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# **DISCONTINUOUS COLLOCATION METHODS** FOR SELF-FORCE APPLICATIONS



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#### **TEUKOLSKY (OR REGGE-WHEELER) EQUATION**

- Strongly hyperbolic PDE describing linear perturbations in Kerr (or Schwarzschild)
- Typically solved as ODE in frequency domain
- Used to compute GSF in radiation (or Regge-Wheeler) gauge
- Time-domain Teukolsky solvers have been increasing in sofistication

```
\left[\left(\nabla^{\mu} + A^{\mu}\right)\nabla_{\mu} + V\right]\psi_{0}(r) = S
```

(Shah, Friedman, Whiting; Barack, Pound, Merlin; Kohen, Kegeles; Barack, Ori et al) For highly eccentric or unbound orbits, time-domain self-consistent evolution is ideal (Burko, Khanna, Pullin, Hughes, Poisson, Lousto, Zenginoglu, Harms, Bernuzzi, Brügmann

#### METHOD OF LINES

- Reduce to first-order in time  $\partial_t U_A = L^A$
- for smooth methods

$$^{AB}U_B + S_B, \quad U_A = \left(\psi_0, \dot{\psi}_0\right)$$

Discretize U, L in space, evolved as coupled ODE system in time  $\frac{d\mathbf{U}}{dt} = \mathbf{L}\mathbf{U} + \mathbf{S}$ 

L contains spatial differentiation matrices, which are typically valid only

### **TEUKOLSKY EQUATION WITH PARTICLE SOURCE**

- Approximate δ as Gaussian pulse (Harms, Bernuzzi and Brügmann)
- Construct finite-difference representation of  $\delta$  (Hughes et al)
- "Lagrangian" picture: Domain decomposition + time-dependent grid + jump conditions across particle (Canizares and Sopuerta; Field, Hesthaven and Lau, Diener et al).
- "Eulerian" picture: Discontinuous collocation method + fixed grid + jump conditions across particle (Markakis, Barack et al).

#### C. Markakis and L. Barack [arXiv:1406.4865]

```
[(\nabla^{\mu} + A^{\mu})\nabla_{\mu} + V]\psi_{0}(r) = G\delta(r - \xi(t)) + F\delta'(r - \xi(t))
```

## **DISCONTINUOUS COLLOCATION METHOD**

- Jumps in solution  $\psi$  and its derivatives across particle known a priori from field equation  $\psi^{(k)}(\xi^+) - \psi^{(k)}(\xi^-) = J_k, \quad k = 0, 1, 2, ...$
- Collocation (finite-difference and pseudo-spectral) methods are based on Lagrange interpolation
- Construct discontinuous generalization to Lagrange interpolation that uses  $\{J_i\}$  as input.

#### LAGRANGE INTERPOLATION

N<sup>th</sup> order polynomial

p(x)

Collocation conditions

p(x)

Solution: Lagrange interpolating polynomial

$$p(x) = \sum_{j=0}^{N} f_j(x) \pi_j(x), \quad \pi_j = \prod_{k=0}^{N} \frac{x - x_k}{x_j - x_k}$$

$$z) = \sum_{j=0}^{N} c_j x^j$$

$$= f_i, \quad i = 0, 1, ..., N$$

#### **DISCONTINUOUS LAGRANGE INTERPOLATION**

- N<sup>th</sup> order piecewise polynomial
- $p(x) = \sum \left[ e^{-\frac{1}{2}} \right]$ *j*=0 Collocation conditions
  - p(x) =

N

Jump conditions

- $p^{(k)}(\xi^+) p^{(k)}$
- Solution: interpolating piecewise polynomial

$$p(x) = \sum_{j=0}^{N} [j]$$

$$\theta(x-\xi)c_j^+ + \theta(\xi-x)c_j^-]x^j$$

$$= f_i, \quad i = 0, 1, ..., N$$

$$J^{(k)}(\xi^{-}) = J_k, \quad k = 0, 1, 2, \dots$$

 $f_j(x) + \overline{\Delta(x_j - \xi; x - \xi)} ]\pi_j(x)$ 

#### **DISCONTINUOUS INTERPOLATION, DIFFERENTIATION, INTEGRATION**

Interpolation

$$p(x) = \sum_{j=0}^{N} [f_j(x) + \Delta(x_j - \xi; x - \xi)] \pi_j(x)$$

Differentiation

$$p^{(n)}(x_i) = \sum_{j=0}^N D_{ij}^{(n)}[f_j(x) + \Delta(x_j - \xi; x_i - \xi)], \quad D_{ij}^{(n)} = \pi_j^{(n)}(x_i)$$

$$\Delta(x_j - \xi; x_i - \xi) = [\theta(x_i - \xi) - \theta(x_j - \xi)] \sum_k \frac{J_k}{k!} (x_j - \xi)^k$$

Integration...

#### **SELF-FORCE ON STATIC SCALAR PARTICLE (SCHWARZSCHILD)**



Equidistant nodes

exact function f(x)

- Lagrange interpolation (M = -1)
  - Discontinuous interpolation (M = 6)

= Discontinuous interpolation (M = 12)

#### **SELF-FORCE ON STATIC SCALAR PARTICLE (SCHWARZSCHILD)**



Chebyschev-Gauss-Lobatto nodes

## SELF-FORCE ON STATIC SCALAR PARTICLE (SCHWARZSCHILD)



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Chebyschev-Gauss-Lobatto nodes

#### TIME INTEGRATION

- Explicit methods (e.g. Runge-Kutta) are not time-symmetric.
  Energy not conserved. Conditionally stable (CFL limited).
- Implicit methods (e.g. trapezium rule, Hermite rule, Lotkin rule and their higher compositions, Suzuki-Yoshida).
   Energy conserved. Unconditionally stable (no CFL limit).
- For linear PDEs, implicit methods do not incur extra cost (matrix inversion).
- For GSF computations in the time domain, avoid RK. Opt for Crank-Nicolson and higher order generalizations. Invert and store matrices to make scheme 'explicit'.

#### C. Markakis et al. [arXiv:1901.09967]

#### CONCLUSIONS

- High accuracy for pseudospectral nodes
- 'Eulerian' picture. Grid fixed/stored, particle can move freely inside domain
- Avoids domain decomposition and coordinate mapping
- Can use a single hyperboloidal slice for the whole spacetime
- Efficient and easy to implement
- Must use relatively large number of jumps
- Must update  $\Delta(x_j,\xi;x_i,\xi)$  as particle moves
- Self-consistent evolution under influence of GSF possible
  - C. Markakis and L. Barack [arXiv:1406.4865]
  - L. Barack and P. Giudice, rXiv:1702.04204
  - C. Markakis et al. [arXiv:1901.09967]

Solving Teukolsky equation for Hertz potential Ψ can be used to directly reconstruct GSF (Barack and Giudice)