

The Interpolating Function

Masazumi Honda



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References: MH,JHEP1412 019 (1408.2960),
MH-DPJ,NPB900 (2015) 533 (1504.02276),
AC-MH-ST and AC-MH, to appear,
MH, work in progress

based on collaborations with

Abhishek Chowdhury (HRI), Dileep P. Jatkar (HRI),
Somyadip Thakur (TIFR)

Perturbative expansion

- ubiquitous
- does not often give satisfactory understanding of physics...
(unless it has nice properties)
- even if it has nice property,
higher order computation is usually hard task

Today I consider somewhat limited situation:

We know

perturbative expansions around 2-points

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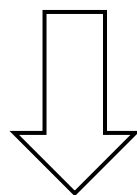
e.g. theories with S-duality, theories with gravity dual,
lattice field theory with weak & strong coupling expansions,
statistical systems with high & low temperature expansions, etc...

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perturbative expansions around **2**-points

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Approximation at **finite** values of parameters

How do we **interpolate** these two expansions?

Tool : Interpolating function

Single function consistent with the 2 perturbative expansions

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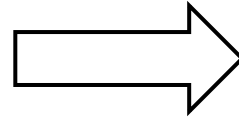
$$\left\{ \begin{array}{l} \text{Small-g exp.: } F(g) = g + \mathcal{O}(g^2) \\ \text{Large-g exp.: } F(g) = 2 + \mathcal{O}(g^{-1}) \end{array} \right.$$

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Interpolation



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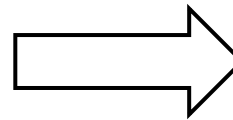
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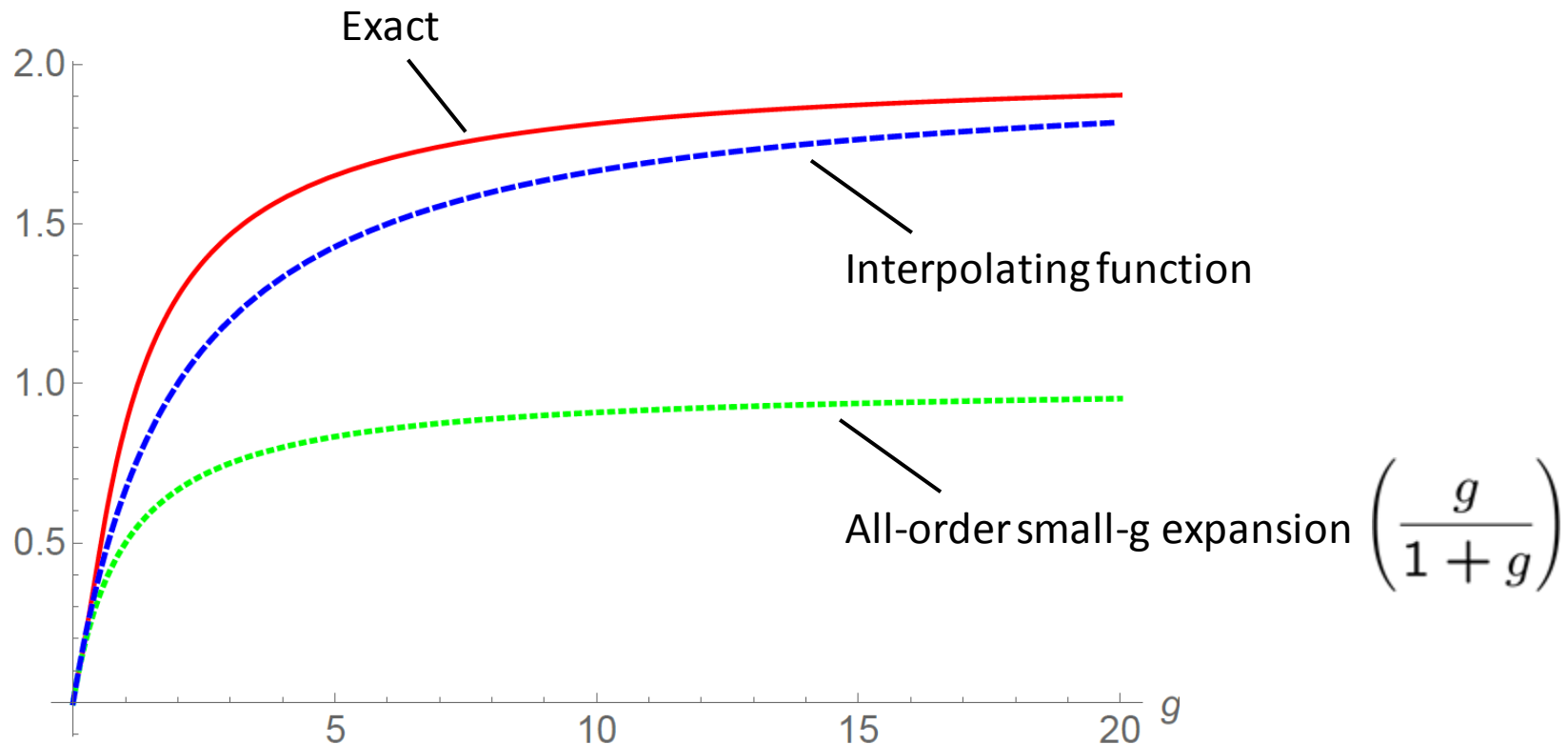
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Outline

- introduce a class of **interpolating functions** [MH'14]
 - generalization of Pade and Sen's interpolating function

- **Analytic property** of interpolating function & dimensions of **twist operators in planar ABJM**
 - indirect evidence for a recent conjecture on unknown function often called $h(\lambda)$, which appears in the context of integrability [Chowdhury-MH, to appear]

Introduction to Interpolating function

Setup

Suppose that **we know small-g and large-g expansions** of a function $F(g)$:

$$F(g) = \begin{cases} g^a(s_0 + s_1g + s_2g^2 + \dots), \\ g^b(l_0 + l_1g^{-1} + l_2g^{-2} + \dots), \end{cases}$$

Then we would like to find approximation of $F(g)$ at finite g .

