes

- Care only about 2d family of linearly growing gauge vectors that produce bounded metric perturbations. (but not today)
- Can be fixed by a generalization of the methods of [Merlin et al., 2016].
- Not done (yet) for generic orbits.



Step 3: Completion

[Wald 1968]

$$\ker \hat{\psi}_4 = (\delta M \text{ and } \delta J \text{ perturbations} + \text{gauge modes})$$

δM and δJ perturbations [Merlin et al., 2016][MvdM, 2017]

- Abbott-Deser integrals measure mass/angular momentum inside closed surfaces.
- $\text{im}(S_{\text{CCK}}^\dagger)$ does not contribute to Abbott-Deser integrals [MvdM, 2017].

Gauge modes

- Care only about 2d family of linearly growing gauge vectors that produce bounded metric perturbations. (but not today)
- Can be fixed by a generalization of the methods of [Merlin et al., 2016].
- Not done (yet) for generic orbits.



Step 3: Completion

[Wald 1968]

$$\ker \hat{\psi}_4 = (\delta M \text{ and } \delta J \text{ perturbations} + \text{gauge modes})$$

δM and δJ perturbations [Merlin et al., 2016][MvdM, 2017]

- Abbott-Deser integrals measure mass/angular momentum inside closed surfaces.
- $\text{im}(\mathcal{S}_{\text{CCK}}^\dagger)$ does not contribute to Abbott-Deser integrals [MvdM, 2017].

Gauge modes

- Care only about 2d family of linearly growing gauge vectors that produce bounded metric perturbations. (but not today)
- Can be fixed by a generalization of the methods of [Merlin et al., 2016].
- Not done (yet) for generic orbits.



Step 4: Mode sum regularization

Mode sum [Pound, Merlin & Barack, 2013]

$$F^\mu = \left(\sum_{\ell=0}^{\infty} \frac{F_{\ell, \text{Rad}}^{\mu, +} + F_{\ell, \text{Rad}}^{\mu, -}}{2} - B_{\text{Lor}}^\mu - \frac{C_{\text{Lor}}^\mu}{L} \right) - D_{\text{Lor}}^\mu.$$

Re-expansion to spherical harmonics

$$\mathcal{F}_{\text{Rad}}^{\mu, \pm} = \sum_{\substack{m\omega l \\ ij}} \mathcal{C}_{lm\omega ij}^{\mu, \pm}(r, \textcolor{red}{z}) {}_2R_{lm\omega}^{\pm, (\text{i})}(r) {}_2S_{lm\omega}^{(\text{j})}(z) e^{im\phi - i\omega t} + c.c.$$

becomes

$$= \sum_{\substack{m\omega ijk \\ l_1 l_2 l_3 l}} \left(\tilde{\mathcal{C}}_{lm\omega ijk}^{\mu, \pm} {}_2R_{lm\omega}^{\pm, (\text{i})} ({}_2b_{m\omega})_{l_1}^{\text{l}} {}_j^m \mathcal{A}_{l_2}^{l_1 k} \mathcal{B}_{\ell}^{l_2} e^{im\phi - i\omega t} + c.c. \right) Y_{\ell m}(z)$$

