

Calculating the scalar self-force during $r\theta$ -resonances

Zachary Nasipak¹ and Charles R. Evans¹

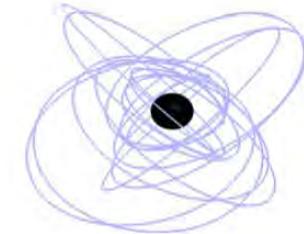
22 June 2020

23rd Capra Meeting on Radiation Reaction in General Relativity
University of Texas at Austin



Motivation

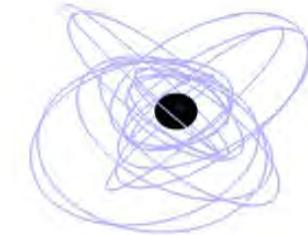
- Model evolution of EMRIs
 - Parameterize motion with slowly evolving ‘constants’
 - Example: $(p, e, x); (E, L_z, Q); (\Upsilon_r, \Upsilon_\theta, \Upsilon_\phi)$
 - Orbital $r\theta$ -resonances in EMRIs
 - Radial and polar motion are commensurate
$$\Upsilon_r/\Upsilon_\theta = \beta_r/\beta_\theta \quad (\beta_i \in \mathbb{Z})$$
 - Example: $(\beta_r, \beta_\theta) = (1,2) \Rightarrow 1:2$ $r\theta$ -resonance
 - Almost all EMRIs encounter at least one strong (1:3, 1:2, 2:3) $r\theta$ -resonance [Ruanes & Hughes (2014)]



Motivation

- Model evolution of EMRIs
 - Parameterize motion with slowly evolving ‘constants’
 - Example: $(p, e, x); (E, L_z, Q); (\Upsilon_r, \Upsilon_\theta, \Upsilon_\phi)$
 - Orbital $r\theta$ -resonances in EMRIs
 - Radial and polar motion are commensurate

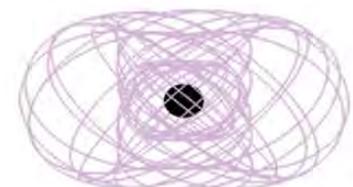
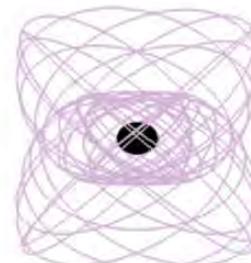
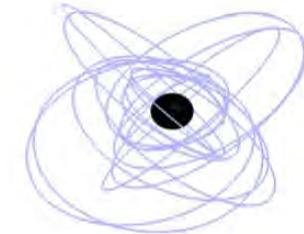
$$\Upsilon_r/\Upsilon_\theta = \beta_r/\beta_\theta \quad (\beta_i \in \mathbb{Z})$$



- Example: $(\beta_r, \beta_\theta) = (1,2) \Rightarrow 1:2 r\theta\text{-resonance}$
- Almost all EMRIs encounter at least one strong (1:3, 1:2, 2:3) $r\theta$ -resonance [Kocsayon & Hughes (2014)]

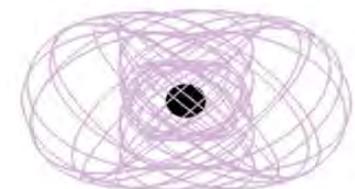
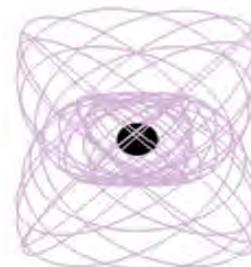
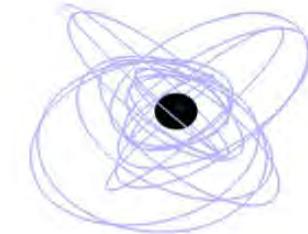
Motivation

- Model evolution of EMRIs
 - Parameterize motion with slowly evolving ‘constants’
 - Example: $(p, e, x); (E, L_z, Q); (\Upsilon_r, \Upsilon_\theta, \Upsilon_\phi)$
 - Orbital $r\theta$ -resonances in EMRIs
 - Radial and polar motion are commensurate
$$\Upsilon_r/\Upsilon_\theta = \beta_r/\beta_\theta \quad (\beta_i \in \mathbb{Z})$$
 - Example: $(\beta_r, \beta_\theta) = (1, 2) \Rightarrow 1:2$ $r\theta$ -resonance
 - Almost all EMRIs encounter at least one strong (1:3, 1:2, 2:3) $r\theta$ -resonance [Rughesi & Hughes (2014)]



Motivation

- Model evolution of EMRIs
 - Parameterize motion with slowly evolving ‘constants’
 - Example: $(p, e, x); (E, L_z, Q); (\Upsilon_r, \Upsilon_\theta, \Upsilon_\phi)$
 - Orbital $r\theta$ -resonances in EMRIs
 - Radial and polar motion are commensurate
$$\Upsilon_r/\Upsilon_\theta = \beta_r/\beta_\theta \quad (\beta_i \in \mathbb{Z})$$
 - Example: $(\beta_r, \beta_\theta) = (1, 2) \Rightarrow 1:2$ $r\theta$ -resonance
 - Almost all EMRIs encounter at least one strong (1:3, 1:2, 2:3) $r\theta$ -resonance [Ruangsriv & Hughes (2014)]



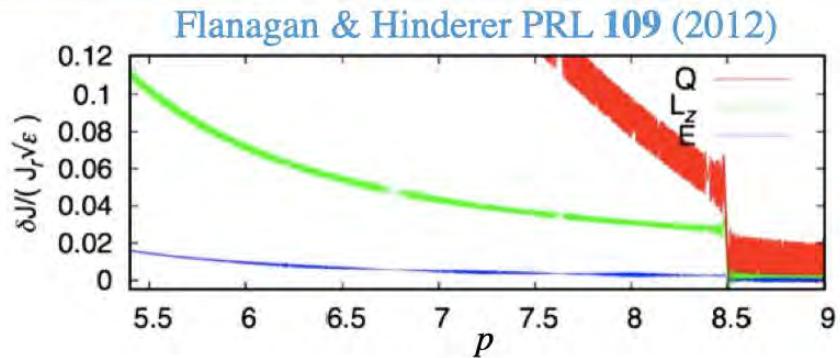
Motivation

- Importance of transient orbital $r\theta$ -resonances

- ‘Kicks’ in the orbital ‘constants’, e.g., (E, L_z, Q)
 - Scale as $\sim \epsilon^{1/2}$ ($\epsilon \equiv \mu/M$)
 - Sensitive to initial conditions/phases at resonance
- Impact EMRI waveform phase ϕ @ 1/2-order
$$\phi = \phi^{(-1)}\epsilon^{-1} + \phi^{(-1/2)}\epsilon^{-1/2} + \phi^{(0)}\epsilon^0 + \dots$$

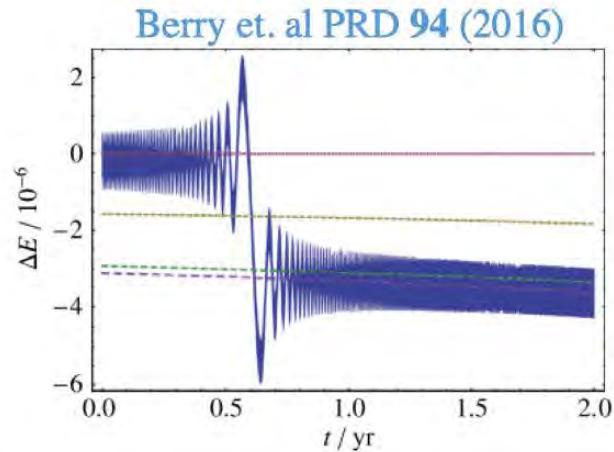
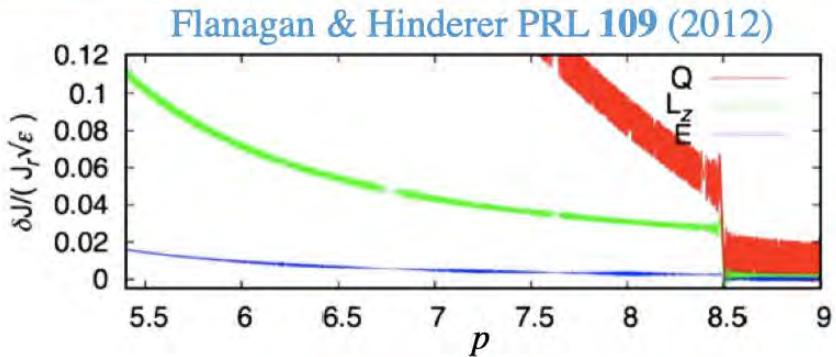
Motivation

- Importance of transient orbital $r\theta$ -resonances
 - ‘Kicks’ in the orbital ‘constants’, e.g., (E, L_z, Q)
 - Scale as $\sim \epsilon^{1/2}$ ($\epsilon \equiv \mu/M$)
 - Sensitive to initial conditions/phases at resonance
 - Impact EMRI waveform phase ϕ @ 1/2-order
$$\phi = \phi^{(-1)}\epsilon^{-1} + \phi^{(-1/2)}\epsilon^{-1/2} + \phi^{(0)}\epsilon^0 + \dots$$



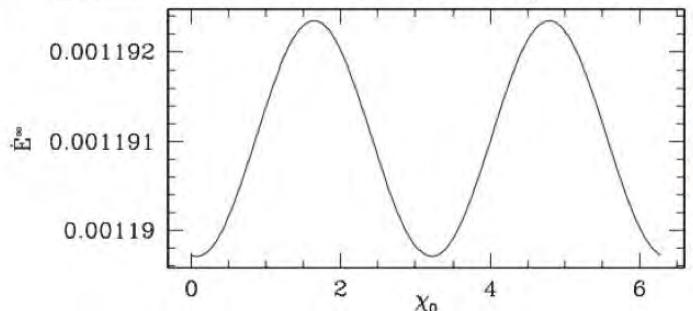
Motivation

- Importance of transient orbital $r\theta$ -resonances
 - ‘Kicks’ in the orbital ‘constants’, e.g., (E, L_z, Q)
 - Scale as $\sim \epsilon^{1/2}$ ($\epsilon \equiv \mu/M$)
 - Sensitive to initial conditions/phases at resonance
 - Impact EMRI waveform phase ϕ @ 1/2-order
$$\phi = \phi^{(-1)}\epsilon^{-1} + \phi^{(-1/2)}\epsilon^{-1/2} + \phi^{(0)}\epsilon^0 + \dots$$