(When we have expansions around $g=g_1$ and $g=g_2$,
changing the variable as $x=(g-g_1)/(g-g_2)$ gives small-x and large-x expansions)

(Two-point) Pade approximant

$$\mathcal{P}_{m,n}(g) = s_0 g^a \frac{1 + \sum_{k=1}^p c_k g^k}{1 + \sum_{k=1}^q d_k g^k},$$

$$\left(p = \frac{m + n + 1 + (b - a)}{2} \in \mathbf{Z}, \quad q = \frac{m + n + 1 - (b - a)}{2} \in \mathbf{Z} \right)$$

The coefficients are determined to reproduce the small-g exp. up to $\mathcal{O}(g^{a+m+1})$ and large-g exp. up to $\mathcal{O}(g^{b-n-1})$

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Some properties:

▪ can construct only for $(b-a) \in \mathbf{Z}$ (although avoidable by a change of variable)

$$\left[(b-a) : \text{even} \rightarrow (m+n) : \text{odd}, (b-a) : \text{odd} \rightarrow (m+n) : \text{even} \right]$$

▪ No branch cut

Fractional Power of Polynomial (FPP)

[Sen '13]

$$F_{m,n}(g) = s_0 g^a \left[1 + \sum_{k=1}^m c_k g^k + \sum_{k=0}^n d_k g^{m+n+1-k} \right]^{\frac{b-a}{m+n+1}}$$

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Some properties:

- can construct for arbitrary (a,b,m,n)
- Type of branch cut is uniquely determined by (a,b,m,n)

Fractional Power of Rational function (FPR)

$$F_{m,n}^{(\alpha)}(g) = s_0 g^a \left[\frac{1 + \sum_{k=1}^p c_k g^k}{1 + \sum_{k=1}^q d_k g^k} \right]^\alpha,$$

[MH'14]

$$\left(p = \frac{1}{2} \left(m + n + 1 + \frac{b-a}{\alpha} \right) \in \mathbf{Z}, \quad q = \frac{1}{2} \left(m + n + 1 - \frac{b-a}{\alpha} \right) \in \mathbf{Z} \right)$$

$$\alpha = \begin{cases} \frac{a-b}{2\ell+1} & \text{for } m+n : \text{ even} \\ \frac{a-b}{2\ell} & \text{for } m+n : \text{ odd} \end{cases}, \quad \text{with } \ell \in \mathbf{Z}.$$

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There are many cases where FPR gives very precise approximation.

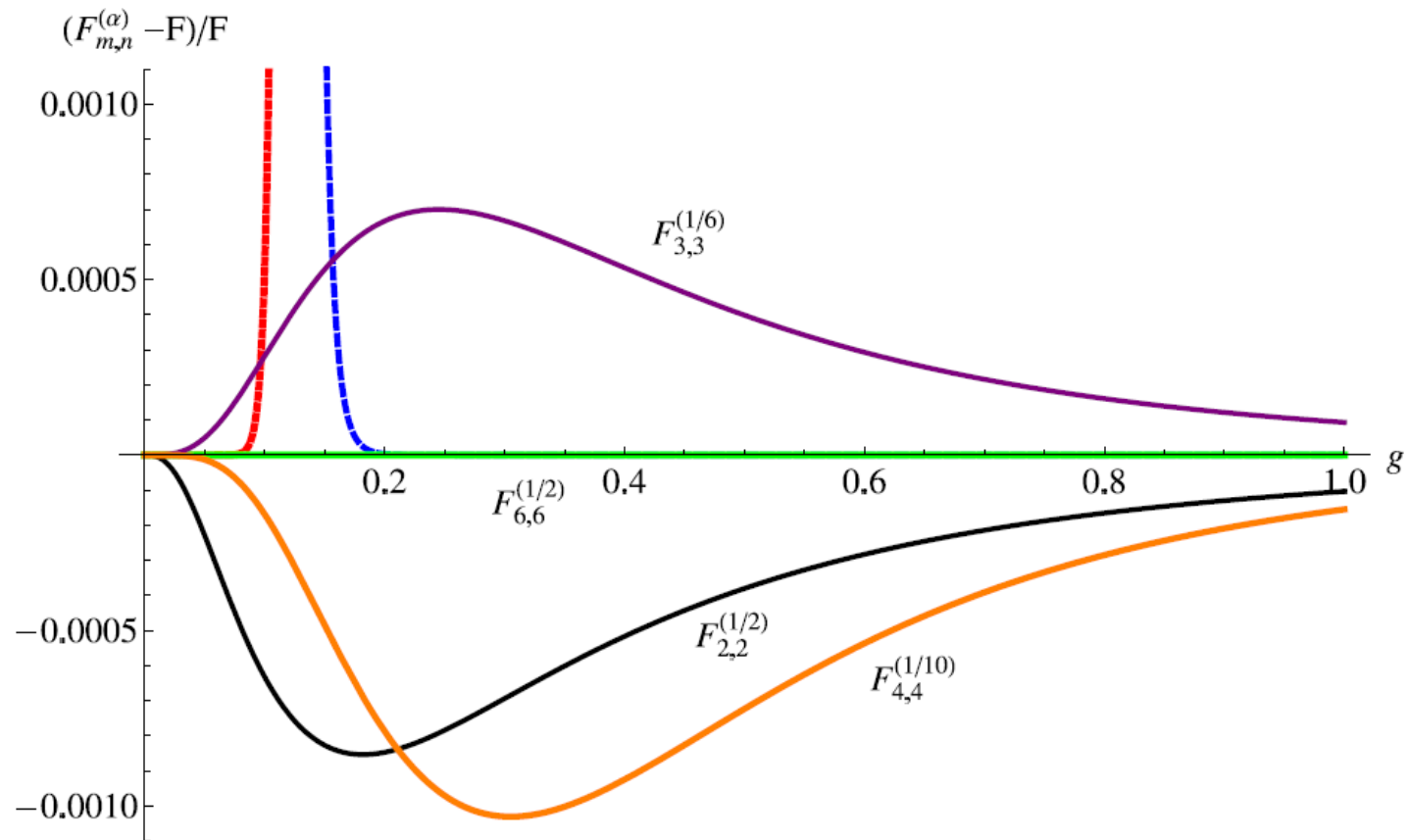
(although there are also many unsuccessful cases)