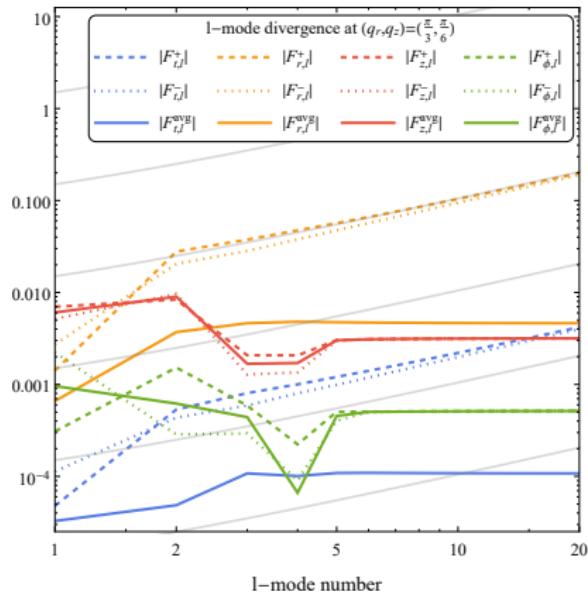


Overview of Results



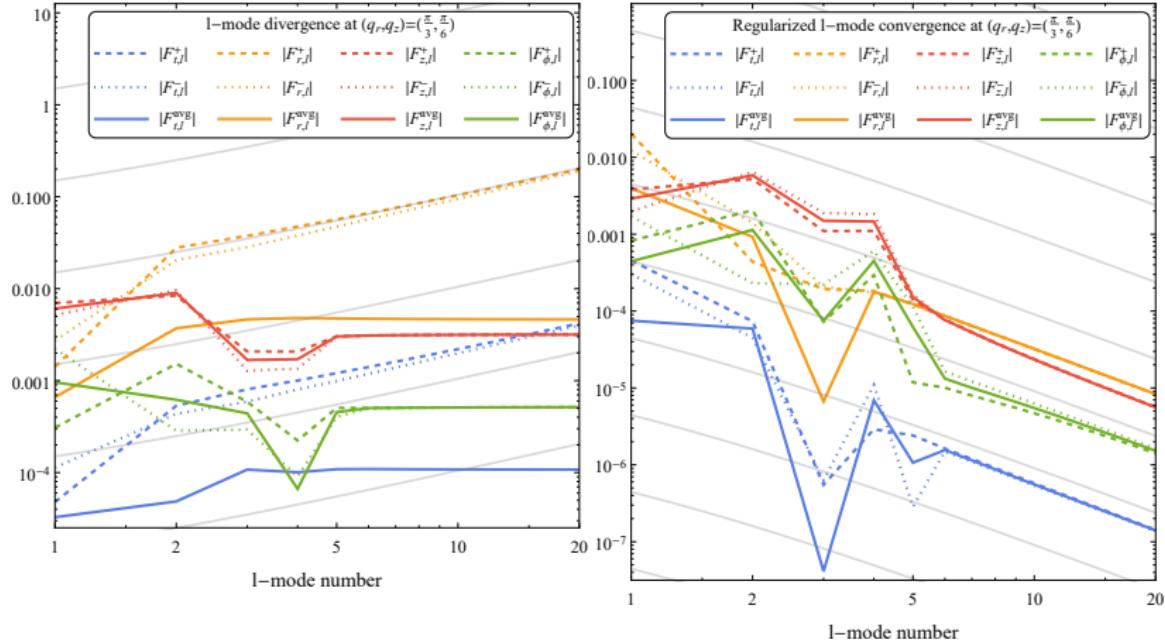
l -mode convergence test

$a = 0.9M$, $p = 10$, $e = 0.1$, and $z_{\max} = \cos \theta_{\min} = 0.1$

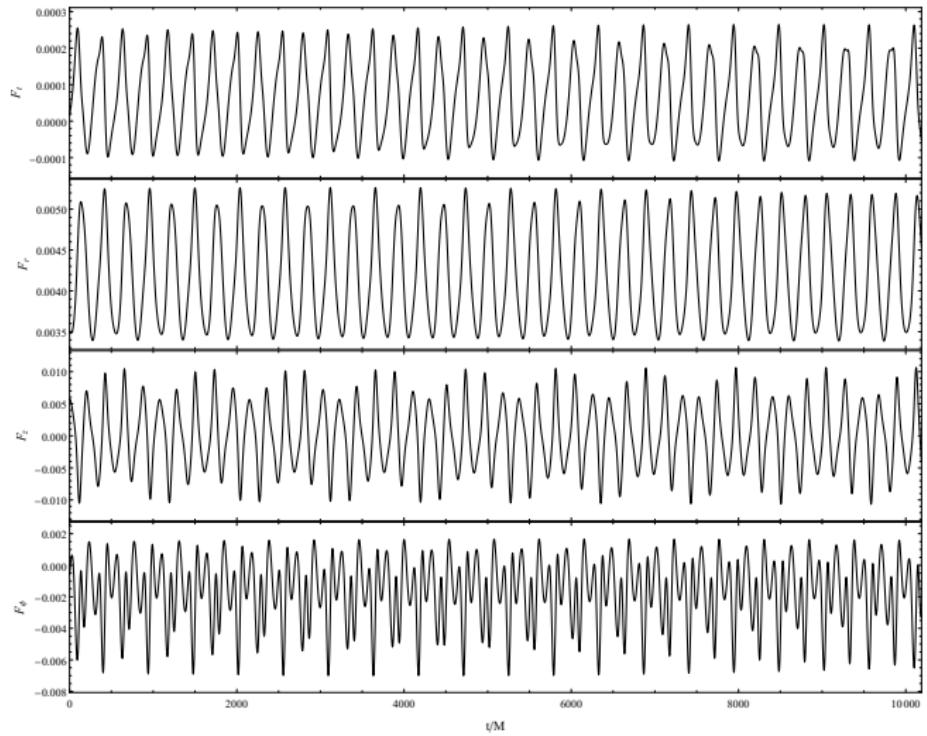


l -mode convergence test

$a = 0.9M$, $p = 10$, $e = 0.1$, and $z_{\max} = \cos \theta_{\min} = 0.1$



Timeseries

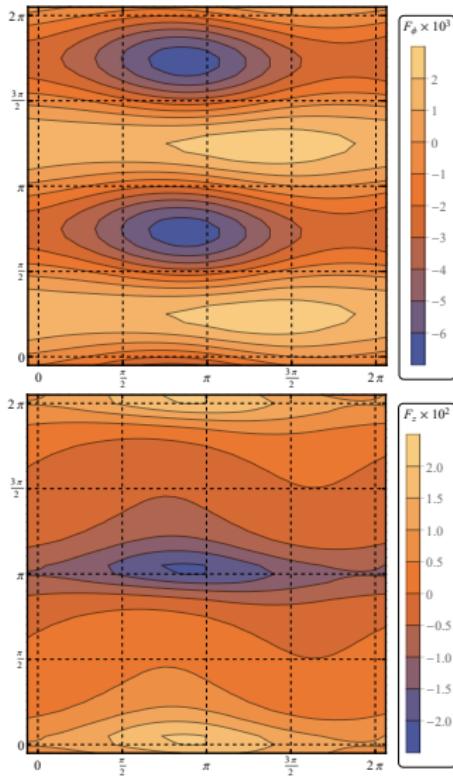
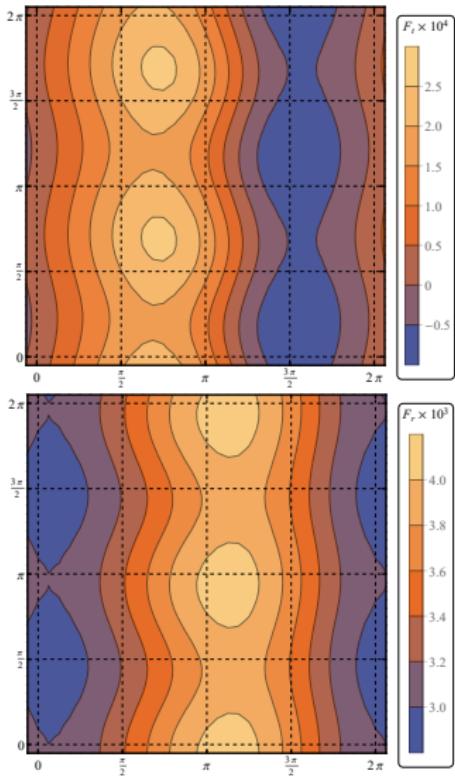


$a = 0.9M$,
 $p = 10$,
 $e = 0.1$, and
 $z_{\max} = \cos \theta_{\min} = 0.5$

- F_μ is biperiodic
- not very informative



"Torus" plots



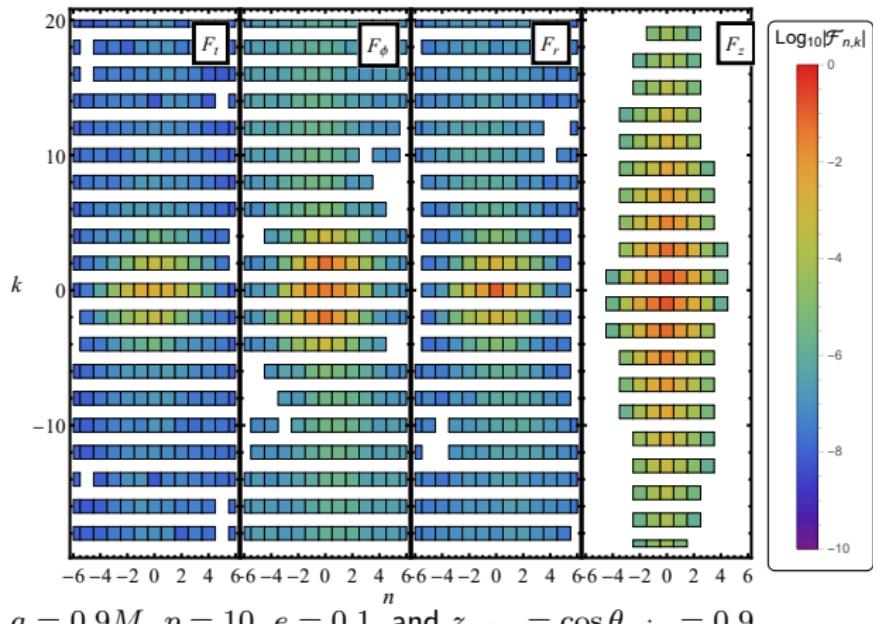
$a = 0.9M$,
 $p = 10$,
 $e = 0.1$, and
 $z_{\max} = 0.9$

- q_r vs q_z
- angles conjugate to radial and polar action
- F_t , F_r , and F_ϕ are π periodic in q_z



Spectra

$$F(q_r, q_z) = \sum_{nk} \mathcal{F}_{nk} e^{inq_r + ikq_z}$$



$$a = 0.9M, p = 10, e = 0.1, \text{ and } z_{\max} = \cos \theta_{\min} = 0.9$$

- Half of the modes vanish due to up/down symmetry.



The good, the bad, and the ugly

The good ...

- We can calculate the self-force on generic orbits in Kerr spacetime.
- High inclinations pose little problems.

The bad ...

- Slow ($\sim 10^4$ CPU hours/geodesic)
- Memory intensive.

The ugly ...

- Re-expansion to spherical harmonic modes is unpleasant.



Summary and Outlook

Summary

Calculation of the GSF corrections to generic orbits is a reality.

To do...

- spherical orbits
- explore resonances
- invariants (redshift, spin precession, ISSO shift)
- fill parameter space

Thank you for listening!

Acknowledgments



This work has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 705229.