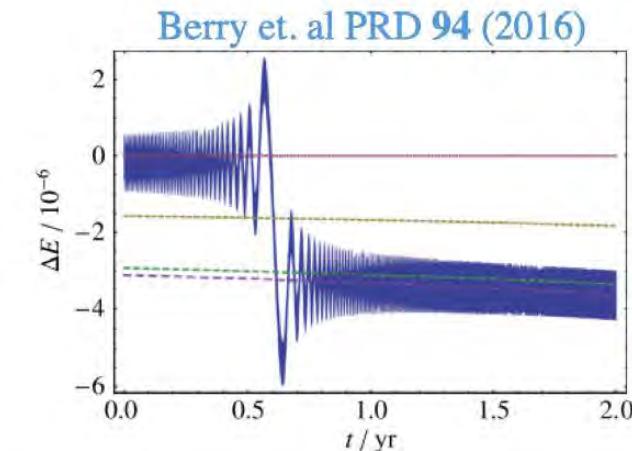
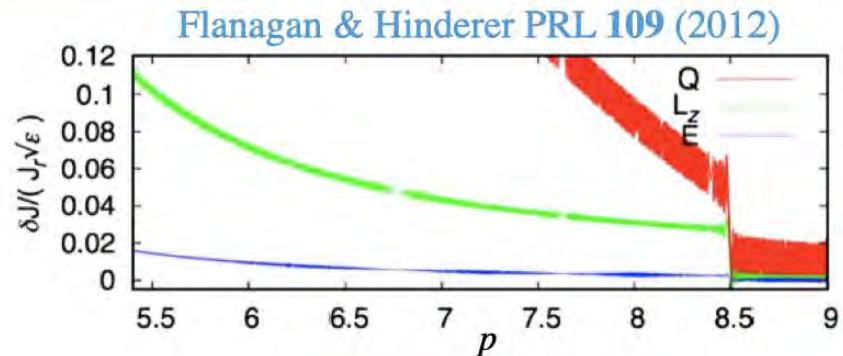
Motivation

- Importance of transient orbital $r\theta$ -resonances
 - ‘Kicks’ in the orbital ‘constants’, e.g., (E, L_z, Q)
 - Scale as $\sim \epsilon^{1/2}$ ($\epsilon \equiv \mu/M$)
 - Sensitive to initial conditions/phases at resonance



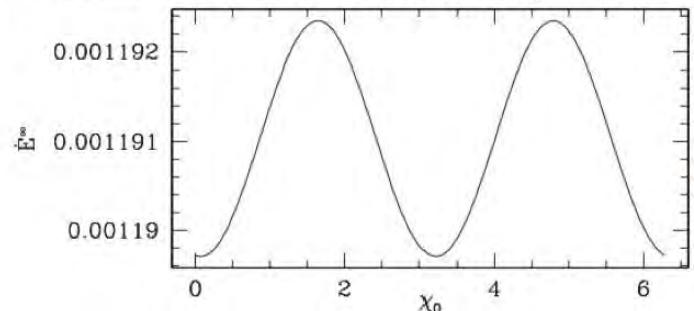
Flanagan et. al
PRD 89 (2014)

- Impact EMRI waveform phase ϕ @ 1/2-order
$$\phi = \phi^{(-1)}\epsilon^{-1} + \phi^{(-1/2)}\epsilon^{-1/2} + \phi^{(0)}\epsilon^0 + \dots$$



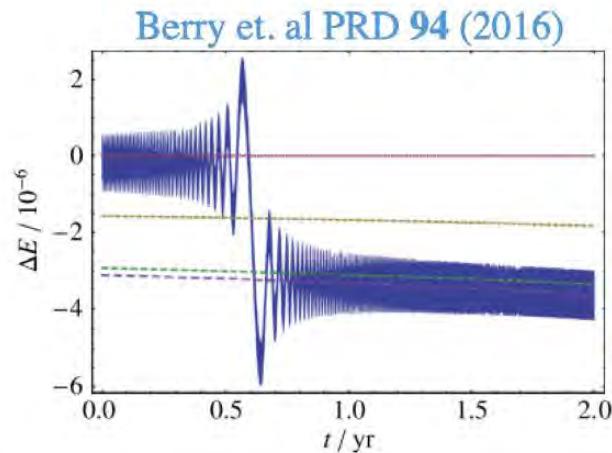
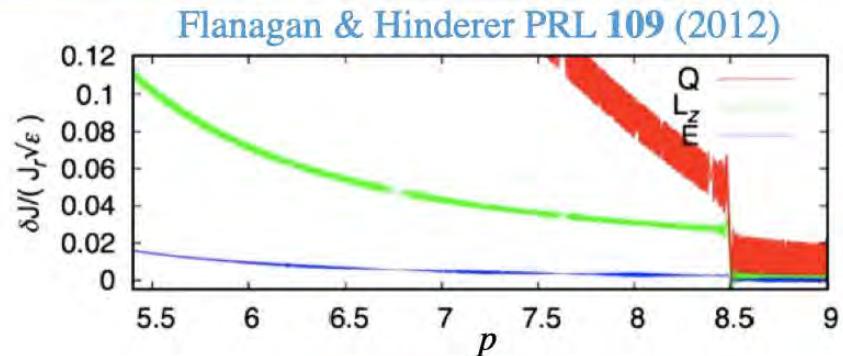
Motivation

- Importance of transient orbital $r\theta$ -resonances
 - ‘Kicks’ in the orbital ‘constants’, e.g., (E, L_z, Q)
 - Scale as $\sim \epsilon^{1/2}$ ($\epsilon \equiv \mu/M$)
 - Sensitive to initial conditions/phases at resonance



Flanagan et. al
PRD 89 (2014)

- Impact EMRI waveform phase ϕ @ 1/2-order
$$\phi = \phi^{(-1)}\epsilon^{-1} + \underline{\phi^{(-1/2)}\epsilon^{-1/2}} + \phi^{(0)}\epsilon^0 + \dots$$

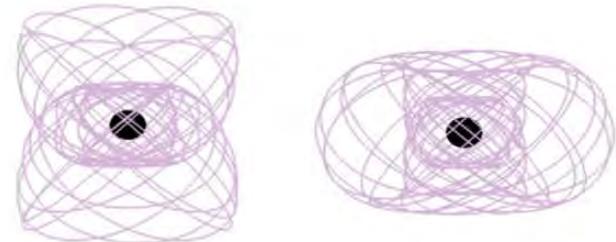


Motivation

- Unanswered questions, including...
 - Incorporating initial conditions in self-force calculations?
 - Conservative contributions? [Flanagan & Hinderer PRL 109 (2012)]
 - Hamiltonian formulation of $\langle \dot{Q} \rangle$ includes potential conservative terms [Iacopini et. al PTEP 013E01 (2019)]
 - *Integrability conjecture:* conservative dynamics will not contribute to adiabatic evolution during resonances
 - And more..., e.g., likelihood of sustained resonances? [van de Meent PRD 89 (2014)]
- Goal: study self-force during $r\theta$ -resonances using scalar self-force (SSF) model [Nasipak et. al PRD 100 (2019)]

Motivation

- Unanswered questions, including...
 - Incorporating initial conditions in self-force calculations?
 - Conservative contributions? [Flanagan & Hinderer PRD 109 (2012)]
 - Hamiltonian formulation of $\langle \dot{Q} \rangle$ includes potential conservative terms [Izquierdo et. al PTEP 013E01 (2019)]
 - *Integrability conjecture:* conservative dynamics will not contribute to adiabatic evolution during resonances
 - And more..., e.g., likelihood of sustained resonances? [van de Meent PRD 89 (2014)]
- Goal: study self-force during $r\theta$ -resonances using scalar self-force (SSF) model [Nasipak et. al PRD 100 (2019)]



Motivation

- Unanswered questions, including...
 - Incorporating initial conditions in self-force calculations?
 - Conservative contributions? [Flanagan & Hinderer PRL 109 (2012)]
 - Hamiltonian formulation of $\langle \dot{Q} \rangle$ includes potential conservative terms [Iacopini et. al PTEP 013E01 (2019)]
 - *Integrability conjecture*: conservative dynamics will not contribute to adiabatic evolution during resonances
 - And more..., e.g., likelihood of sustained resonances? [van de Meent PRD 89 (2014)]
- Goal: study self-force during $r\theta$ -resonances using scalar self-force (SSF) model [Nasipak et. al PRD 100 (2019)]

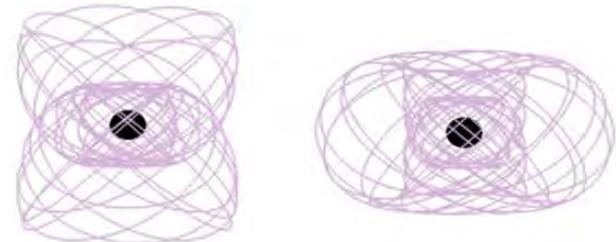
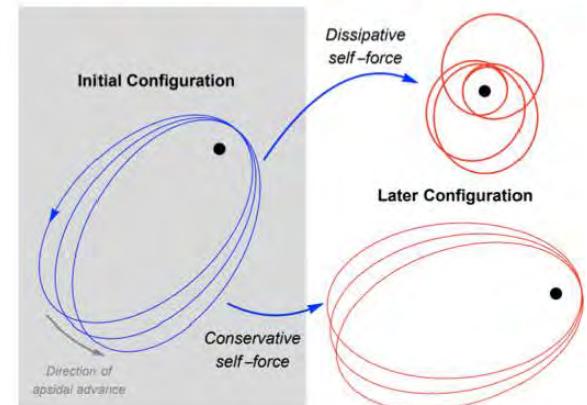


Image Credit: Osburn et. al (2016)



Motivation

- Unanswered questions, including...
 - Incorporating initial conditions in self-force calculations?
 - Conservative contributions? [Flanagan & Hinderer PRL 109 (2012)]
 - Hamiltonian formulation of $\langle \dot{Q} \rangle$ includes potential conservative terms [Isoyama et. al PTEP 013E01 (2019)]
 - *Integrability conjecture*: conservative dynamics will not contribute to adiabatic evolution during resonances
 - And more..., e.g., likelihood of sustained resonances? [van de Meent PRD 89 (2014)]
- Goal: study self-force during $r\theta$ -resonances using scalar self-force (SSF) model [Nasipak et. al PRD 100 (2019)]

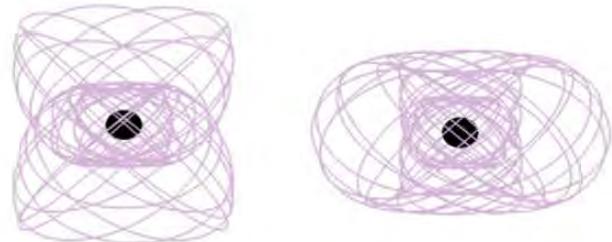
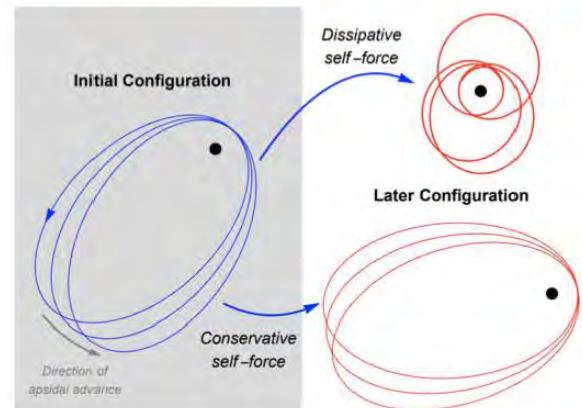


Image Credit: Osburn et. al (2016)



Motivation

- Unanswered questions, including...
 - Incorporating initial conditions in self-force calculations?
 - Conservative contributions? [Flanagan & Hinderer PRL 109 (2012)]
 - Hamiltonian formulation of $\langle \dot{Q} \rangle$ includes potential conservative terms [Isoyama et. al PTEP 013E01 (2019)]
 - *Integrability conjecture*: conservative dynamics will not contribute to adiabatic evolution during resonances
 - And more..., e.g., likelihood of sustained resonances? [van de Meent PRD 89 (2014)]
- Goal: study self-force during $r\theta$ -resonances using scalar self-force (SSF) model [Nasipak et. al PRD 100 (2019)]