Ex.) Partition function of 0d ϕ^4 theory

[Sen'13, MH'14]

$$F(g) = \int_{-\infty}^{\infty} dx e^{-\frac{x^2}{2} - g^2 x^4} = \frac{e^{\frac{1}{32g^2}}}{2\sqrt{2}g} K_{\frac{1}{4}}\left(\frac{1}{32g^2}\right),$$

$$F(g) = F_{m,n}^{(\alpha)}(g) + \mathcal{O}(g^{a+m+1}, g^{b-n-1})$$

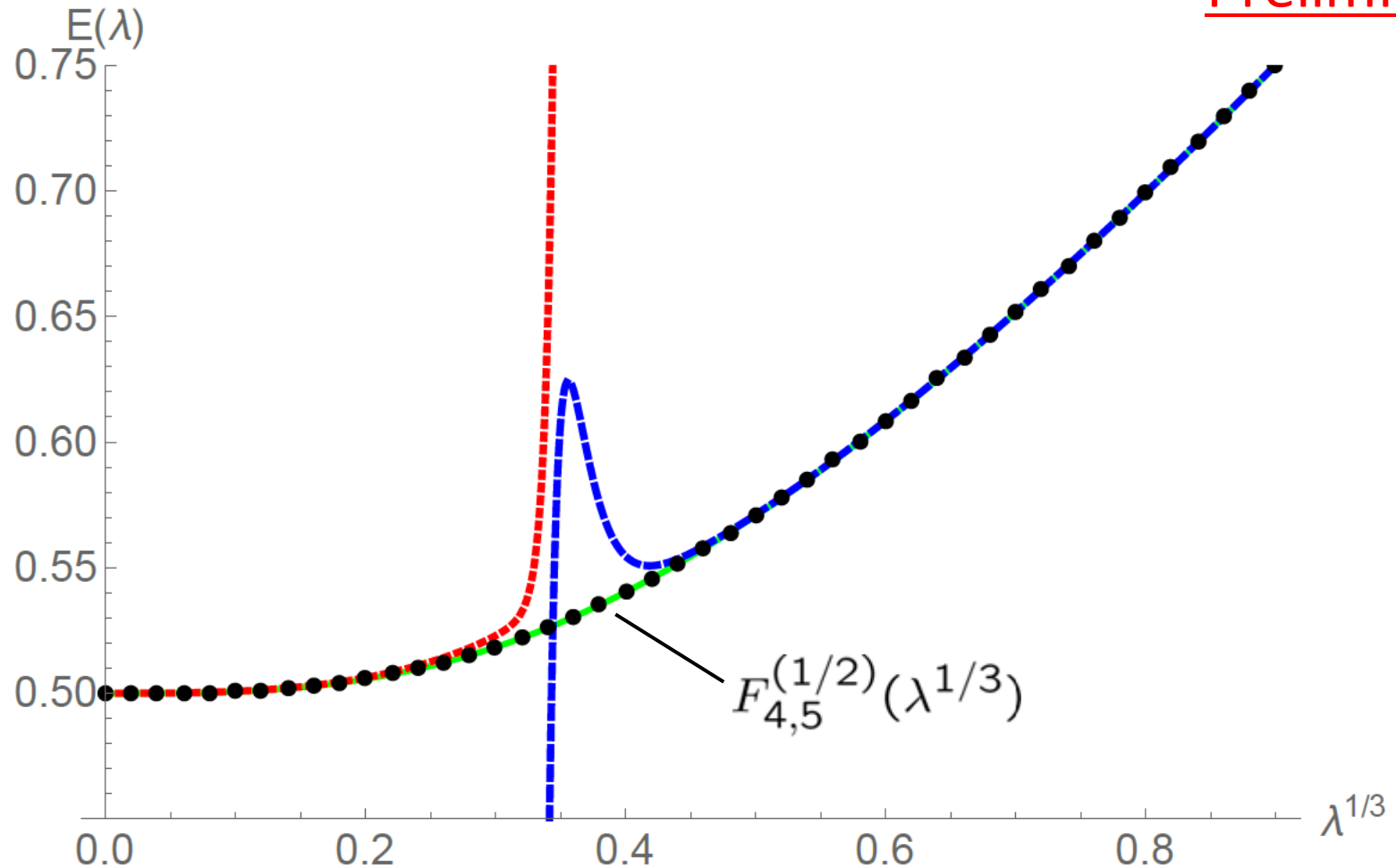


Ex.) Grand state Energy in anharmonic oscillator

[MH, work in progress]

$$\left[-\frac{d^2}{dx^2} + \frac{1}{4}x^2 + \frac{1}{4}\lambda x^4 \right] \psi(x) = E(\lambda)\psi(x)$$

