Universality Crossover in Parametric (ch)RMT and Lattice Dirac Spectrum

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#### PLAN

- Intro/Motivation

   universality crossover
   Dirac spectrum vs chRMT
   NLσM and universality
- 2. Parametric (chiral) RMT level spacing / smallest EV
- 3. SU(2)×U(1) Dirac Spectrum random flux × noise model (check) pure LGT scaling and low-energy constant



··· analytical

numerical

## 1. Universality of spectral fluctuation

energy levels of excited nuclei

Wigner 1950s



## Universality of spectral fluctuation



# Universality classes

 $\sigma$ : **C**-herm.  $\rightarrow$  **C**-herm.,  $\sigma^2 = 1$ Involutive symmetry  $H = \sigma(H)$ 

. . .

none	C-herm.	GUE
$\sigma(H) = H^*$	<b>R</b> -sym.	GOE
$\sigma(H) = \tau_2 H^* \tau_2^{-1}$	H-selfdual	GSE
$\sigma(H) = -\gamma_5 H \gamma_5^{-1}$		chG*E

classification by Riemann symm. spaces Zimbauer 96

. . .

# Universality classes

 $\sigma$ : C-herm.  $\rightarrow$  C-herm.,  $\sigma^2 = 1$ Involutive symmetry  $H = \sigma(H)$ 

$$D_{\mathrm{U}(1)} = \gamma_{\mu} \left( i\partial_{\mu} + A_{\mu} \right) = -\gamma_{5} D_{\mathrm{U}(1)} \gamma_{5}^{-1} \quad \text{same symm. as chGUE}$$
$$D_{\mathrm{SU}(2) \text{ fnd}} = \gamma_{\mu} \left( i\partial_{\mu} + A_{\mu}^{a} \tau_{a} \right) = (\tau_{2} C) D_{\mathrm{SU}(2) \text{ fnd}}^{*} (\tau_{2} C)^{-1} \quad \text{chGOE}$$
$$D_{\mathrm{SU}(N) \text{ adj}} = \gamma_{\mu} \left( i\partial_{\mu} + A_{\mu}^{a} f_{a} \right) = C D_{\mathrm{SU}(N) \text{ adj}}^{*} C^{-1} \quad \text{chGSE}$$

expect to agree w/ chG#E Verbaarschot 94

## Dirac spectrum

#### SU(2) fund KS

Berbenni-Bitsch et al 97



## Smallest Dirac EV

SU(2 or 3) overlap Edwards et al 99



match w/ chGOE, chGUE, chGSE

#### Origin of universality

spectral correlator NLoM

$$\left\langle \det(\lambda - H) \cdots \right\rangle \xrightarrow{\text{RSS}} DQ \exp\left\{ -\frac{\pi}{4\Delta V} \int d\mathbf{r} \left( \frac{D}{2} \operatorname{str} |\nabla Q|^2 - i \operatorname{str} \Lambda Q \right) \right\}$$
$$\downarrow \quad \epsilon \text{ regime}$$
$$\int_{\text{RSS}} dQ \exp\left\{ i \frac{\pi}{4\Delta} \operatorname{str} \Lambda Q \right\}$$

---> QCD : flavor SSB Weinberg 60s, Leutwyler-Smiga 92

- AH : H-S transf, sdl pt Wegner 81, Efetov 83
   QCD : strong coupling Kawamoto-Smit 81, Nagao-Nishigaki 01 gauge/gravity ?
- RMT : H-S transf Efetov 83
   QCh : sum over periodic orbit pairs Muller et al 09

### Origin of universality

spectral correlator NLoM

$$\left\langle \det(\lambda - D) \cdots \right\rangle \xrightarrow{\text{SU}(N_{\text{F}})} DU \exp\left\{-\int dx \left(\frac{f_{\pi}^{2}}{2} \operatorname{tr} \left|\partial_{\mu}U\right|^{2} - i \Sigma \operatorname{Re} \operatorname{tr} \Lambda U\right)\right\}$$
$$\downarrow \quad \epsilon \operatorname{regime}$$
$$\int_{\operatorname{SU}(N_{\text{F}})} dU \exp\left\{i V\Sigma \operatorname{Re} \operatorname{tr} \Lambda U\right\}$$

- ---> QCD : flavor SSB Weinberg 60s, Leutwyler-Smiga 92
- AH : H-S transf, sdl pt Wegner 81, Efetov 83
  QCD : strong coupling Kawamoto-Smit 81, Nagao-Nishigaki 01
  gauge/gravity ?
- RMT technical advantage (orth. polyn.)
  QCh : sum over periodic orbit pairs Muller et al 09

## 2. Parametric RMT

symmetry breaking perturb.

$$H = H_0 + \alpha H_1, \quad \sigma(H_0) = H_0, \quad \sigma(H_1) \neq H_1$$
  

$$\alpha = O(\Delta \sqrt{V})^{<1} \Rightarrow H_0 \text{ dominates}$$
  

$$\alpha = O(\Delta \sqrt{V})^{>1} \Rightarrow H_1 \text{ dominates}$$
  

$$\alpha = O(\Delta \sqrt{V})^1 \Rightarrow \text{ crossover, parametrized by } \rho = \frac{\alpha}{\Delta \sqrt{V}}$$

nonchiralGOE, GSE $\rightarrow$ GUEPandey Mehta 83chiralchGOE, chGSE $\rightarrow$ chGUENagao 03supercondGUE, CI $\rightarrow$ CKoziy Skvortsov 11

## Parametric RMT

spectral correlator NL<sub>0</sub>M

$$\left\langle \det(\lambda - H) \cdots \right\rangle \xrightarrow{\qquad} \int_{\text{RSS}} dU \ e^{-S[U]}$$
$$S[U] = i \frac{\pi}{2} \ \text{tr} \underbrace{X}_{\text{unf. EV}} U \quad - \underbrace{\frac{\pi^2 \rho^2}{4}}_{\text{Xover parameter}} \ \text{tr} \left(U\tau_3\right)^2$$

• the same NL $\sigma$ M as 2C chiral L at finite density

- fit the spectrum to pRMT = measure the LEC  $\rho$
- ---> QCD : flavor SSB Kogut-Stephanov-Toublan 00, Dunne-Nishigaki 03
- → AH · RMT : H-S transf Altland-Iida-Efetov 93 QCD [SC, gauge/gravity] ?