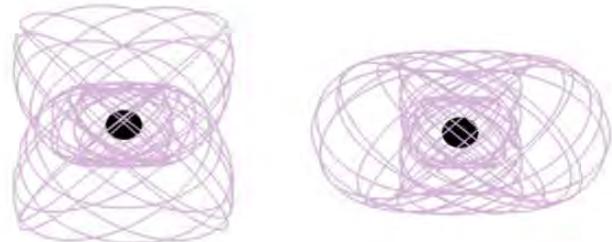
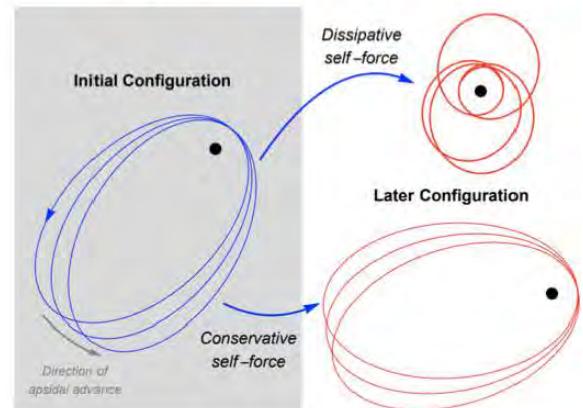


Image Credit: Osburn et. al (2016)



Motivation

- Unanswered questions, including...
 - Incorporating initial conditions in self-force calculations?
 - Conservative contributions? [Flanagan & Hinderer PRL 109 (2012)]
 - Hamiltonian formulation of $\langle \dot{Q} \rangle$ includes potential conservative terms [Isoyama et. al PTEP 013E01 (2019)]
 - *Integrability conjecture*: conservative dynamics will not contribute to adiabatic evolution during resonances
 - And more..., e.g., likelihood of sustained resonances? [van de Meent PRD 89 (2014)]
 - Goal: study self-force during $r\theta$ -resonances using scalar self-force (SSF) model [Nasipak et. al PRD 100 (2019)]

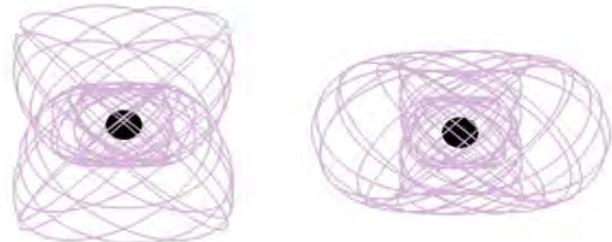
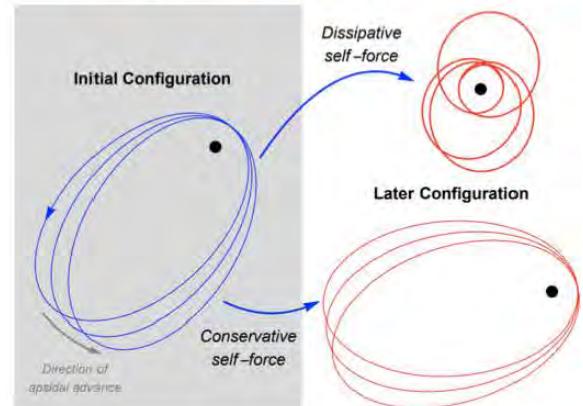


Image Credit: Osburn et. al (2016)



Motivation

- Unanswered questions, including...
 - Incorporating initial conditions in self-force calculations?
 - Conservative contributions? [Flanagan & Hinderer PRL 109 (2012)]
 - Hamiltonian formulation of $\langle \dot{Q} \rangle$ includes potential conservative terms [Isoyama et. al PTEP 013E01 (2019)]
 - *Integrability conjecture*: conservative dynamics will not contribute to adiabatic evolution during resonances
 - And more..., e.g., likelihood of sustained resonances? [van de Meent PRD 89 (2014)]
- Goal: study self-force during $r\theta$ -resonances using scalar self-force (SSF) model [Nasipak et. al PRD 100 (2019)]

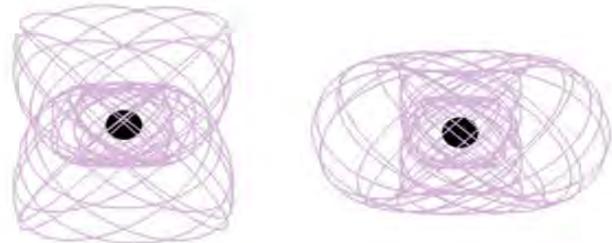
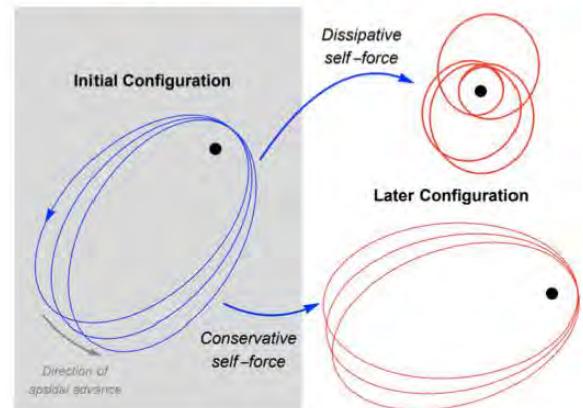


Image Credit: Osburn et. al (2016)



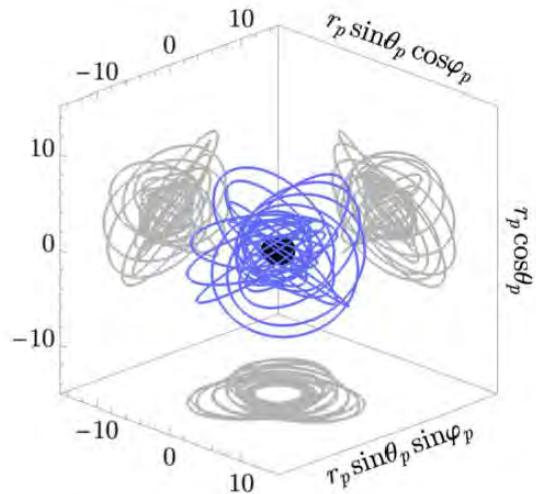
Impact of initial phases

- Consider non-resonant inclined, eccentric geodesic
 - Radial & polar motion separate w/ (Carter-)Mino time λ
 - Define angle variables: $q_r \equiv \Upsilon_r \lambda$, $q_\theta \equiv \Upsilon_\theta \lambda$



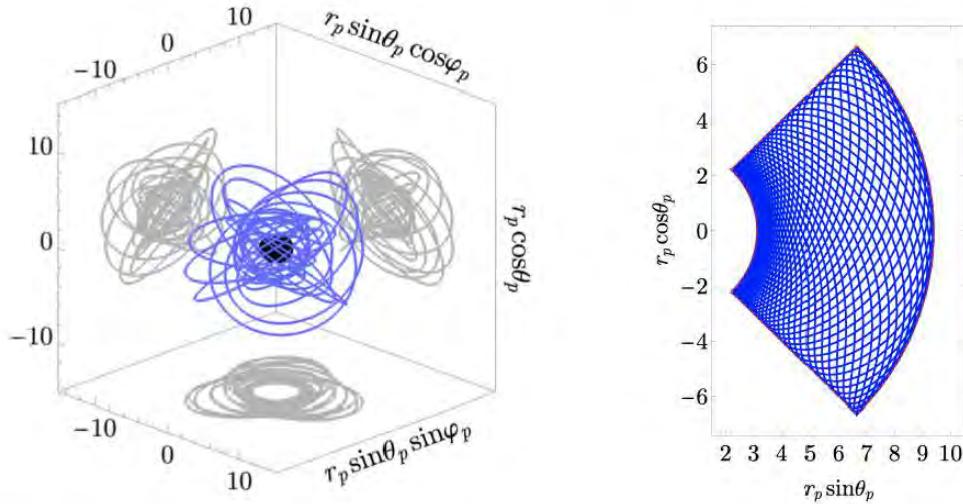
Impact of initial phases

- Consider non-resonant inclined, eccentric geodesic
 - Radial & polar motion separate w/ (Carter-)Mino time λ
 - Define angle variables $q_r \equiv \Upsilon_r \lambda, q_\theta \equiv \Upsilon_\theta \lambda$



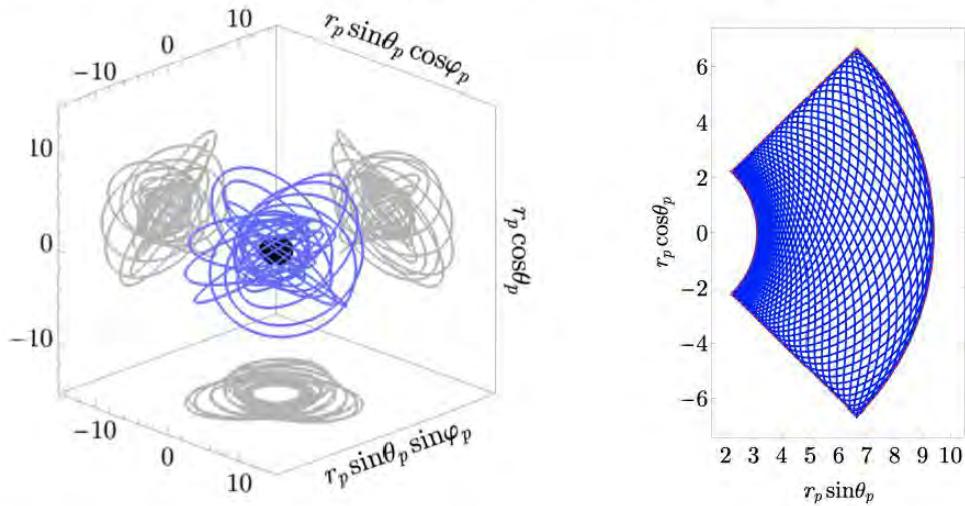
Impact of initial phases

- Consider non-resonant inclined, eccentric geodesic
 - Radial & polar motion separate w/ (Carter-)Mino time λ
 - Define angle variables $q_r \equiv \Upsilon_r \lambda, q_\theta \equiv \Upsilon_\theta \lambda$