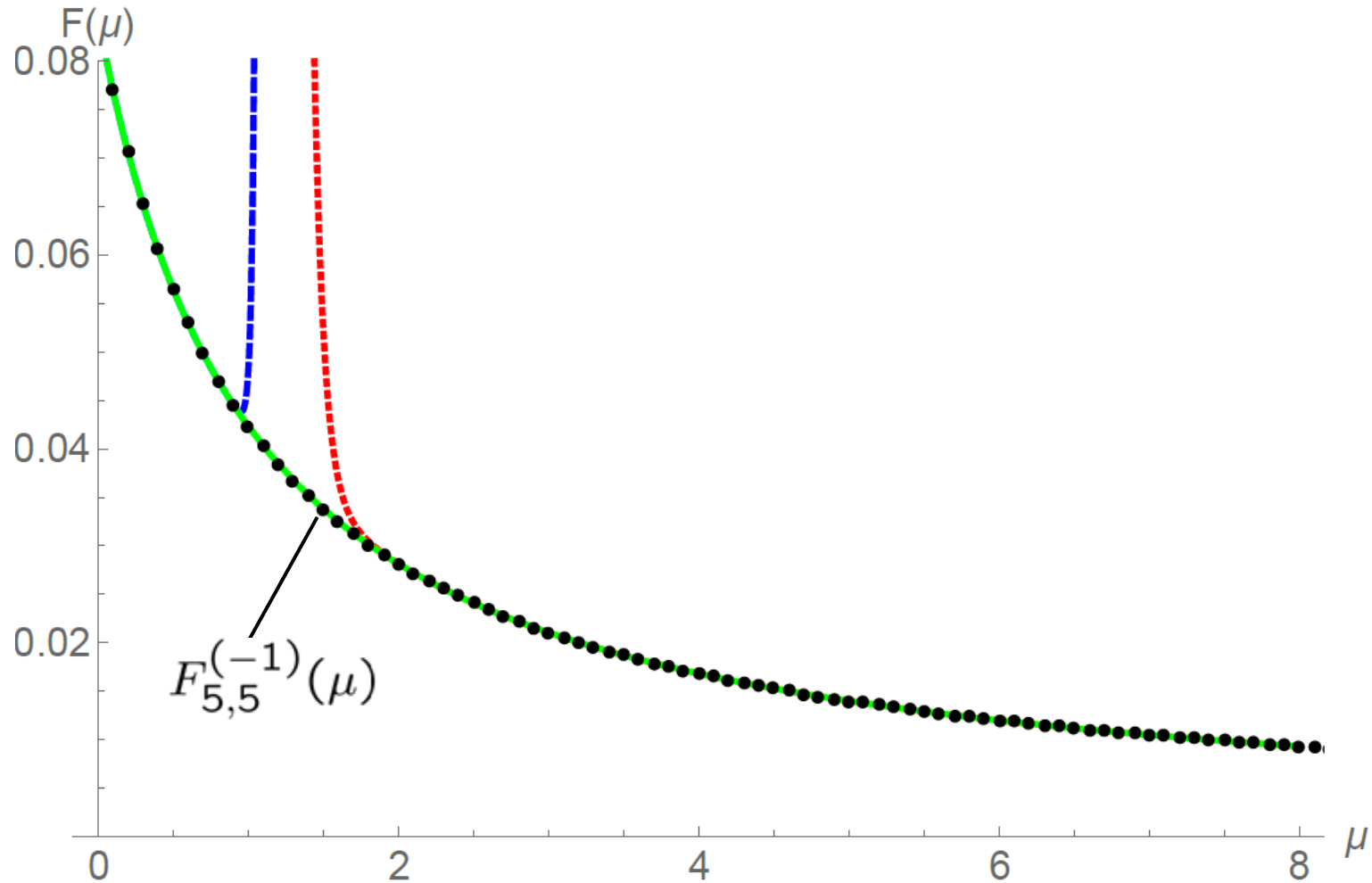
Preliminary



Ex.) Free energy of $c=1$ non-critical string

[MH'14]