QCh : sum over p.o. pairs Saito-Nagao-Muller-Braun 09

### Fit to param. RMT

Dupuis-Montambaux 91



caution: P(s) curve is just Wigner-like approximant

counterpart in SU(2)×U(1) Dirac spectra?



## Parametric chRMT : the results

by using Nystrøm-type method with Gauss-Legendre quadratures of computing Fredholm Pfaffian of the matrix-valued dynamical Bessel kernels, I have computed and plotted the 1<sup>st</sup> EV distributions for chGOE-chGUE, chGSE-chGUE xover

$$H = \begin{pmatrix} \mathbf{R} \\ \mathbf{R}^T \end{pmatrix} + \alpha \begin{pmatrix} \mathbf{C} \\ \mathbf{C}^+ \end{pmatrix}, \begin{pmatrix} \mathbf{H} \\ \mathbf{H}^D \end{pmatrix} + \alpha \begin{pmatrix} \mathbf{C} \\ \mathbf{C}^+ \end{pmatrix}$$

**R**, **C**, **H**: Gauss random  $N \times N$ 

#### chGOE-chGUE

#### chGSE-chGUE



#### **GOE-GUE**

**GSE-GUE** 



# 3. SU(2)×U(1) Lattice Dirac Spectrum

Open question : does the chG#E-chGUE xover occur in physics?

partly inspired by the recent trend on isospin-related physics : QCD+QED simulation, measure Dirac spectra of 2C-QCD+QED

$$D = \gamma_{\mu} \left( i \partial_{\mu} + g A^{a}_{\mu} T_{a} + e A_{\mu} \right)$$

U(1) : symmetry breaker

- can be fitted to the parametric chRMT at all?
- can determine the symmetry-breaking LEC  $\rho$  accurately?
- if yes, does  $\rho$  scale as expected?

## SU(2)×U(1) Lattice Dirac Spectrum

Strategy : smaller lattice  $V=4^4 \sim 6^4$  , more samples  $O(10^4)$ 

• toy model for scaling check: SU(2) random x U(1) noise  $\theta_{x,\mu} \in [-p, \dot{p}]$ .d.  $\beta_2 = 0, \quad p = 0.01 \sim 10$ • quenched SU(2) x U(1)  $\{\beta_2 = 0, 0.5, 1\} \times \{\beta_1 = 1000 \sim 3000\}$ strong cpl extremely weak cpl

#### global SU(2) Dirac spectrum



- global spectrum of SU(2) intact w/ weak U(1)
- parameter  $\rho$  depend on mean spacing  $\Delta \Rightarrow$  choose a plateau from the global spectrum for best bulk result

SU(2)×U(1) noise model - bulk

$$V = 4^4$$
,  $N_{\text{conf}} = 4e4$ ,  $\beta_2 = 0$ ,  $p = 0.02 \sim 10$ 



SU(2)×U(1) noise model - origin

$$V = 4^4$$
,  $N_{\text{conf}} = 4e4$ ,  $\beta_2 = 0$ ,  $p = 0.02 \sim 10$ 



SU(2)×U(1) noise model - bulk

$$V = 12^2 \sim 24^2$$
,  $N_{\text{conf}} = 8e3$ ,  $\beta_2 = 0$ ,  $p = 0.03$ 



#### SU(2)×U(1) noise model



- $\cdot$  perfectly fittable, accurate determine  $\rho$
- $\cdot$  Xover parameter ho scales as expected :

$$\rho = \operatorname{const} \frac{p}{\Delta \sqrt{V}}$$

SU(2)×U(1) LGT - bulk

 $V = 6^4$ ,  $N_{\text{conf}} = 1800$ ,  $\beta_2 = 1$ ,  $\beta_1 = 1000 \sim 3000$ 



SU(2)×U(1) LGT - origin  $V = 4^4$ ,  $N_{conf} = 40000$ ,  $\beta_2 = 1$ ,  $\beta_1 = 1000, 1500, 2000$ 



### $SU(2) \times U(1) LGT$



- $\cdot$  fittable, fair determin.  $\rho$
- $\cdot$  Xover parameter ho scales as expected :

$$\rho = \operatorname{const} \frac{\beta_1^{-1/2}}{\Delta \sqrt{V}}$$
  
LEC

# 4. Summary

- computed/plotted
- smallest EV distr for  $chG(O,S)E \rightarrow chGUE$  crossover
- level spacing distr for  $G(O,S)E \rightarrow GUE$  crossover
- $\Box$  measured the SU(2)×U(1) KS Dirac spectra for
- random SU(2) flux × U(1) noise toy model
- pure LGT at strong  $\beta_2 \times \text{very weak } \beta_1$
- fitting and scaling
- parametric (ch)RMT fits SU(2)×U(1) well in all cases
- accurate measurement of crossover parameter  $\rho$  expected scaling  $\Rightarrow$  nontrivial continuum limit
  - $\Rightarrow$  determine LEC for the symm break. term in  $L_{ch}$

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