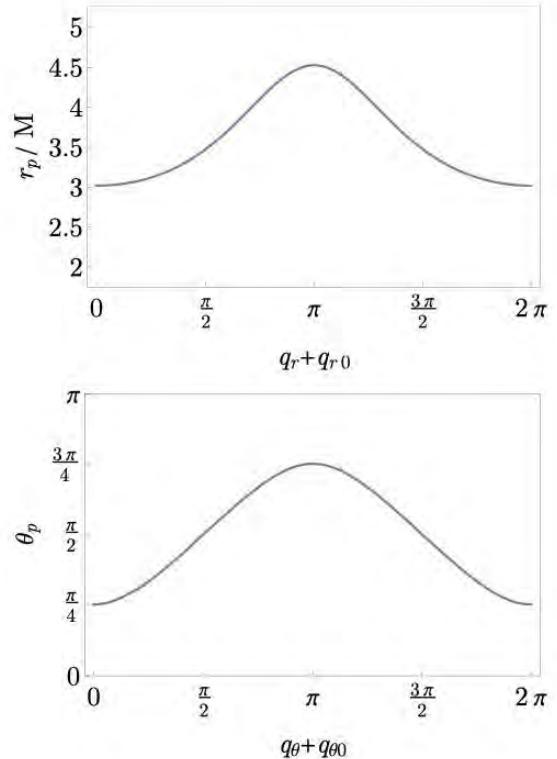
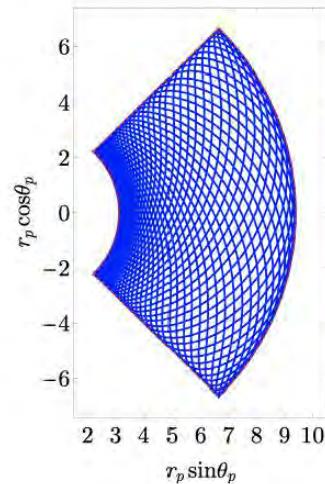
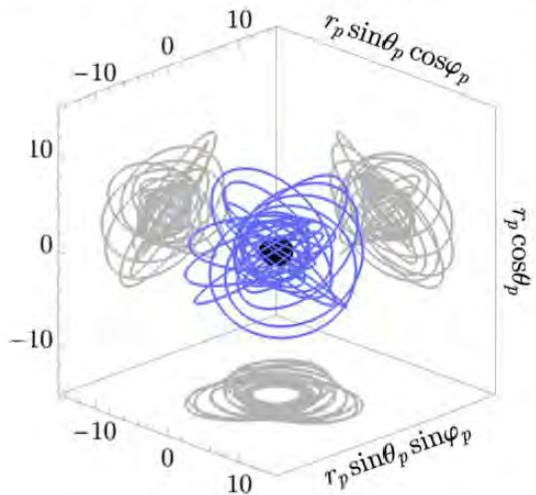
Impact of initial phases

- Consider non-resonant inclined, eccentric geodesic
 - Radial & polar motion separate w/ (Carter-)Mino time λ
 - Define angle variables $q_r \equiv \Upsilon_r \lambda, q_\theta \equiv \Upsilon_\theta \lambda$



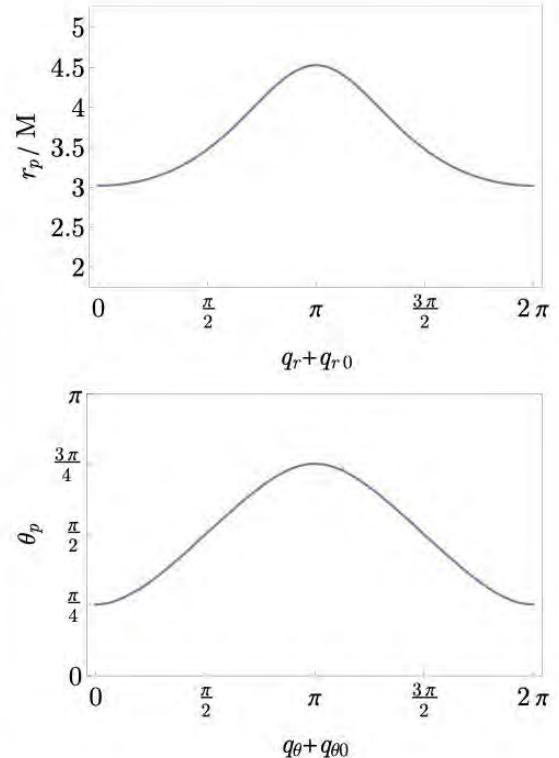
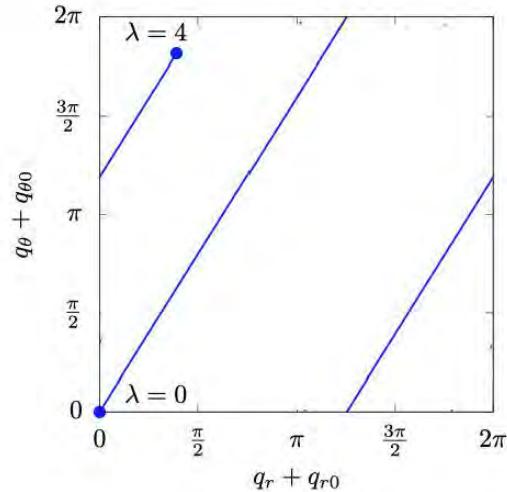
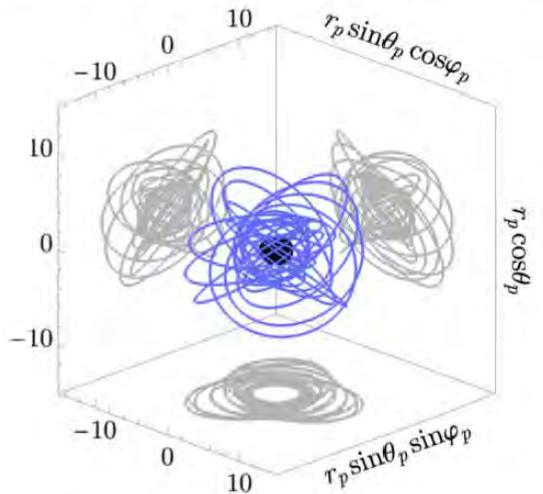
Impact of initial phases

- Consider non-resonant inclined, eccentric geodesic
 - Radial & polar motion separate w/ (Carter-)Mino time λ
 - Define angle variables $q_r \equiv \Upsilon_r \lambda, q_\theta \equiv \Upsilon_\theta \lambda$



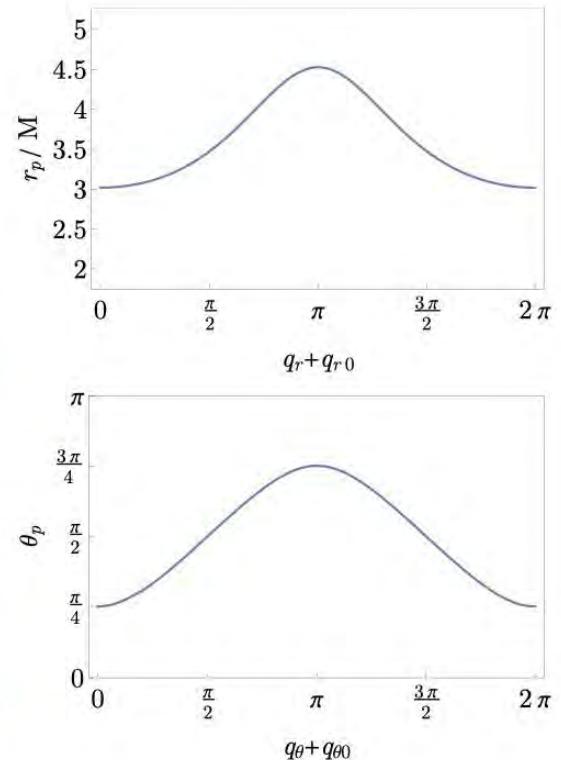
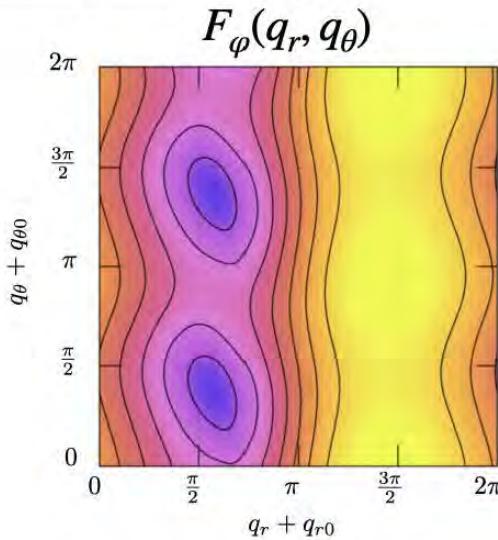
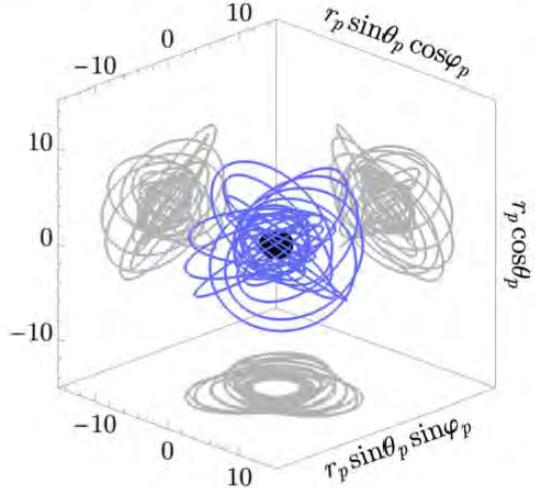
Impact of initial phases

- Consider non-resonant inclined, eccentric geodesic
 - Radial & polar motion separate w/ (Carter-)Mino time λ
 - Define angle variables $q_r \equiv \Upsilon_r \lambda, q_\theta \equiv \Upsilon_\theta \lambda$



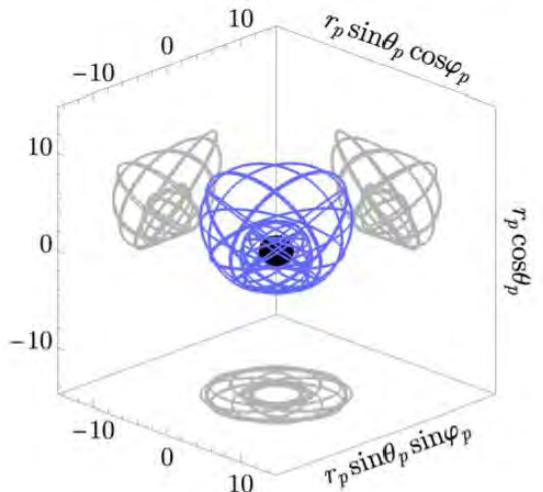
Impact of initial phases

- Consider non-resonant inclined, eccentric geodesic
 - Radial & polar motion separate w/ (Carter-)Mino time λ
 - Define angle variables $q_r \equiv \Upsilon_r \lambda, q_\theta \equiv \Upsilon_\theta \lambda$



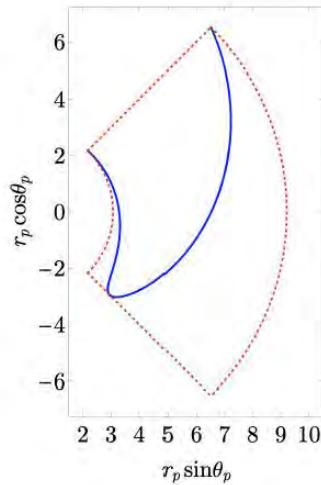
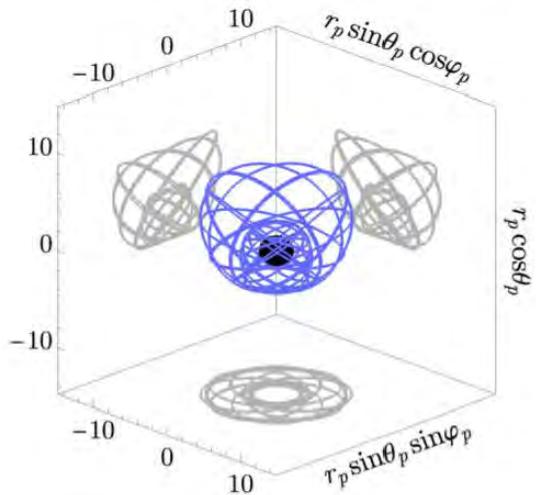
Impact of initial phases