μ : cosmological constant

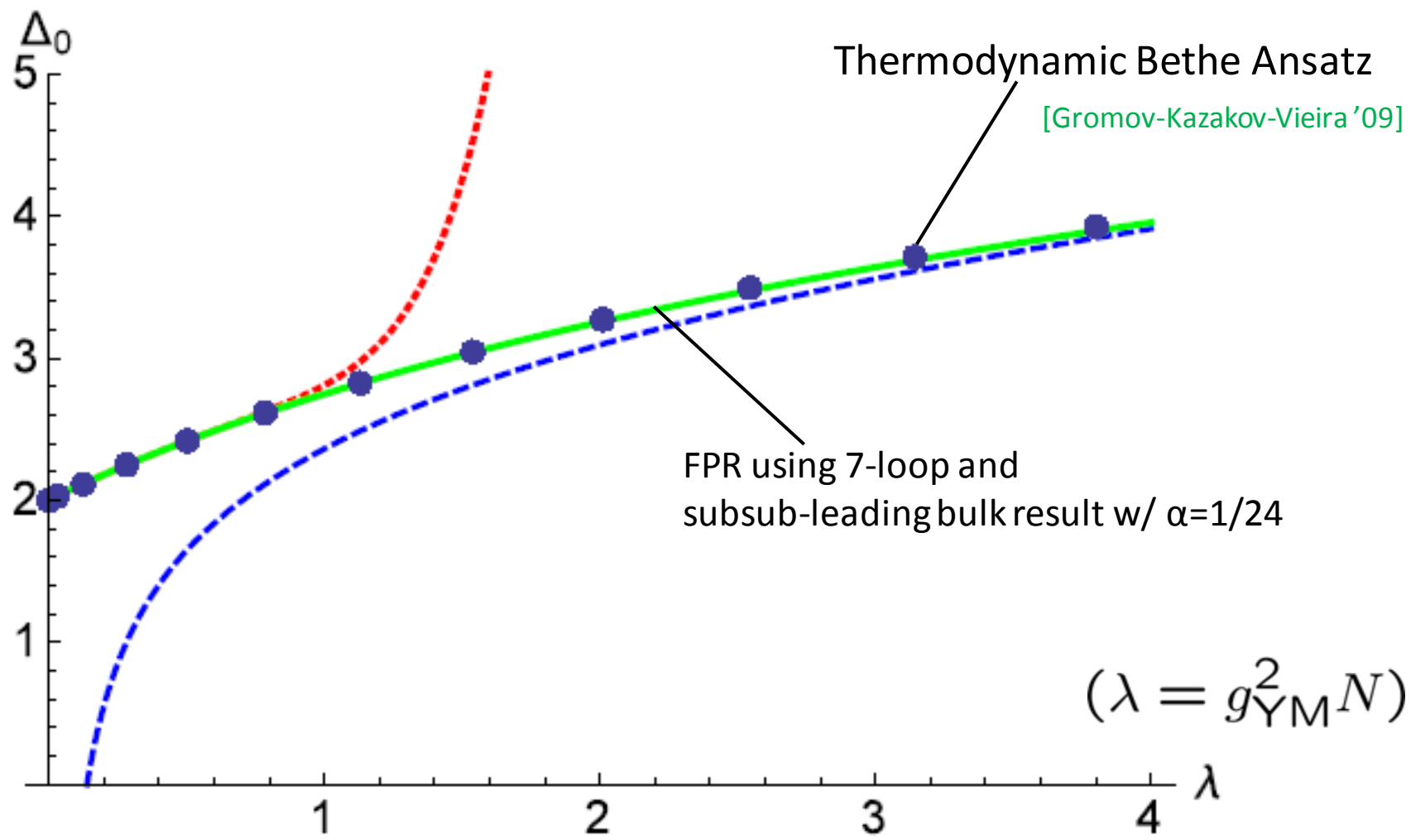


Ex.) Dimension of Konishi op. in planar $\mathcal{N} = 4$ SYM

[Chowdhury-MH-Thakur, to appear]

$$\mathcal{O}_{\text{Konishi}} = \text{Tr} \phi^I \phi^I$$

Preliminary



Analytic property of interpolating function & Twist operators in planar ABJM

[Chowdhury-MH, to appear]

Twist-operators in ABJM

ABJM theory:

[Aharony-Bergman-Jafferis-Maldacena '08]

$$3d \mathcal{N} = 6 \text{ U}(\mathbf{N})_k \times \text{ U}(\mathbf{N})_{-k}$$

(k: CS level)

superconformal Chern-Simons theory

$$\mathcal{O}_{L,S} = \text{Tr} \left[D_+^S (Y^1 Y_4^\dagger)^L \right] \quad Y^1, Y_4^\dagger : \text{(anti-)bi-fundamental scalar}$$

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The dimension of this operator is anomalous (unless $S=0$):

$$\Delta_{L,S}(k, N) = L + S + \gamma_{L,S}(k, N)$$

Here we focus on the **planar** limit:

$$\Delta_{L,S}(k, N) = \Delta_{L,S}^{(0)}(\lambda) + \mathcal{O}(N^{-2})$$

$$\lambda = \frac{N}{k}$$

Dressed coupling constant $h(\lambda)$

In the context of **integrability** analysis,
the dimension is described in terms of an **unknown function $h(\lambda)$** .

[Giombi-Gaiotto-Yin]

$h(\lambda) \propto$ (Central charge of $SU(2|2)$ sub-superconformal algebra)

$$h(\lambda) = \lambda + \mathcal{O}(\lambda^3) = \sqrt{\frac{\lambda}{2}} + \mathcal{O}(\lambda^0)$$

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Recently, **exact form of $h(\lambda)$ has been conjectured** as

[Gromov-Sizov]

$$\lambda = \frac{\sinh(2\pi h)}{2\pi} {}_3F_2\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}; 1, \frac{3}{2}; -\sinh^2(2\pi h)\right).$$