- Consider resonant inclined, eccentric geodesic
 - Single resonant frequency $\rightarrow \Upsilon \equiv \Upsilon_r/\beta_r = \Upsilon_\theta/\beta_\theta$
 - Define $\bar{q} \equiv \Upsilon\lambda = q_r/\beta_r = q_\theta/\beta_\theta$, $\bar{q}_0 \equiv q_{\theta 0}/\beta_\theta - q_{r0}/\beta_r$



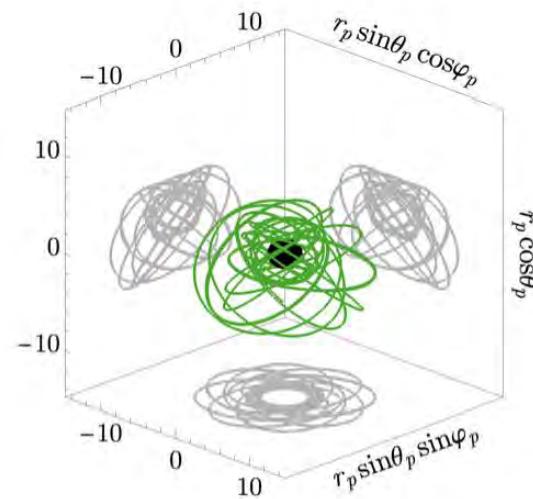
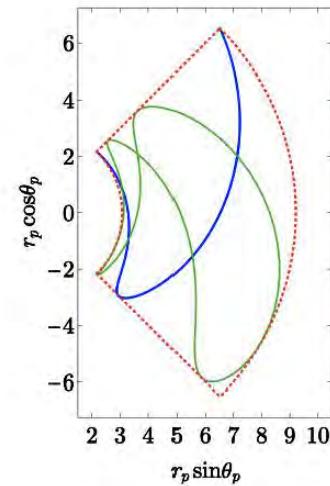
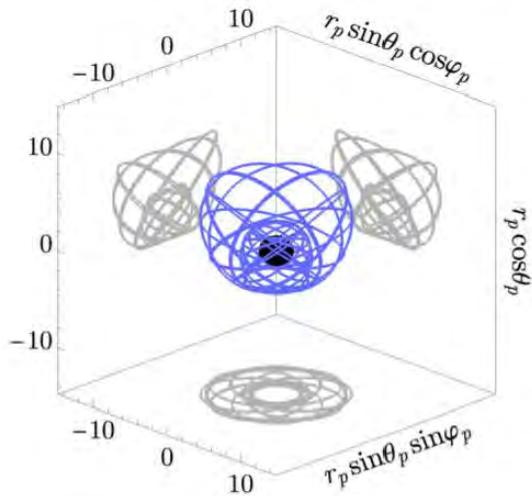
Impact of initial phases

- Consider resonant inclined, eccentric geodesic
 - Single resonant frequency $\rightarrow \Upsilon \equiv \Upsilon_r/\beta_r = \Upsilon_\theta/\beta_\theta$
 - Define $\bar{q} \equiv \Upsilon\lambda = q_r/\beta_r = q_\theta/\beta_\theta$, $\bar{q}_0 \equiv q_{\theta 0}/\beta_\theta - q_{r0}/\beta_r$



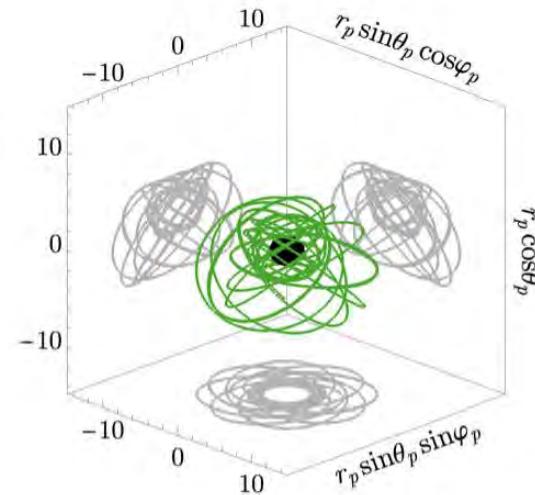
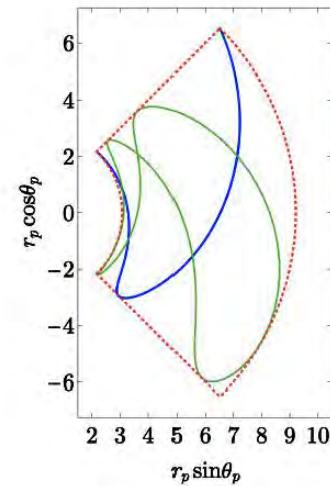
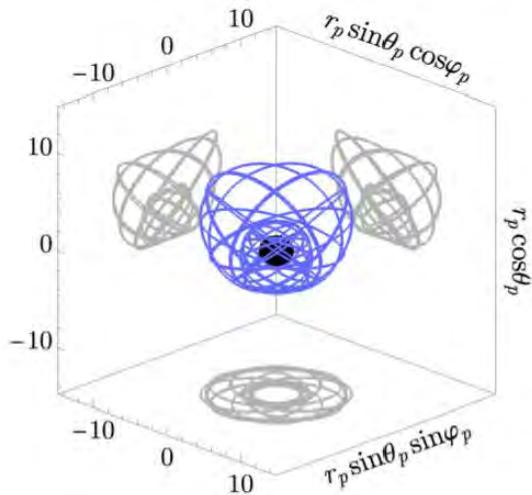
Impact of initial phases

- Consider resonant inclined, eccentric geodesic
 - Single resonant frequency $\rightarrow \Upsilon \equiv \Upsilon_r/\beta_r = \Upsilon_\theta/\beta_\theta$
 - Define $\bar{q} \equiv \Upsilon\lambda = q_r/\beta_r = q_\theta/\beta_\theta$, $\bar{q}_0 \equiv q_{\theta 0}/\beta_\theta - q_{r0}/\beta_r$



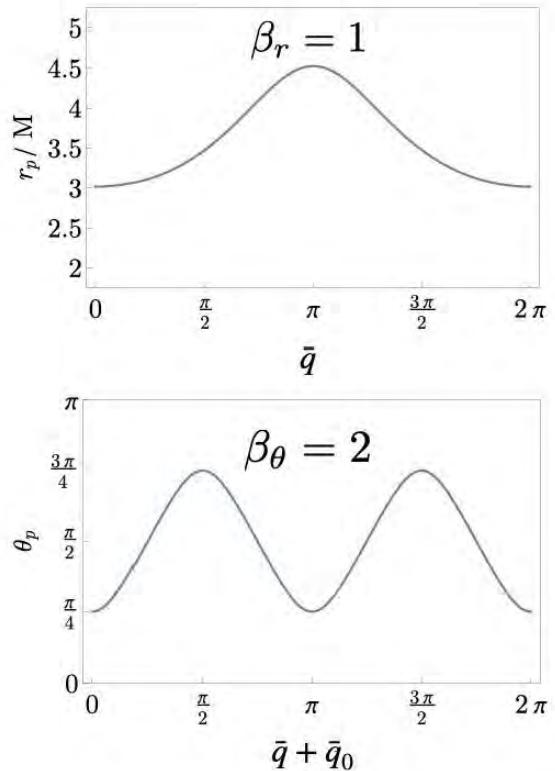
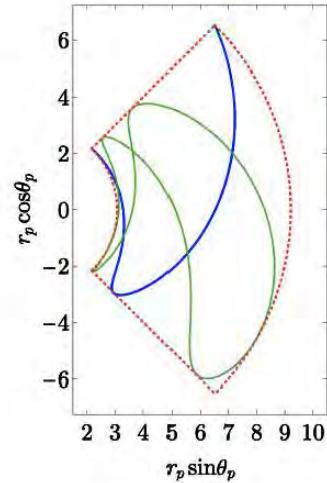
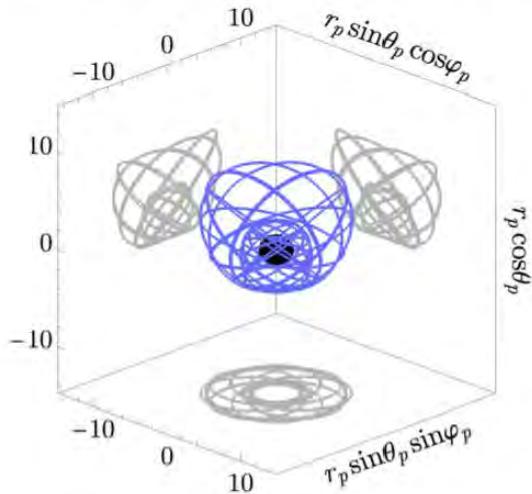
Impact of initial phases

- Consider resonant inclined, eccentric geodesic
 - Single resonant frequency $\rightarrow \Upsilon \equiv \Upsilon_r/\beta_r = \Upsilon_\theta/\beta_\theta$
 - Define $\bar{q} \equiv \Upsilon\lambda = q_r/\beta_r = q_\theta/\beta_\theta$, $\bar{q}_0 \equiv q_{\theta 0}/\beta_\theta - q_{r0}/\beta_r$



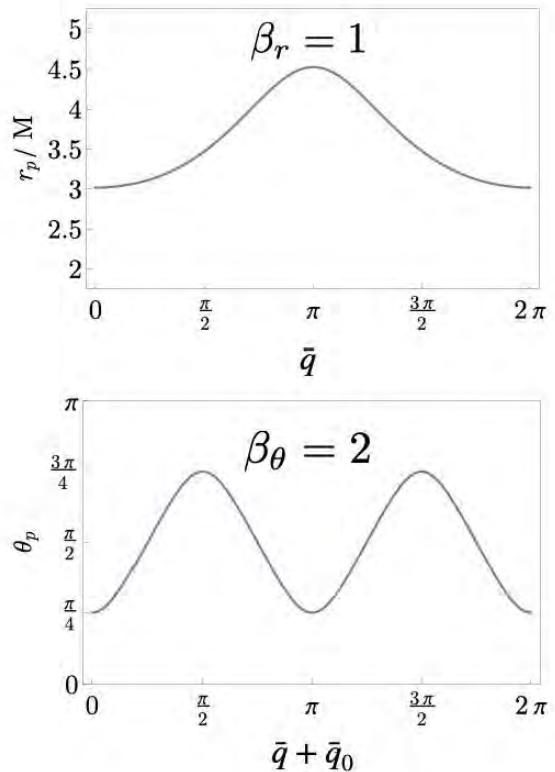
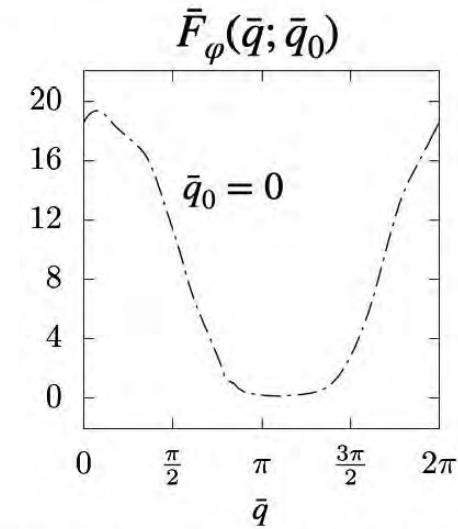
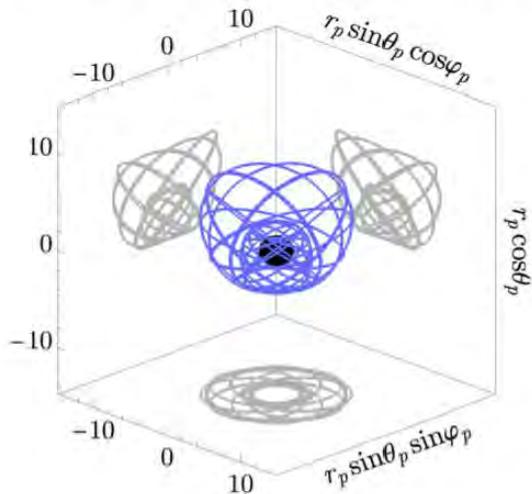
Impact of initial phases

- Consider resonant inclined, eccentric geodesic
 - Single resonant frequency $\rightarrow \Upsilon \equiv \Upsilon_r/\beta_r = \Upsilon_\theta/\beta_\theta$
 - Define $\bar{q} \equiv \Upsilon\lambda = q_r/\beta_r = q_\theta/\beta_\theta, \bar{q}_0 \equiv q_{\theta 0}/\beta_\theta - q_{r0}/\beta_r$



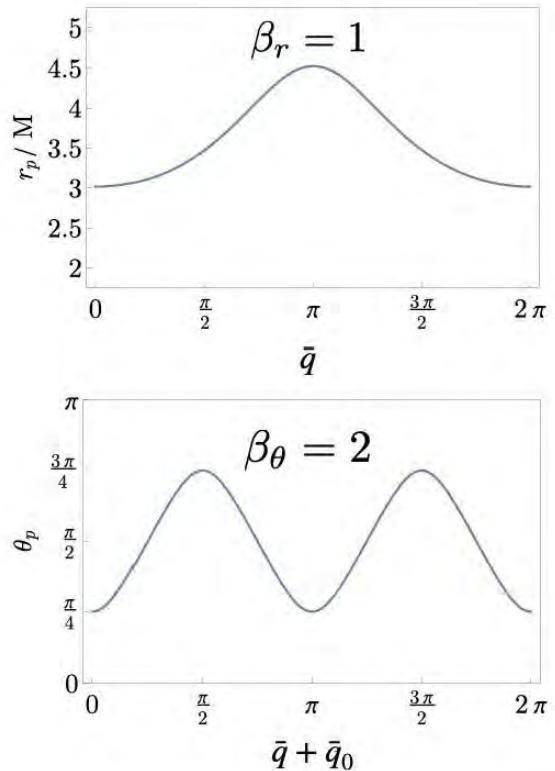
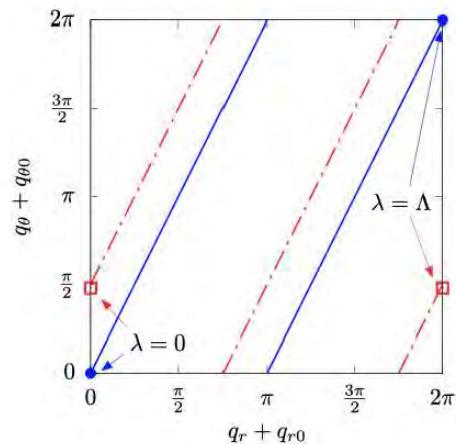
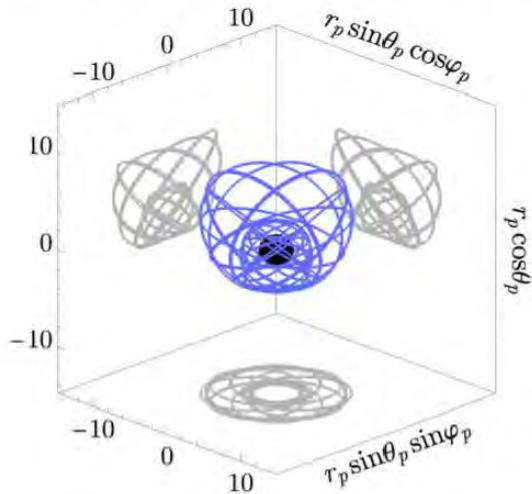
Impact of initial phases

- Consider resonant inclined, eccentric geodesic
 - Single resonant frequency $\rightarrow \Upsilon \equiv \Upsilon_r/\beta_r = \Upsilon_\theta/\beta_\theta$
 - Define $\bar{q} \equiv \Upsilon\lambda = q_r/\beta_r = q_\theta/\beta_\theta, \bar{q}_0 \equiv q_{\theta 0}/\beta_\theta - q_{r0}/\beta_r$



Impact of initial phases

- Consider resonant inclined, eccentric geodesic
 - Single resonant frequency $\rightarrow \Upsilon \equiv \Upsilon_r/\beta_r = \Upsilon_\theta/\beta_\theta$
 - Define $\bar{q} \equiv \Upsilon\lambda = q_r/\beta_r = q_\theta/\beta_\theta, \bar{q}_0 \equiv q_{\theta 0}/\beta_\theta - q_{r0}/\beta_r$



Scalar self-force for resonant orbits

- Constructing SSF

- Calculate SSF just as if resonance is a non-resonant geodesic with $q_{r0} = q_{\theta0} = 0$
- Choose value of \bar{q}_0 to choose a particular resonant orbit
- Sample SSF along flow of this orbit on the torus



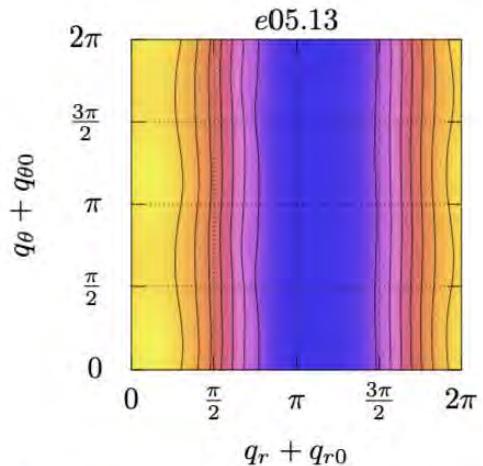
Scalar self-force for resonant orbits

- Constructing SSF
 - Calculate SSF just as if resonance is a non-resonant geodesic with $q_{r0} = q_{\theta0} = 0$
 - Choose value of \bar{q}_0 to choose a particular resonant orbit
 - Sample SSF along flow of this orbit on the torus



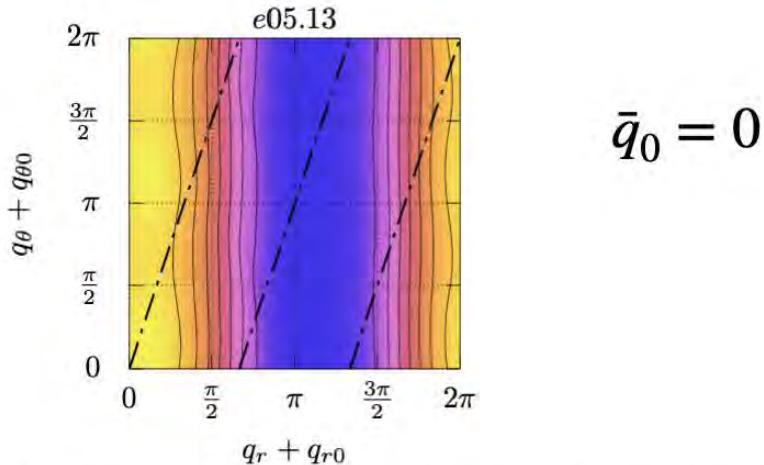
Scalar self-force for resonant orbits

- Constructing SSF
 - Calculate SSF just as if resonance is a non-resonant geodesic with $q_{r0} = q_{\theta0} = 0$
 - Choose value of \bar{q}_0 to choose a particular resonant orbit
 - Sample SSF along flow of this orbit on the torus



Scalar self-force for resonant orbits

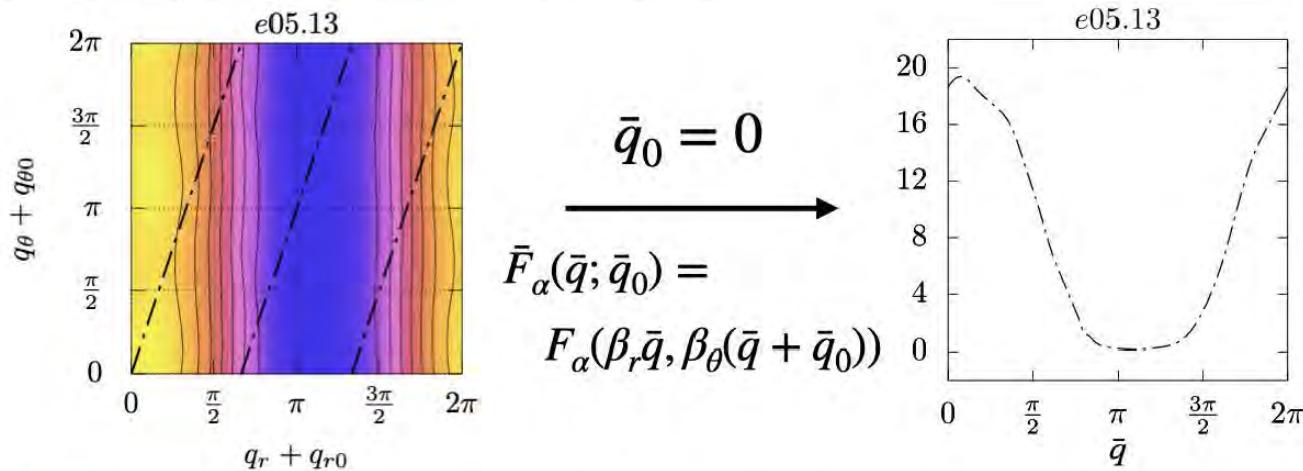
- Constructing SSF
 - Calculate SSF just as if resonance is a non-resonant geodesic with $q_{r0} = q_{\theta0} = 0$
 - Choose value of \bar{q}_0 to choose a particular resonant orbit
 - Sample SSF along flow of this orbit on the torus



Scalar self-force for resonant orbits

- Constructing SSF

- Calculate SSF just as if resonance is a non-resonant geodesic with $q_{r0} = q_{\theta0} = 0$
- Choose value of \bar{q}_0 to choose a particular resonant orbit
- Sample SSF along flow of this orbit on the torus