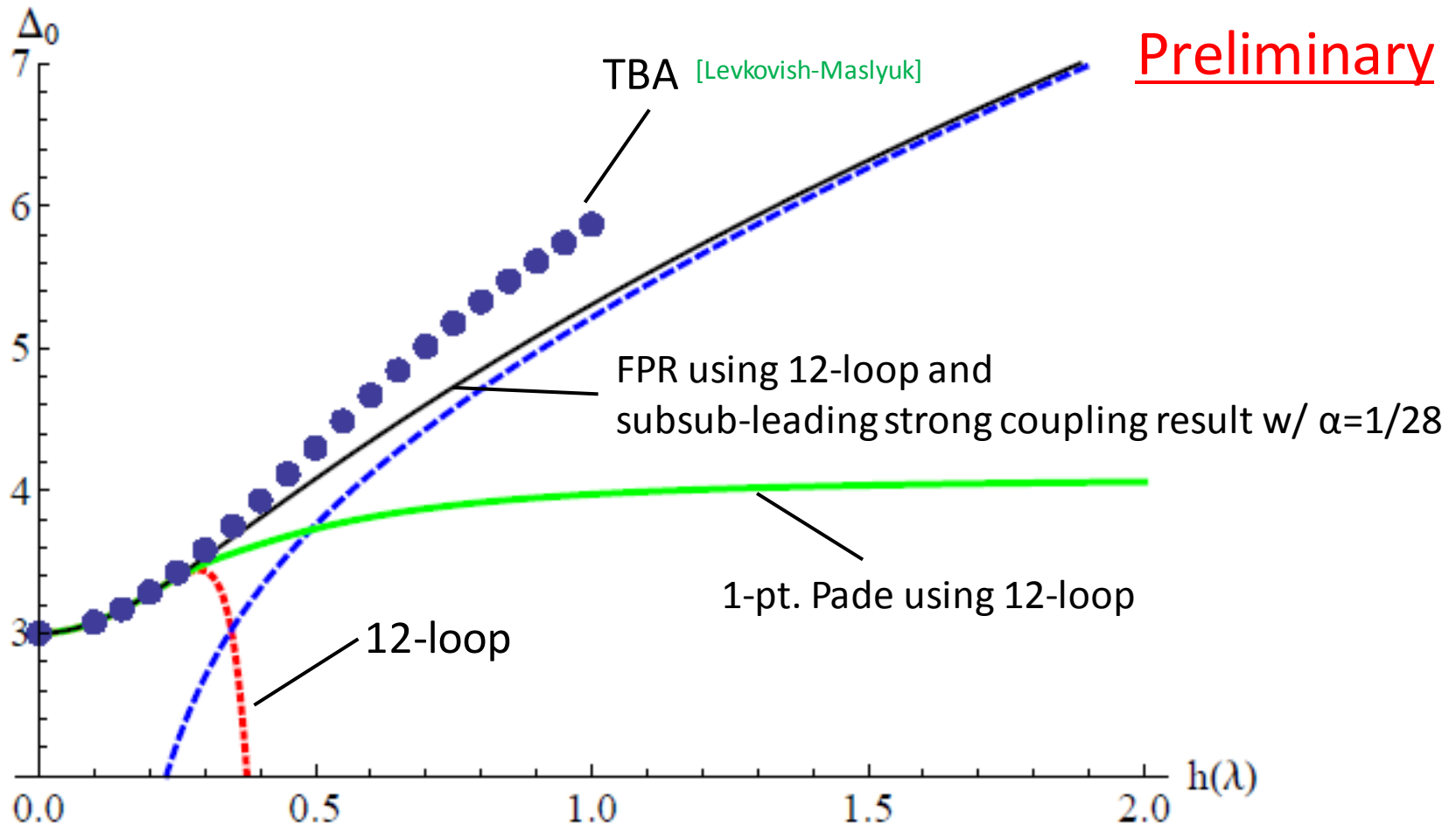
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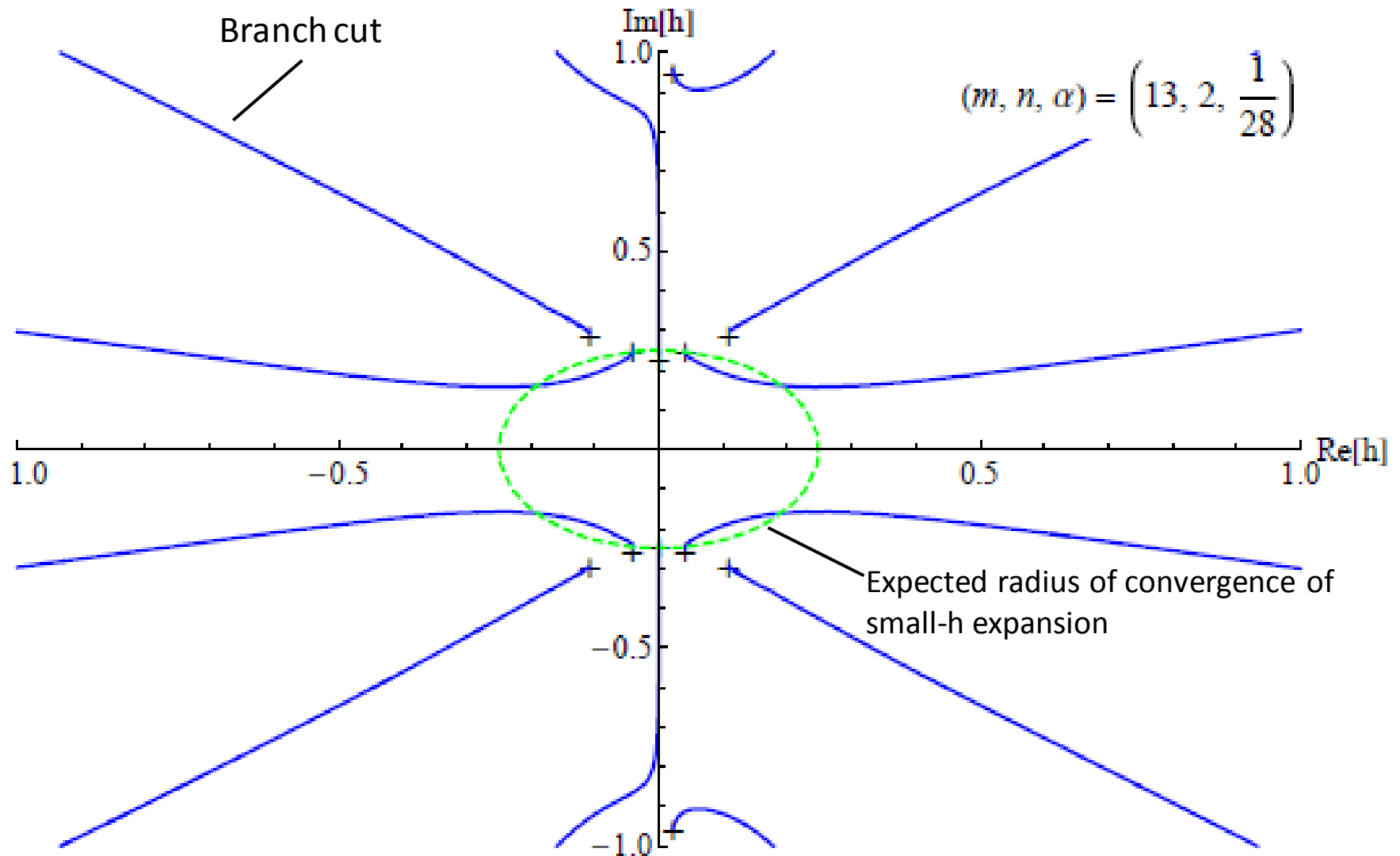