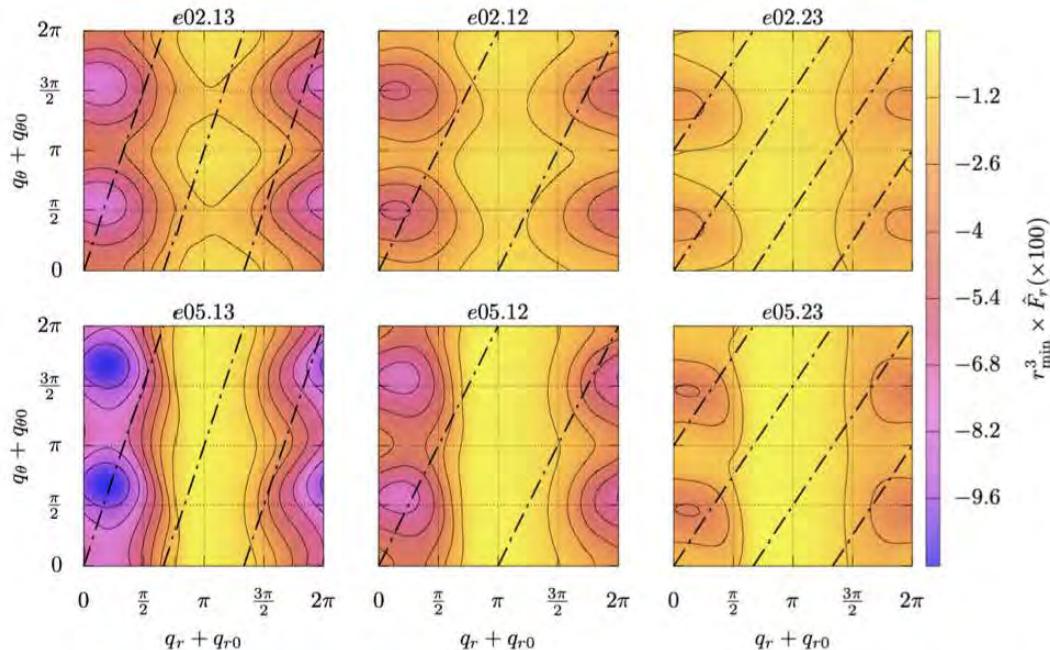
Scalar self-force for resonant orbits

- Analyzed 1:3, 1:2, & 2:3 resonances for $e = 0.2$ & $e = 0.5$ ($a = 0.9, x = \cos \pi/4$)



Scalar self-force for resonant orbits

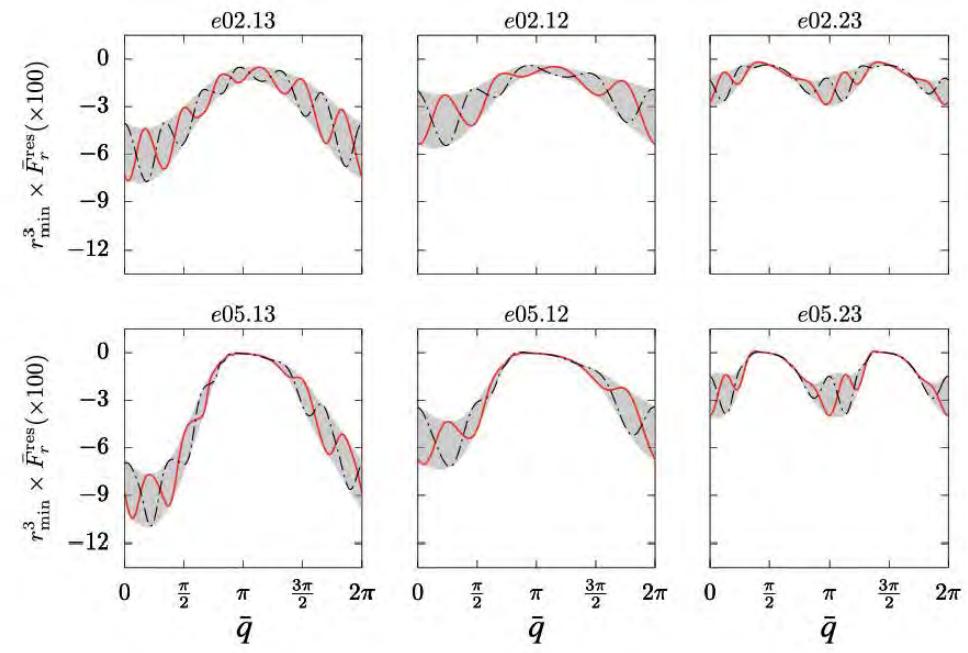
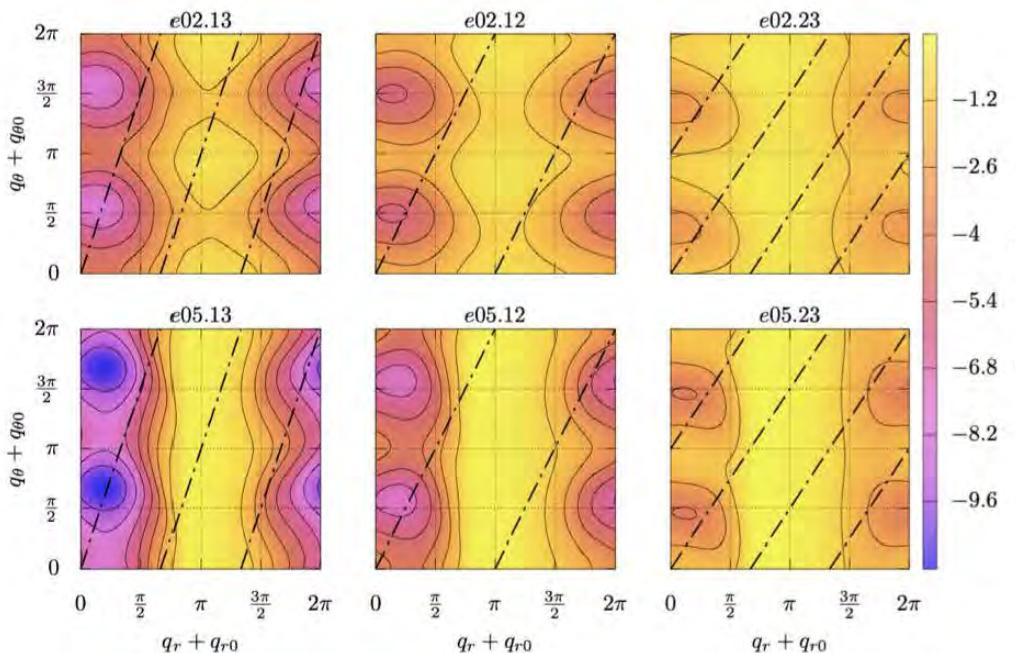
- Analyzed 1:3, 1:2, & 2:3 resonances for $e = 0.2$ & $e = 0.5$ ($a = 0.9, x = \cos \pi/4$)



THE UNIVERSITY
of NORTH CAROLINA
at CHAPEL HILL

Scalar self-force for resonant orbits

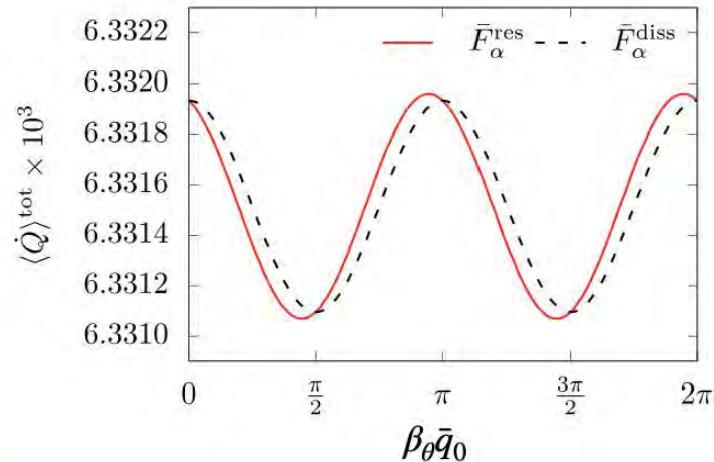
- Analyzed 1:3, 1:2, & 2:3 resonances for $e = 0.2$ & $e = 0.5$ ($a = 0.9$, $x = \cos \pi/4$)



THE UNIVERSITY
of NORTH CAROLINA
at CHAPEL HILL

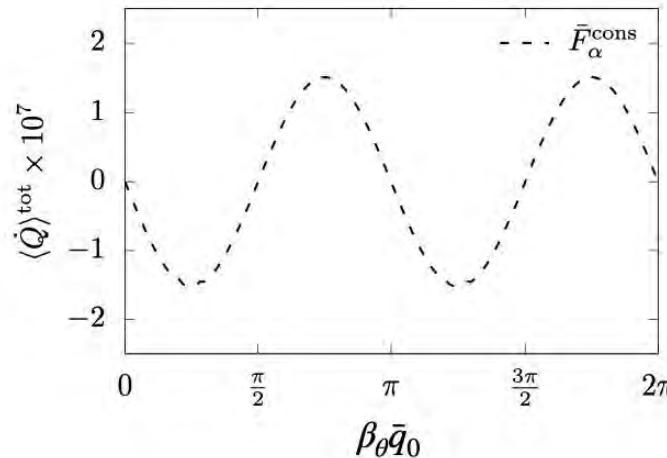
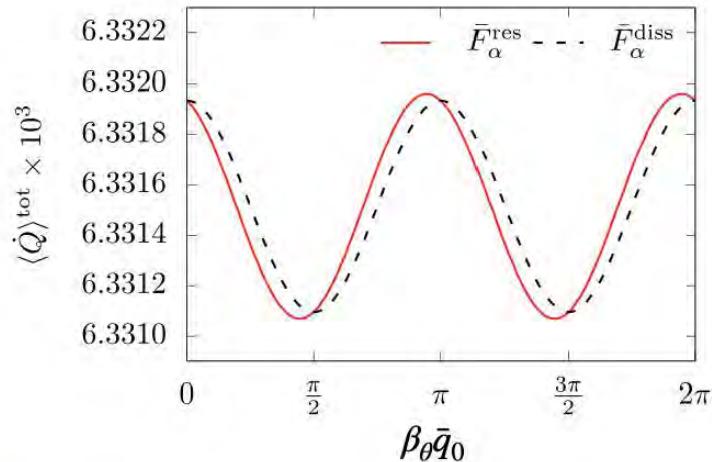
Testing the integrability conjecture

- Conservative contributions to $\langle \dot{Q} \rangle^{\text{tot}}$
 - Potential conservative contribution suppressed by 4 orders of magnitude
 - Numerical error estimated to be 5×10^{-7} ; consistent w/ zero



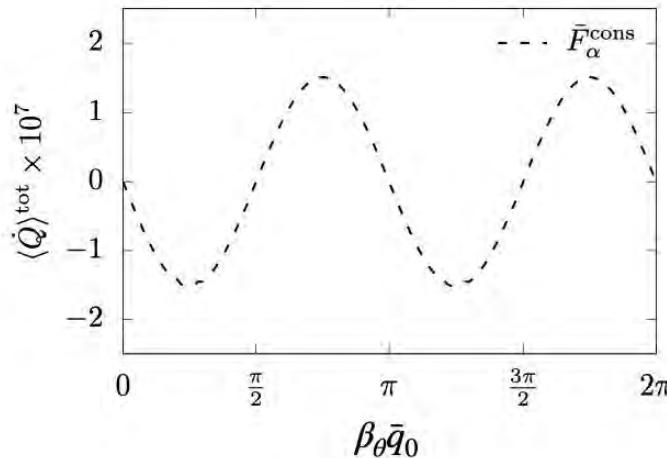
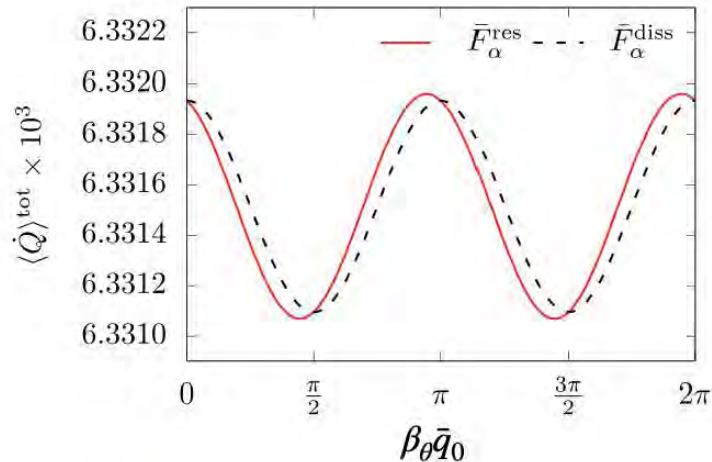
Testing the integrability conjecture

- Conservative contributions to $\langle \dot{Q} \rangle^{\text{tot}}$
 - Potential conservative contribution suppressed by 4 orders of magnitude
 - Numerical error estimated to be 5×10^{-7} ; consistent w/ zero



Testing the integrability conjecture

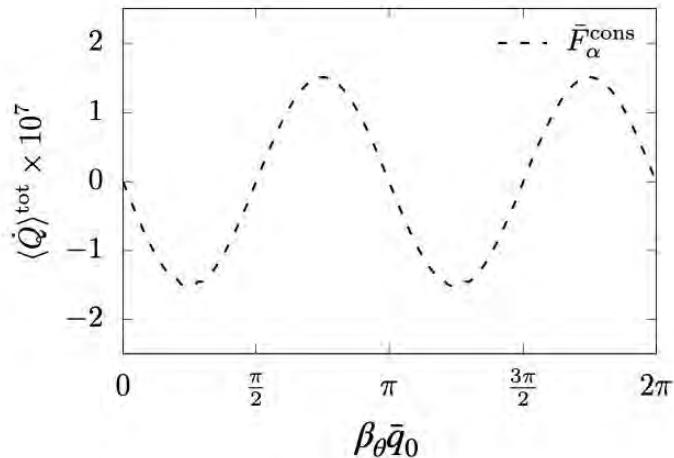
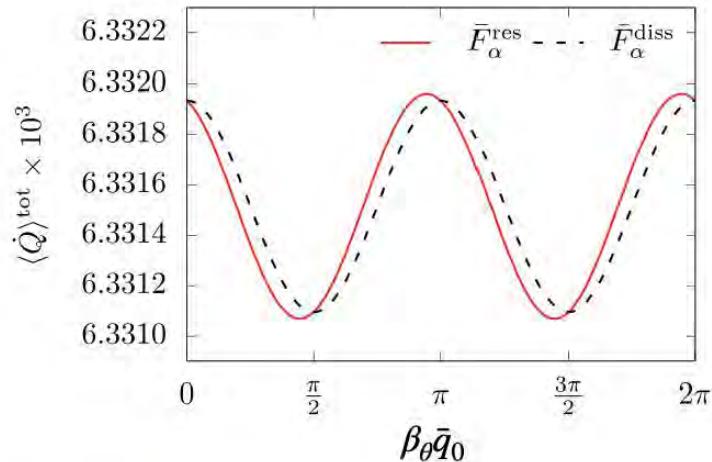
- Conservative contributions to $\langle \dot{Q} \rangle^{\text{tot}}$
 - Potential conservative contribution suppressed by 4 orders of magnitude
 - Numerical error estimated to be 5×10^{-7} ; consistent w/ zero



Testing the integrability conjecture

- Conservative contributions to $\langle \dot{Q} \rangle^{\text{tot}}$
 - Potential conservative contribution suppressed by 4 orders of magnitude
 - Numerical error estimated to be 5×10^{-7} ; consistent w/ zero

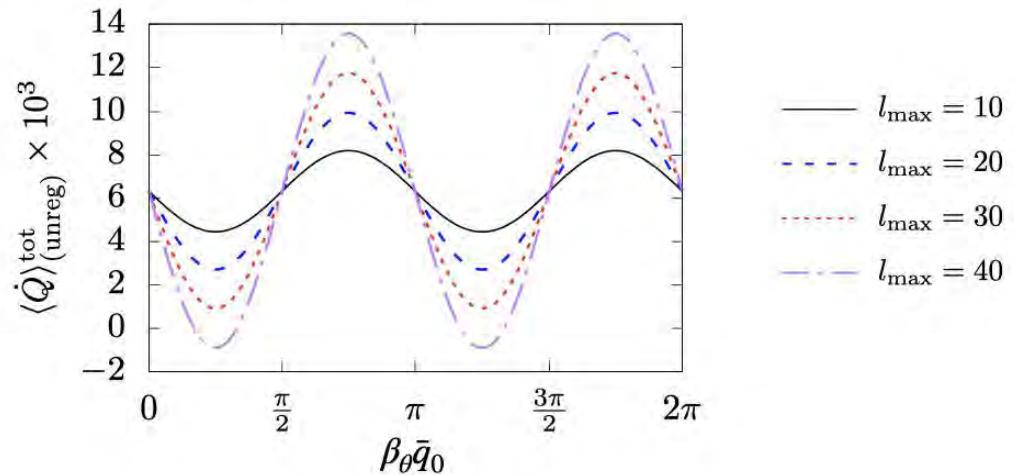
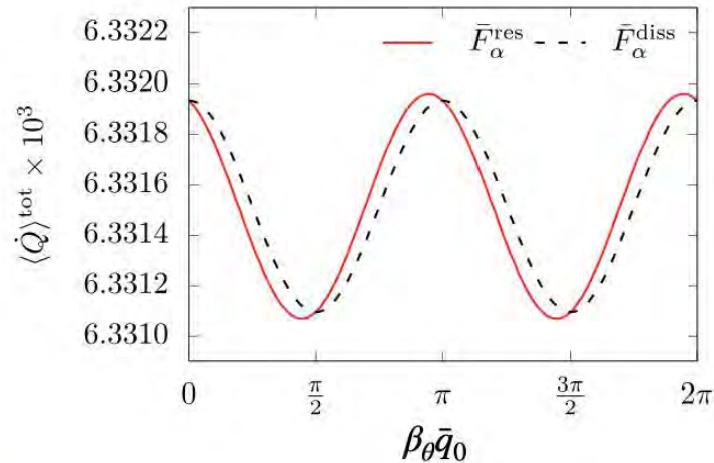
Conservative contribution is negligible



Testing the integrability conjecture

- Conservative contributions to $\langle \dot{Q} \rangle^{\text{tot}}$
 - Potential conservative contribution suppressed by 4 orders of magnitude
 - Numerical error estimated to be 5×10^{-7} ; consistent w/ zero

Conservative contribution is negligible



Summary

- Key results from scalar model
 - SSF & GSF can be calculated exact same way for resonant sources as non-resonant (in FD)
 - In fact, sampling in q_r and q_θ more efficient than sampling in \bar{q} and \bar{q}_0
 - Integrability conjecture still holds - conservative contribution to $\langle \dot{Q} \rangle^{\text{tot}}$ negligible within errors of code
 - Need improved regularization scheme to improve numerical analysis



Summary

- Key results from scalar model
 - SSF & GSF can be calculated **exact same way** for resonant sources as non-resonant (in FD)
 - In fact, sampling in q_r and q_θ more efficient than sampling in \bar{q} and \bar{q}_0
 - Integrability conjecture still holds - conservative contribution to $\langle \dot{Q} \rangle^{\text{tot}}$ negligible within errors of code
 - Need improved regularization scheme to improve numerical analysis



Summary

- Key results from scalar model
 - SSF & GSF can be calculated **exact same way** for resonant sources as non-resonant (in FD)
 - In fact, sampling in q_r and q_θ more efficient than sampling in \bar{q} and \bar{q}_0
 - Integrability conjecture still holds - conservative contribution to $\langle \dot{Q} \rangle^{\text{tot}}$ negligible within errors of code
 - Need improved regularization scheme to improve numerical analysis

Summary

- Key results from scalar model
 - SSF & GSF can be calculated **exact same way** for resonant sources as non-resonant (in FD)
 - In fact, sampling in q_r and q_θ more efficient than sampling in \bar{q} and \bar{q}_0
 - Integrability conjecture still holds - conservative contribution to $\langle \dot{Q} \rangle^{\text{tot}}$ negligible within errors of code
 - Need improved regularization scheme to improve numerical analysis

Questions?

Scalar self-force (SSF) problem

- Gravitational case ($g_{\mu\nu} = g_{\mu\nu}^{\text{Kerr}} + \epsilon h_{\mu\nu}$)

- Equations of motion

$$\mu u^\beta \nabla_\beta u^\alpha = \epsilon F_{\text{GSF}}^\alpha$$

- Gravitational self-force (GSF) equations

$$F_{\text{GSF}}^\alpha = \mu \nabla^{\alpha\beta\gamma} h_{\beta\gamma}^R$$

- Field equations & metric reconstruction

$$\hat{\mathcal{O}}_{\pm 2} \Psi_{\pm 2} = T_{\pm 2} \implies h_{\alpha\beta}^{\text{ret}} = \hat{\mathcal{S}}_{\pm 2}^\dagger \Psi_{\pm 2}^{\text{Hertz}} + h_{\alpha\beta}^{\text{comp}}$$

- Regularization:

$$h_{\alpha\beta}^R = h_{\alpha\beta}^{\text{ret}} - h_{\alpha\beta}^S$$

- Scalar case ($h_{\mu\nu} \rightarrow 0, q$)

- Equations of motion

$$u^\beta \nabla_\beta (\mu u^\alpha) = (q/M) F^\alpha$$

- Scalar self-force (SSF) equations

$$F^\alpha = q \nabla^\alpha \Phi^R$$

- Field equations

$$\hat{\mathcal{O}}_0 \Psi_0 = T_0 \implies \square_g \Phi = -4\pi T_0$$

- Regularization:

$$\Phi^R = \Phi^{\text{ret}} - \Phi^S$$

Impact of initial phases

- Self-force for generic orbits

- Tri-periodic discrete spectrum

$$\Upsilon_{mkn} = m\Upsilon_\varphi + k\Upsilon_\theta + n\Upsilon_r$$

- Separate radial and polar dependence

$$F^\alpha(q_r, q_\theta) = \sum_{\hat{l}lmkn} Z_{\hat{l}lmkn} f_{\hat{l}lmkn}^\alpha(q_r) g_{\hat{l}lmkn}^\alpha(q_\theta)$$

- $Z_{\hat{l}lmkn}$ related to source integration
 - Can express self-force as function of λ

$$F^\alpha(\lambda) = F^\alpha(\Upsilon_r \lambda, \Upsilon_\theta \lambda)$$

- Self-force for generic orbits

- Bi-periodic discrete spectrum

$$\Upsilon_{mkn} = m\Upsilon_\varphi + N\Upsilon$$

- Separate radial and polar dependence

$$\bar{F}^\alpha(\bar{q}; \bar{q}_0) = \sum_{\hat{l}lmN} \bar{Z}_{\hat{l}lmN}(\bar{q}_0) \bar{f}_{\hat{l}lmN}^\alpha(\bar{q} + \bar{q}_0) \bar{g}_{\hat{l}lmN}^\alpha(\bar{q})$$

- Can express resonant self-force in terms of non-resonant case

$$\bar{F}^\alpha(\bar{q}; \bar{q}_0) = F^\alpha(\beta_r \bar{q}, \beta_\theta \bar{q} + \beta_\theta \bar{q}_0)$$