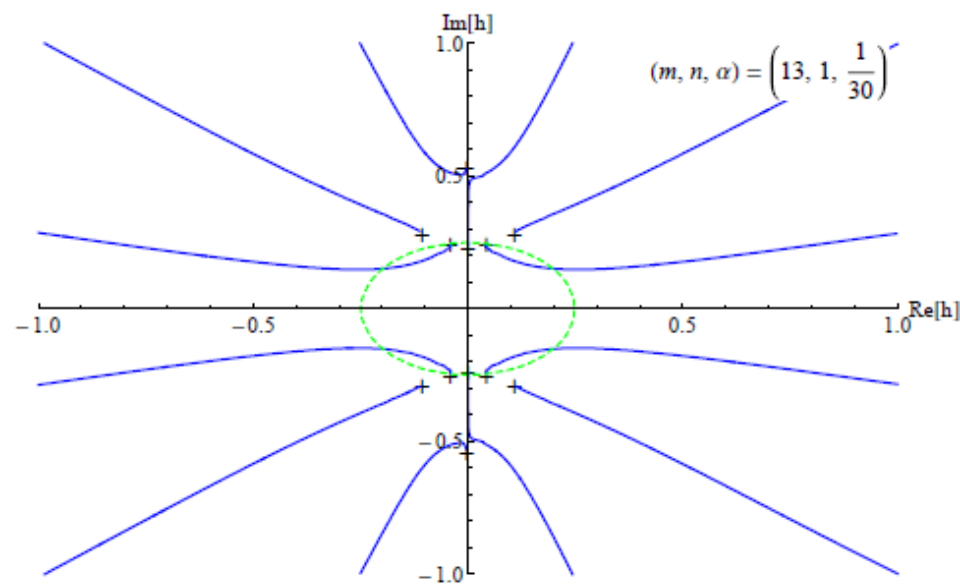
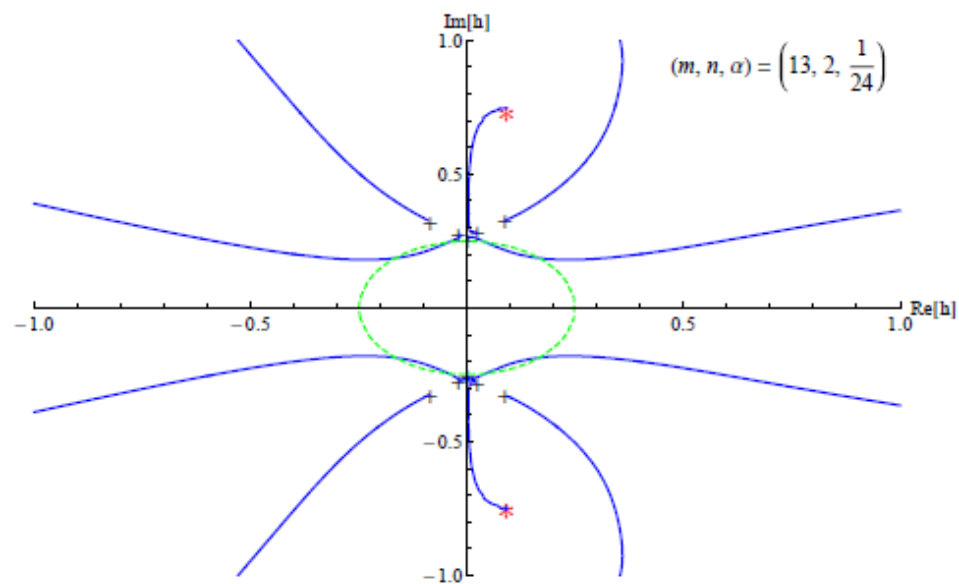
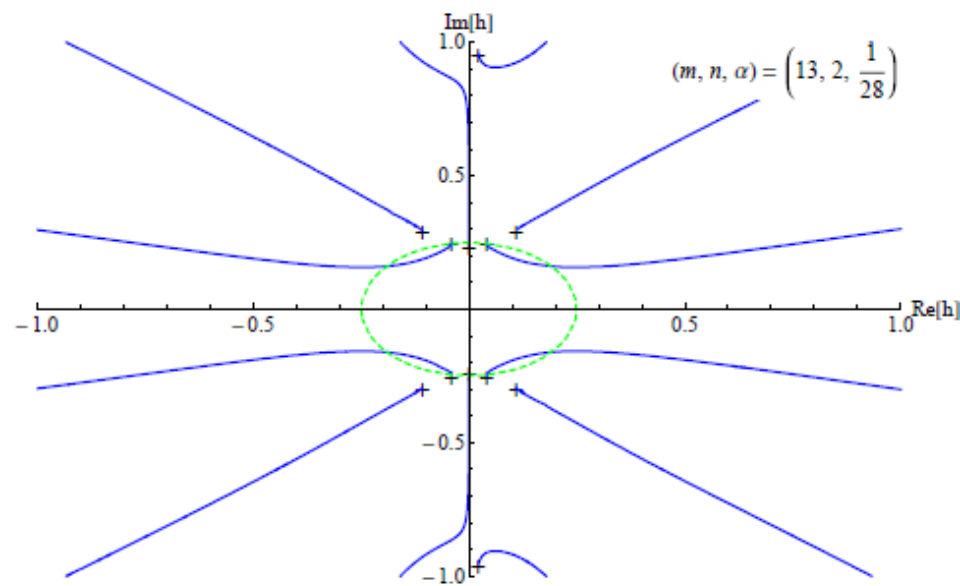
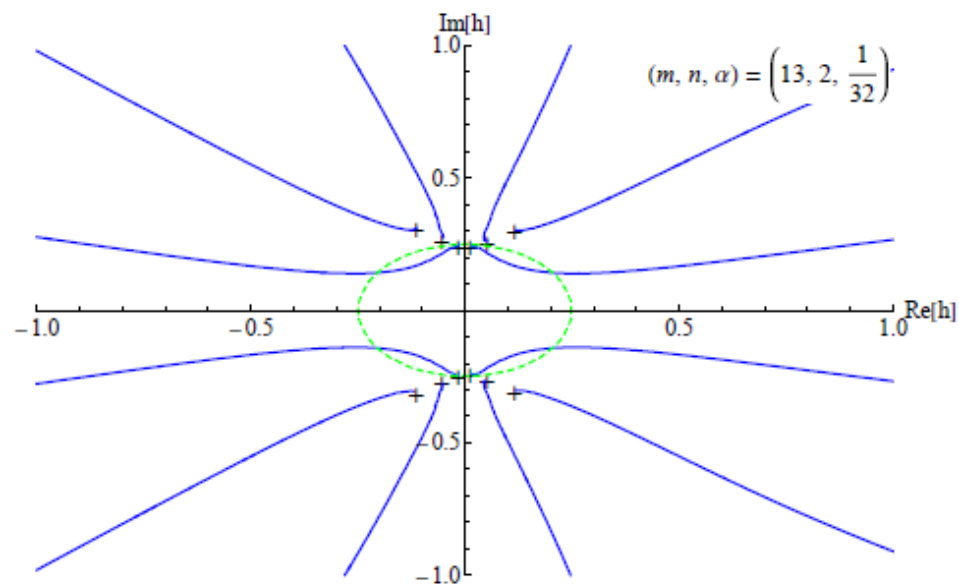
Preliminary

(Similar results hold for many other interpolating functions)

Analytic property of FPR for $(L,S)=(1,2)$

Preliminary





Many interpolating functions have singularities around $h = \pm i/4$!!

Similar results hold also for other (L, S) .

Physical Interpretation

Many interpolating functions have **singularity** around $h = \pm i/4$.

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If the conjecture $\lambda = \frac{\sinh(2\pi h)}{2\pi} {}_3F_2\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}; 1, \frac{3}{2}; -\sinh^2(2\pi h)\right)$ is correct,

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(where ABJM free energy behaves as the one of $c=1$ non-critical string.)

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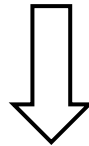
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Indirect evidence for the conjecture on $h(\lambda)$

Summary

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- We have introduced a class of interpolating functions (FPR)

$$F_{m,n}^{(\alpha)}(g) = s_0 g^a \left[\frac{1 + \sum_{k=1}^p c_k g^k}{1 + \sum_{k=1}^q d_k g^k} \right]^\alpha,$$

which includes Pade and FPP as the special cases

- **Analytic property** of FPR gives physical information on the dimensions of **twist operators in the planar ABJM**
 - indirect evidence for the recent conjecture on the dressed coupling constant $h(\lambda)$

Results on which I didn't talk (due to time)

- We can construct **many** interpolating functions
 - Which does give the best approximation?
 - **Criterion for the best approximation** [MH'14]
- Analytic property of interpolating function and **Stokes Phenomena** [MH-Jatkar '15]
- Interpolating function **in Borel plane?** [MH-Jatkar '15]
 - Naïve idea is failed.
- Comparison with **resurgence** approach [MH-Jatkar '15]

- **S-duality invariant** interpolating function for twist op. in N=4 SYM

$$F_m^{(s,\alpha)}(\tau) = \left[\frac{\sum_{k=1}^p c_k E_{s+k}(\tau)}{\sum_{k=1}^q d_k E_{s+k}(\tau)} \right]^\alpha$$

[Chowdhury-MH-Thakur, to appear]

[Generalization of
Beem- Rastelli-Sen-van Rees, Alday-Bissi]

- Compare with conformal bootstrap and draw conformal manifold
- Applications to 2d Ising model, lattice YM, W-loop in N=4 SYM etc...

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